

Model Independent bounds on NP in Tree Level (constraints from flavour observables)

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Based on:

*A. Lenz, GTX : 1912.07621 /
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Operator Basis

$$\hat{\mathcal{H}}_{eff}^{|\Delta B|=1} = \frac{G_F}{\sqrt{2}} \left\{ \sum_{p,p'=u,c} \lambda_{pp'}^{(q)} \sum_{i=1,2} C_i^{q,pp'}(\mu) \hat{Q}_i^{q,pp'} \right. \\ \left. + \sum_{p=u,c} \lambda_p^{(q)} \left[\sum_{i=3}^{10} C_i^q(\mu) \hat{Q}_i^q + C_{7\gamma}^q \hat{Q}_{7\gamma}^q + C_{8g}^q \hat{Q}_{8g}^q \right] \right\} + h. \quad \begin{aligned} \lambda_p^{(q)} &= V_{pb} V_{pq}^* , \\ \lambda_{pp'}^{(q)} &= V_{pb} V_{p'q}^* . \end{aligned}$$

$$\begin{aligned} \hat{Q}_1^{q,pp'} &= (\bar{\hat{p}}_\beta \hat{b}_\alpha)_{V-A} (\bar{\hat{q}}_\alpha \hat{p}'_\beta)_{V-A} , & \hat{Q}_2^{q,pp'} &= (\bar{\hat{p}} \hat{b})_{V-A} (\bar{\hat{q}} \hat{p}')_{V-A} , \\ \hat{Q}_3^q &= (\bar{\hat{q}} \hat{b})_{V-A} \sum_k (\bar{\hat{k}} \hat{k})_{V-A} , & \hat{Q}_4^q &= (\bar{\hat{q}}_\alpha \hat{b}_\beta)_{V-A} \sum_k (\bar{\hat{k}}_\beta \hat{k}_\alpha)_{V-A} , \\ \hat{Q}_5^q &= (\bar{\hat{q}} \hat{b})_{V-A} \sum_k (\bar{\hat{k}} \hat{k})_{V+A} , & \hat{Q}_6^q &= (\bar{\hat{q}}_\alpha \hat{b}_\beta)_{V-A} \sum_k (\bar{\hat{k}}_\beta \hat{k}_\alpha)_{V+A} , \\ \hat{Q}_7^q &= (\bar{\hat{q}} \hat{b})_{V-A} \sum_k \frac{3}{2} e_k (\bar{\hat{k}} \hat{k})_{V+A} , & \hat{Q}_8^q &= (\bar{\hat{q}}_\alpha \hat{b}_\beta)_{V-A} \sum_k \frac{3}{2} e_k (\bar{\hat{k}}_\beta \hat{k}_\alpha)_{V+A} , \\ \hat{Q}_9^q &= (\bar{\hat{q}} \hat{b})_{V-A} \sum_k \frac{3}{2} e_k (\bar{\hat{k}} \hat{k})_{V-A} , & \hat{Q}_{10}^q &= (\bar{\hat{q}}_\alpha \hat{b}_\beta)_{V-A} \sum_k \frac{3}{2} e_k (\bar{\hat{k}}_\beta \hat{k}_\alpha)_{V-A} , \\ \hat{Q}_{7\gamma}^q &= \frac{e}{8\pi^2} m_b \bar{\hat{q}} \sigma_{\mu\nu} (1 + \gamma_5) \hat{F}^{\mu\nu} \hat{b} , & \hat{Q}_{8g}^q &= \frac{g_s}{8\pi^2} m_b \bar{\hat{q}} \sigma_{\mu\nu} (1 + \gamma_5) \hat{G}^{\mu\nu} \hat{b} . \end{aligned}$$

New Physics at Tree Level

Strategy:

$$C_1(M_W) := C_1^{SM}(M_W) + \Delta C_1(M_W),$$
$$C_2(M_W) := C_2^{SM}(M_W) + \Delta C_2(M_W),$$

$$\vec{C}(\mu) = U(\mu, M_W, \alpha) \vec{C}(M_W),$$

$$\mu = m_b$$

Evolution matrix

$$U(\mu, M_W, \alpha) = \left[U_0 + \frac{\alpha_s(\mu)}{4\pi} \mathbf{J} U_0 - \frac{\alpha_s(M_W)}{4\pi} U_0 \mathbf{J} + \frac{\alpha}{4\pi} \left(\frac{4\pi}{\alpha_s(\mu)} \mathbf{R}_0 + \mathbf{R}_1 \right) \right],$$

$$C_1^{SM}(m_b) \simeq -0.30$$

$$C_2^{SM}(m_b) \simeq 1.0$$

Chi² Fit

$$\chi^2(\vec{\omega}) = \sum_i \left(\frac{\tilde{O}_{i,\text{exp}} - \tilde{O}_{i,\text{theo}}(\vec{\omega})}{\sigma_{i,\text{exp}}} \right)^2,$$

parameters to be fitted

$$\vec{\omega} = \left(\Delta C_1(M_W), \Delta C_2(M_W), \vec{\lambda} \right).$$

nuisance parameters

decay constants, form factors, masses, other QCD quantities, etc.

Use a likelihood ratio test.

Observables

$$b \rightarrow u \bar{u} d$$

$$R_{\pi\pi} = \frac{\Gamma(B^+ \rightarrow \pi^+ \pi^0)}{d\Gamma(\bar{B}_d^0 \rightarrow \pi^+ \ell^- \bar{\nu}_\ell)/dq^2|_{q^2=0}}$$

$$\begin{array}{l} \bar{B}_d^0 \rightarrow \pi^+ \pi^- \\ B_d^0 \rightarrow \pi^+ \pi^- \end{array} \Rightarrow S_{\pi\pi} = \frac{2 \operatorname{Im}(\lambda_{\pi\pi}^d)}{1 + |\lambda_{\pi\pi}^d|^2}, \quad \lambda_{\pi\pi}^d = \left[\frac{V_{td} V_{tb}^*}{|V_{td} V_{tb}^*|} \right]^2 \frac{\bar{\mathcal{A}}_{\pi^+ \pi^-}}{\mathcal{A}_{\pi^+ \pi^-}}.$$

$$\begin{array}{l} \bar{B}_d^0 \rightarrow \pi^+ \rho^- \\ \bar{B}_d^0 \rightarrow \rho^+ \pi^- \end{array} \Rightarrow \tilde{S}_{\pi\rho} = \frac{2 \operatorname{Im}(\lambda_{\pi\rho}^d)}{1 + |\lambda_{\pi\rho}^d|^2}, \quad \tilde{S}_{\rho\pi} = \frac{2 \operatorname{Im}(\lambda_{\rho\pi}^d)}{1 + |\lambda_{\rho\pi}^d|^2},$$

$$R_{\rho\rho} = \frac{\mathcal{B}r(B^- \rightarrow \rho_L^- \rho_L^0)}{\mathcal{B}r(\bar{B}_d^0 \rightarrow \rho_L^+ \rho_L^-)} = \frac{|\mathcal{A}_{\rho^- \rho^0}|^2}{|\mathcal{A}_{\rho^+ \rho^-}|^2},$$

Observables

$$b \rightarrow c \bar{u} d$$

$$R_{D^*\pi} = \frac{\Gamma(\bar{B}^0 \rightarrow D^{*+} \pi^-)}{d\Gamma(\bar{B}^0 \rightarrow D^{*+} l^- \bar{\nu}_l) / dq^2 |_{q^2=m_\pi^2}}$$

$$b \rightarrow c \bar{c} s$$

$$\mathcal{B}r(\bar{B} \rightarrow X_s \gamma) \quad S_{J/\psi\phi} = \frac{2 \operatorname{Im}(\lambda_{J/\psi\phi}^s)}{1 + |\lambda_{J/\psi\phi}^s|^2} :$$

$$\Delta\Gamma_s$$

$$\frac{\tau_{B_s}}{\tau_{B_d}}$$

Observables

$$b \rightarrow c \bar{c} d$$

$$M_{12}^d \quad \longrightarrow \quad \sin(2\beta_d) \quad \mathcal{B}_r (\bar{B} \rightarrow X_d \gamma)$$

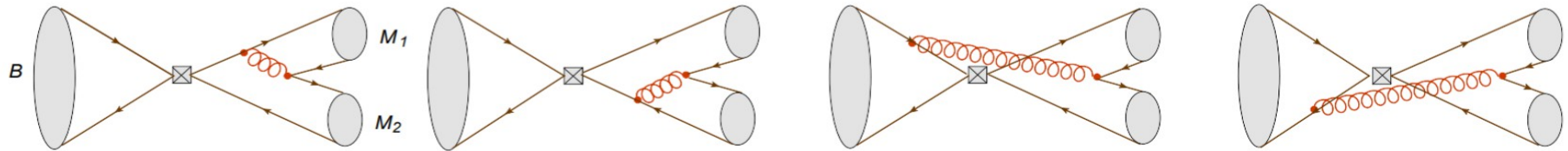
Multi-channel Observables

$$a_{sl}^s \quad a_{sl}^d \quad \Delta\Gamma_s$$

Useful for obtaining extra constraints in each transition channel and to determine maximal constraints

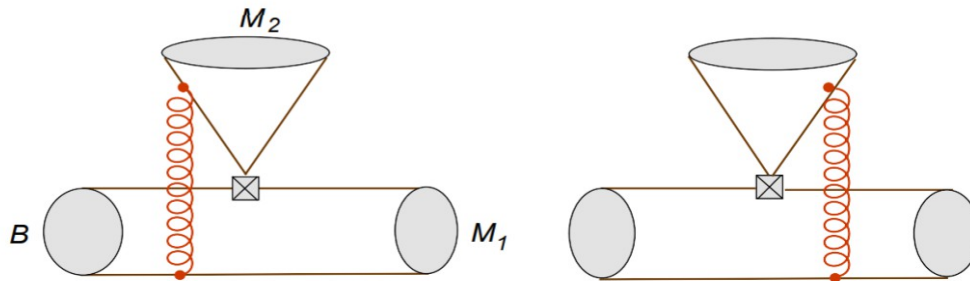
Important uncertainties in QCD-factorization

Annihilation: end-point singularities



$$X_A = \left(1 + \rho_A e^{i\phi_A}\right) \ln \frac{m_B}{\Lambda_h} \quad 0 < \rho_A < 2, \quad 0 < \phi_A < 2\pi$$

Spectator-scattering: end-point singularities



$$X_H = \left(1 + \rho_H e^{i\phi_H}\right) \ln \frac{m_B}{\Lambda_h} \quad 0 < \rho_H < 2, \quad 0 < \phi_H < 2\pi$$

Important uncertainties in QCD-factorization

$$\begin{aligned}
 \bar{A}_{\pi^+\rho^-} = & A_{\pi\rho} \left(\lambda_u^{(d)} \alpha_2^{\pi\rho} + \sum_{p=u,c} \lambda_p^{(d)} \left[\tilde{\alpha}_4^{p,\pi\rho} + \tilde{\alpha}_{4,EW}^{p,\pi\rho} \right. \right. \\
 & \left. \left. + \beta_3^{p,\pi\rho} + \beta_4^{p,\pi\rho} - \frac{1}{2} \beta_{3,EW}^{p,\pi\rho} - \frac{1}{2} \beta_{4,EW}^{p,\pi\rho} \right] \right) \\
 & + A_{\rho\pi} \left(\lambda_u^{(d)} \beta_1^{\rho\pi} + \sum_{p=u,c} \lambda_p^{(d)} \left[\beta_4^{p,\rho\pi} + \beta_{4,EW}^{p,\rho\pi} \right] \right),
 \end{aligned}$$

Individual annihilation topological amplitudes are of O(5%)

The total effect in the uncertainty of the observables can be Important if the Amplitudes depend on several of them.

Example: $R_{\rho\rho}$

$$R_{\rho\rho} = \frac{\mathcal{B}r(B^- \rightarrow \rho_L^- \rho_L^0)}{\mathcal{B}r(\bar{B}_d^0 \rightarrow \rho_L^+ \rho_L^-)}$$

$$\mathcal{A}_{\rho^- \rho^0} = \frac{A_{\rho\rho}}{\sqrt{2}} \left[\lambda_u^{(d)} (\alpha_1^{\rho\rho} + \alpha_2^{\rho\rho}) + \frac{3}{2} \sum_{p=u,c} \lambda_p^{(d)} (\alpha_7^{p,\rho\rho} + \alpha_9^{p,\rho\rho} + \alpha_{10}^{p,\rho\rho}) \right],$$

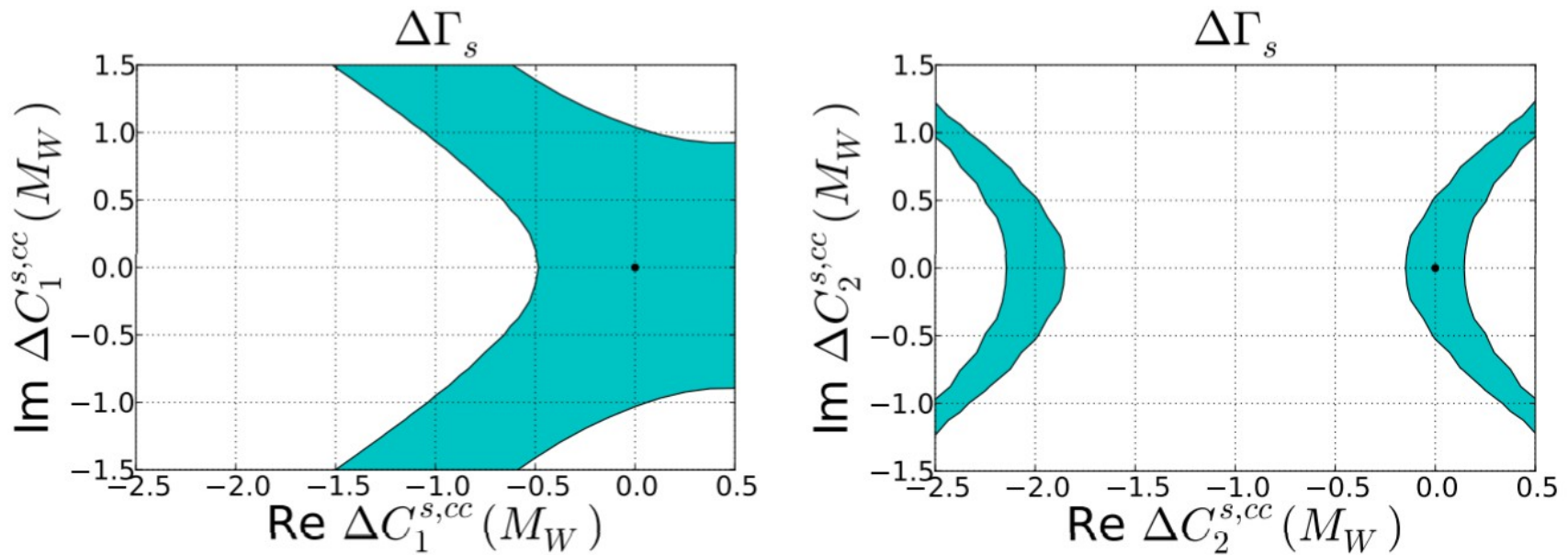
$$\begin{aligned} \mathcal{A}_{\rho^+ \rho^-} = & A_{\rho\rho} \left[\lambda_u^{(d)} (\alpha_2^{\rho\rho} + \beta_2^{\rho\rho}) + \sum_{p=u,c} \lambda_p^{(d)} (\alpha_4^{p,\rho\rho} + \alpha_{10}^{p,\rho\rho} \right. \\ & \left. + \beta_3^{p,\rho\rho} + 2\beta_4^{p,\rho\rho} - \frac{1}{2}\beta_{3,EW}^{p,\rho\rho} + \frac{1}{2}\beta_{4,EW}^{p,\rho\rho}) \right]. \end{aligned}$$

| Parameter | Relative Error |
|----------------------------|----------------|
| $\bar{\chi}_A$ | 26.40% |
| X_H | 23.33% |
| λ_B | 12.32% |
| μ | 6.78% |
| $A_0^{B \rightarrow \rho}$ | 2.54% |
| a_2^{ρ} | 2.24% |
| f_{ρ} | 0.46% |
| Λ_5^{QCD} | 0.45% |
| γ | 0.38% |
| m_b | 0.27% |
| f_B | 0.15% |
| f_{ρ}^{\perp} | 0.15% |
| m_c | 0.12% |
| Total | 38.09% |

$$R_{\rho\rho}^{\text{SM}} = (67.5 \pm 25.7) \cdot 10^{-2}$$

$$R_{\rho\rho}^{\text{Exp}} = (83.14 \pm 8.98) \cdot 10^{-2}$$

Determination of the allowed NP regions per observable

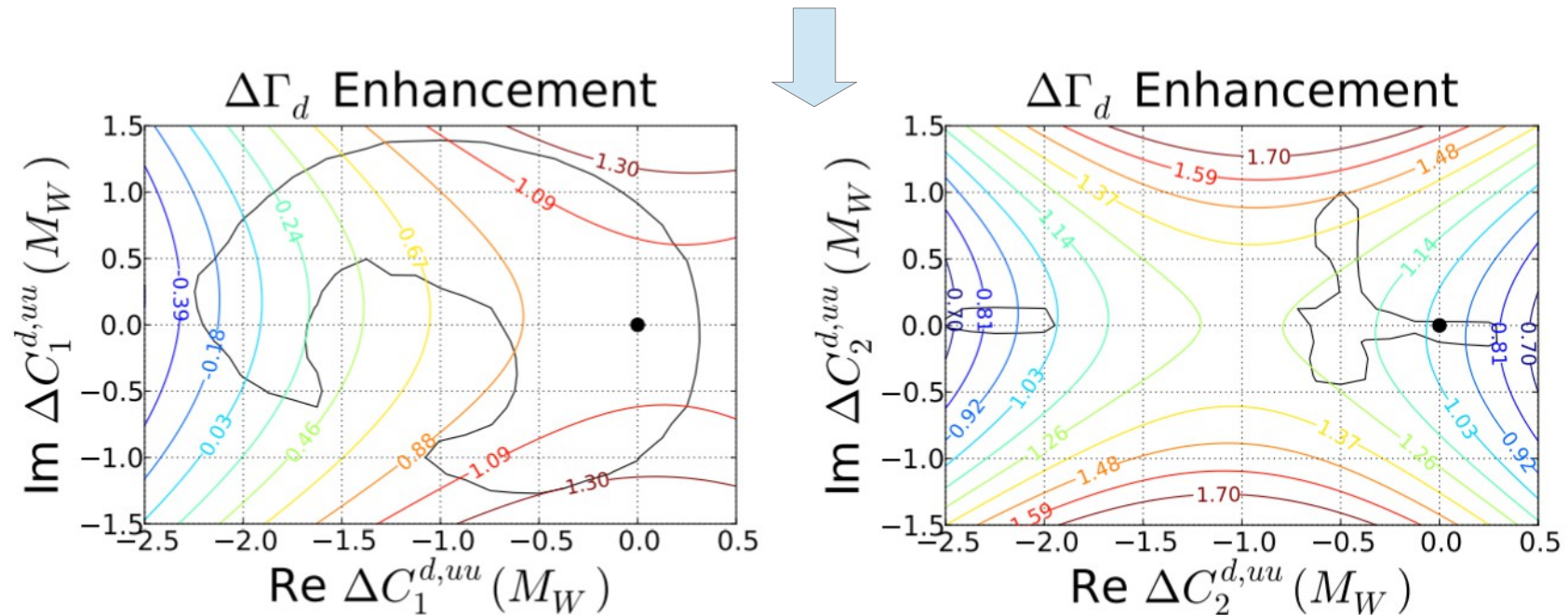


Turn on one coefficient at a time.

*Assume complex NP contribution
to the tree-level Wilson coefficients.*

Perform the fit per b quark level transition

$b \rightarrow u \bar{u} d$



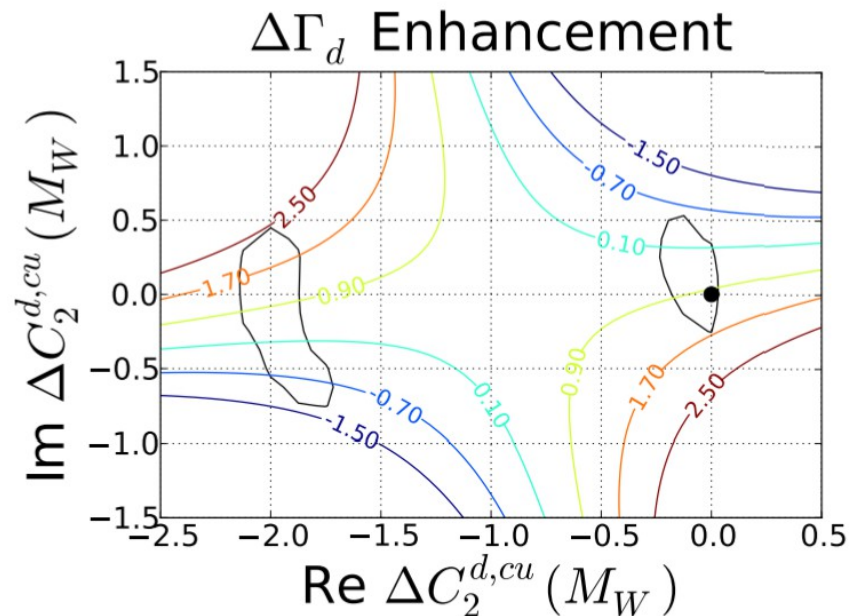
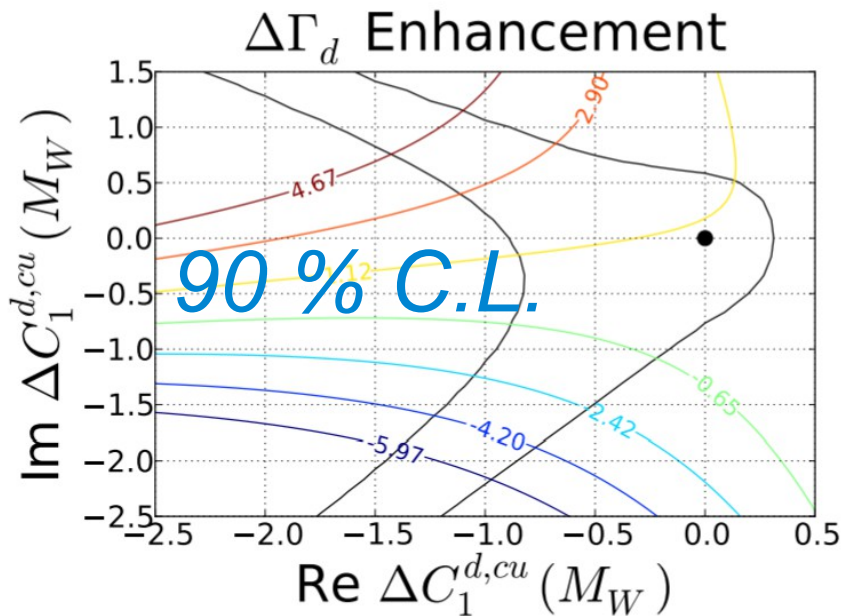
In general we find variations of $O(1)$ in the NP contributions per quark level transition

Implications in mixing

$$\Delta\Gamma_d^{\text{Exp}} = \left(-1.3 \pm 6.6\right) \cdot 10^{-3} \text{ ps}^{-1} \quad \Delta\Gamma_d^{\text{SM}} = \left(2.6 \pm 0.4\right) \cdot 10^{-3} \text{ ps}^{-1}$$

HFLAV: 1612.07233

NP in Tree-Level can lead to enhancements in $\Delta\Gamma_d$



for $\Delta C_1^{d,cu}(M_W)$: $-5.97 < \Delta\Gamma_d/\Delta\Gamma_d^{\text{SM}} < 4.67$, \longrightarrow

for $\Delta C_2^{d,cu}(M_W)$: $-1.5 < \Delta\Gamma_d/\Delta\Gamma_d^{\text{SM}} < 2.50$.

*Saturates bounds ,
Use it to constrain NP in Tree*

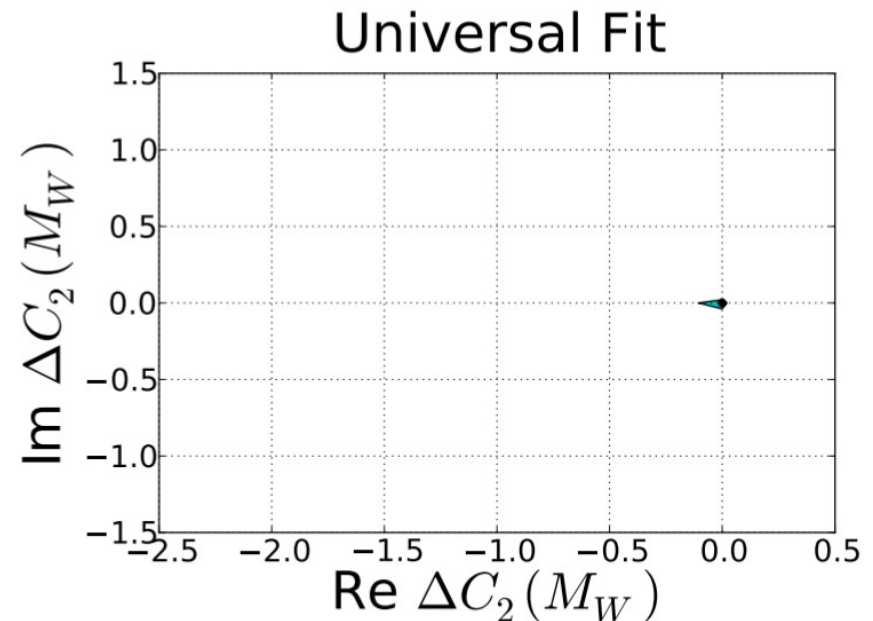
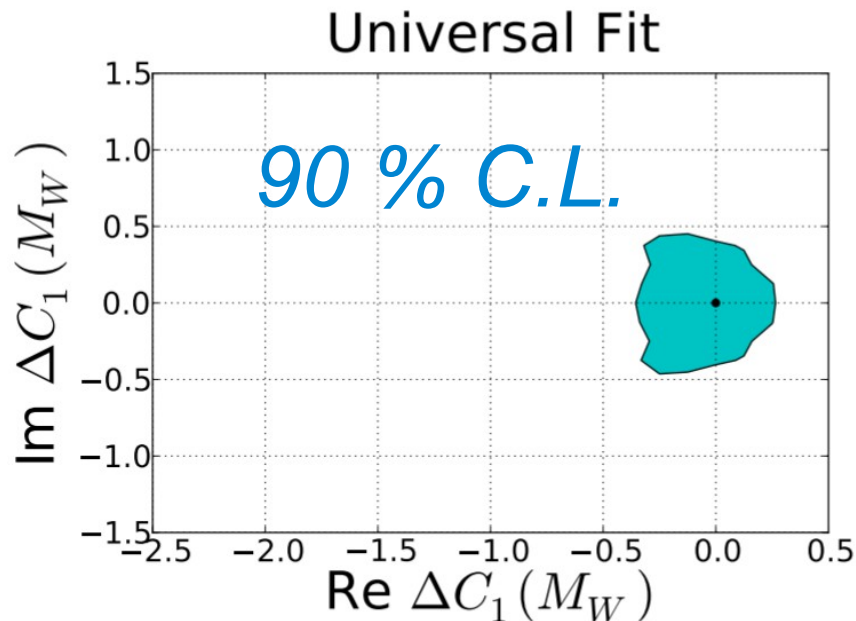
Universal Fit

Obtain maximal NP constraints by taking

$$\Delta C_1^{s,ab}(M_W) = \Delta C_1^{d,ab}(M_W) = \Delta C_1(M_W)$$

$$\Delta C_2^{s,ab}(M_W) = \Delta C_2^{d,ab}(M_W) = \Delta C_2(M_W)$$

for $a = u, c$ and $b = u, c$



Universal Fit

90 % C.L. Regions

$$\begin{aligned} \operatorname{Re} \left[\Delta C_1(M_W) \right] \Big|_{\min} &= -0.36, & \operatorname{Im} \left[\Delta C_1(M_W) \right] \Big|_{\min} &= -0.47, \\ \operatorname{Re} \left[\Delta C_1(M_W) \right] \Big|_{\max} &= 0.26, & \operatorname{Im} \left[\Delta C_1(M_W) \right] \Big|_{\max} &= 0.45, \\ \\ \operatorname{Re} \left[\Delta C_2(M_W) \right] \Big|_{\min} &= -0.11, & \operatorname{Im} \left[\Delta C_2(M_W) \right] \Big|_{\min} &= -0.04, \\ \operatorname{Re} \left[\Delta C_2(M_W) \right] \Big|_{\max} &= 0.02, & \operatorname{Im} \left[\Delta C_2(M_W) \right] \Big|_{\max} &= 0.02. \end{aligned}$$

Impact on the CKM angle γ

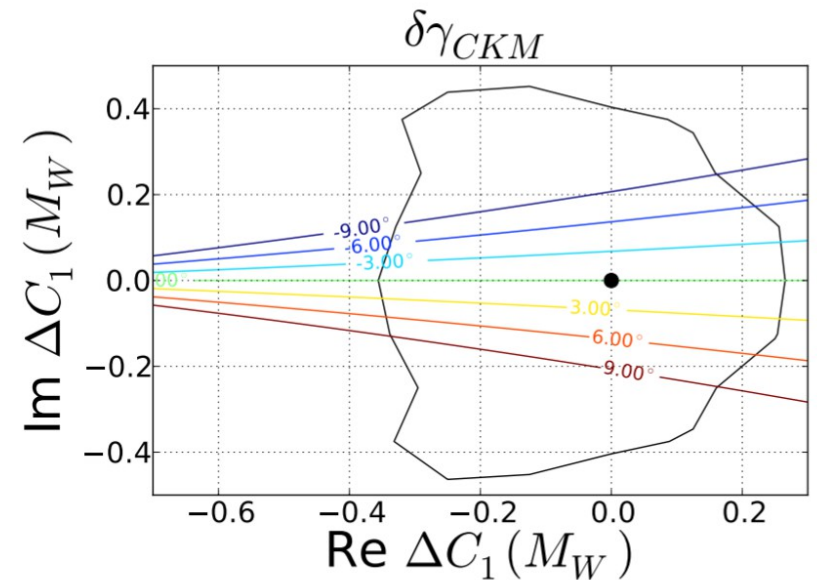
Sensitivity to NP in Tree-Level

$$\delta\gamma = (r_A - r_{A'}) \frac{\text{Im} [\Delta C_1]}{C_2}$$

Large uncertainties in the ratios

$$r_{A'} = \frac{\langle \bar{D}^0 K^- | Q_1^{\bar{u}cs} | B^- \rangle}{\langle \bar{D}^0 K^- | Q_2^{\bar{u}cs} | B^- \rangle},$$

$$r_A = \frac{\langle D^0 K^- | Q_1^{\bar{c}us} | B^- \rangle}{\langle D^0 K^- | Q_2^{\bar{c}us} | B^- \rangle}$$



Reduce $\text{Im} [\Delta C_1]$ to minimize potential effects in γ

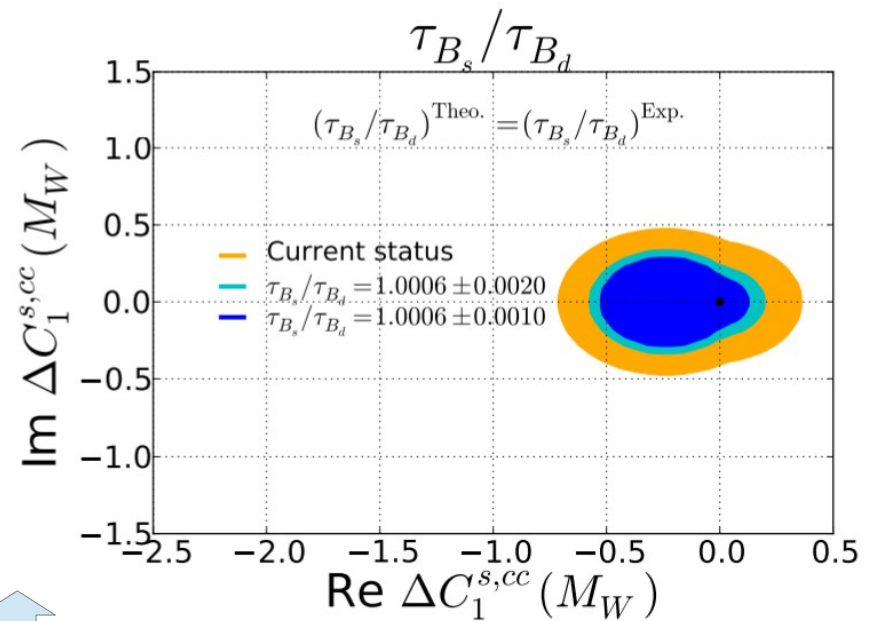
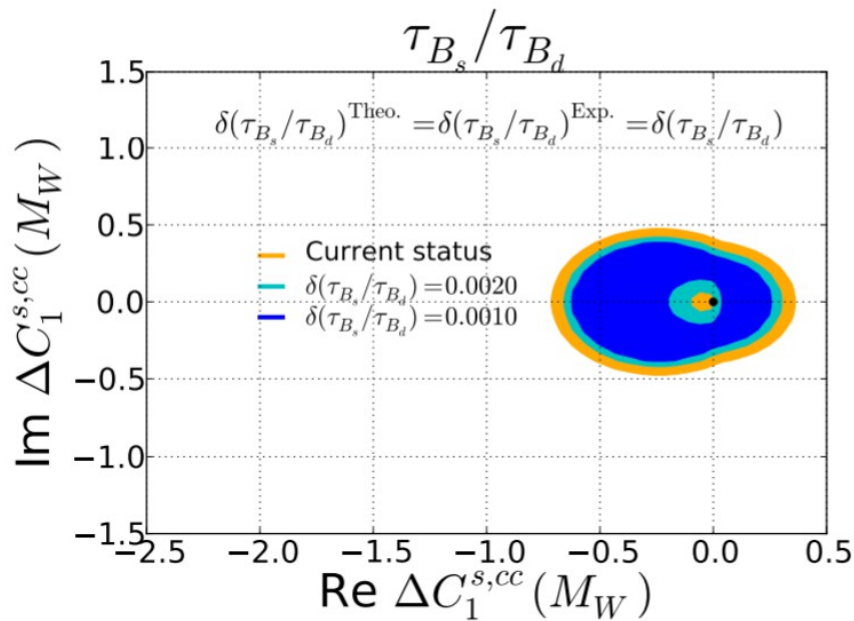
Effect of the Life-time ratio

$$\left(\frac{\tau_{B_s}}{\tau_{B_d}}\right)^{\text{Exp}} = 0.994 \pm 0.004$$

HFLAV: 1612.07233

$$\left(\frac{\tau_{B_s}}{\tau_{B_d}}\right)^{\text{SM}} = 1.0006 \pm 0.0020$$

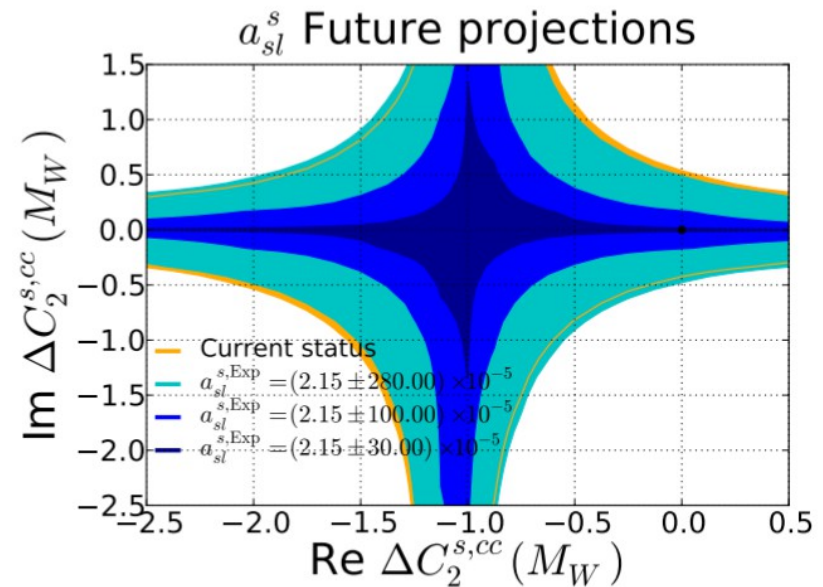
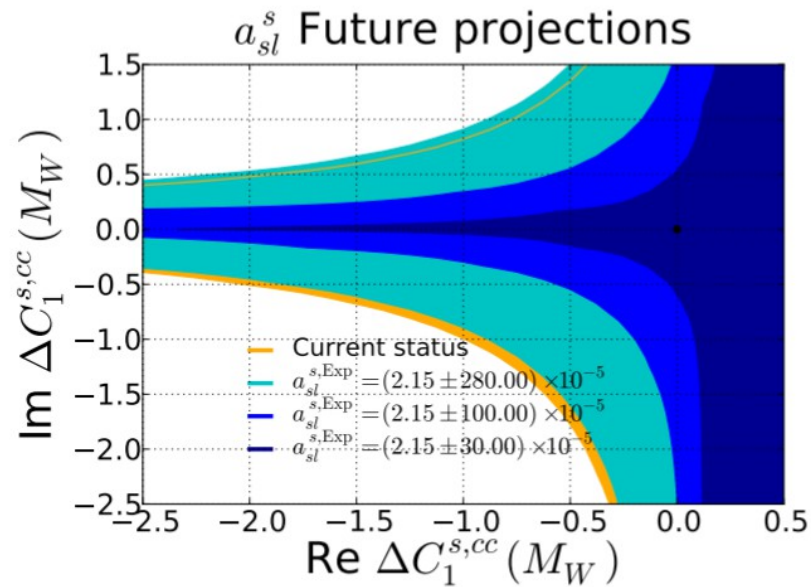
Kirk, Lenz and Rauh: 1711.02100



Impact of the uncertainty

To be updated soon....

Semileptonic Asymmetries



$$a_{sl}^{s,SM} = (2.06 \pm 0.18) \cdot 10^{-5}$$

$$a_{sl}^{s,Exp} = (60 \pm 280) \cdot 10^{-5} \quad \text{HFLAV: 1612.07233}$$

projections for the uncertainty

$$\delta(a_{sl}^s) = 1 \cdot 10^{-3} \quad \text{LHCb 2025}$$

$$\delta(a_{sl}^s) = 3 \cdot 10^{-4} \quad \text{Upgrade II}$$

Closing Remarks

- *Deviations with respect to the SM on the tree-level Wilson coefficients C_1 and C_2 are allowed by current theoretical and experimental results in *flavour observables*.*
- The deviations on C_1 can be as sizeable as 40%.
- Potential enhancements in the observable $\Delta\Gamma_d$
- Potential sizeable effects on the uncertainty of CKM γ .
- Updates on *mixing observables and the neutral B meson lifetime ratio* play a central role in reducing the size of the deviations on C_1 and C_2 .