

Precision Higgs Physics

CRC Annual Meeting 2021

Robert Harlander, 28 May 2021

A MODEL OF LEPTONS*

Steven Weinberg†

Laboratory for Nuclear Science and Physics Department,
Massachusetts Institute of Technology, Cambridge, Massachusetts

(Received 17 October 1967)

Nobel 1979

Therefore, we shall construct our Lagrangian out of L and R , plus gauge fields \vec{A}_μ and B_μ coupled to \vec{T} and Y , plus a spin-zero doublet

$$\varphi = \begin{pmatrix} \varphi^0 \\ \varphi^- \end{pmatrix} \quad (3)$$

whose vacuum expectation value will break \vec{T} and Y and give the electron its mass. The on-

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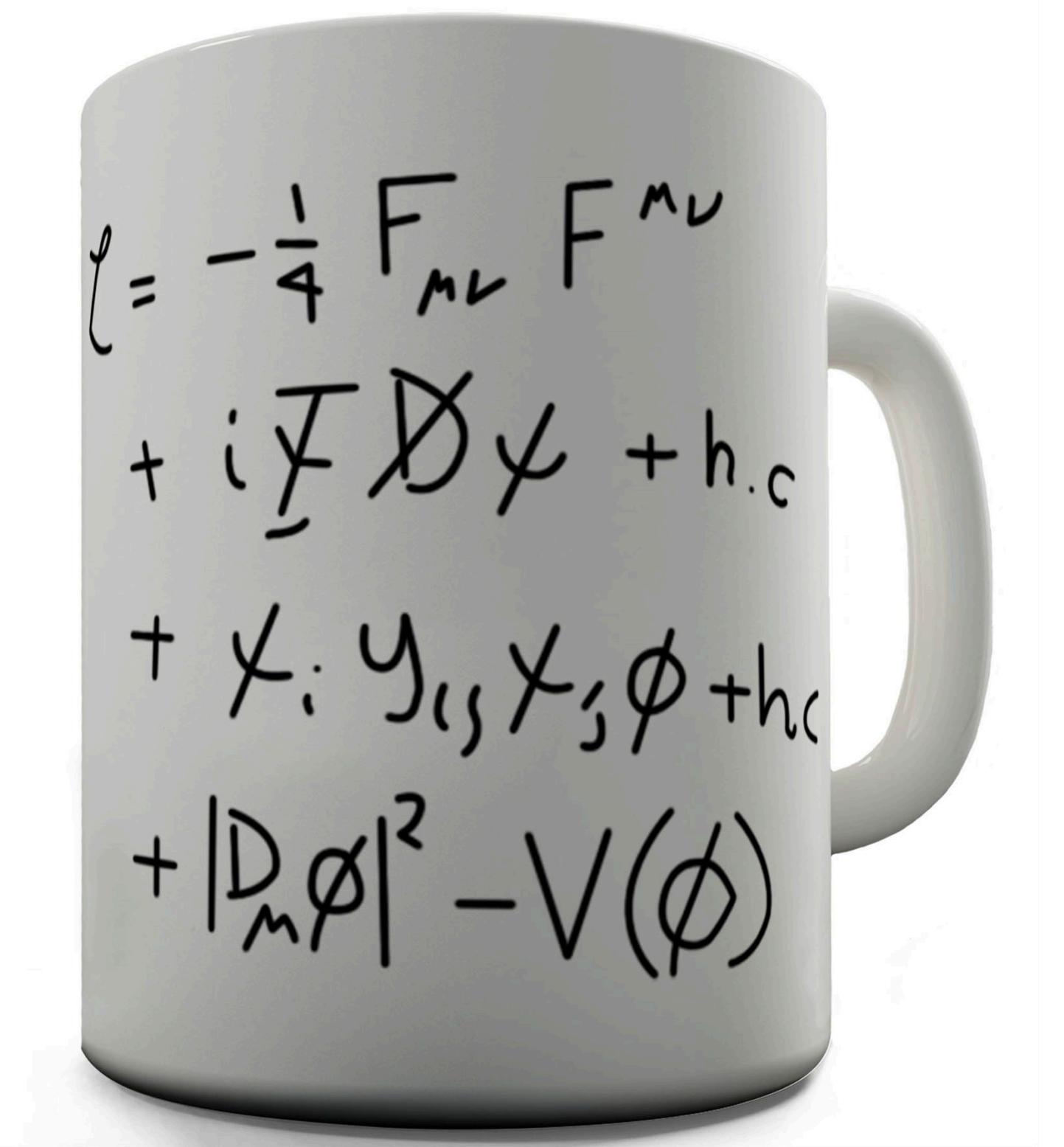
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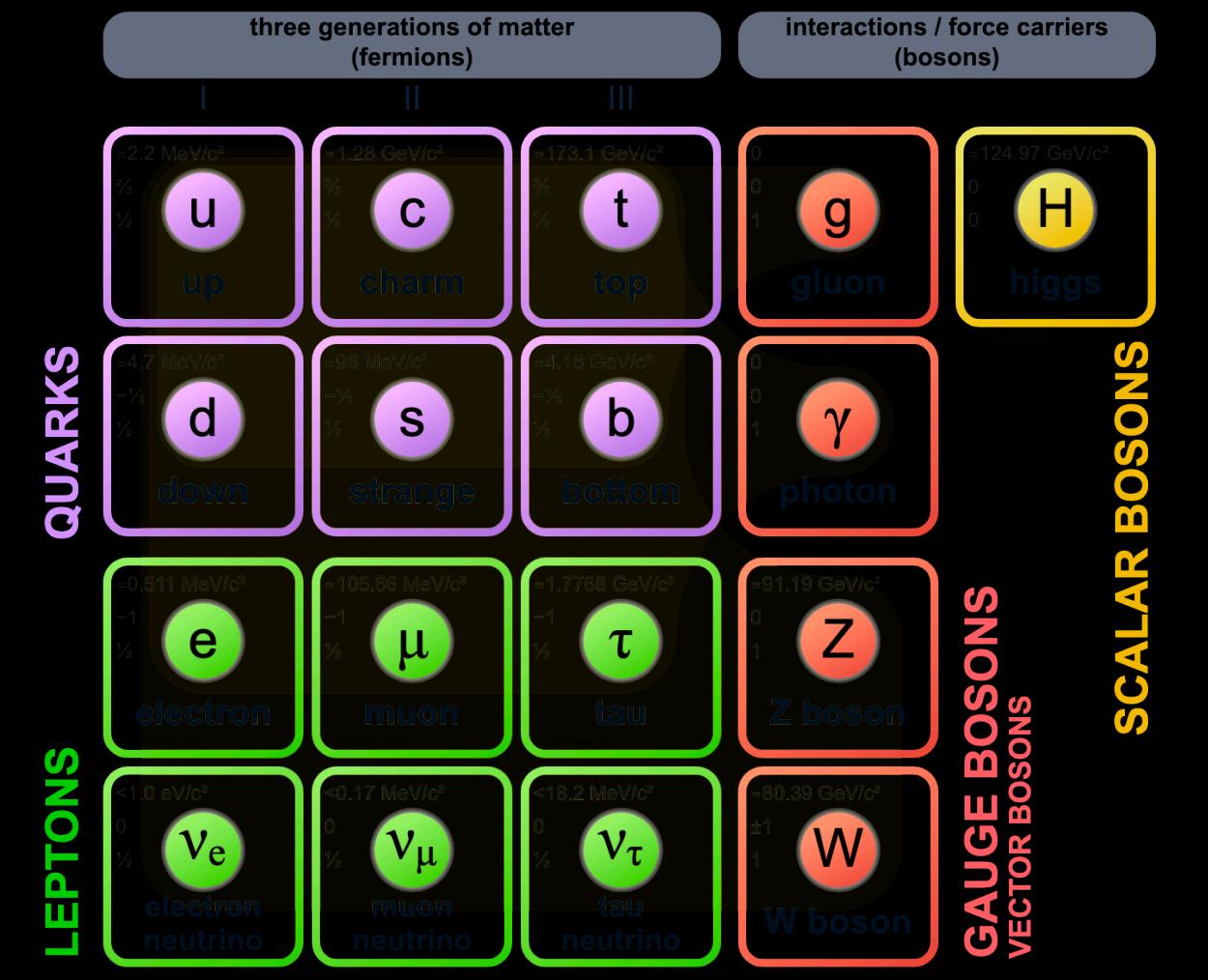
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interaction is multiplied by a factor $-\frac{1}{2}$ rather than $\frac{3}{2}$. Of course our model has too many arbitrary features for these predictions to be taken very seriously, but it is worth keeping in mind that the standard calculation⁸ of the

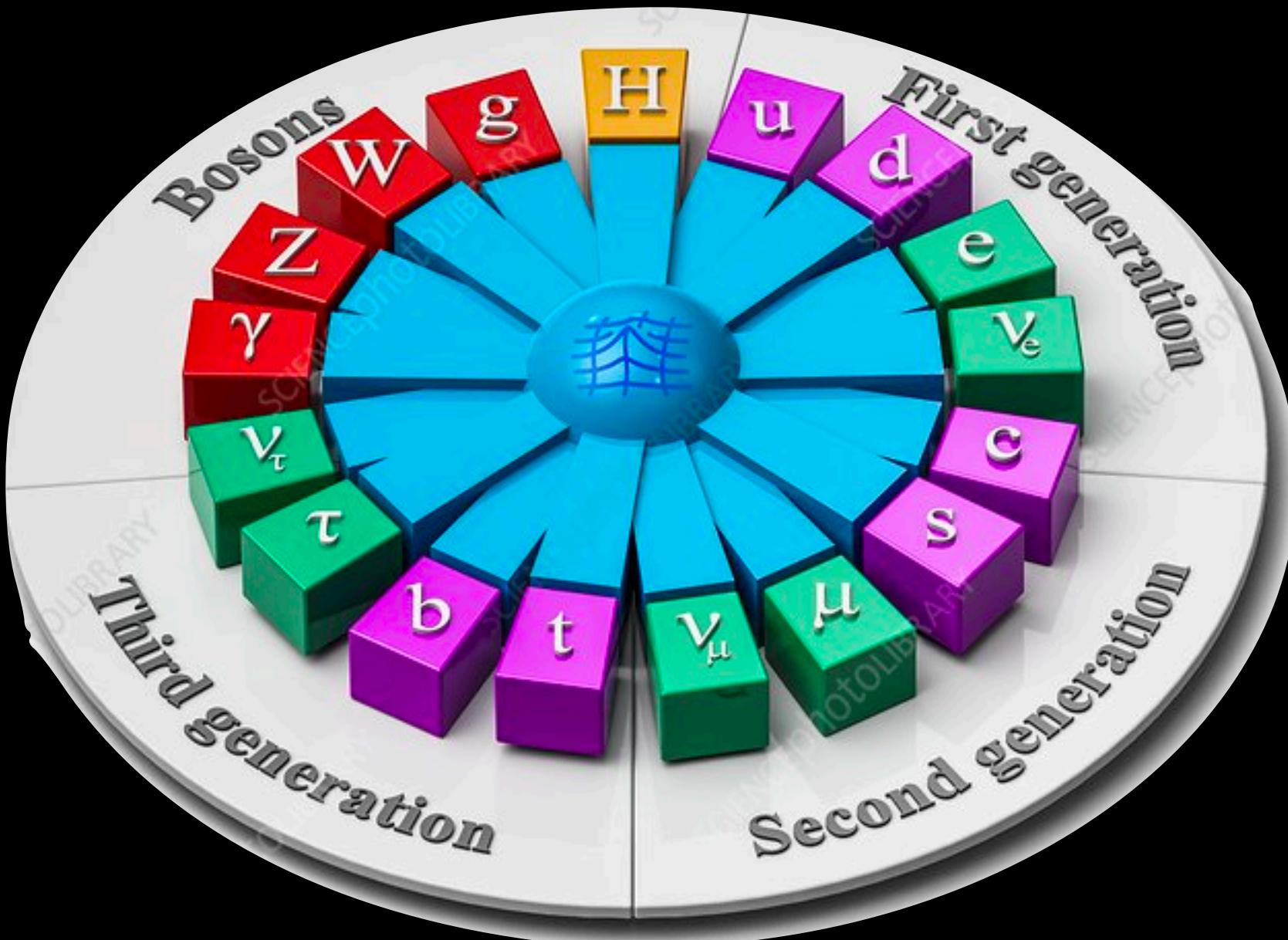




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three generations of matter (fermions)			interactions / force carriers (bosons)	
I	II	III		
 =2.2 MeV/c² $\frac{1}{2}$ $\frac{1}{2}$ u up	 =1.28 GeV/c² $\frac{1}{2}$ $\frac{1}{2}$ c charm	 =173.1 GeV/c² $\frac{1}{2}$ $\frac{1}{2}$ t top	 0 0 1 g gluon	 =124.97 GeV/c² 0 0 H higgs
 =4.7 MeV/c² $-\frac{1}{2}$ $\frac{1}{2}$ d down	 =98 MeV/c² $-\frac{1}{2}$ $\frac{1}{2}$ s strange	 =4.18 GeV/c² $-\frac{1}{2}$ $\frac{1}{2}$ b bottom	 0 0 1 γ photon	
 =0.511 MeV/c² -1 $\frac{1}{2}$ e electron	 =105.66 MeV/c² -1 $\frac{1}{2}$ μ muon	 =1.7768 GeV/c² -1 $\frac{1}{2}$ τ tau	 =91.19 GeV/c² 0 1 Z Z boson	
 <1.0 eV/c² 0 $\frac{1}{2}$ ν_e electron neutrino	 =0.17 MeV/c² 0 $\frac{1}{2}$ ν_μ muon neutrino	 <18.2 MeV/c² 0 $\frac{1}{2}$ ν_τ tau neutrino	 =80.39 GeV/c² ± 1 1 W W boson	
LEPTONS		GAUGE BOSONS VECTOR BOSONS		

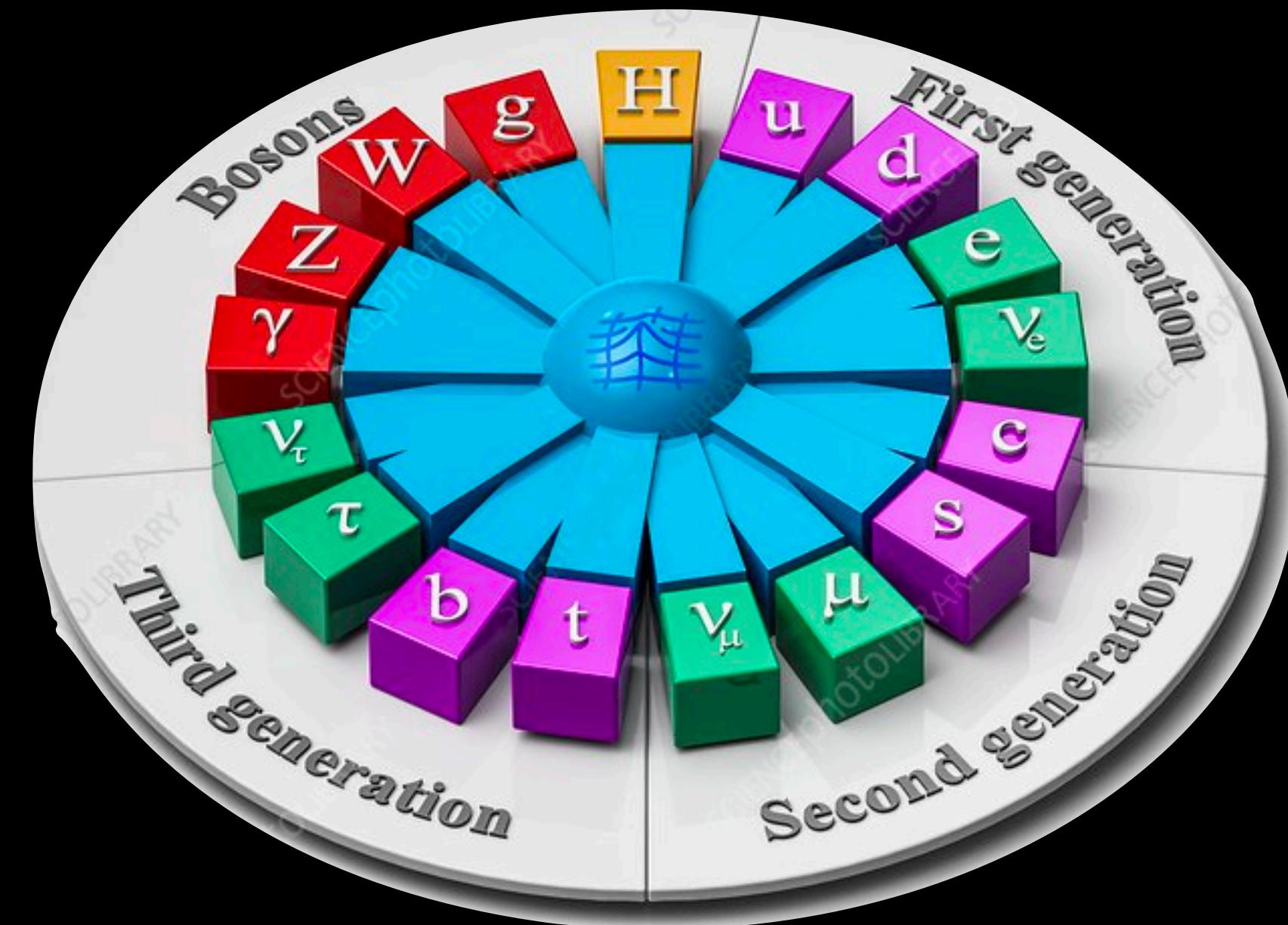
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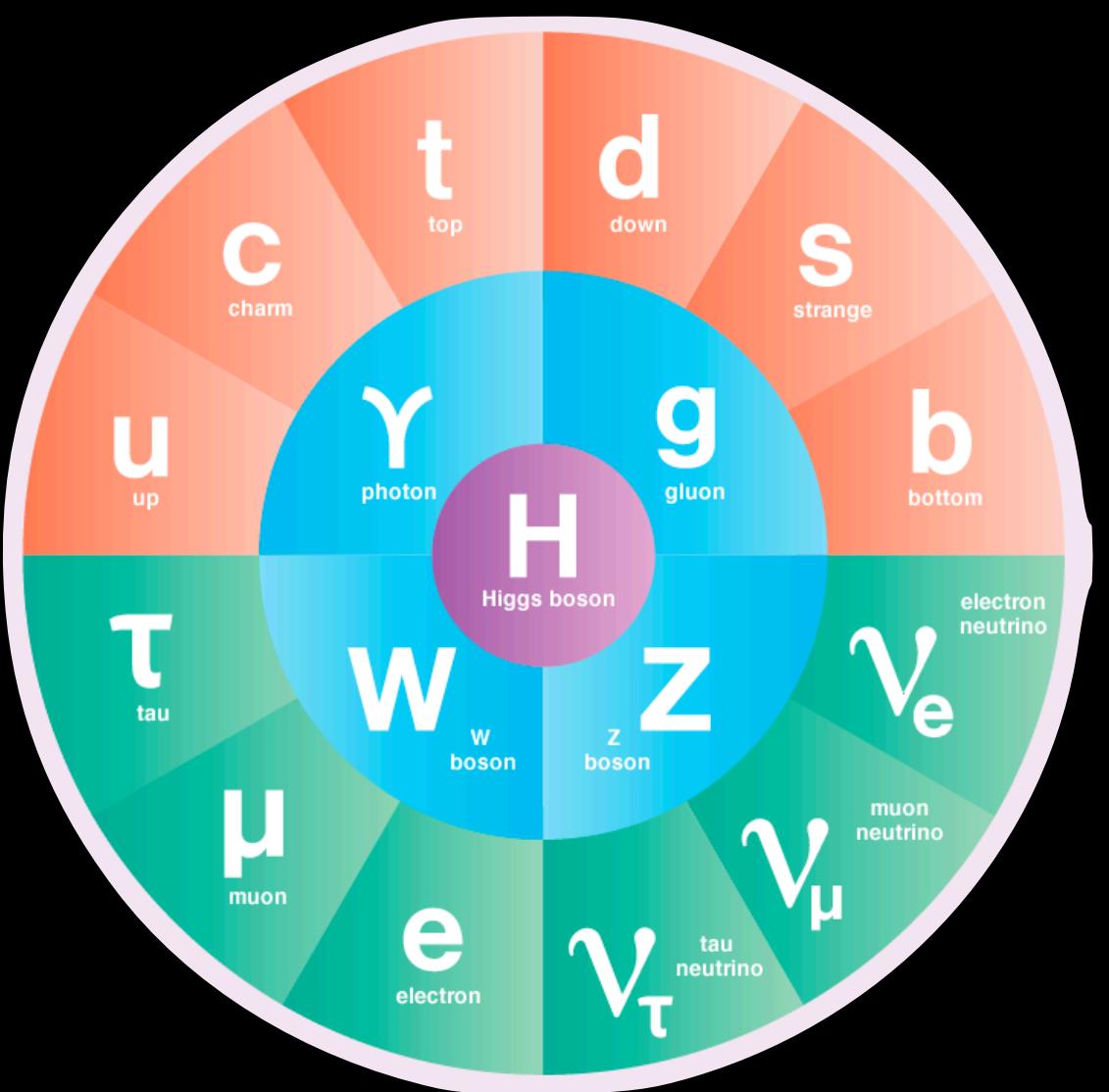
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three generations of matter (fermions)			interactions / force carriers (bosons)	
I	II	III	g	H
1/2 MeV/c ² u up	1.26 GeV/c ² c charm	173.1 GeV/c ² t top	124.97 GeV/c ² gluon	Higgs
1/2 MeV/c ² d down	98 MeV/c ² s strange	4.18 GeV/c ² b bottom	0.0 γ photon	
0.911 MeV/c ² e electron	105.06 MeV/c ² μ muon	1.7708 GeV/c ² τ tau	81.19 GeV/c ² Z boson	
-1.0 eV/c ² ν _e electron neutrino	-0.17 MeV/c ² ν _μ muon neutrino	-18.2 MeV/c ² ν _τ tau neutrino	-80.39 GeV/c ² W boson	

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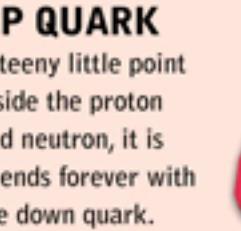
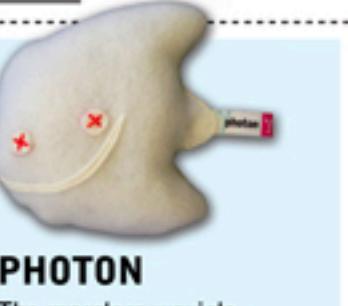
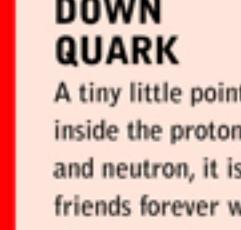
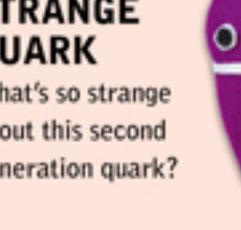
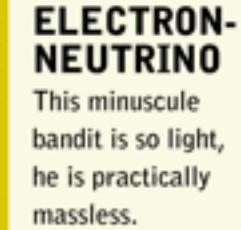
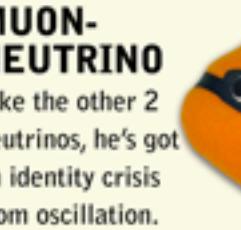
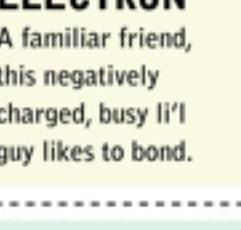
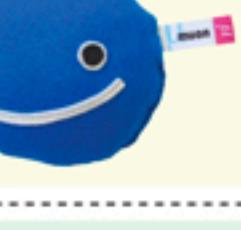
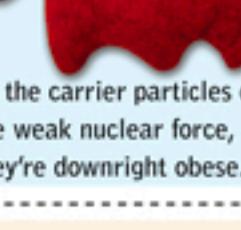
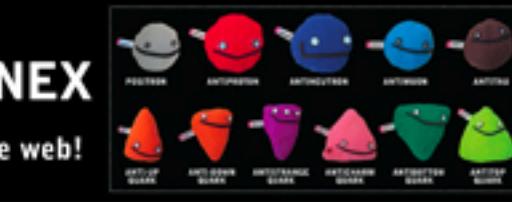
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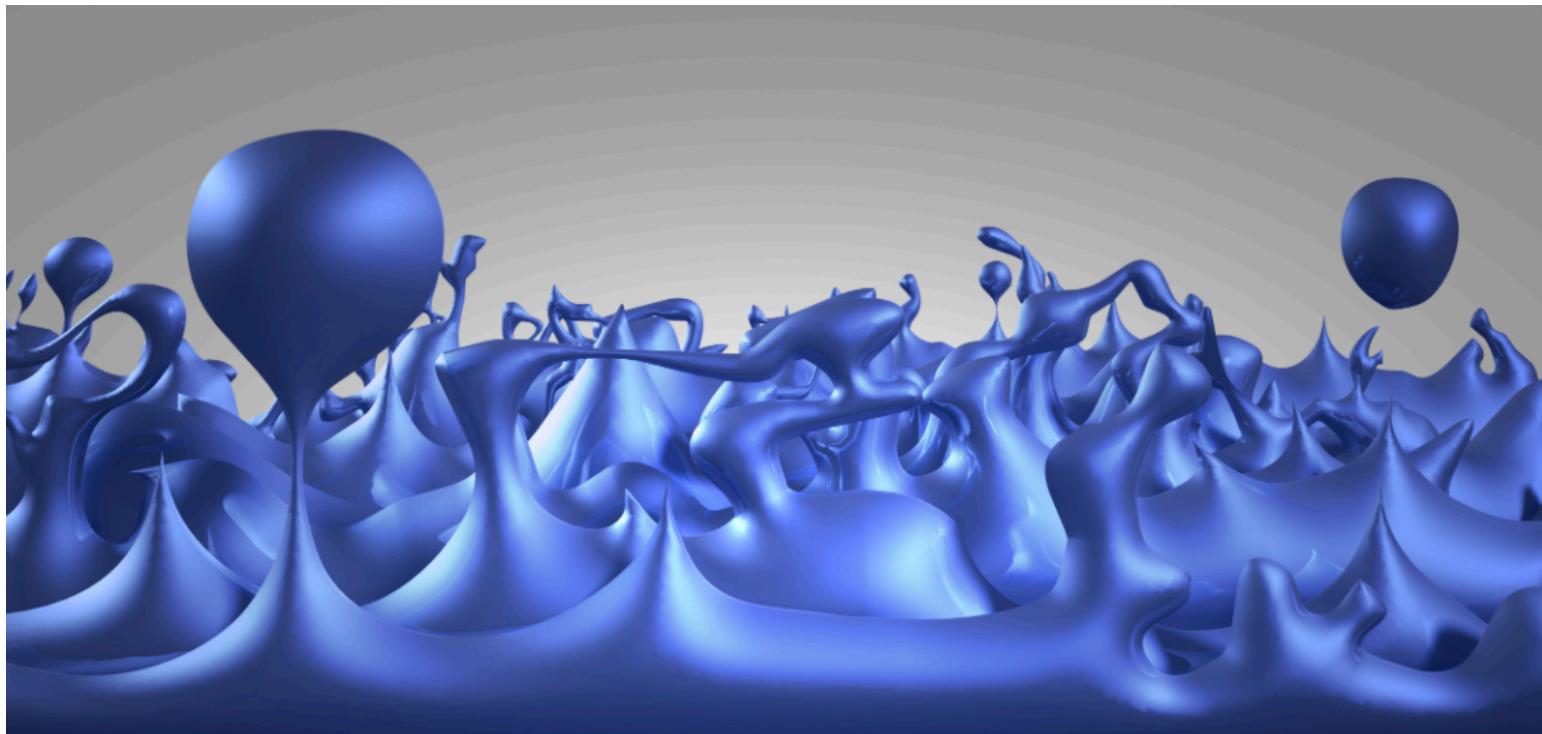
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The PARTICLE ZOO

Subatomic Particle Plush Toys FROM THE STANDARD MODEL OF PHYSICS & beyond!

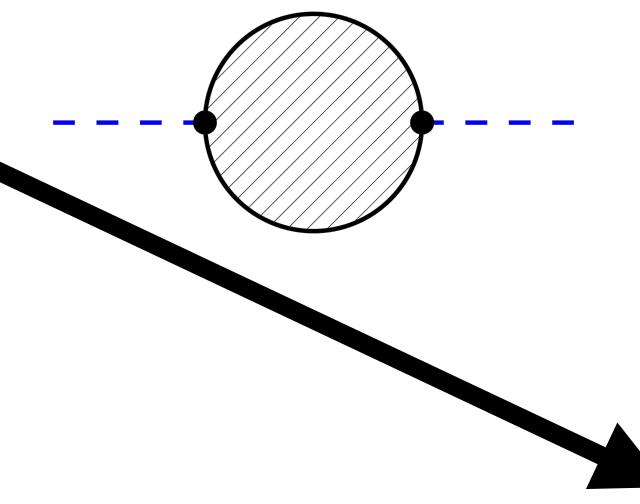
QUARKS		FORCE CARRIERS									
	UP QUARK A teeny little point inside the proton and neutron, it is friends forever with the down quark.		CHARM QUARK A second generation quark, he is charmed, indeed.		TOP QUARK This heavyweight champion doesn't live long enough to make friends with anyone.		PHOTON The massless wavelike we know and love.				
	DOWN QUARK A tiny little point inside the proton and neutron, it is friends forever with the up quark.		STRANGE QUARK What's so strange about this second generation quark?		BOTTOM QUARK This third generation quark is puttin' on the pounds.		GLUON The "glue" of the strong nuclear force.				
LEPTONS			ELECTRON-NEUTRINO This minuscule bandit is so light, he is practically massless.		MUON-NEUTRINO Like the other 2 neutrinos, he's got an identity crisis from oscillation.		TAU-NEUTRINO He's a tau now, but what type of neutrino will he be next?		W BOSON		Z BOSON
	ELECTRON A familiar friend, this negatively charged, busy li'l guy likes to bond.		MUON A "heavy electron" who lives fast and dies young.		TAU A "heavy muon" who could stand to lose a little weight.		HIGGS BOSON He's the one everyone wants to meet, but for now he's playing hard to get. You'd be smiling too if everyone was looking to interview you.		GRAVITON Still unobserved, yet theoretically everywhere, he's got big legs for jumping branes.		PROTON We would not be here without her positivity.
	TACHYON Can this devious and clever particle really travel faster than light?		DARK MATTER The mysterious missing mass. Difficult to see because he's so dark.		NEUTRON He insists on remaining neutral.		NEW! GIFT CARDS		STAMPSHEET Twenty-three particles on one 8.5x11" sheet of perforated "stamps"		
<p>Visit the ANTIPARTICLE ANNEX You can now buy antimatter on the web!</p> 											

© Universe Today

10^{-35}m 

*Image Credit: NASA/CXC/
M. Weiss*

RGE



NATURALNESS, CHIRAL SYMMETRY, AND SPONTANEOUS
CHIRAL SYMMETRY BREAKING

G. 't Hooft

1980

magnitude smaller than another then the smallest must satisfy our "dogma" separately. As we will see, naturalness will put the severest restriction on the occurrence of scalar particles in renormalizable theories. In fact we conjecture that this is the reason why light, weakly interacting scalar particles are not seen.

 10^{-18}m

Standard Model of Elementary Particles

three generations of matter (fermions)			interactions / force carriers (bosons)	
I	II	III		
mass $\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$		
charge $\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$		
spin $\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$		
u up	c charm	t top	g gluon	H higgs
d down	s strange	b bottom	γ photon	
e electron	μ muon	τ tau	Z Z boson	
ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

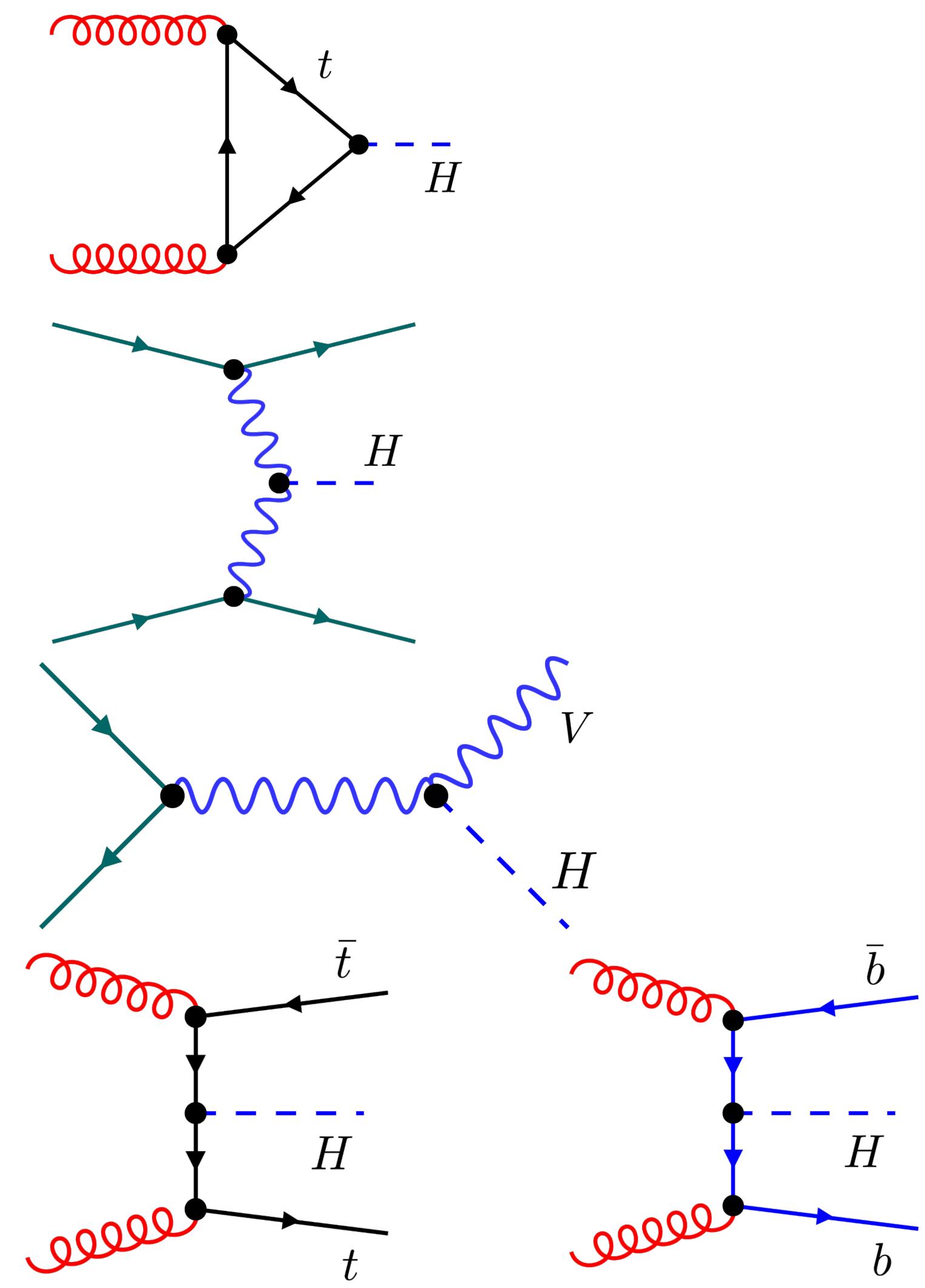
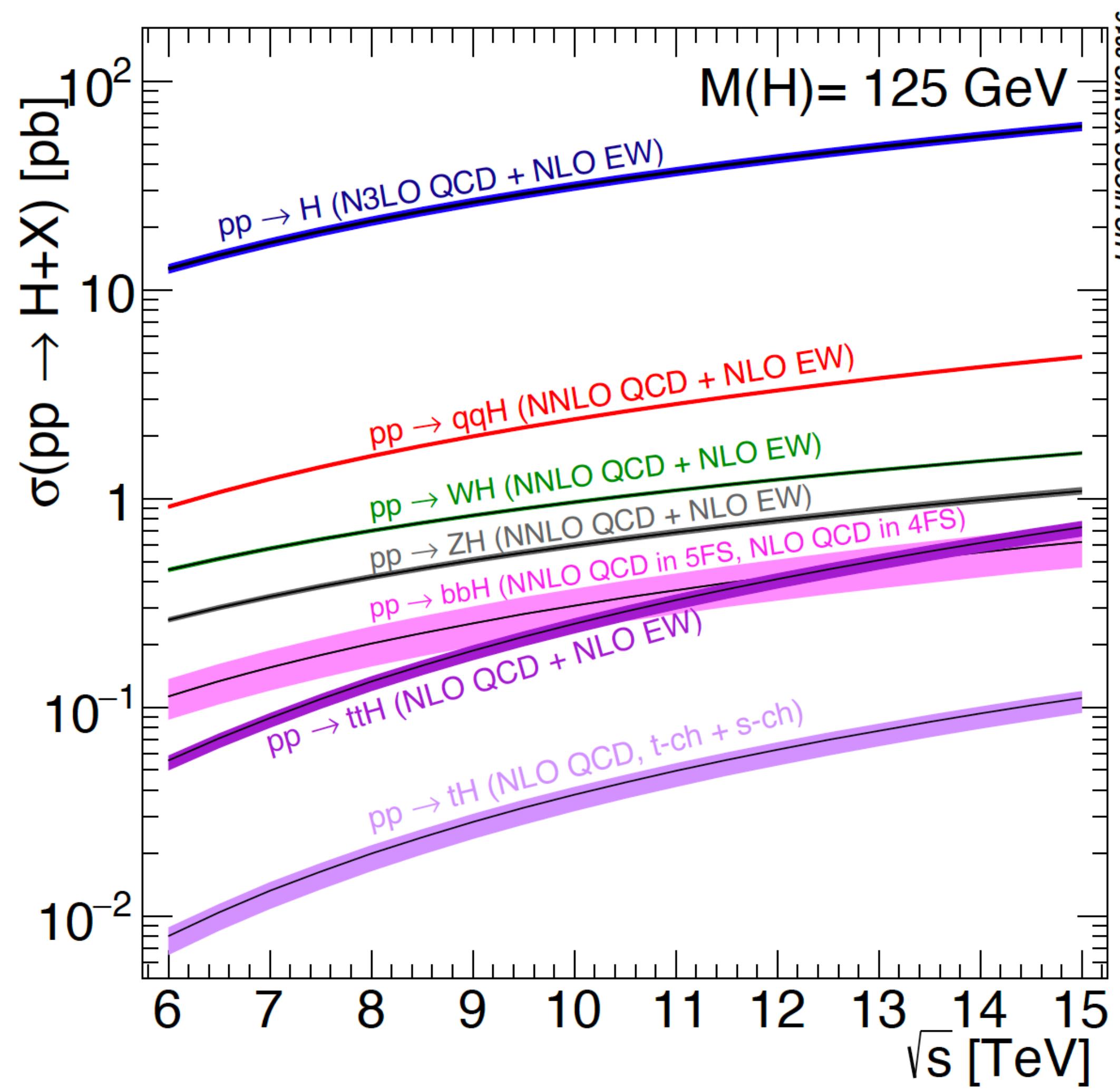
QUARKS

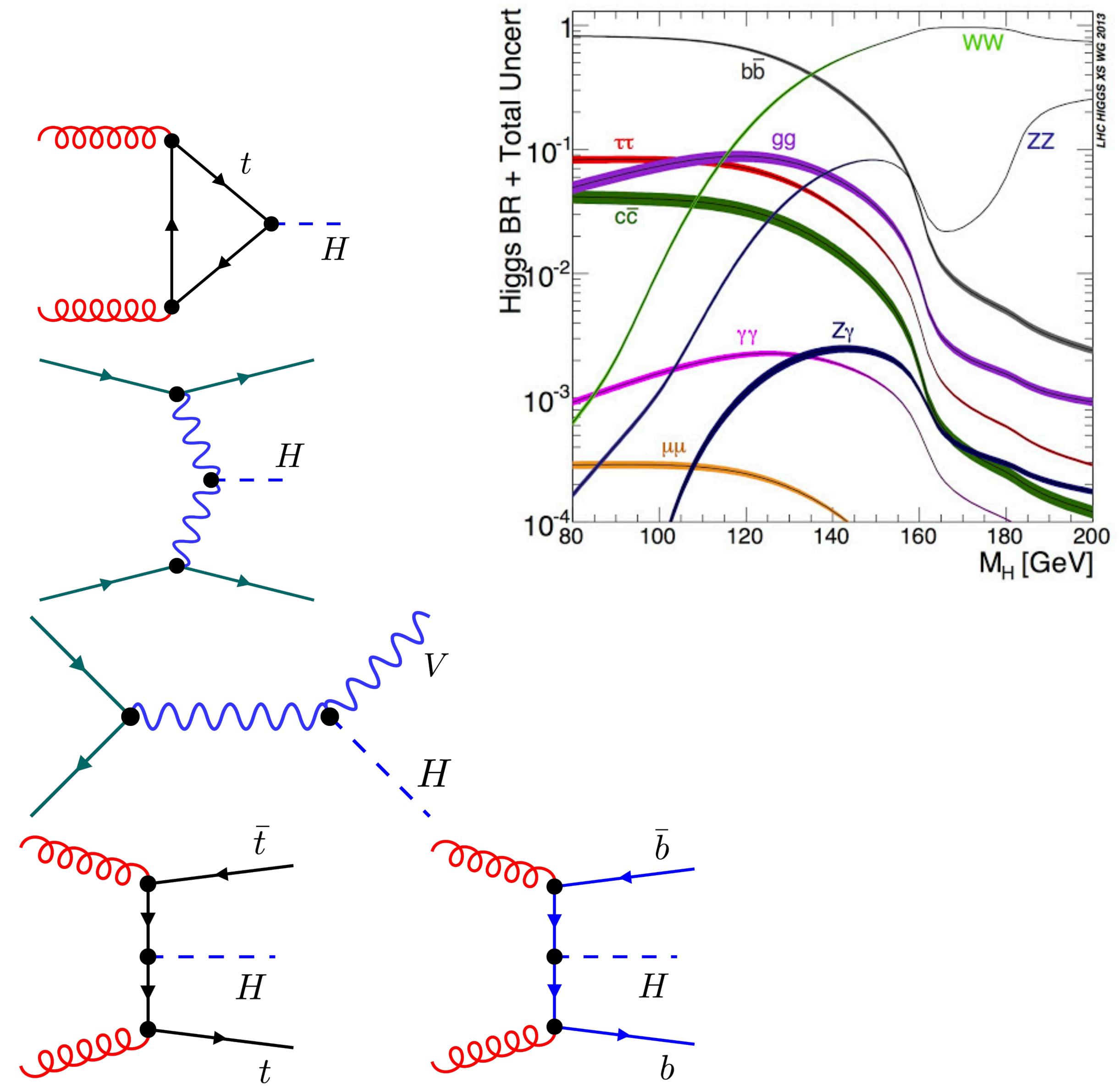
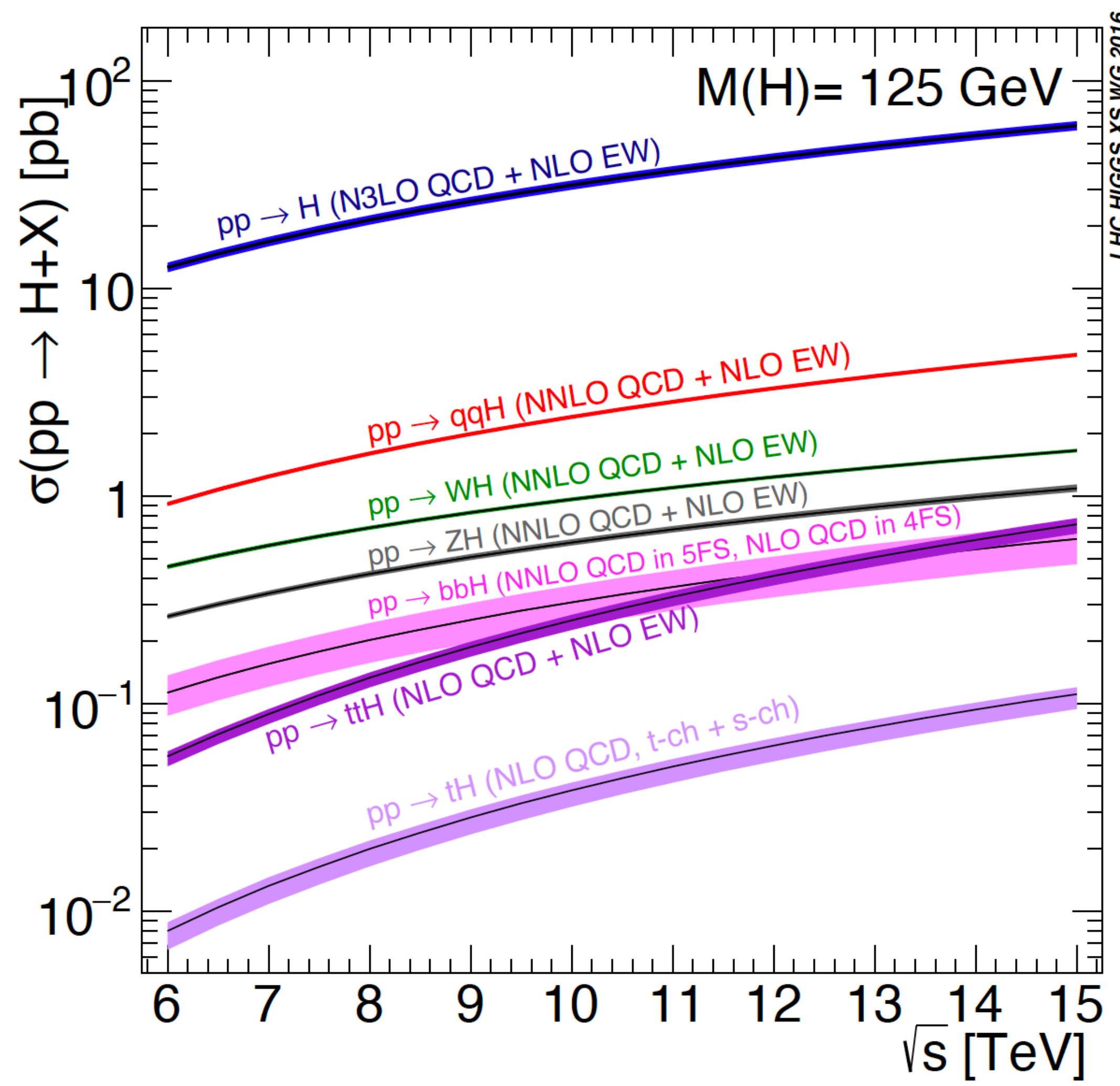
SCALAR BOSONS

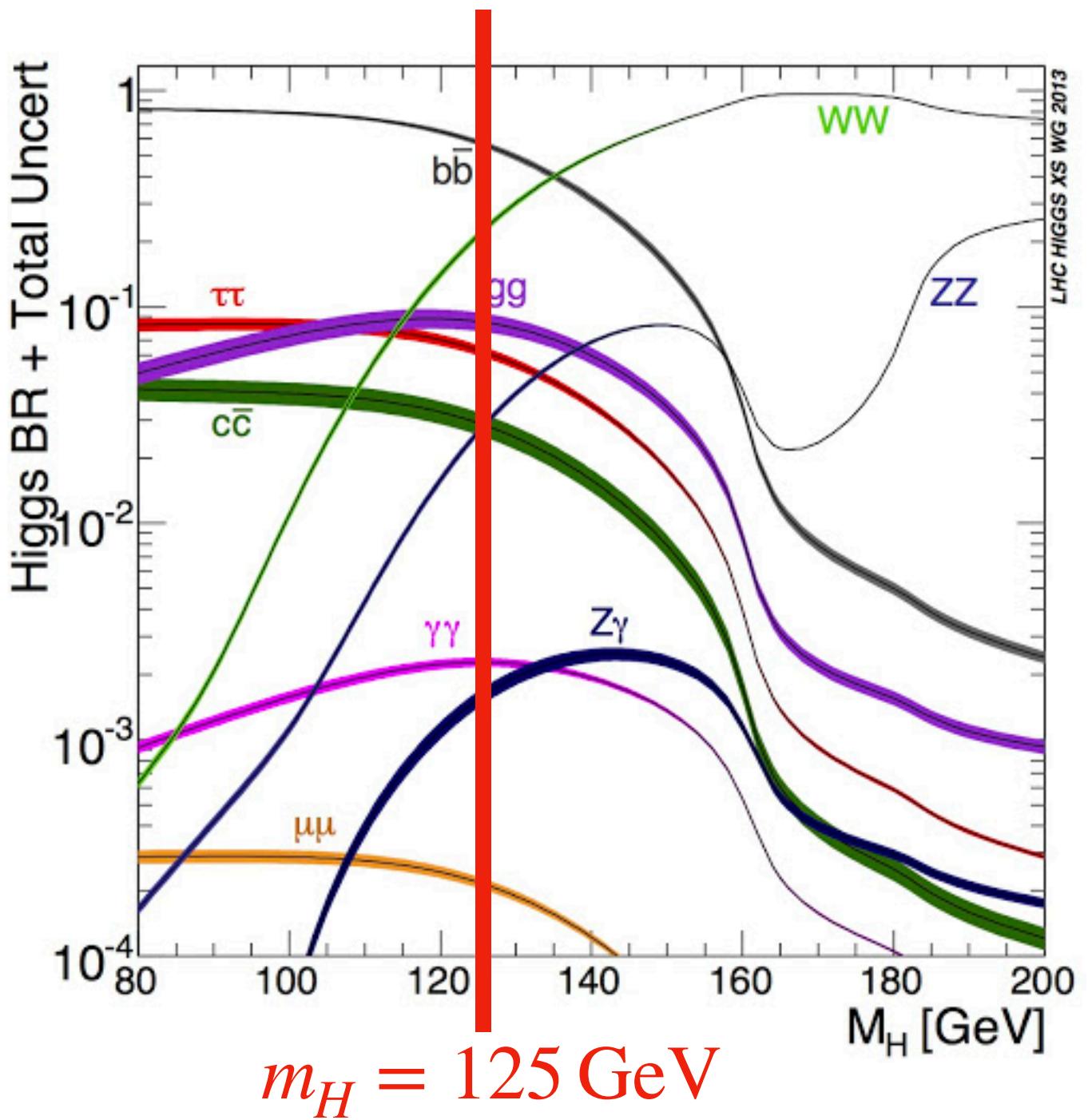
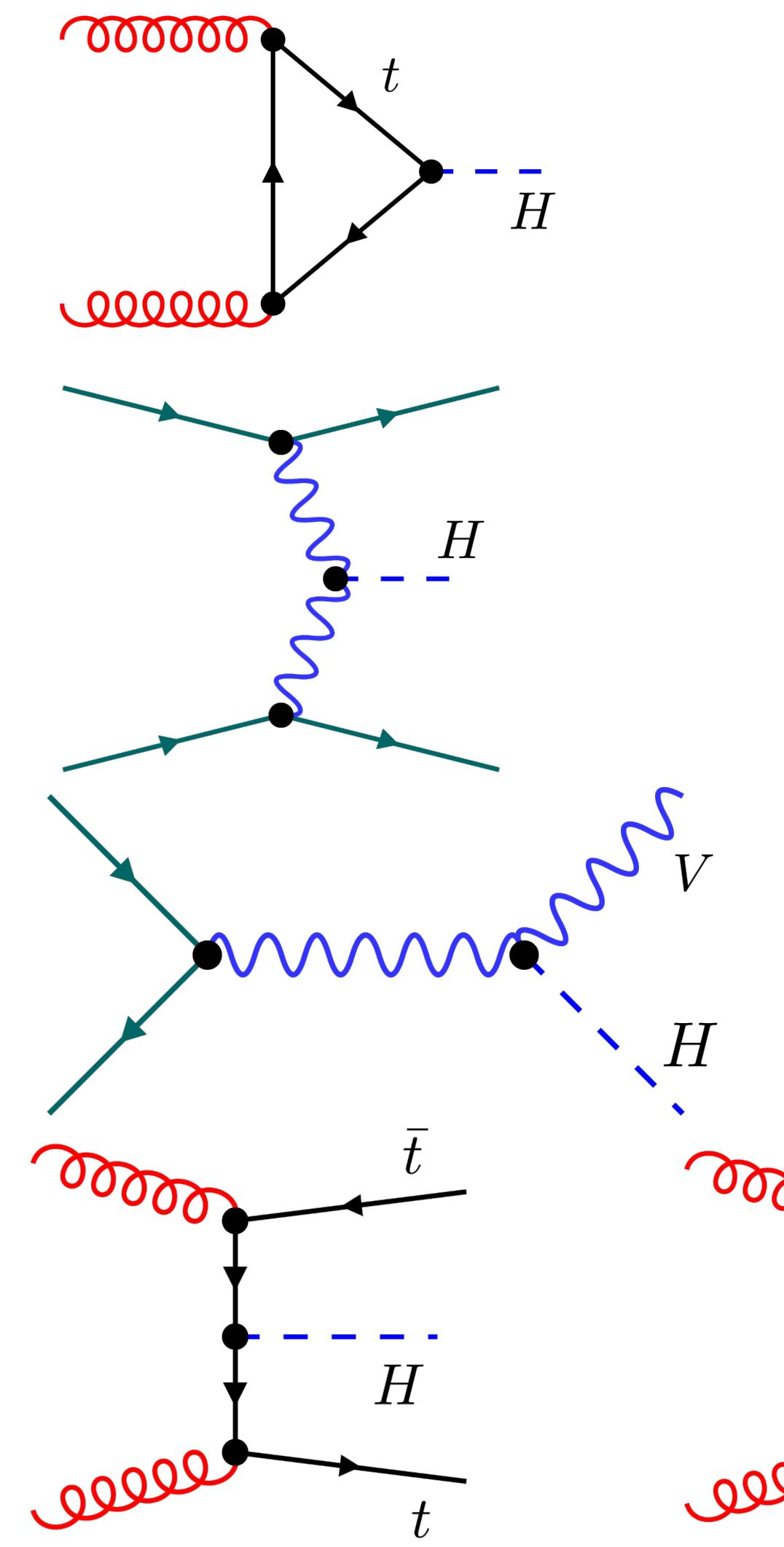
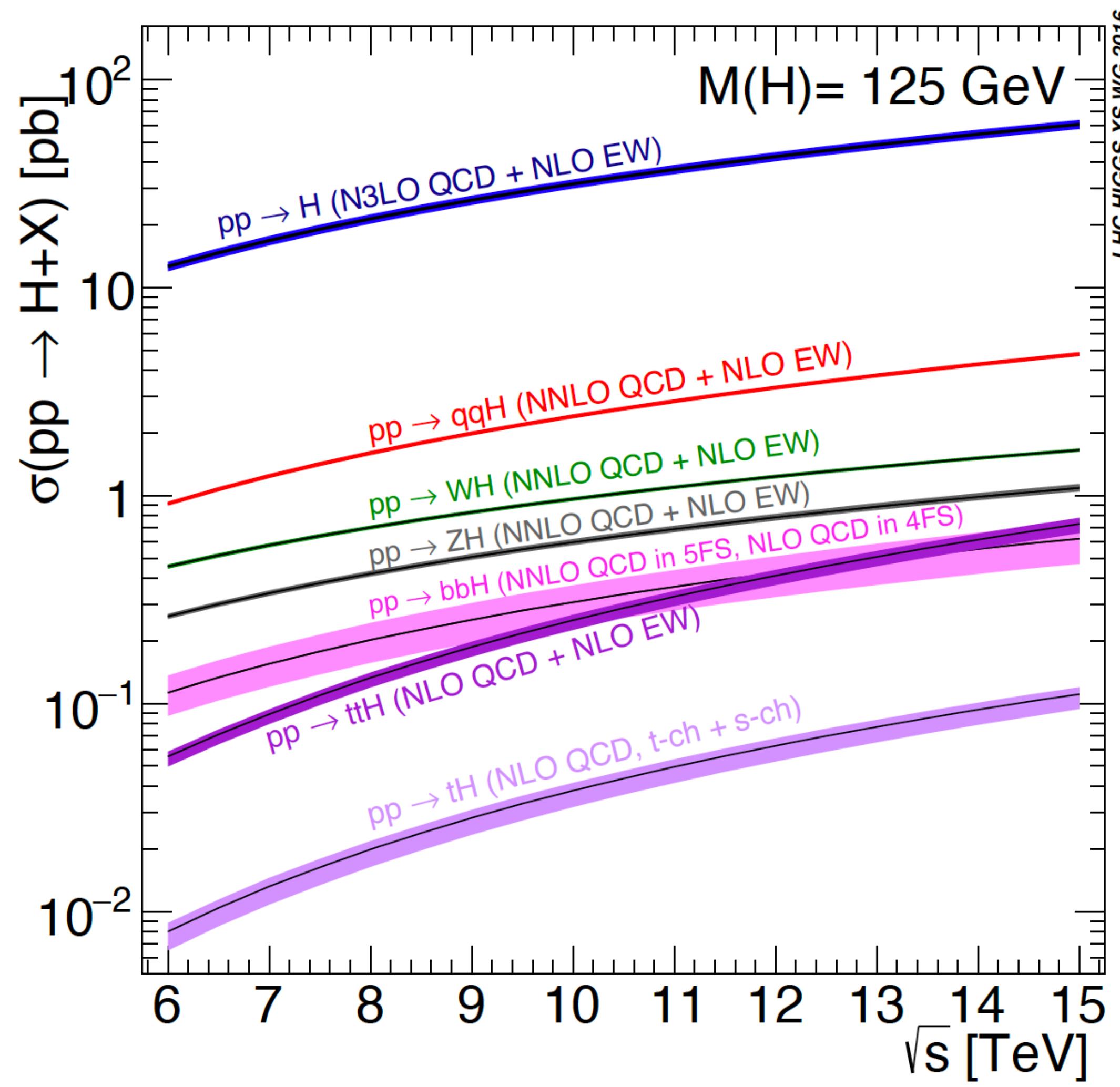
LEPTONS

GAUGE BOSONS
VECTOR BOSONS

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Collider Physics at the Precision Frontier

Gudrun Heinrich

P3H-20-044

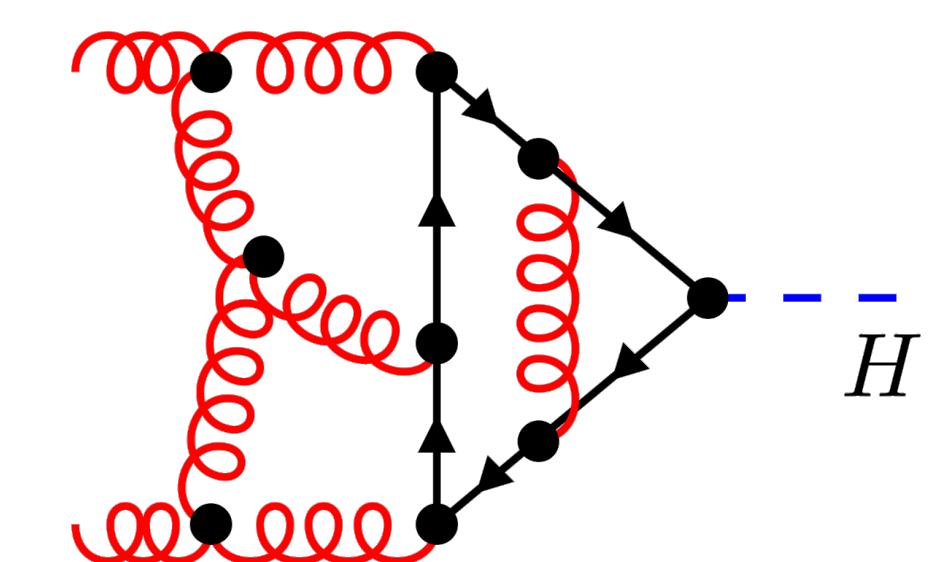
*Karlsruhe Institute of Technology, Institute for Theoretical Physics, Wolfgang-Gaede-Str. 1,
76131 Karlsruhe, Germany*

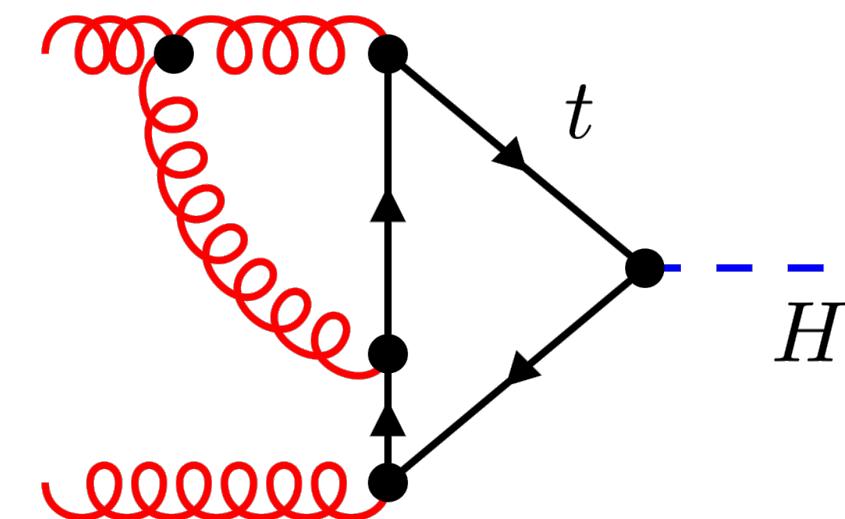
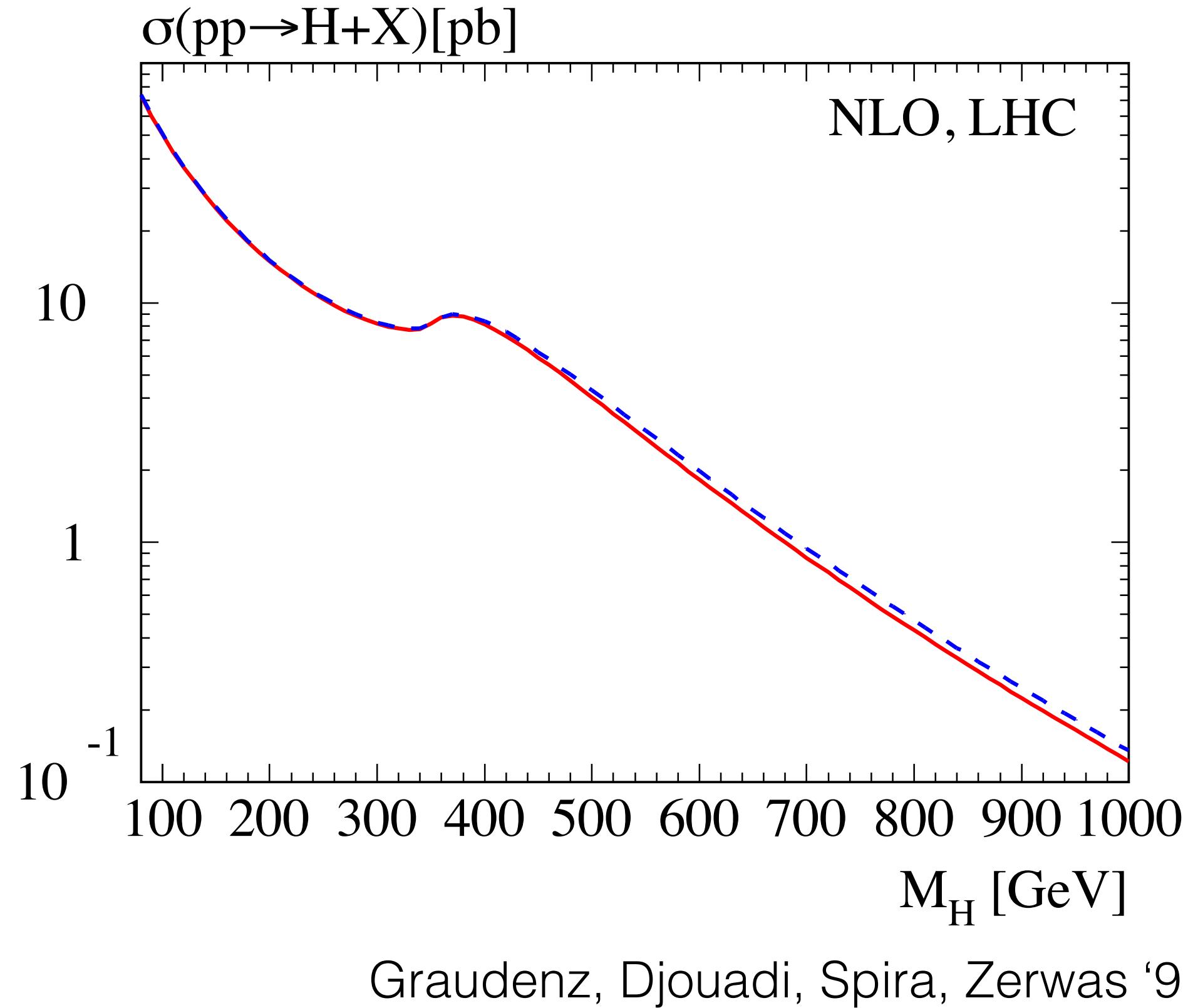
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$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.37\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

LHCH(XS)WG YR4 '16

Inclusive gluon fusion cross section:

$$48.58 \text{ pb} = \begin{array}{lll} 16.00 \text{ pb} & (+32.9\%) & (\text{LO, rEFT}) \\ + 20.84 \text{ pb} & (+42.9\%) & (\text{NLO, rEFT}) \\ - 2.05 \text{ pb} & (-4.2\%) & ((t, b, c), \text{exact NLO}) \\ + 9.56 \text{ pb} & (+19.7\%) & (\text{NNLO, rEFT}) \\ + 0.34 \text{ pb} & (+0.7\%) & (\text{NNLO}, 1/m_t) \\ + 2.40 \text{ pb} & (+4.9\%) & (\text{EW, QCD-EW}) \\ + 1.49 \text{ pb} & (+3.1\%) & (\text{N}^3\text{LO, rEFT}) \end{array}$$

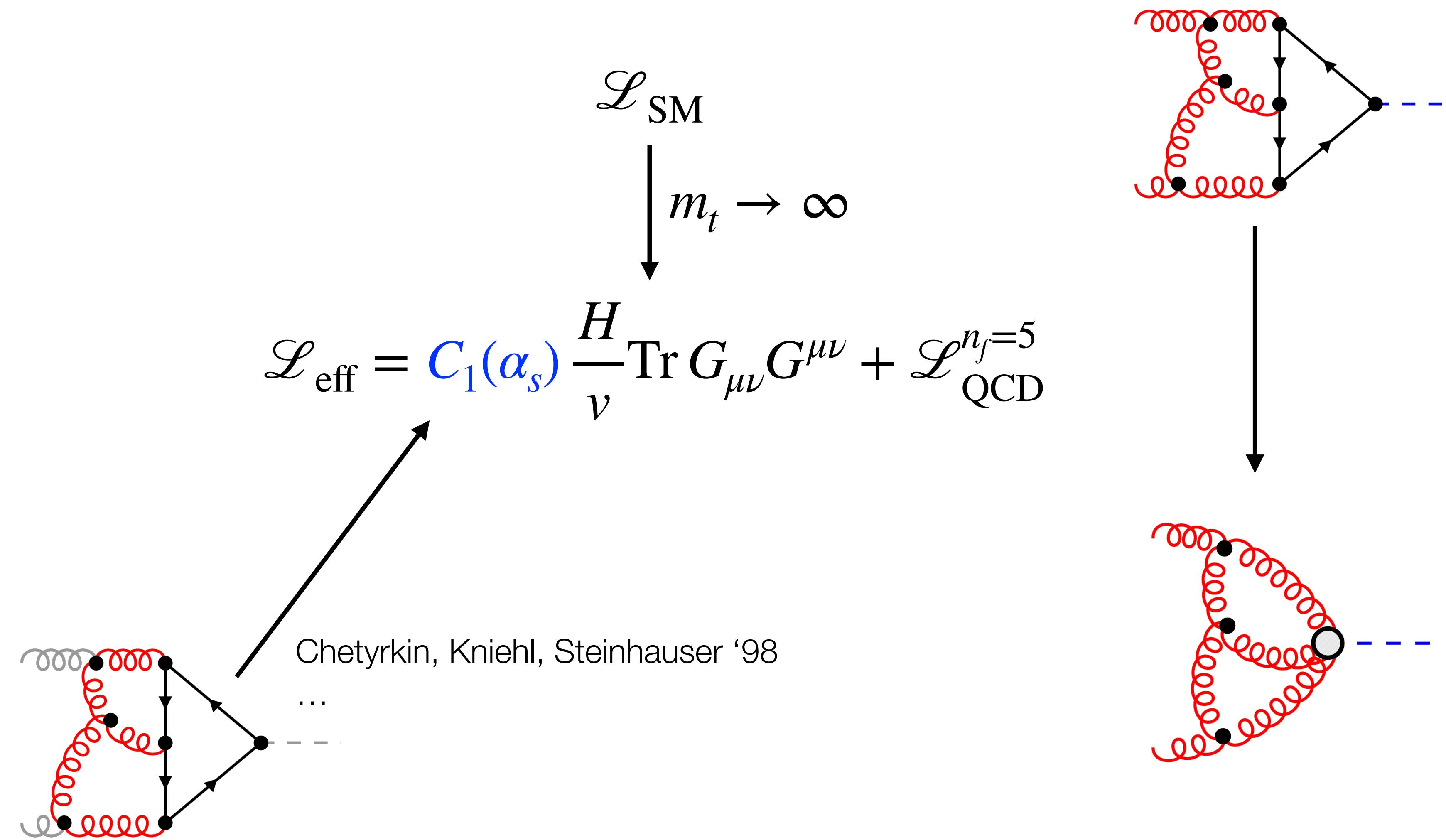


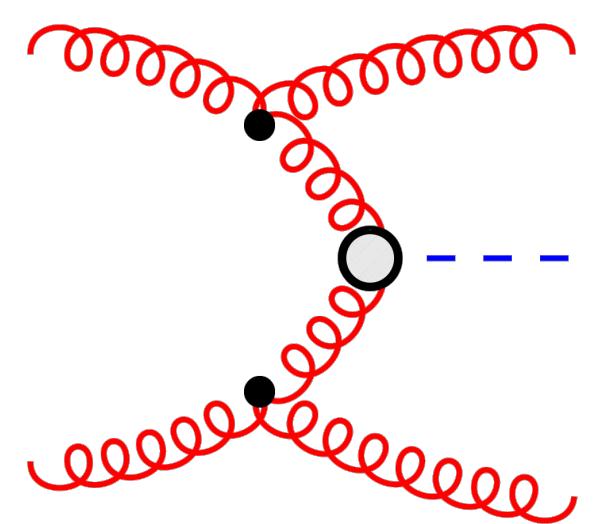
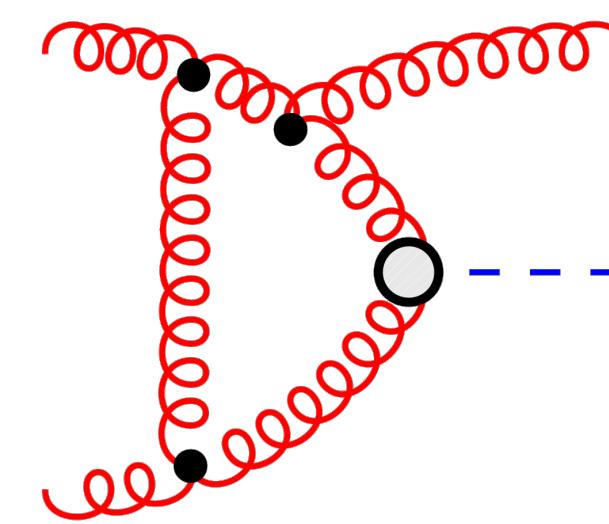
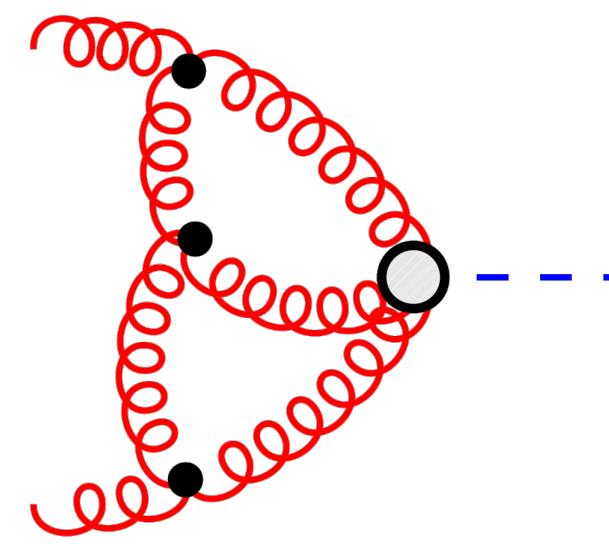


$$\sigma_{\text{HEFT}}^{\text{HO}} = \sigma^{\text{LO}} \left(\frac{\sigma^{\text{HO}}}{\sigma^{\text{LO}}} \right)_{m_t \rightarrow \infty}$$

$$\sigma_{\text{exact}}^{\text{NLO}} = (1 - 0.008) \cdot \sigma_{\text{HEFT}}^{\text{NLO}}$$

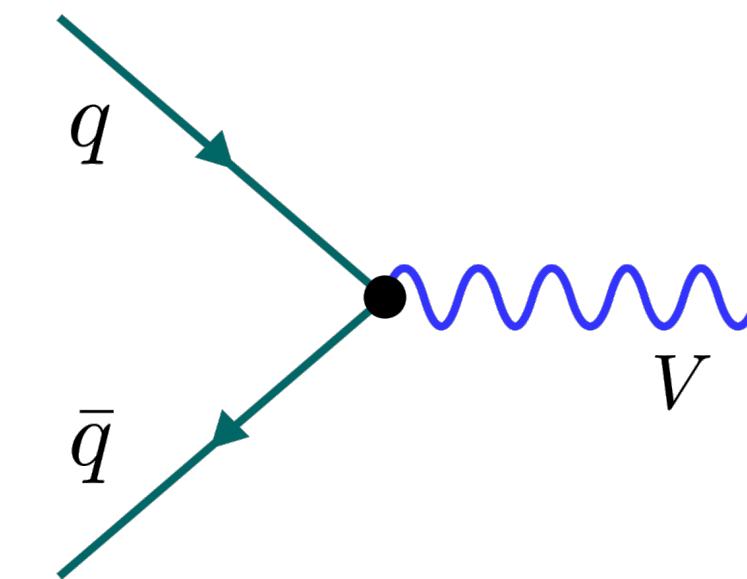
for $m_H = 125 \text{ GeV}$





- need to cancel IR divergences
- around '01: only NLO methods
- but: NNLO Drell-Yan 1990

v. Neerven et al.

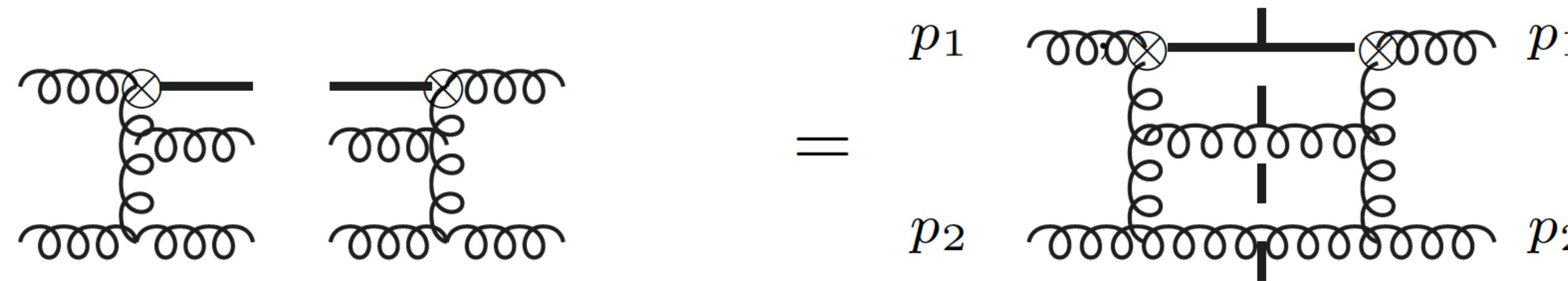


perform phase space integrals analytically
in D dimensions

IR poles cancel analytically (no “subtraction scheme”)

still: phase space integration complicated (back then...)

Reverse unitarity:



$$2i\pi\delta(p^2 - m^2) \rightarrow \frac{1}{p^2 - m^2 + i0} - \frac{1}{p^2 - m^2 - i0}$$

Anastasiou, Melnikov '02
Anastasiou, Dixon, Melnikov '03

Treat phase space integrals like loop integrals:
integration-by-parts \rightarrow master integrals \rightarrow differential equations

soft expansion:

$$z = \hat{s}/M_H^2 \rightarrow 1$$

$$\begin{aligned}\hat{\sigma}_{ij} &= \sum_{n \geq 0} \left(\frac{\alpha_s}{\pi} \right)^n \hat{\sigma}_{ij}^{(n)}, \\ \hat{\sigma}_{ij}^{(n)} &= a^{(n)} \delta(1-x) + \sum_{k=0}^{2n-1} b_k^{(n)} \left[\frac{\ln^k(1-x)}{1-x} \right]_+ \\ &\quad + \sum_{l=0}^{\infty} \sum_{k=0}^{2n-1} c_{lk}^{(n)} (1-x)^l \ln^k(1-x),\end{aligned}$$

soft expansion:

$$z = \hat{s}/M_H^2 \rightarrow 1$$

$$\hat{\sigma}_{ij} = \sum_{n \geq 0} \left(\frac{\alpha_s}{\pi} \right)^n \hat{\sigma}_{ij}^{(n)},$$

soft-virtual approximation

$$\begin{aligned} \hat{\sigma}_{ij}^{(n)} &= a^{(n)} \delta(1-x) + \sum_{k=0}^{2n-1} b_k^{(n)} \left[\frac{\ln^k(1-x)}{1-x} \right]_+ \\ &+ \sum_{l=0}^{\infty} \sum_{k=0}^{2n-1} c_{lk}^{(n)} (1-x)^l \ln^k(1-x), \end{aligned}$$

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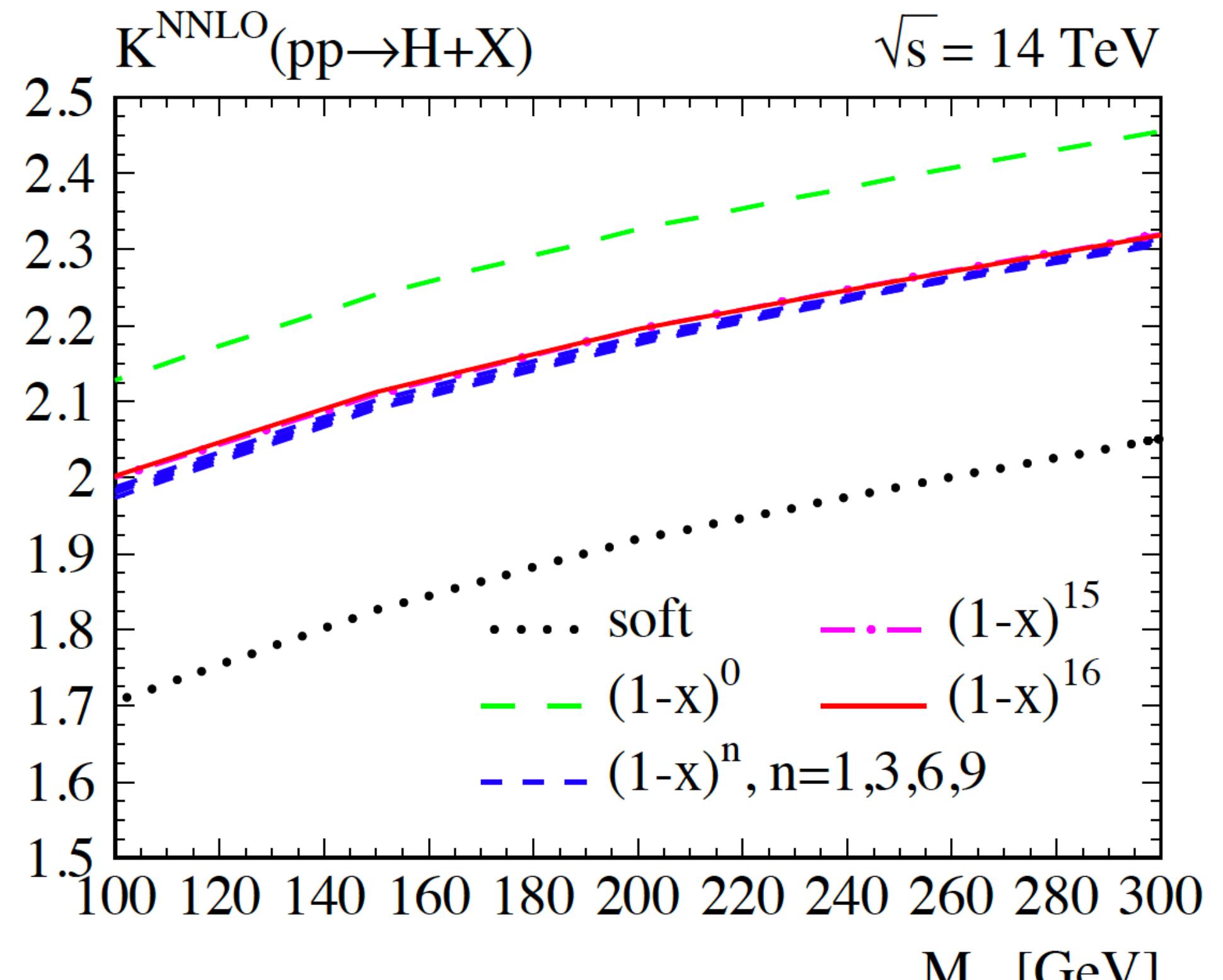
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$$+ \sum_{l=0}^{\infty} \sum_{k=0}^{2n-1} c_{lk}^{(n)} (1-x)^l \ln^k(1-x),$$

$l = 0, k > 0$: next-to-soft approximation

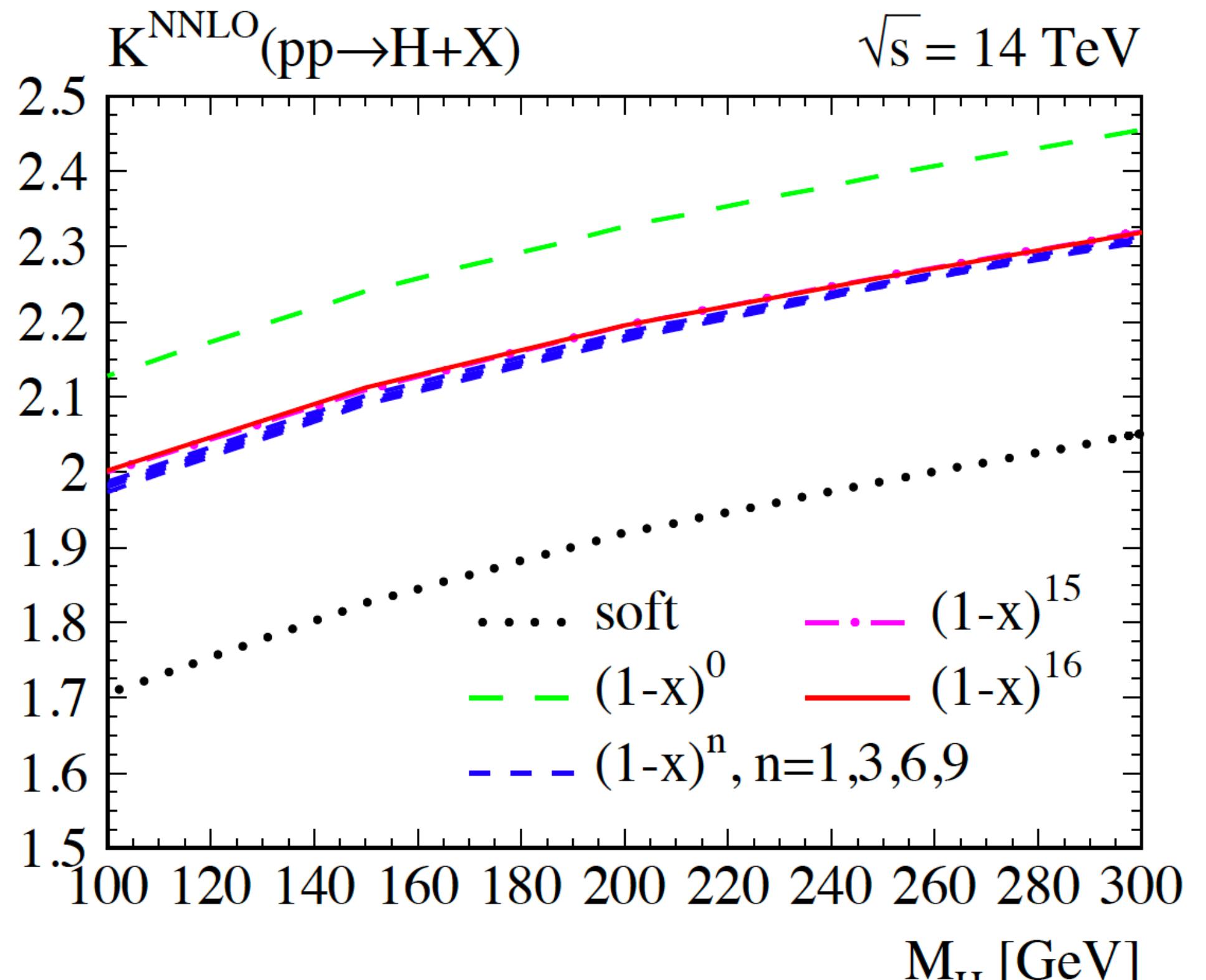
Krämer, Laenen, Spira '96

NNLO



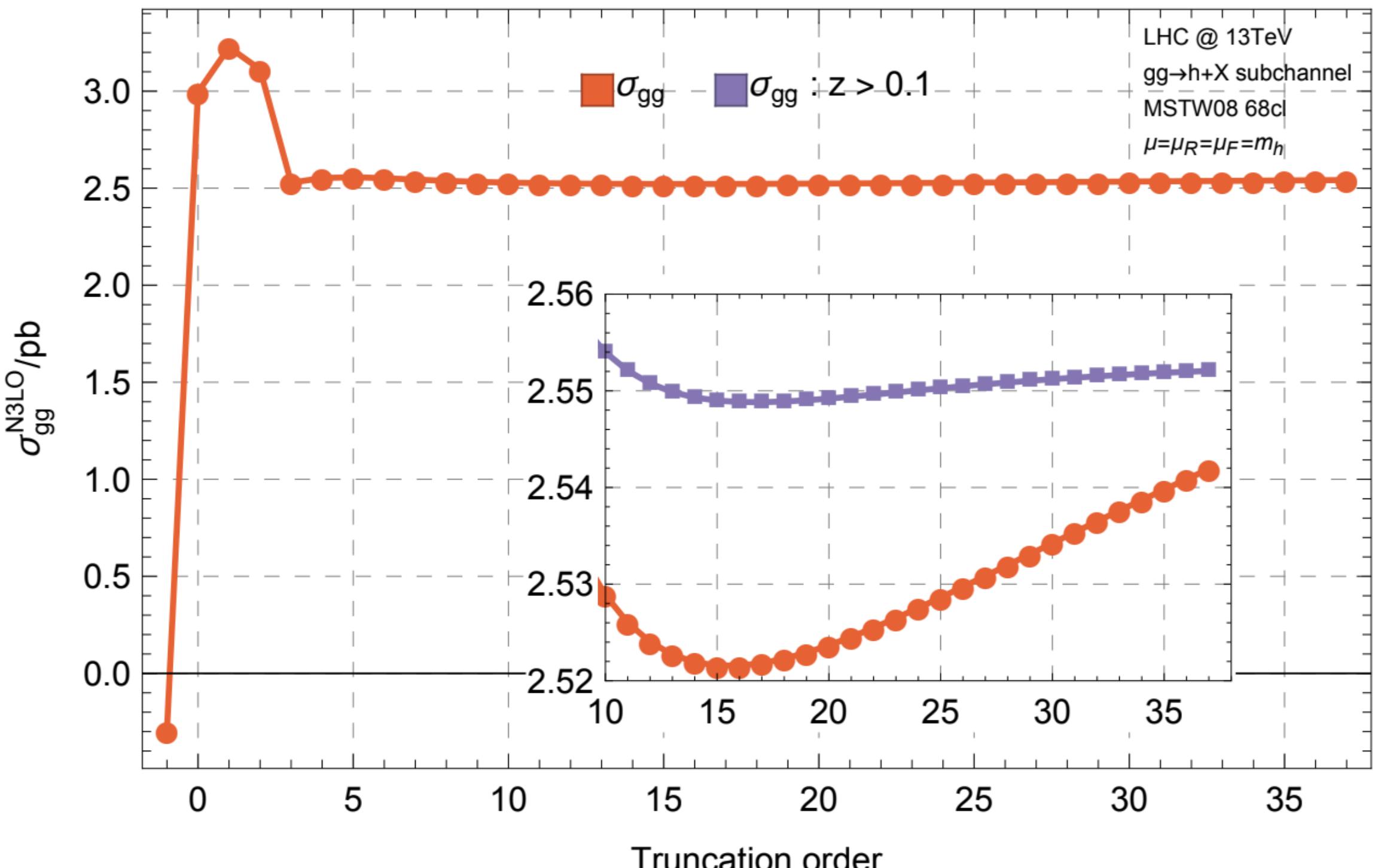
RH, Kilgore '02

NNLO



RH, Kilgore '02

N³LO



Anastasiou, Duhr, Dulat, Herzog, Mistlberger '15

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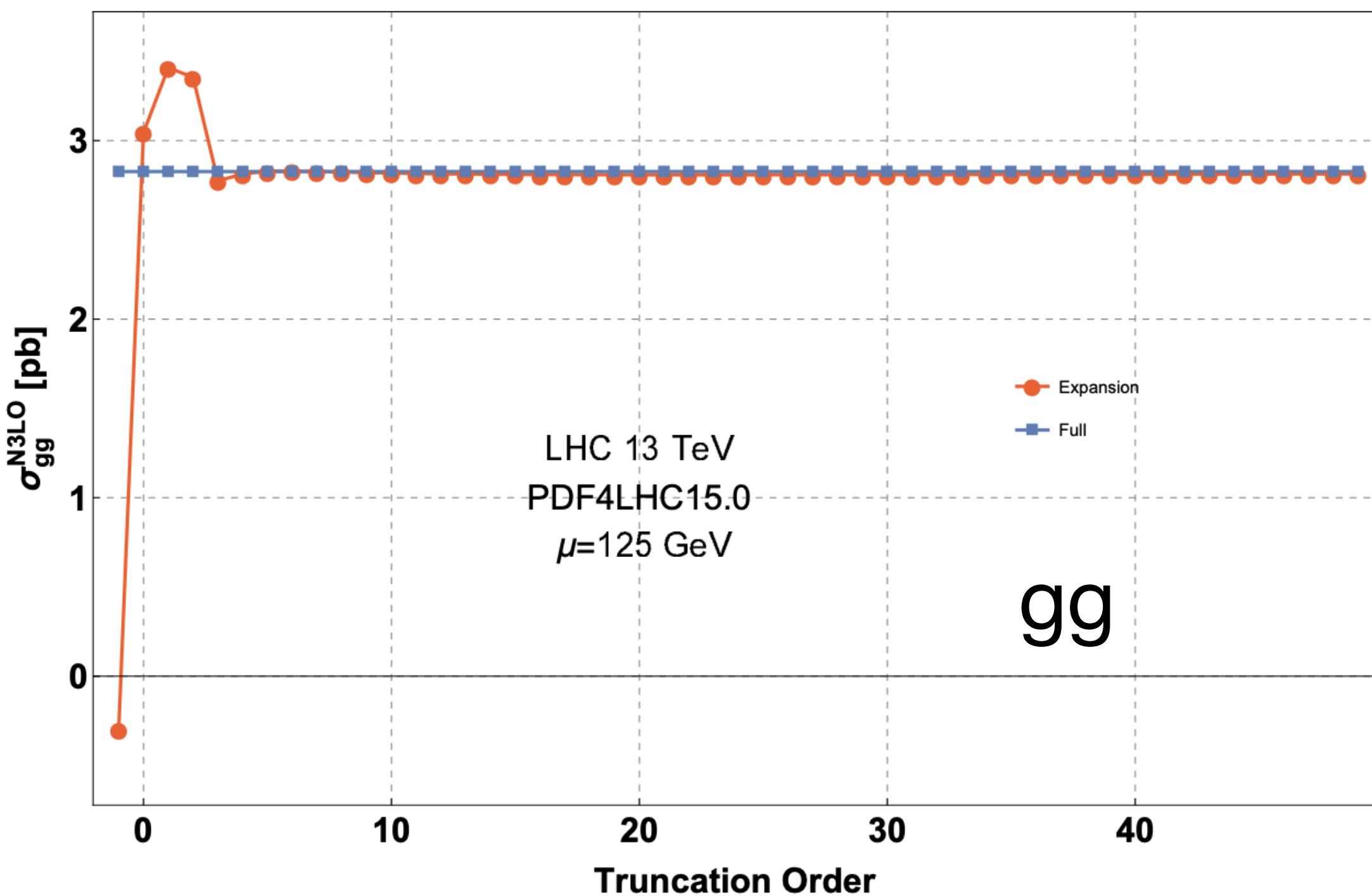
LHCHWG YR4 '16

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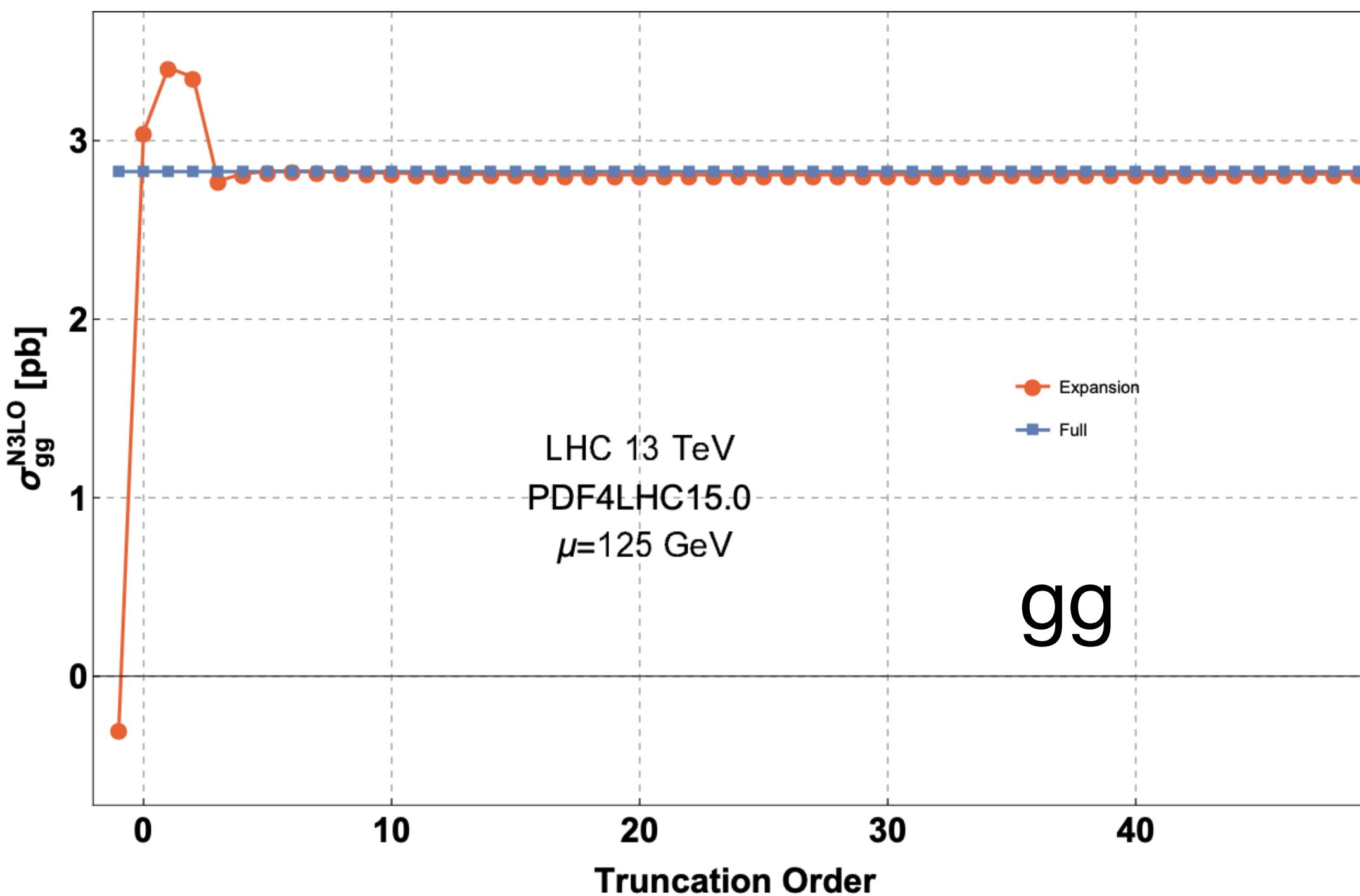
LHCHWG YR4 '16



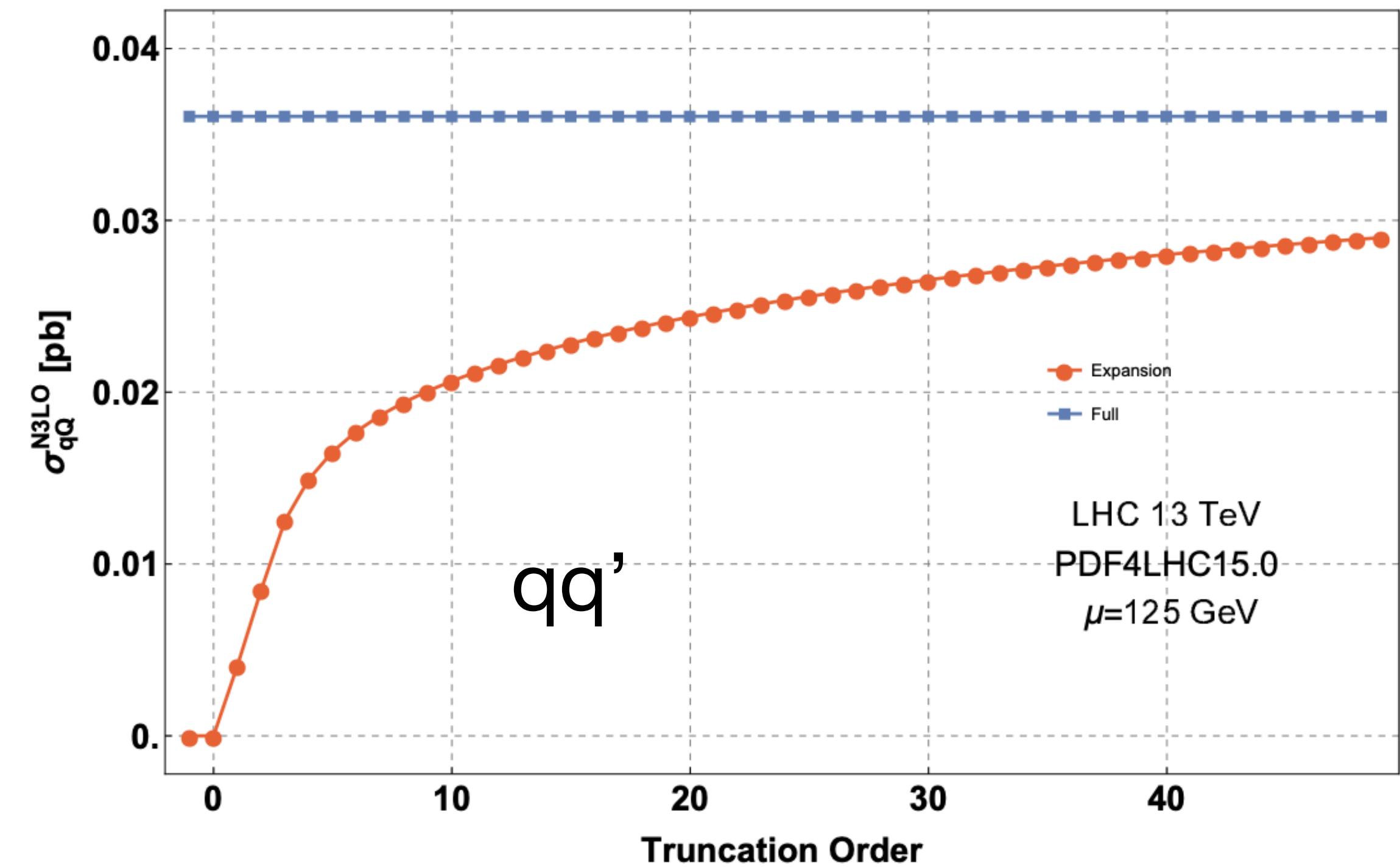
Mistlberger '18

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LHCHWG YR4 '16

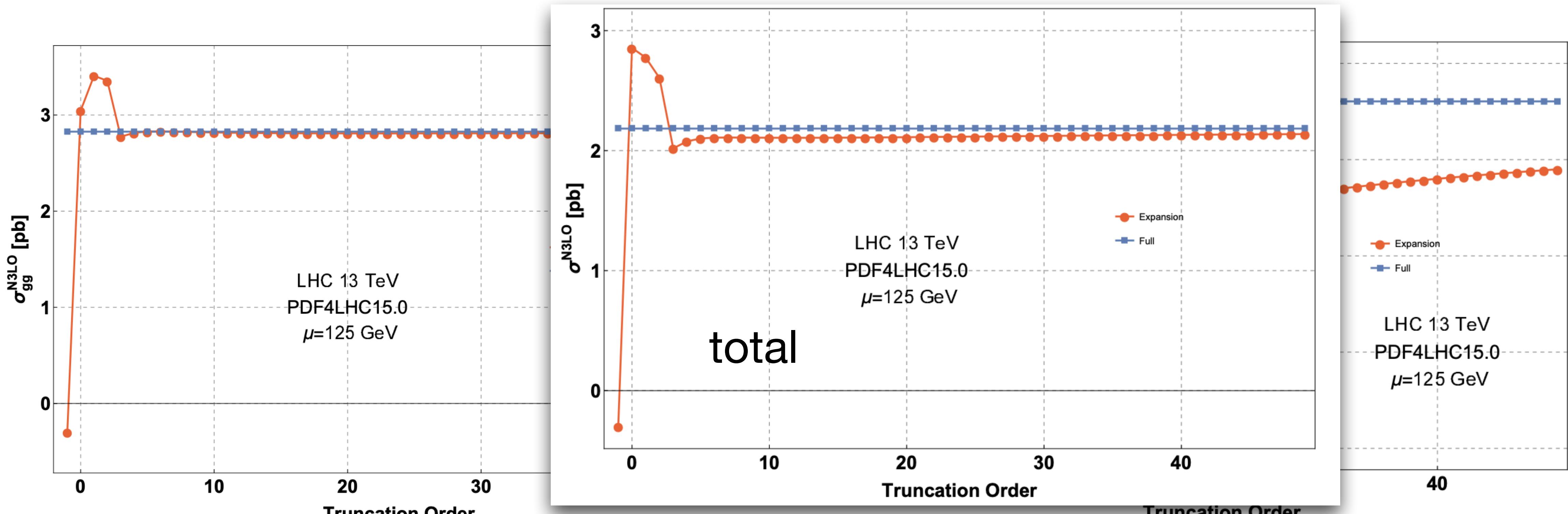


Mistlberger '18



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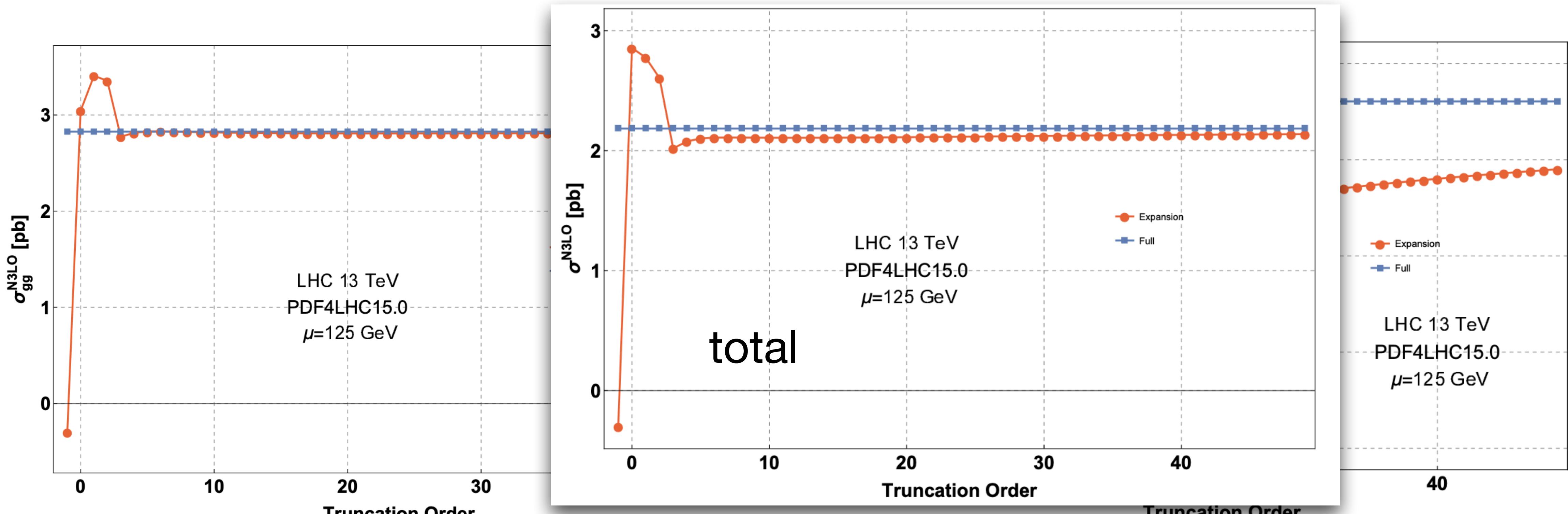
LHCHWG YR4 '16



Mistlberger '18

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

LHCHWG YR4 '16



$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$ 	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

LHCHWG YR4 '16

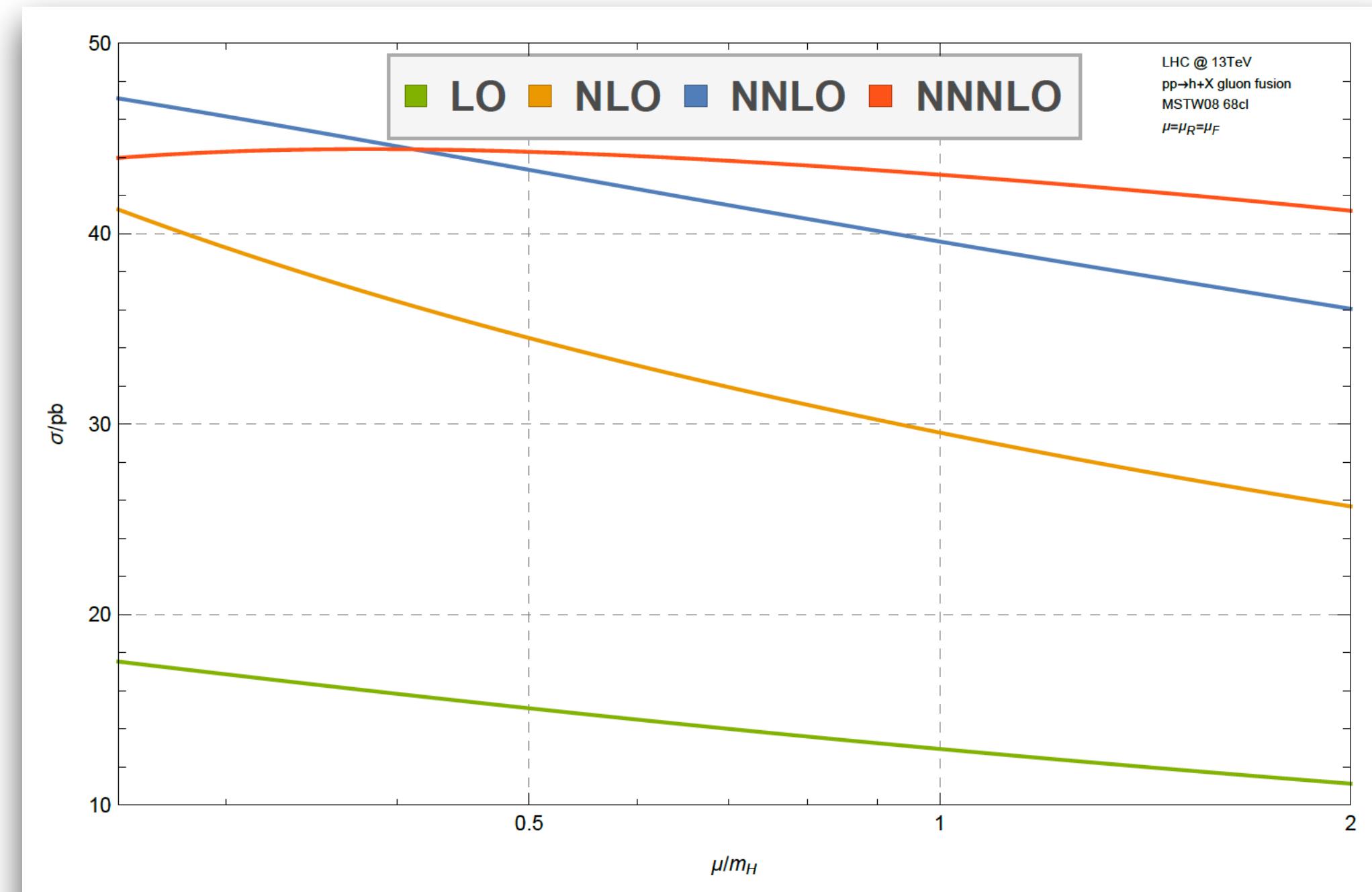
$$48.58 \text{ pb} = \begin{array}{lll} 16.00 \text{ pb} & (+32.9\%) & (\text{LO, rEFT}) \\ + 20.84 \text{ pb} & (+42.9\%) & (\text{NLO, rEFT}) \\ - 2.05 \text{ pb} & (-4.2\%) & ((t, b, c), \text{exact NLO}) \\ + 9.56 \text{ pb} & (+19.7\%) & (\text{NNLO, rEFT}) \\ + 0.34 \text{ pb} & (+0.7\%) & (\text{NNLO}, 1/m_t) \\ + 2.40 \text{ pb} & (+4.9\%) & (\text{EW, QCD-EW}) \\ + 1.49 \text{ pb} & (+3.1\%) & (\text{N}^3\text{LO, rEFT}) \end{array}$$

$K \approx 3$

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

LHCHWG YR4 '16

$$48.58 \text{ pb} = \begin{aligned} & 16.00 \text{ pb} \\ & + 20.84 \text{ pb} \\ & - 2.05 \text{ pb} \\ & + 9.56 \text{ pb} \\ & + 0.34 \text{ pb} \\ & + 2.40 \text{ pb} \\ & + 1.49 \text{ pb} \end{aligned}$$



$K \approx 3$

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.37\%$	$\pm 1.16\%$			

Approximate four-loop QCD corrections to the Higgs-boson production cross section

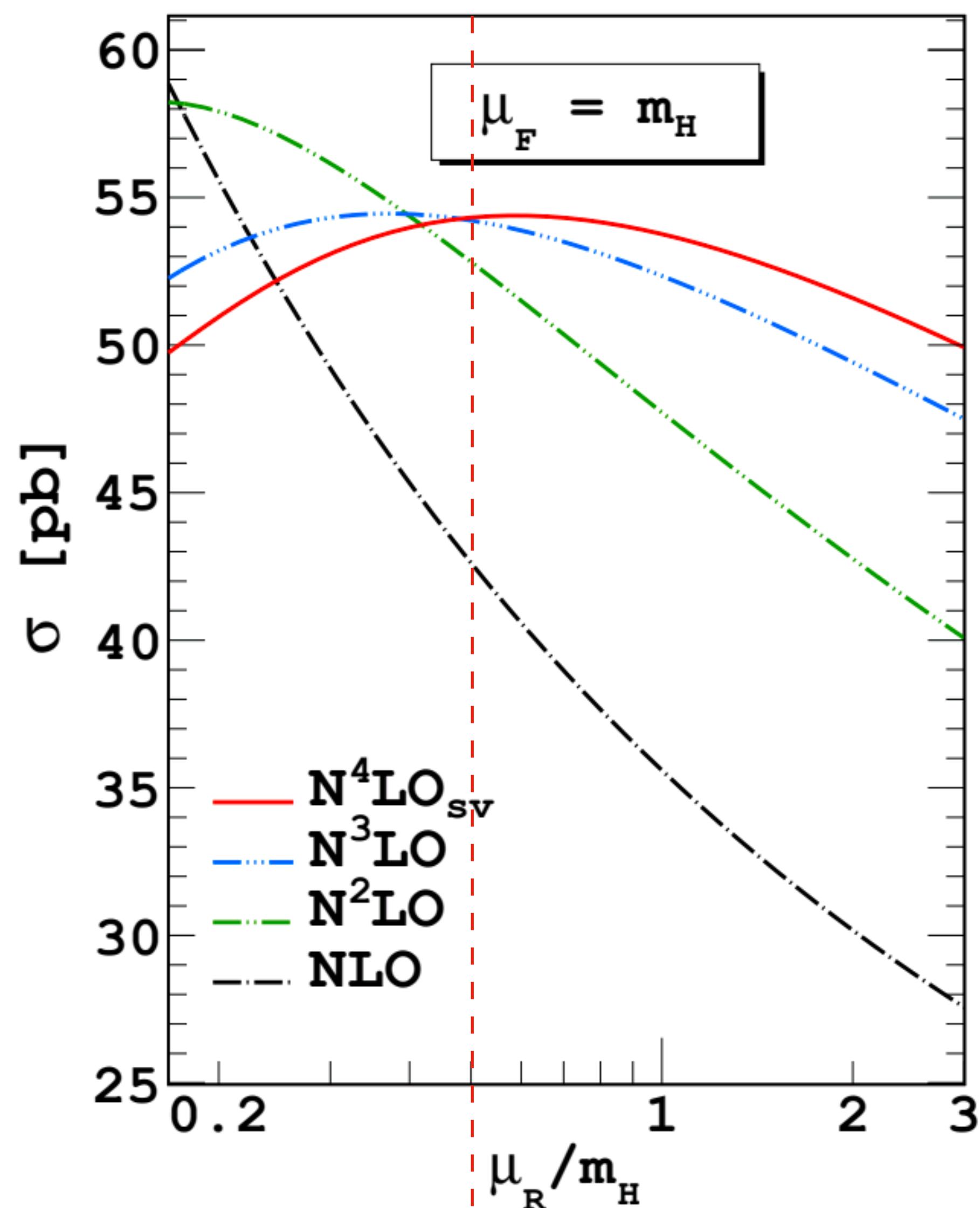
P3H-20-054

$$\hat{\sigma}_{ij} = \sum_{n \geq 0} \left(\frac{\alpha_s}{\pi} \right)^n \hat{\sigma}_{ij}^{(n)},$$

G. Das ^{*a}, S. Moch ^{†b} and A. Vogt ^{‡c}
soft-virtual approximation

$$\begin{aligned} \hat{\sigma}_{ij}^{(n)} &= a^{(n)} \delta(1-x) + \sum_{k=0}^{2n-1} b_k^{(n)} \left[\frac{\ln^k(1-x)}{1-x} \right]_+ \\ &+ \sum_{l=0}^{\infty} \sum_{k=0}^{2n-1} c_{lk}^{(n)} (1-x)^l \ln^k(1-x), \end{aligned}$$

$\delta(\text{scale})$ $\mu_R = m_H/2$



$\delta(\text{PDF-TH})$

$\delta(\text{EW})$

$\delta(t, b, c)$

$\delta(1/m_t)$

.56 pb

± 0.49 pb

± 0.40 pb

± 0.49 pb

Approximate four-loop QCD corrections to the Higgs-boson production cross section

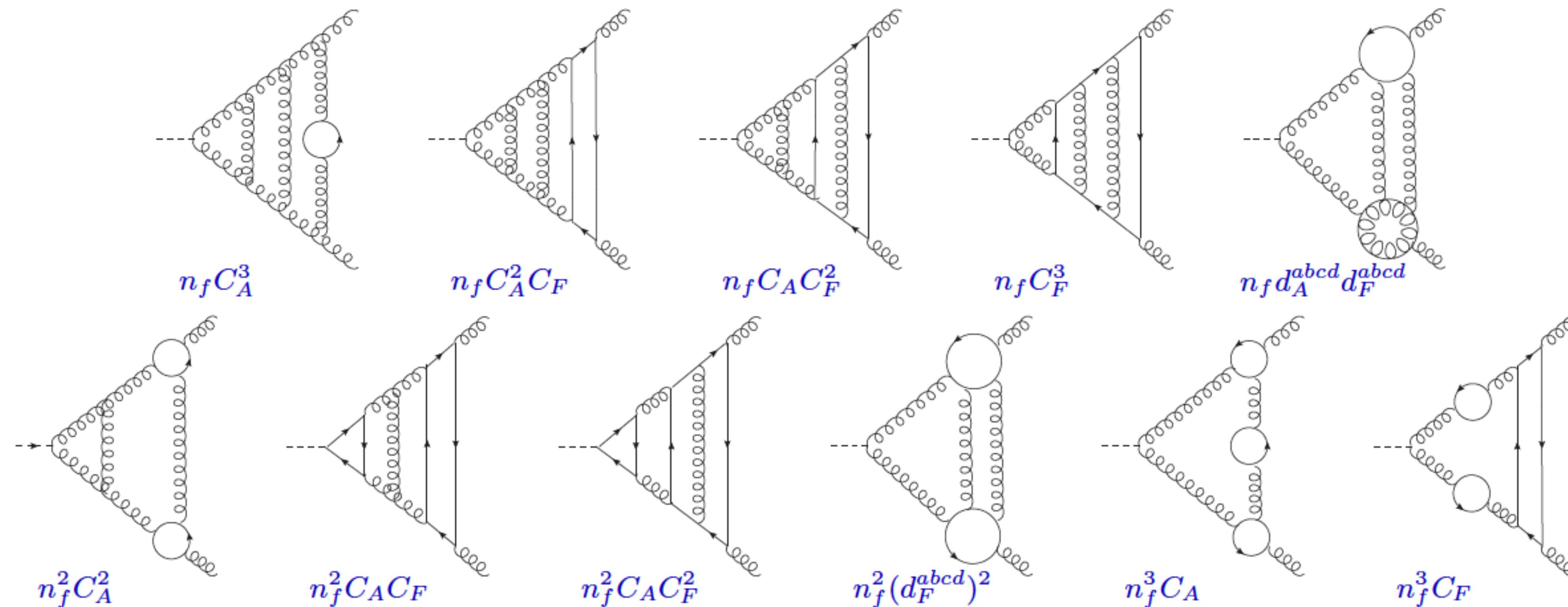
P3H-20-054

G. Das^{*a}, S. Moch^{†b} and A. Vogt^{‡c}
soft-virtual approximation

$$\frac{\ln^k(1-x)}{1-x} + k(1-x),$$

Fermionic corrections to quark and gluon form factors in four-loop QCD

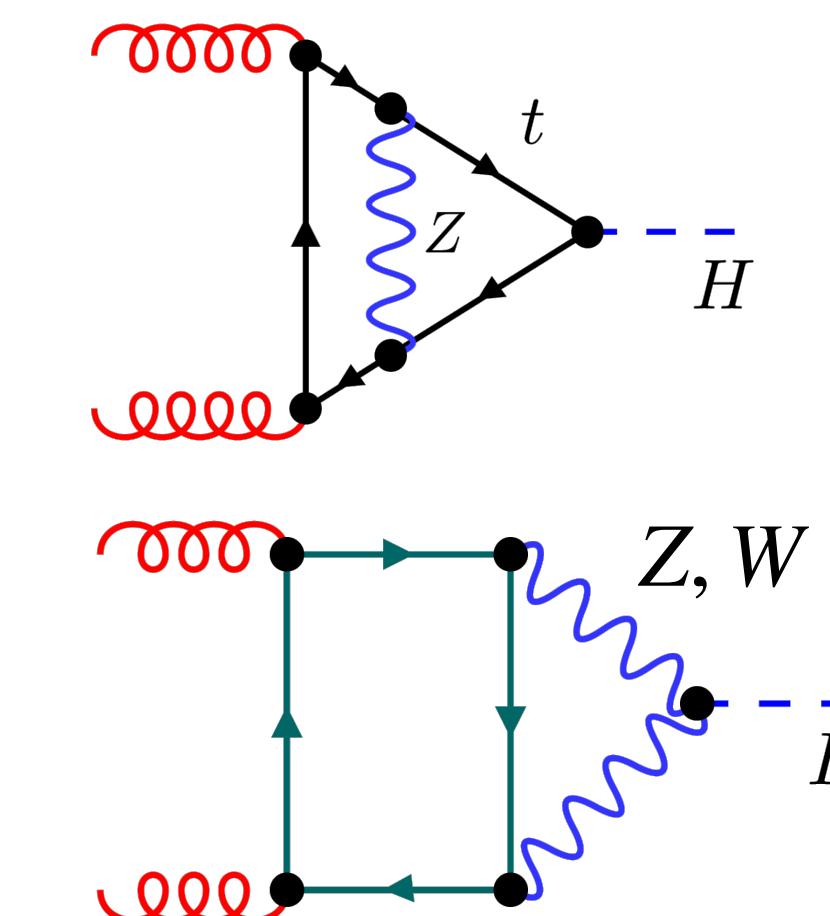
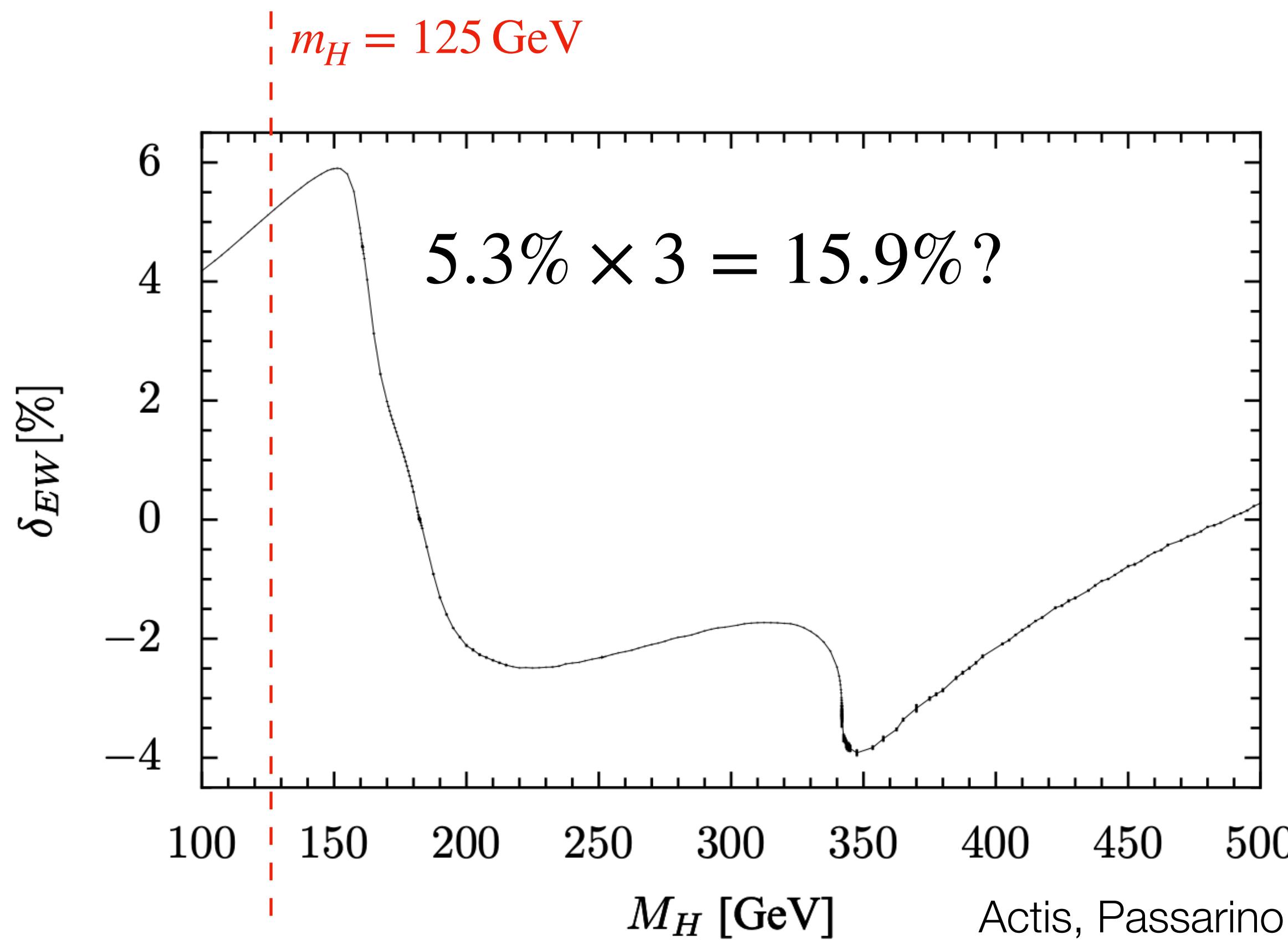
Roman N. Lee,¹ Andreas von Manteuffel,² Robert M. Schabinger,² P3H-21-013
Alexander V. Smirnov,³ Vladimir A. Smirnov,⁴ and Matthias Steinhauser⁵



$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
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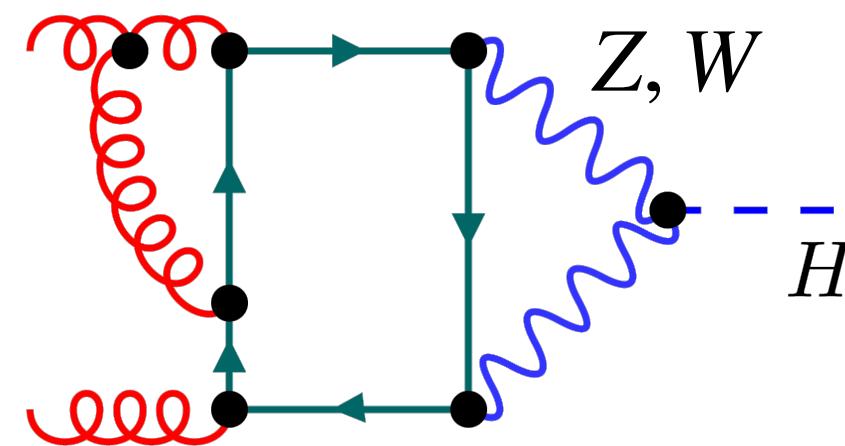
LHCHWG YR4 '16

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
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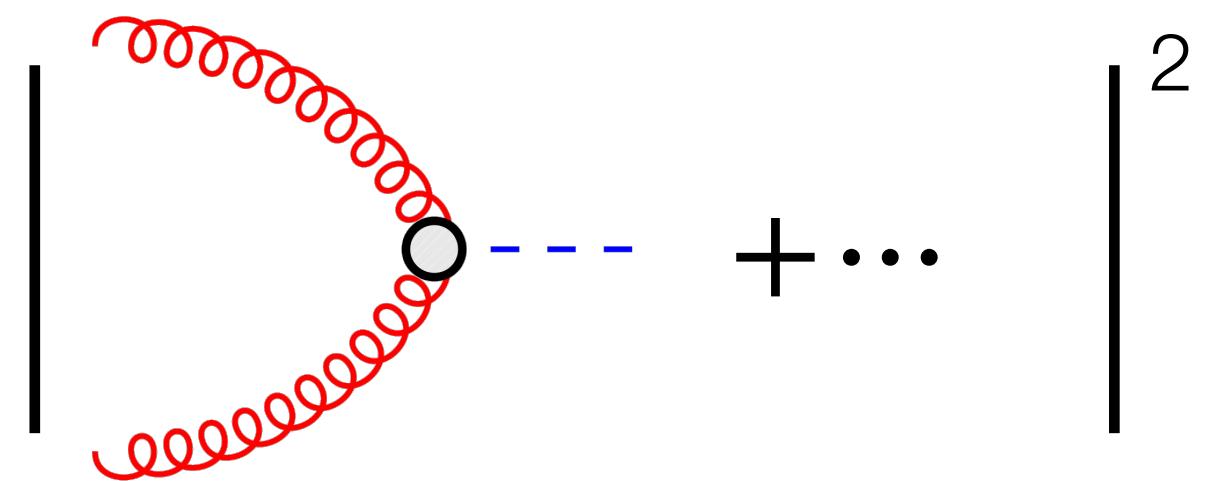
Actis, Passarino, Sturm, Uccirati '08

Mixed QCD/EW effects



Factorization hypothesis:

$$\sigma \sim C^2 \cdot$$



$$C \sim (1 + \lambda_{\text{EW}}) \left(1 + a_s C_{\text{QCD}}^{(1)} + \dots \right) = 1 + \lambda_{\text{EW}} + a_s C_{\text{QCD}}^{(1)} + \lambda_{\text{EW}} a_s C_{\text{QCD}}^{(1)} + \dots$$

\uparrow

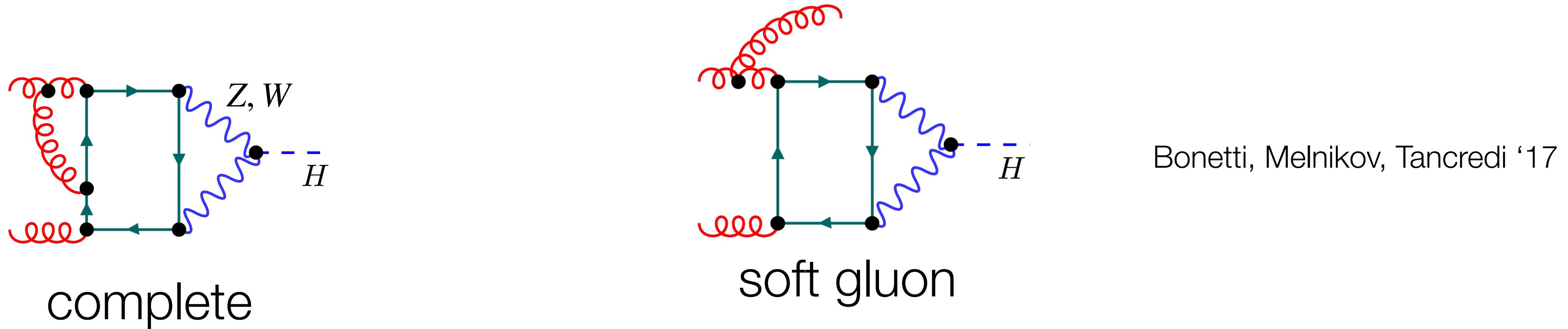
$$0.026 \qquad \qquad \qquad \frac{11}{4} = 2.75$$

$$= 1 + \lambda_{\text{EW}} + a_s C_{\text{QCD}}^{(1)} + \lambda_{\text{EW}} a_s \textcolor{red}{C}_w^{(1)} + \dots$$

$$M_W, M_Z \gg M_H : \quad \textcolor{red}{C}_w^{(1)} = \frac{7}{6} \approx 1.17 \quad \text{[Anastasiou, Boughezal, Petriello '09]}$$

varying $\textcolor{red}{C}_w^{(1)}$: 1% uncertainty estimate

Mixed QCD/EW effects



$$\lim_{E_4 \rightarrow 0} |A_{\text{NLO}}^{\text{real}}|^2 = \frac{\alpha_S}{4\pi} N_C \frac{2p_1 \cdot p_2}{p_1 \cdot p_4 p_2 \cdot p_4} |A_{\text{LO}}|^2 + \mathcal{O}(p_4^{-1}),$$

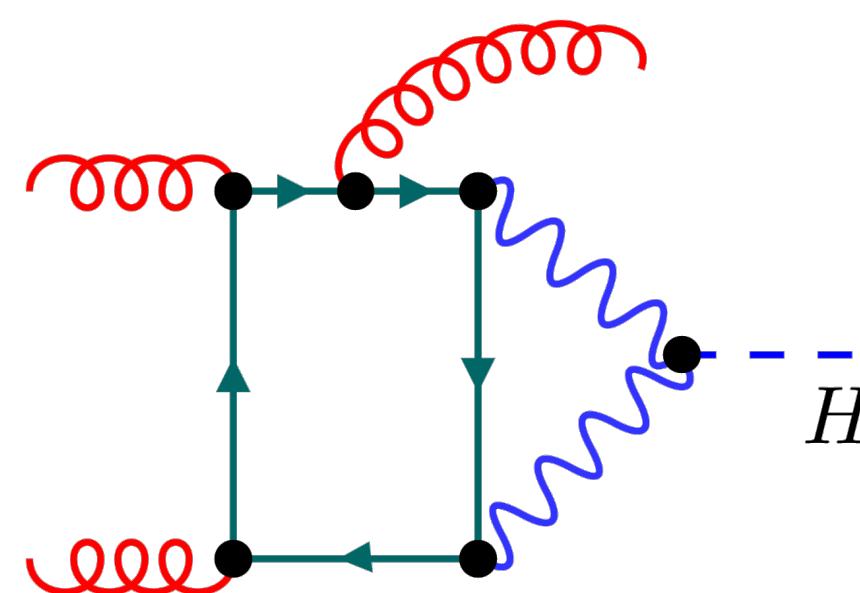
$$\begin{aligned}\sigma_{\text{LO}}^{\text{QCD}} &= 20.6 \text{ pb}, \\ \sigma_{\text{NLO}}^{\text{QCD}} &= 32.7 \text{ pb},\end{aligned}$$

$$\begin{aligned}\sigma_{\text{LO}}^{\text{QCD-EW}} &= 21.7 \text{ pb} &\Rightarrow &+5.3\% \text{ at LO} \\ \sigma_{\text{NLO}}^{\text{QCD-EW}} &= 34.4 \text{ pb} &\Rightarrow &+5.2\% \text{ at NLO}\end{aligned}$$

Anastasiou, Del Duca, Furlan, Mistlberger, Moriello, Schweitzer, Specchia '18

$m_W, m_Z \rightarrow 0$

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
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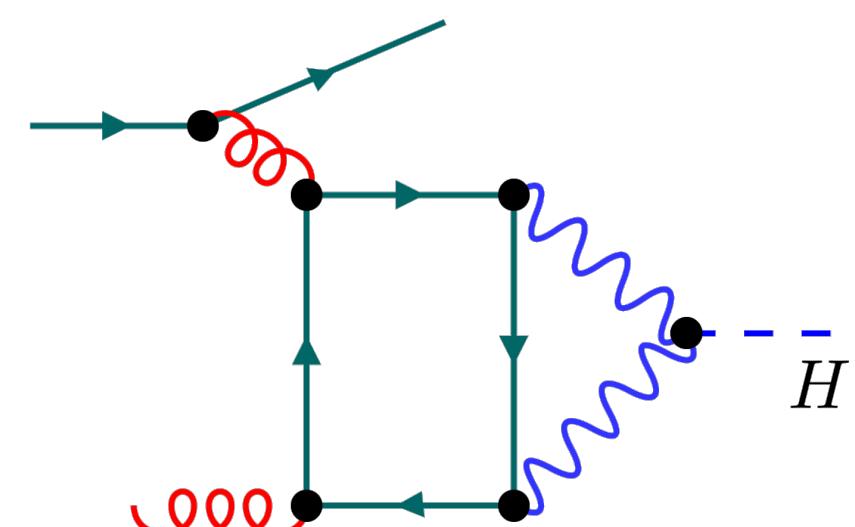
Two-loop mixed QCD-EW corrections to $gg \rightarrow Hg$

P3H-20-036

Marco Bonetti^{a,b} Erik Panzer^c Vladimir A. Smirnov^{d,e} Lorenzo Tancredi^f

^aInstitute for Theoretical Particle Physics, KIT,
Wolfgang-Gaede-Strasse 1, D-76128 Karlsruhe, Germany

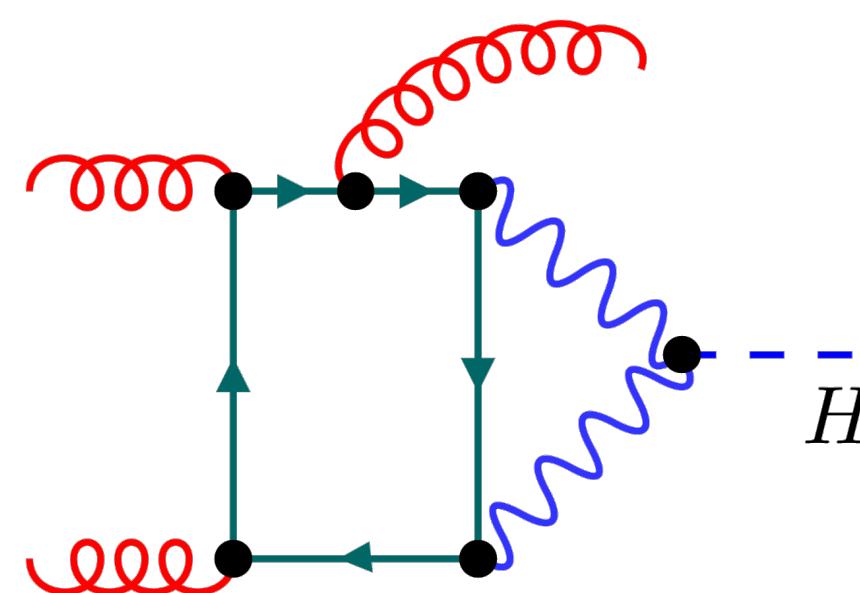
^bInstitute for Theoretical Particle Physics and Cosmology, RWTH Aachen University,
Sommerfeldstrasse 16, D-52056 Aachen, Germany



missing

also: Becchetti, Bonciani, Del Duca, Hirschi, Moriello '21
Hirschi, Lionetti, Schweitzer '19

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
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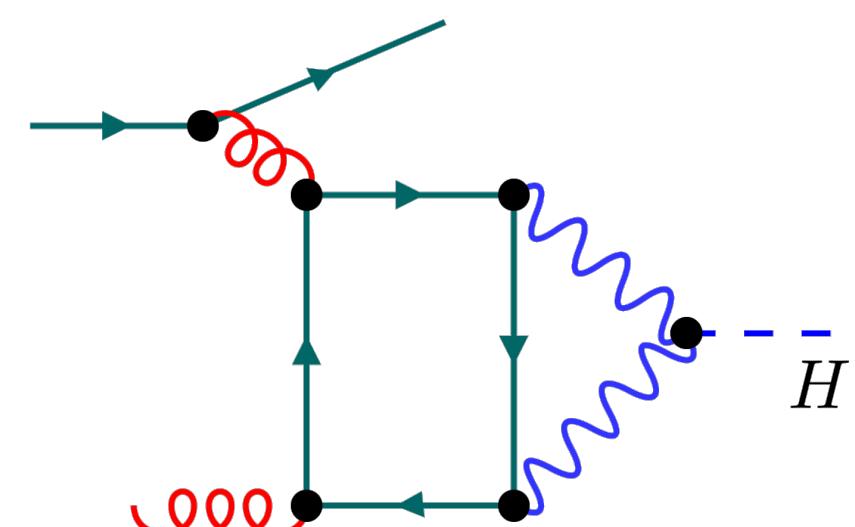
Two-loop mixed QCD-EW corrections to $gg \rightarrow Hg$

P3H-20-036

Marco Bonetti^{a,b} Erik Panzer^c Vladimir A. Smirnov^{d,e} Lorenzo Tancredi^f

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^bInstitute for Theoretical Particle Physics and Cosmology, RWTH Aachen University,
Sommerfeldstrasse 16, D-52056 Aachen, Germany



missing

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$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

LHCHWG YR4 '16

NNLO cross section with

NNLO PDFs

NLO PDFs

$$\frac{1}{2} \left| \frac{\sigma_{EFT}^{(2),NNLO} - \sigma_{EFT}^{(2),NLO}}{\sigma_{EFT}^{(2),NNLO}} \right| = \frac{1}{2} 2.31\% = 1.16\%$$

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
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LHCHWG YR4 '16

$$48.58 \text{ pb} = \begin{array}{ll} 16.00 \text{ pb} & (+32.9\%) \\ + 20.84 \text{ pb} & (+42.9\%) \\ - 2.05 \text{ pb} & (-4.2\%) \\ + 9.56 \text{ pb} & (+19.7\%) \\ + 0.34 \text{ pb} & (+0.7\%) \\ + 2.40 \text{ pb} & (+4.9\%) \\ + 1.49 \text{ pb} & (+3.1\%) \end{array} \begin{array}{l} (\text{LO, rEFT}) \\ (\text{NLO, rEFT}) \\ ((t, b, c), \text{exact NLO}) \\ (\text{NNLO, rEFT}) \\ (\text{NNLO}, 1/m_t) \\ (\text{EW, QCD-EW}) \\ (\text{N}^3\text{LO}, \text{rEFT}) \end{array}$$

$$\sigma_{\text{HEFT}}^{\text{HO}} = \sigma^{\text{LO}} \left(\frac{\sigma^{\text{HO}}}{\sigma^{\text{LO}}} \right)_{m_t \rightarrow \infty}$$

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

NLO:

	α_s^2	α_s^3	Δm_t			
gg	$16.30 + 19.64$	$+0.0345$	$+0.2\%$	$\rightarrow + 0.1\%$		
qg		$1.49 - 0.3696$	-24.8%	$\rightarrow - 1.0\%$		
qq		$0.02 + 0.0322$	$+161\%$	$\rightarrow + 0.1\%$		
total	$16.30 + 21.15$	-0.3029	$+1.4\%$	$\rightarrow - 0.8\%$		

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

NNLO:

	α_s^2	α_s^3	α_s^4	
gg	$16.30 + 19.64 + 8.76$			$0.2\% \rightarrow 0.04\%$
qg		$1.49 + 0.84$		$24.8\% \rightarrow 0.4\%$
qq		$0.02 + 0.10$		$161\% \rightarrow 0.3\%$
total	$16.30 + 21.15 + 9.79$			$3.5\% \rightarrow 0.74\%$

$$\begin{array}{c} \mathcal{L}_{\text{SM}} \\ \downarrow m_t \rightarrow \infty \\ \mathcal{L}_{\text{eff}} = C_1(\alpha_s) \frac{H}{v} \text{Tr } G_{\mu\nu} G^{\mu\nu} + \mathcal{L}_{\text{QCD}}^{n_f=5} \end{array}$$

$$\begin{array}{c}
 \mathcal{L}_{\text{SM}} \\
 \downarrow m_t \rightarrow \infty \\
 \mathcal{L}_{\text{eff}} = C_1(\alpha_s) \frac{H}{v} \text{Tr } G_{\mu\nu} G^{\mu\nu} + \mathcal{L}_{\text{QCD}}^{n_f=5}
 \end{array}$$

higher orders in $1/m_t$:
need higher dimensional operators

$$\begin{array}{c}
 \mathcal{L}_{\text{SM}} \\
 \downarrow m_t \rightarrow \infty \\
 \mathcal{L}_{\text{eff}} = C_1(\alpha_s) \frac{H}{v} \text{Tr } G_{\mu\nu} G^{\mu\nu} + \mathcal{L}_{\text{QCD}}^{n_f=5}
 \end{array}$$

higher orders in $1/m_t$:
need higher dimensional operators

$$C_2(\alpha_s) \langle \frac{H}{v} \text{Tr } G_{\mu\nu} G^{\mu\rho} G_\rho^\nu \rangle$$

dim 6 (0 + 1)

$$\mathcal{O}_{6;1}'' = \frac{1}{3} \text{Tr}(F_{12}F_{13}F_{23}) . \quad (\text{B.1})$$

dim 8 (1 + 1)

$$\mathcal{O}_{8;1}'' = \text{Tr}(D_1F_{23}D_4F_{23}F_{14}); \quad \mathcal{O}_{8;2}'' = \text{Tr}(D_1F_{23}D_1F_{24}F_{34}) . \quad (\text{B.2})$$

dim 10 (3 + 2)

$$\begin{aligned} \mathcal{O}_{10;1}'' &= \text{Tr}(D_{12}F_{34}D_{15}F_{34}F_{25}), \quad \mathcal{O}_{10;2}'' = \text{Tr}(D_{12}F_{34}D_5F_{34}D_1F_{25}), \quad \mathcal{O}_{10;3}'' = \text{Tr}(D_2F_{34}D_{15}F_{34}D_1F_{25}); \\ \mathcal{O}_{10;4}'' &= \text{Tr}(D_{12}F_{34}D_1F_{35}D_2F_{45}), \quad \mathcal{O}_{10;5}'' = \text{Tr}(D_{12}F_{34}D_{12}F_{35}F_{45}) . \end{aligned} \quad (\text{B.3})$$

dim 12 (6 + 4)

$$\begin{aligned} \mathcal{O}_{12;1}'' &= \text{Tr}(D_{123}F_{45}D_{126}F_{45}F_{36}), \quad \mathcal{O}_{12;2}'' = \text{Tr}(D_{123}F_{45}D_{16}F_{45}D_2F_{36}), \quad \mathcal{O}_{12;3}'' = \text{Tr}(D_{13}F_{45}D_{126}F_{45}D_2F_{36}), \\ \mathcal{O}_{12;4}'' &= \text{Tr}(D_{123}F_{45}D_6F_{45}D_{12}F_{36}), \quad \mathcal{O}_{12;5}'' = \text{Tr}(D_{13}F_{45}D_{26}F_{45}D_{12}F_{36}), \quad \mathcal{O}_{12;6}'' = \text{Tr}(D_3F_{45}D_{126}F_{45}D_{12}F_{36}); \\ \mathcal{O}_{12;7}'' &= \text{Tr}(D_{12}F_{45}D_{13}F_{46}D_{23}F_{56}), \quad \mathcal{O}_{12;8}'' = \text{Tr}(D_{12}F_{45}D_{123}F_{46}D_3F_{56}), \quad \mathcal{O}_{12;9}'' = \text{Tr}(D_{123}F_{45}D_{12}F_{46}D_3F_{56}), \\ \mathcal{O}_{12;10}'' &= \text{Tr}(D_{123}F_{45}D_{123}F_{46}F_{56}) . \end{aligned} \quad (\text{B.4})$$

dim 14 (10 + 5)

$$\begin{aligned} \mathcal{O}_{14;1}'' &= \text{Tr}(D_{1234}F_{56}D_{1237}F_{56}F_{47}), \quad \mathcal{O}_{14;2}'' = \text{Tr}(D_{1234}F_{56}D_{127}F_{56}D_3F_{47}), \quad \mathcal{O}_{14;3}'' = \text{Tr}(D_{124}F_{56}D_{1237}F_{56}D_3F_{47}), \\ \mathcal{O}_{14;4}'' &= \text{Tr}(D_{1234}F_{56}D_{17}F_{56}D_{23}F_{47}), \quad \mathcal{O}_{14;5}'' = \text{Tr}(D_{124}F_{56}D_{137}F_{56}D_{23}F_{47}), \quad \mathcal{O}_{14;6}'' = \text{Tr}(D_{14}F_{56}D_{1237}F_{56}D_{23}F_{47}) \\ \mathcal{O}_{14;7}'' &= \text{Tr}(D_{1234}F_{56}D_7F_{56}D_{123}F_{47}), \quad \mathcal{O}_{14;8}'' = \text{Tr}(D_{124}F_{56}D_{37}F_{56}D_{123}F_{47}), \quad \mathcal{O}_{14;9}'' = \text{Tr}(D_{14}F_{56}D_{237}F_{56}D_{123}F_{47}), \\ \mathcal{O}_{14;10}'' &= \text{Tr}(D_4F_{56}D_{1237}F_{56}D_{123}F_{47}); \\ \mathcal{O}_{14;11}'' &= \text{Tr}(D_{123}F_{56}D_{124}F_{57}D_{34}F_{67}), \quad \mathcal{O}_{14;12}'' = \text{Tr}(D_{1234}F_{56}D_{12}F_{57}D_{34}F_{67}), \quad \mathcal{O}_{14;13}'' = \text{Tr}(D_{123}F_{56}D_{1234}F_{57}D_4F_{67}), \\ \mathcal{O}_{14;14}'' &= \text{Tr}(D_{1234}F_{56}D_{123}F_{57}D_4F_{67}), \quad \mathcal{O}_{14;15}'' = \text{Tr}(D_{1234}F_{56}D_{1234}F_{57}F_{67}) . \end{aligned} \quad (\text{B.5})$$

dim 16 (15 + 7)

$$\begin{aligned} \mathcal{O}_{16;1}'' &= \text{Tr}(D_{12345}F_{67}D_{12348}F_{67}F_{58}), \quad \mathcal{O}_{16;2}'' = \text{Tr}(D_{12345}F_{67}D_{1238}F_{67}D_4F_{58}), \quad \mathcal{O}_{16;3}'' = \text{Tr}(D_{1235}F_{67}D_{12348}F_{67}D_4F_{58}), \\ \mathcal{O}_{16;4}'' &= \text{Tr}(D_{12345}F_{67}D_{128}F_{67}D_{34}F_{58}), \quad \mathcal{O}_{16;5}'' = \text{Tr}(D_{1235}F_{67}D_{1248}F_{67}D_{34}F_{58}), \quad \mathcal{O}_{16;6}'' = \text{Tr}(D_{125}F_{67}D_{12348}F_{67}D_{34}F_{58}) \\ \mathcal{O}_{16;7}'' &= \text{Tr}(D_{12345}F_{67}D_{18}F_{67}D_{234}F_{58}), \quad \mathcal{O}_{16;8}'' = \text{Tr}(D_{1235}F_{67}D_{148}F_{67}D_{234}F_{58}), \quad \mathcal{O}_{16;9}'' = \text{Tr}(D_{125}F_{67}D_{1348}F_{67}D_{234}F_{58}), \\ \mathcal{O}_{16;10}'' &= \text{Tr}(D_{15}F_{67}D_{12348}F_{67}D_{234}F_{58}), \quad \mathcal{O}_{16;11}'' = \text{Tr}(D_{12345}F_{67}D_8F_{67}D_{1234}F_{58}), \quad \mathcal{O}_{16;12}'' = \text{Tr}(D_{1235}F_{67}D_{48}F_{67}D_{1234}F_{58}), \\ \mathcal{O}_{16;13}'' &= \text{Tr}(D_{125}F_{67}D_{348}F_{67}D_{1234}F_{58}), \quad \mathcal{O}_{16;14}'' = \text{Tr}(D_{15}F_{67}D_{2348}F_{67}D_{1234}F_{58}), \quad \mathcal{O}_{16;15}'' = \text{Tr}(D_5F_{67}D_{12348}F_{67}D_{1234}F_{58}); \\ \mathcal{O}_{16;16}'' &= \text{Tr}(D_{1234}F_{67}D_{125}F_{68}D_{345}F_{78}), \quad \mathcal{O}_{16;17}'' = \text{Tr}(D_{123}F_{67}D_{12345}F_{68}D_{45}F_{78}), \quad \mathcal{O}_{16;18}'' = \text{Tr}(D_{1234}F_{67}D_{1235}F_{68}D_{45}F_{78}), \\ \mathcal{O}_{16;19}'' &= \text{Tr}(D_{12345}F_{67}D_{123}F_{68}D_{45}F_{78}), \quad \mathcal{O}_{16;20}'' = \text{Tr}(D_{1234}F_{67}D_{12345}F_{68}D_5F_{78}), \quad \mathcal{O}_{16;21}'' = \text{Tr}(D_{12345}F_{67}D_{1234}F_{68}D_5F_{78}), \\ \mathcal{O}_{16;22}'' &= \text{Tr}(D_{12345}F_{67}D_{12345}F_{68}F_{78}) . \end{aligned} \quad (\text{B.6})$$

Jin, Ren, Yang '21

dim 6 (0 + 1)

$$\mathcal{O}_{6;1}'' = \frac{1}{3} \text{Tr}(F_{12}F_{13}F_{23}). \quad (\text{B.1})$$

dim 8 (1 + 1)

$$\mathcal{O}_{8;1}'' = \text{Tr}(D_1F_{23}D_4F_{23}F_{14}); \quad \mathcal{O}_{8;2}'' = \text{Tr}(D_1F_{23}D_1F_{24}F_{34}). \quad (\text{B.2})$$

dim 10 (3 + 2)

$$\begin{aligned} \mathcal{O}_{10;1}'' &= \text{Tr}(D_{12}F_{34}D_{15}F_{34}F_{25}), \quad \mathcal{O}_{10;2}'' = \text{Tr}(D_{12}F_{34}D_5F_{34}D_1F_{25}), \quad \mathcal{O}_{10;3}'' = \text{Tr}(D_2F_{34}D_{15}F_{34}D_1F_{25}); \\ \mathcal{O}_{10;4}'' &= \text{Tr}(D_{12}F_{34}D_1F_{35}D_2F_{45}), \quad \mathcal{O}_{10;5}'' = \text{Tr}(D_{12}F_{34}D_{12}F_{35}F_{45}). \end{aligned} \quad (\text{B.3})$$

dim 12 (6 + 4)

$$\begin{aligned} \mathcal{O}_{12;1}'' &= \text{Tr}(D_{123}F_{45}D_{126}F_{45}F_{36}), \quad \mathcal{O}_{12;2}'' = \text{Tr}(D_{123}F_{45}D_{16}F_{45}D_2F_{36}), \quad \mathcal{O}_{12;3}'' = \text{Tr}(D_{13}F_{45}D_{126}F_{45}D_2F_{36}), \\ \mathcal{O}_{12;4}'' &= \text{Tr}(D_{123}F_{45}D_6F_{45}D_{12}F_{36}), \quad \mathcal{O}_{12;5}'' = \text{Tr}(D_{13}F_{45}D_{26}F_{45}D_{12}F_{36}), \quad \mathcal{O}_{12;6}'' = \text{Tr}(D_3F_{45}D_{126}F_{45}D_{12}F_{36}); \\ \mathcal{O}_{12;7}'' &= \text{Tr}(D_{12}F_{45}D_{13}F_{46}D_{23}F_{56}), \quad \mathcal{O}_{12;8}'' = \text{Tr}(D_{12}F_{45}D_{123}F_{46}D_3F_{56}), \quad \mathcal{O}_{12;9}'' = \text{Tr}(D_{123}F_{45}D_{12}F_{46}D_3F_{56}), \\ \mathcal{O}_{12;10}'' &= \text{Tr}(D_{123}F_{45}D_{123}F_{46}F_{56}). \end{aligned} \quad (\text{B.4})$$

dim 14 (10 + 5)

$$\begin{aligned} \mathcal{O}_{14;1}'' &= \text{Tr}(D_{1234}F_{56}D_{1237}F_{56}F_{47}), \quad \mathcal{O}_{14;2}'' = \text{Tr}(D_{1234}F_{56}D_{127}F_{56}D_3F_{47}), \quad \mathcal{O}_{14;3}'' = \text{Tr}(D_{124}F_{56}D_{1237}F_{56}D_3F_{47}), \\ \mathcal{O}_{14;4}'' &= \text{Tr}(D_{1234}F_{56}D_{17}F_{56}D_{23}F_{47}), \quad \mathcal{O}_{14;5}'' = \text{Tr}(D_{124}F_{56}D_{137}F_{56}D_{23}F_{47}), \quad \mathcal{O}_{14;6}'' = \text{Tr}(D_{14}F_{56}D_{1237}F_{56}D_{23}F_{47}) \\ \mathcal{O}_{14;7}'' &= \text{Tr}(D_{1234}F_{56}D_7F_{56}D_{123}F_{47}), \quad \mathcal{O}_{14;8}'' = \text{Tr}(D_{124}F_{56}D_{37}F_{56}D_{123}F_{47}), \quad \mathcal{O}_{14;9}'' = \text{Tr}(D_{14}F_{56}D_{237}F_{56}D_{123}F_{47}), \\ \mathcal{O}_{14;10}'' &= \text{Tr}(D_4F_{56}D_{1237}F_{56}D_{123}F_{47}); \\ \mathcal{O}_{14;11}'' &= \text{Tr}(D_{123}F_{56}D_{124}F_{57}D_{34}F_{67}), \quad \mathcal{O}_{14;12}'' = \text{Tr}(D_{1234}F_{56}D_{12}F_{57}D_{34}F_{67}), \quad \mathcal{O}_{14;13}'' = \text{Tr}(D_{123}F_{56}D_{1234}F_{57}D_4F_{67}), \\ \mathcal{O}_{14;14}'' &= \text{Tr}(D_{1234}F_{56}D_{123}F_{57}D_4F_{67}), \quad \mathcal{O}_{14;15}'' = \text{Tr}(D_{1234}F_{56}D_{1234}F_{57}F_{67}). \end{aligned} \quad (\text{B.5})$$

dim 16 (15 + 7)

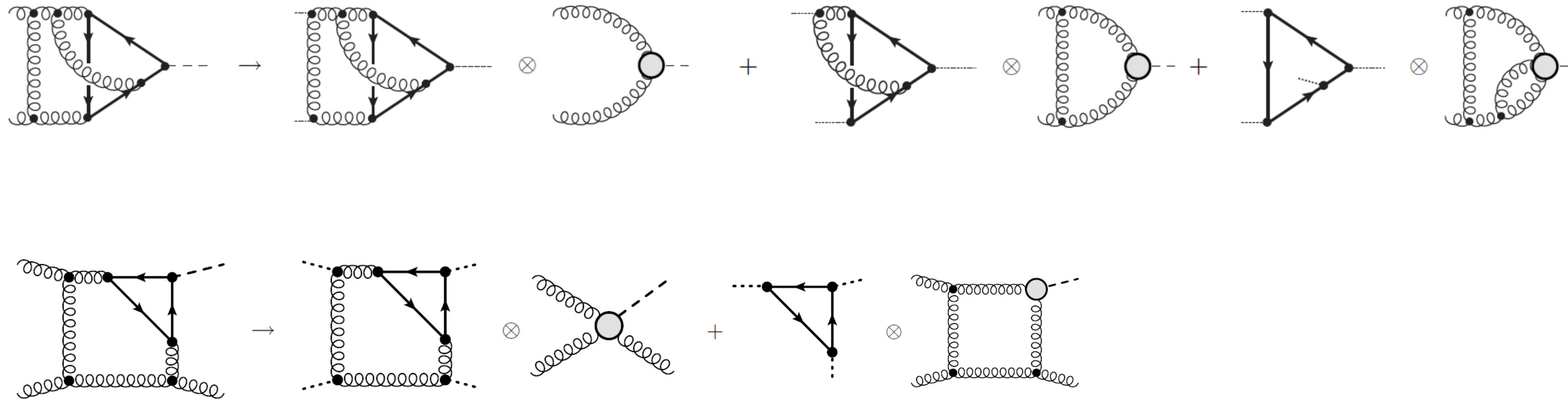
$$\begin{aligned} \mathcal{O}_{16;1}'' &= \text{Tr}(D_{12345}F_{67}D_{12348}F_{67}F_{58}), \quad \mathcal{O}_{16;2}'' = \text{Tr}(D_{12345}F_{67}D_{1238}F_{67}D_4F_{58}), \quad \mathcal{O}_{16;3}'' = \text{Tr}(D_{1235}F_{67}D_{12348}F_{67}D_4F_{58}), \\ \mathcal{O}_{16;4}'' &= \text{Tr}(D_{12345}F_{67}D_{128}F_{67}D_{34}F_{58}), \quad \mathcal{O}_{16;5}'' = \text{Tr}(D_{1235}F_{67}D_{1248}F_{67}D_{34}F_{58}), \quad \mathcal{O}_{16;6}'' = \text{Tr}(D_{125}F_{67}D_{12348}F_{67}D_{34}F_{58}) \\ \mathcal{O}_{16;7}'' &= \text{Tr}(D_{12345}F_{67}D_{18}F_{67}D_{234}F_{58}), \quad \mathcal{O}_{16;8}'' = \text{Tr}(D_{1235}F_{67}D_{148}F_{67}D_{234}F_{58}), \quad \mathcal{O}_{16;9}'' = \text{Tr}(D_{125}F_{67}D_{1348}F_{67}D_{234}F_{58}), \\ \mathcal{O}_{16;10}'' &= \text{Tr}(D_{15}F_{67}D_{12348}F_{67}D_{234}F_{58}), \quad \mathcal{O}_{16;11}'' = \text{Tr}(D_{12345}F_{67}D_8F_{67}D_{1234}F_{58}), \quad \mathcal{O}_{16;12}'' = \text{Tr}(D_{1235}F_{67}D_{48}F_{67}D_{1234}F_{58}), \\ \mathcal{O}_{16;13}'' &= \text{Tr}(D_{125}F_{67}D_{348}F_{67}D_{1234}F_{58}), \quad \mathcal{O}_{16;14}'' = \text{Tr}(D_{15}F_{67}D_{2348}F_{67}D_{1234}F_{58}), \quad \mathcal{O}_{16;15}'' = \text{Tr}(D_5F_{67}D_{12348}F_{67}D_{1234}F_{58}); \\ \mathcal{O}_{16;16}'' &= \text{Tr}(D_{1234}F_{67}D_{125}F_{68}D_{345}F_{78}), \quad \mathcal{O}_{16;17}'' = \text{Tr}(D_{123}F_{67}D_{12345}F_{68}D_{45}F_{78}), \quad \mathcal{O}_{16;18}'' = \text{Tr}(D_{1234}F_{67}D_{1235}F_{68}D_{45}F_{78}), \\ \mathcal{O}_{16;19}'' &= \text{Tr}(D_{12345}F_{67}D_{123}F_{68}D_{45}F_{78}), \quad \mathcal{O}_{16;20}'' = \text{Tr}(D_{1234}F_{67}D_{12345}F_{68}D_5F_{78}), \quad \mathcal{O}_{16;21}'' = \text{Tr}(D_{12345}F_{67}D_{1234}F_{68}D_5F_{78}), \\ \mathcal{O}_{16;22}'' &= \text{Tr}(D_{12345}F_{67}D_{12345}F_{68}F_{78}). \end{aligned} \quad (\text{B.6})$$

Instead: asymptotic expansions

Jin, Ren, Yang '21

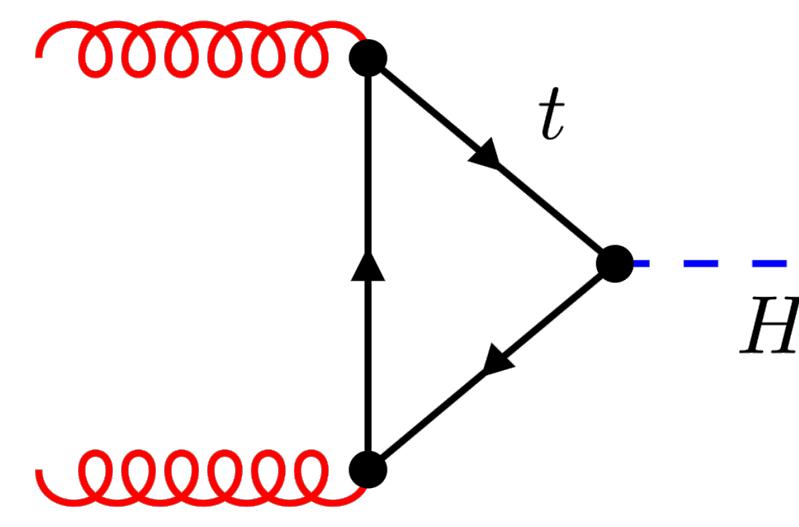
Asymptotic expansions

Calculating Wilson coefficient and matrix element “at once”:

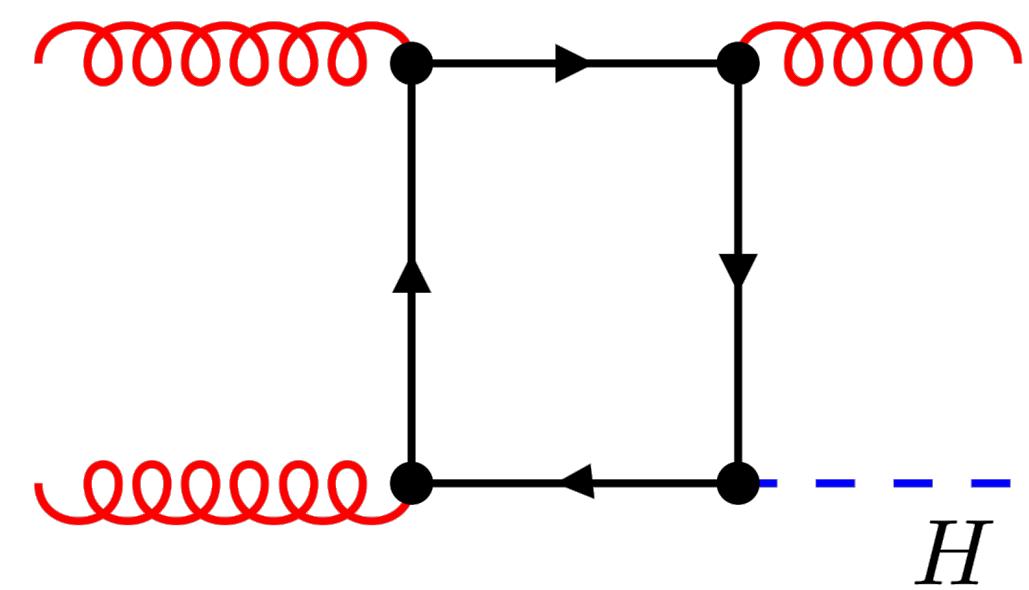


RH, Ozeren '09
Pak, Regal, Steinhauser '09

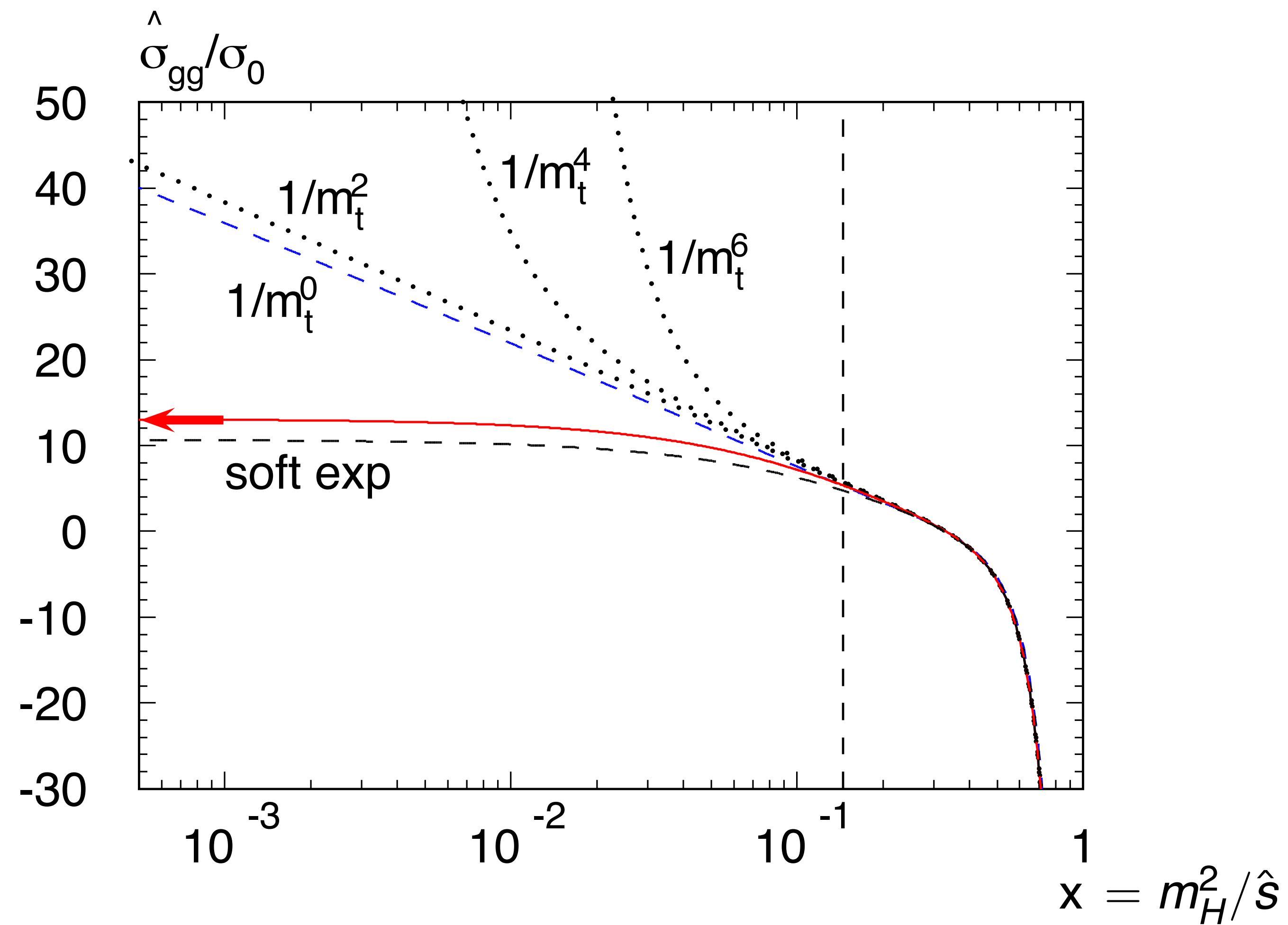
problem:



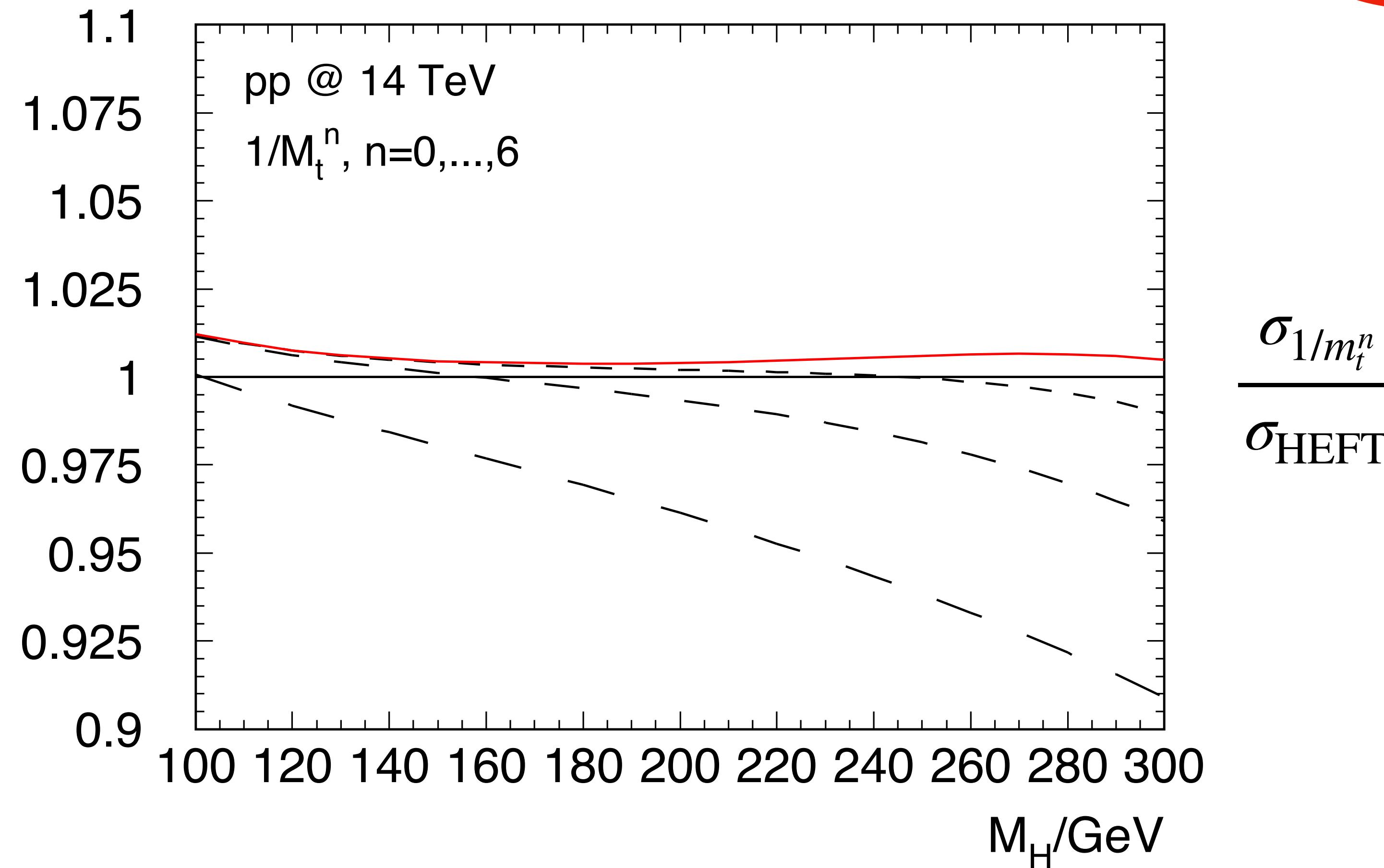
$$\hat{s} = m_H^2 < 4m_t^2$$



$$m_H^2 < \hat{s} < s < 13 \text{ TeV}$$



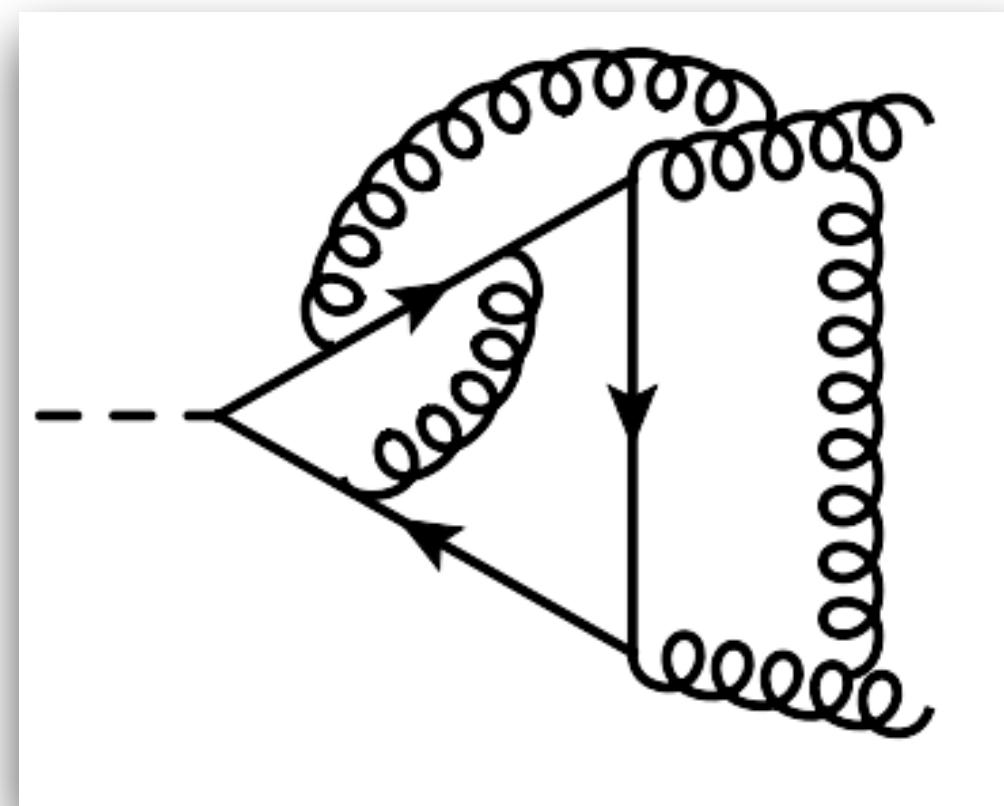
$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$



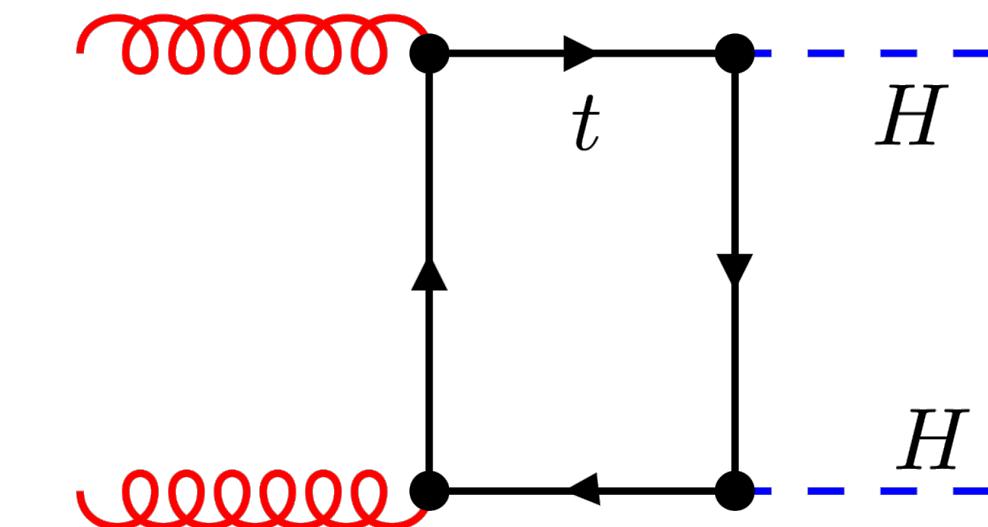
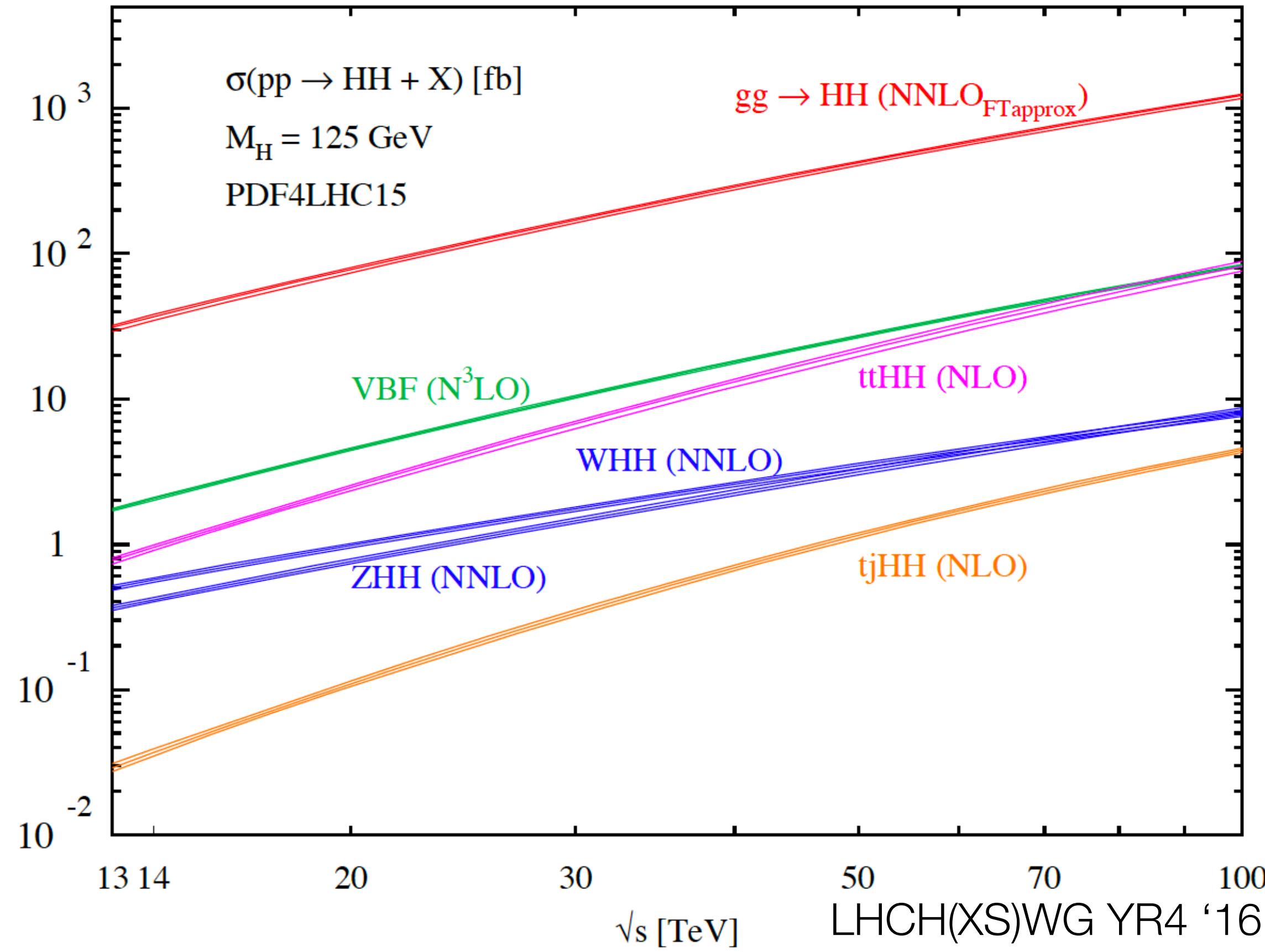
Top Quark Mass Effects in Next-To-Next-To-Next-To-Leading Order Higgs Boson Production: Virtual Corrections

P3H-19-045

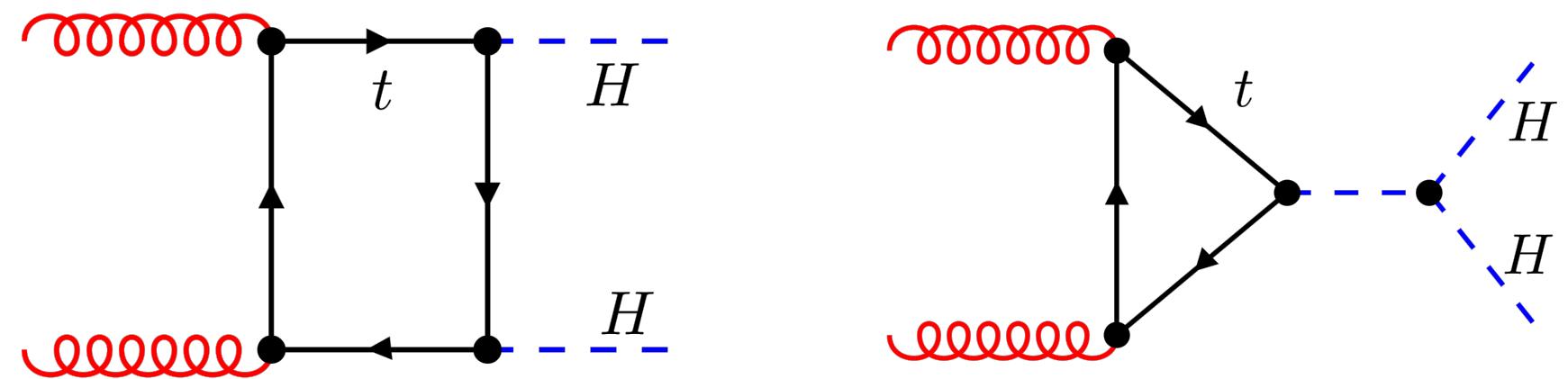
Joshua Davies,^{1,*} Florian Herren,^{1,†} and Matthias Steinhauser^{1,‡}



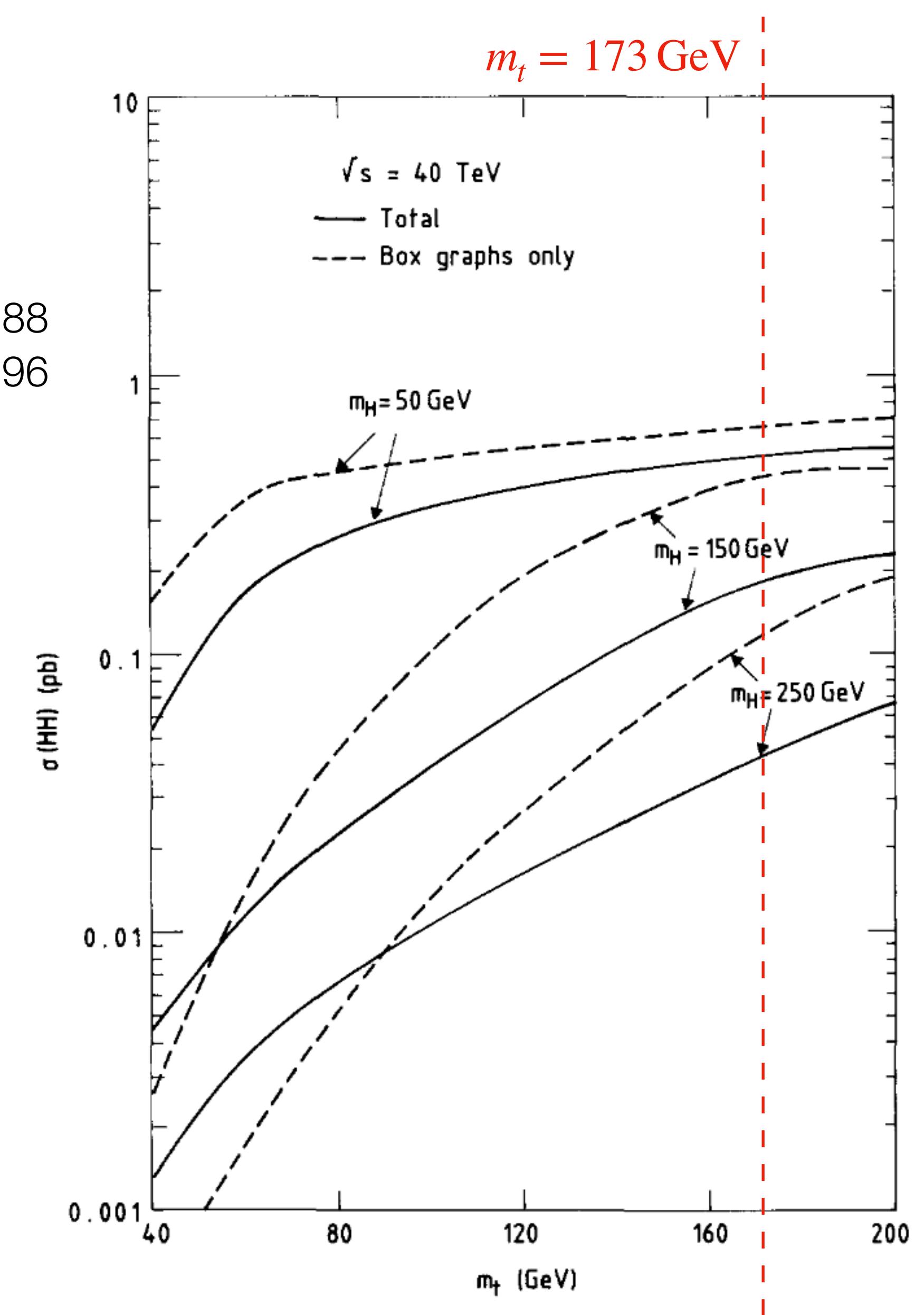
Double-Higgs production

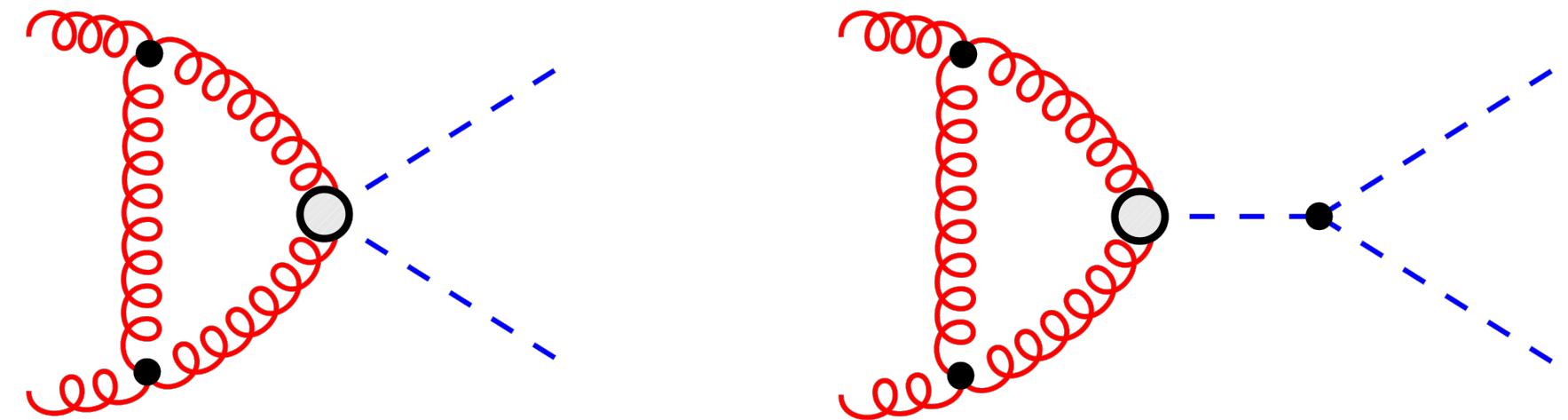


Double-Higgs production

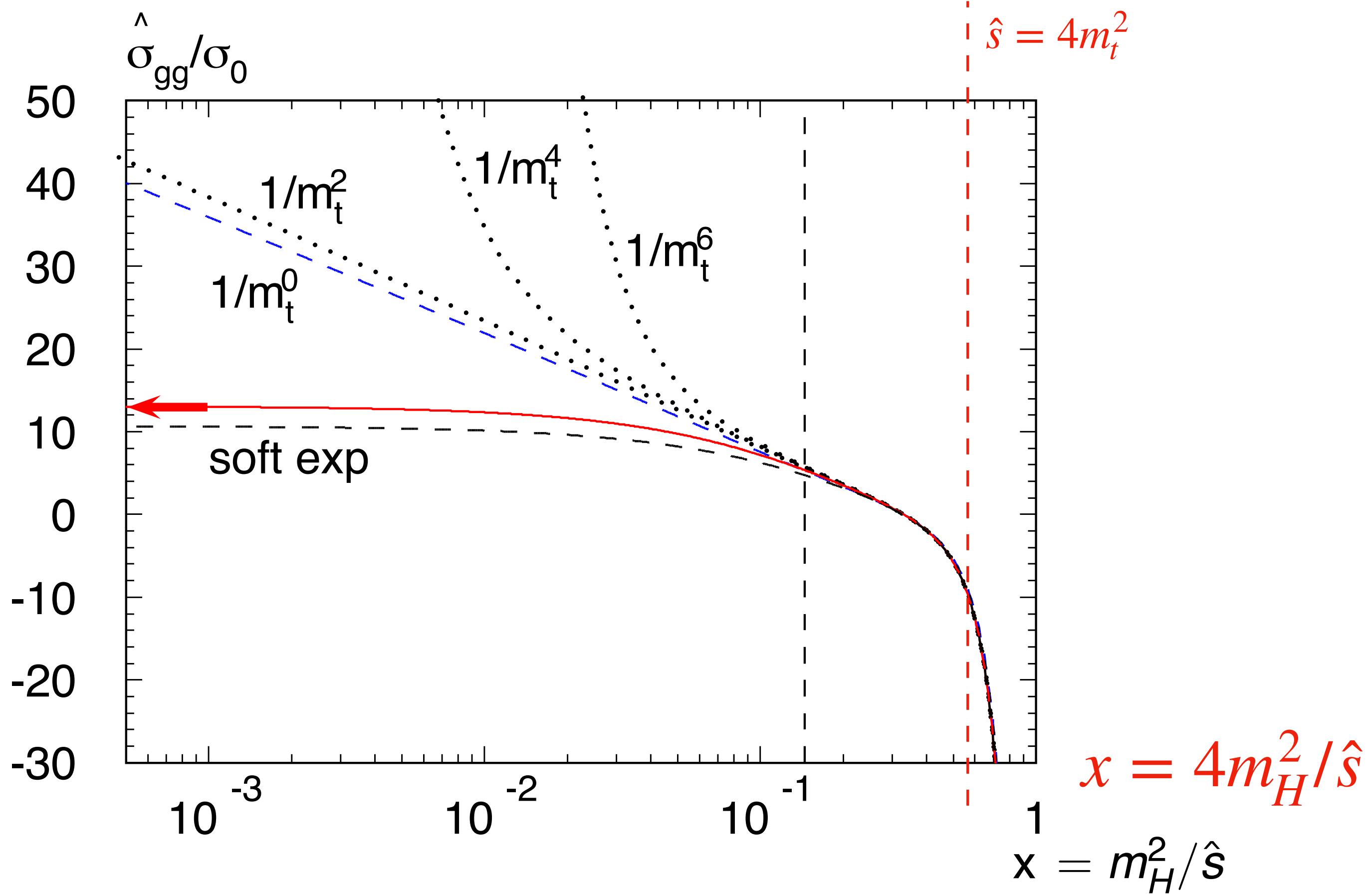


v.d.Bij, Glover '88
Plehn, Spira, Zerwas '96



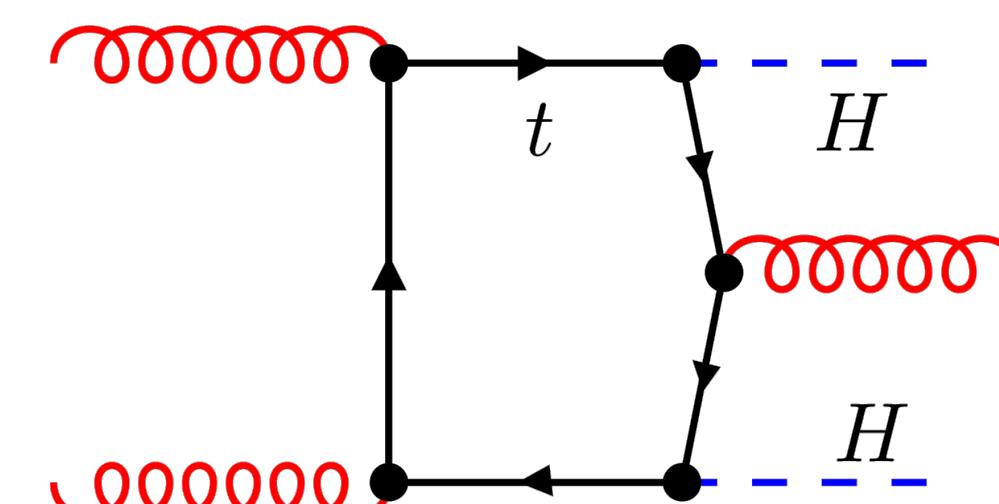


Dawson, Dittmaier, Spira '98



$\text{FT}_{\text{approx}}$ real radiation exact

Maltoni, Vryonidou, Zaro '14



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D. de Florian and J. Mazzitelli, *Higgs Boson Pair Production at Next-to-Next-to-Leading Order in QCD*, *Phys. Rev. Lett.* **111** (2013) 201801, [[arXiv:1309.6594](#)].

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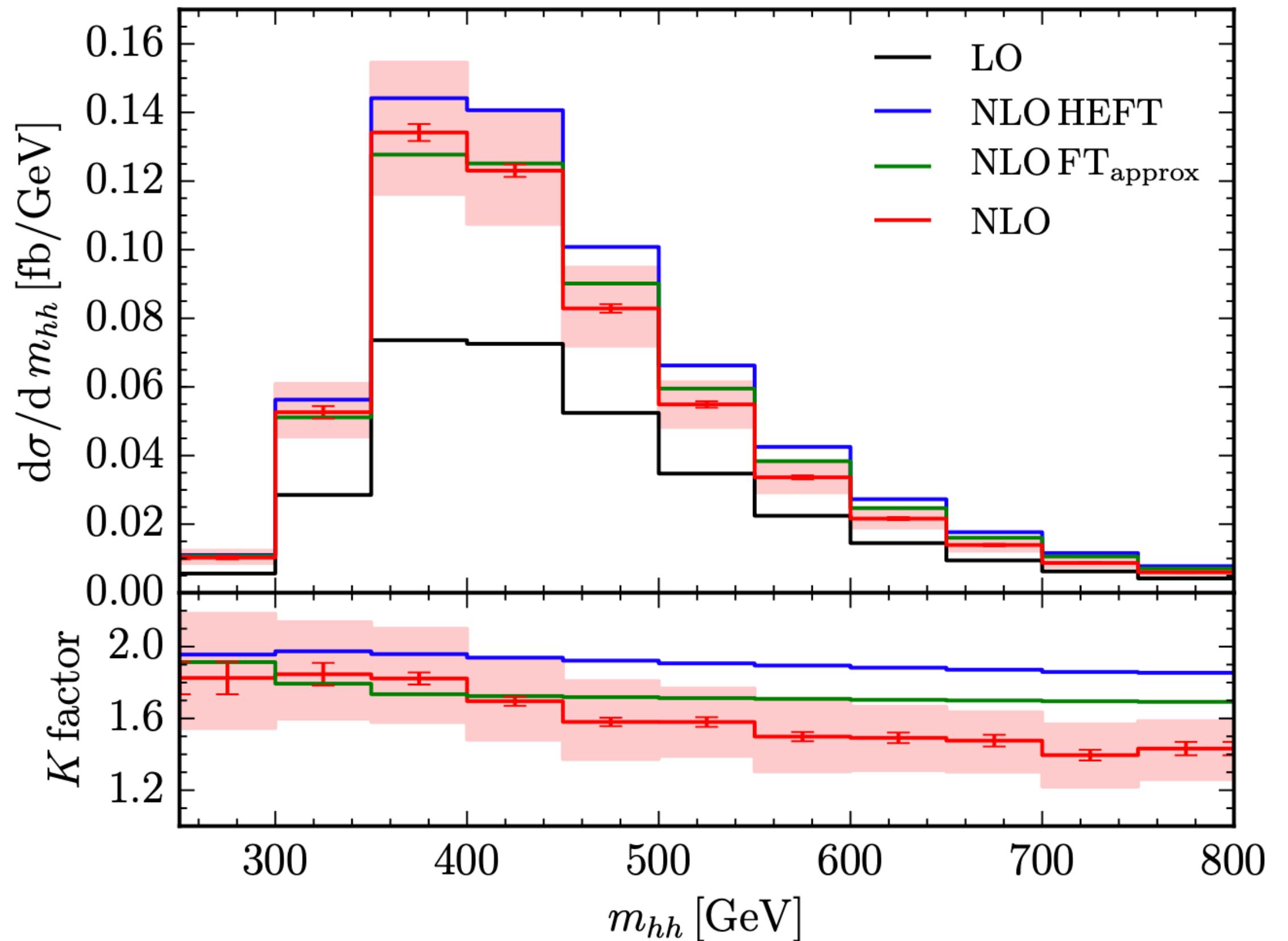
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D. de Florian, M. Grazzini, C. Hanga, S. Kallweit, J. M. Lindert, P. Maierhfer, J. Mazzitelli, and D. Rathlev, *Differential Higgs Boson Pair Production at Next-to-Next-to-Leading Order in QCD*, [arXiv:1606.09519](#).

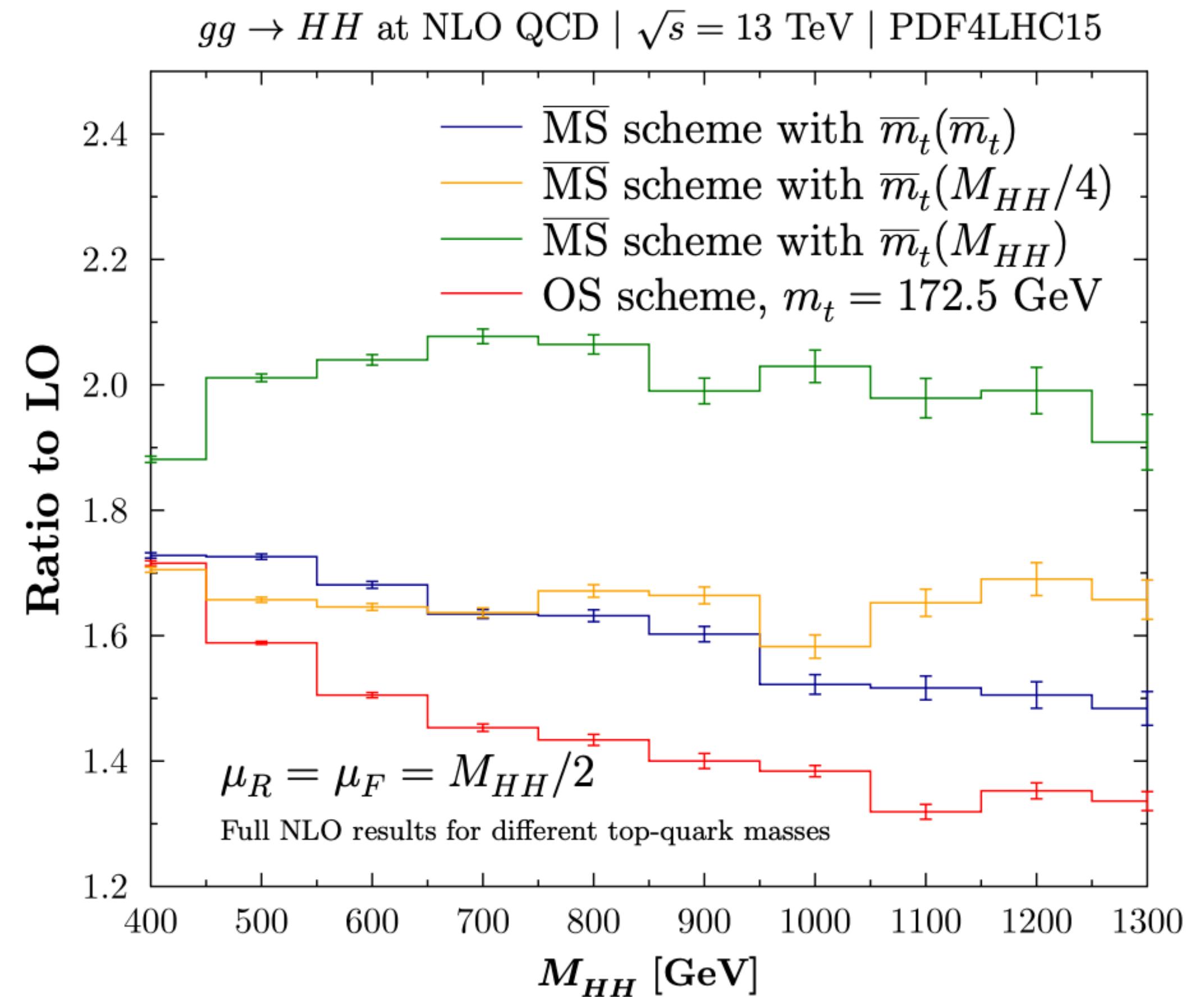
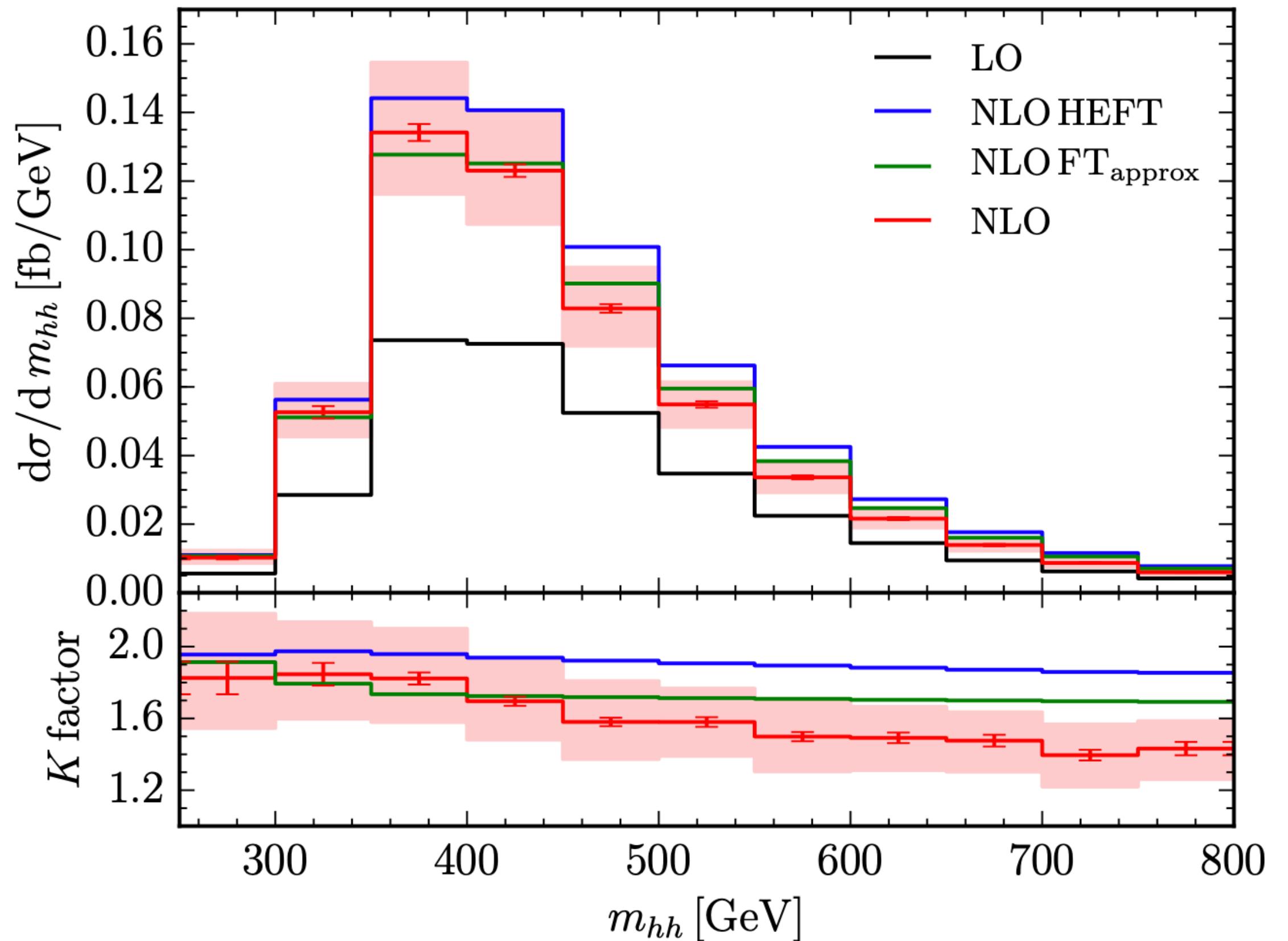
L.-B. Chen, H. T. Li, H.-S. Shao and J. Wang, *Higgs boson pair production via gluon fusion at N^3LO in QCD*, *Phys. Lett.* **B803** (2020) 135292 [[1909.06808](#)].

HEFT and
sub-leading terms
in $1/m_t$
through NLO

HEFT through N^3LO



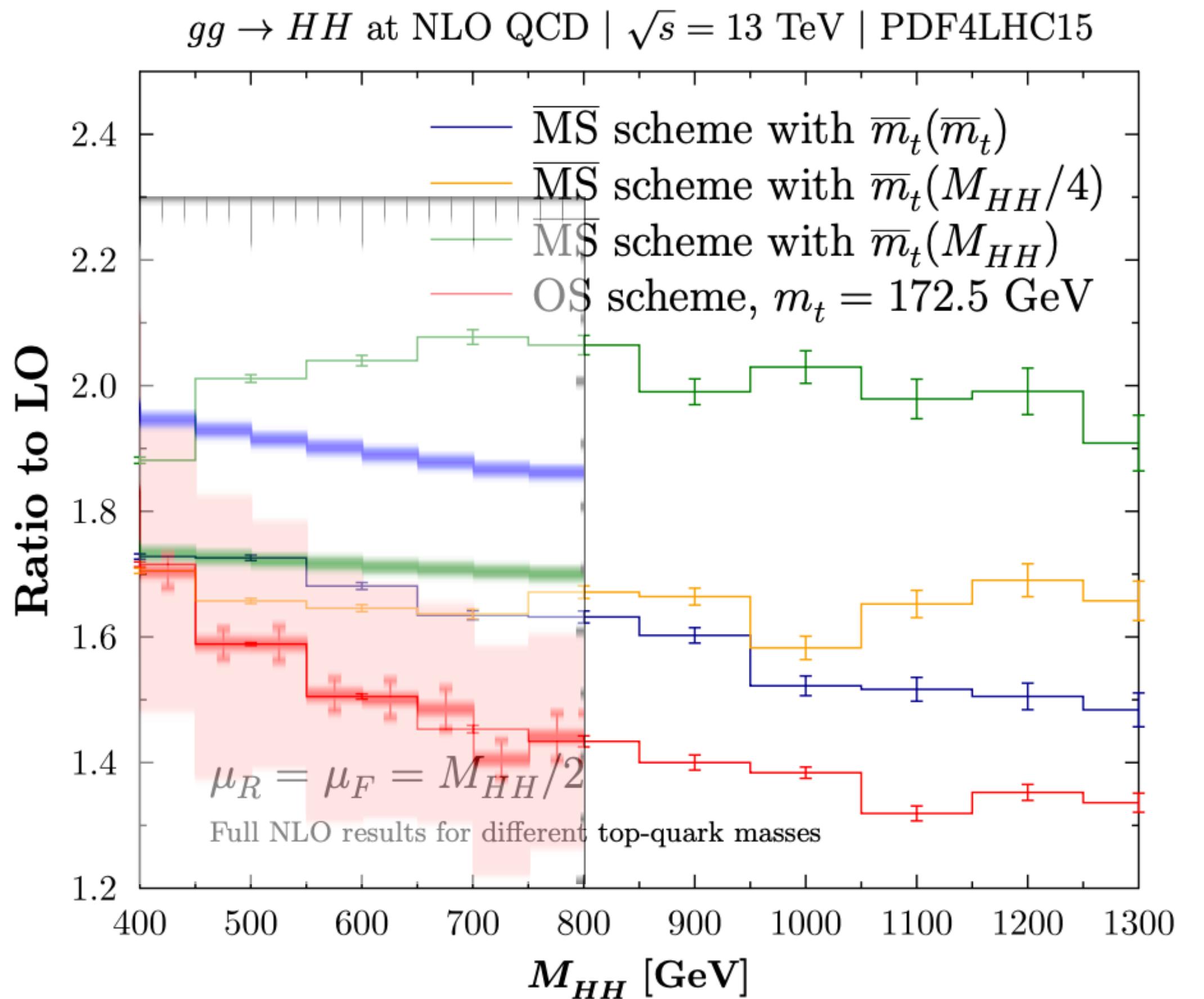
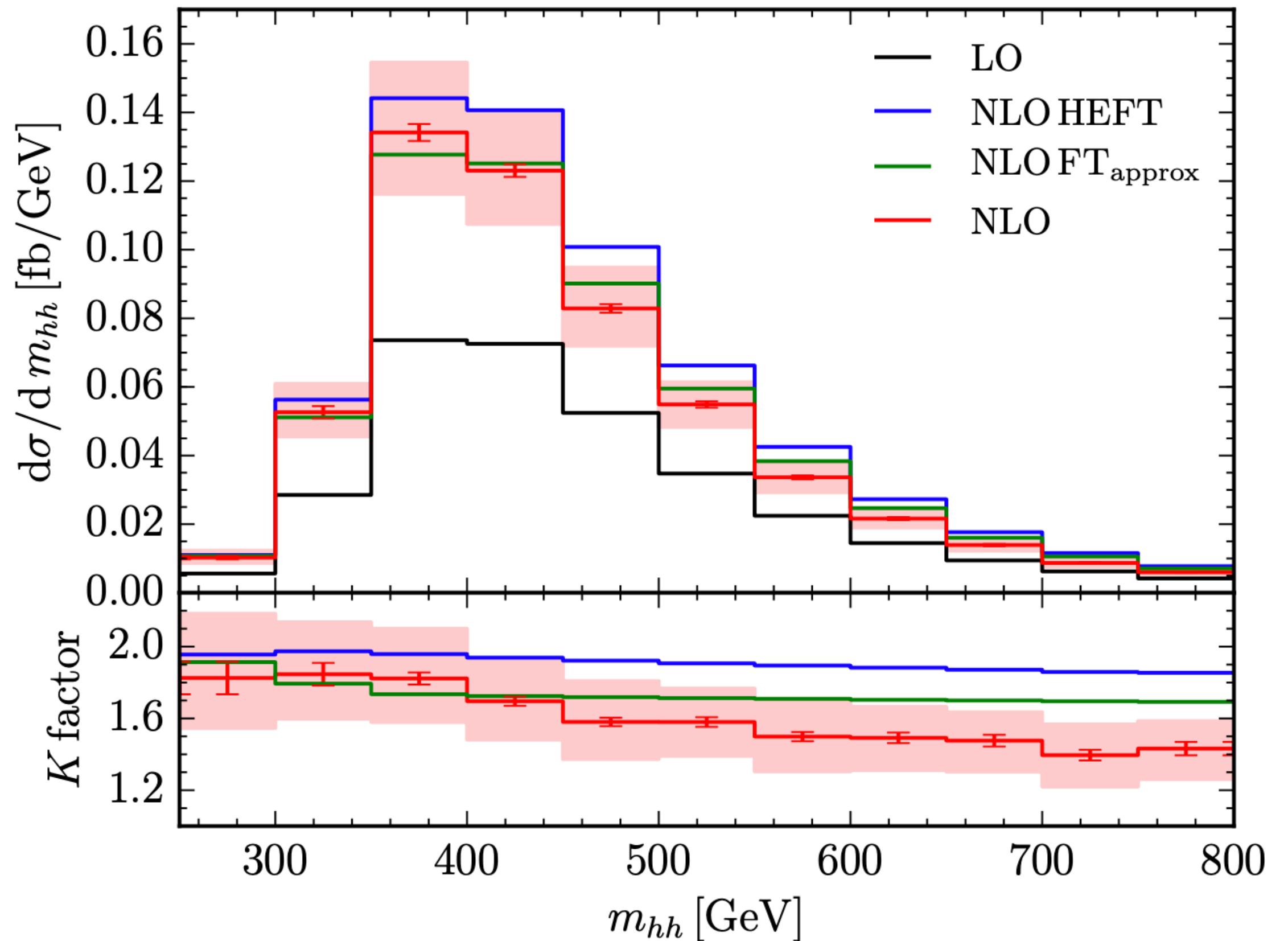
Borowka, Greiner, Heinrich, Jones, Kerner, Schlenk, Schubert, Zirke '16
 Baglio, Campanario, Glaus, Mühlleitner, Spira, Streicher '19



Baglio, Campanario, Glaus, Mühlleitner, Ronca, Spira '20

Borowka, Greiner, Heinrich, Jones, Kerner, Schlenk, Schubert, Zirke '16

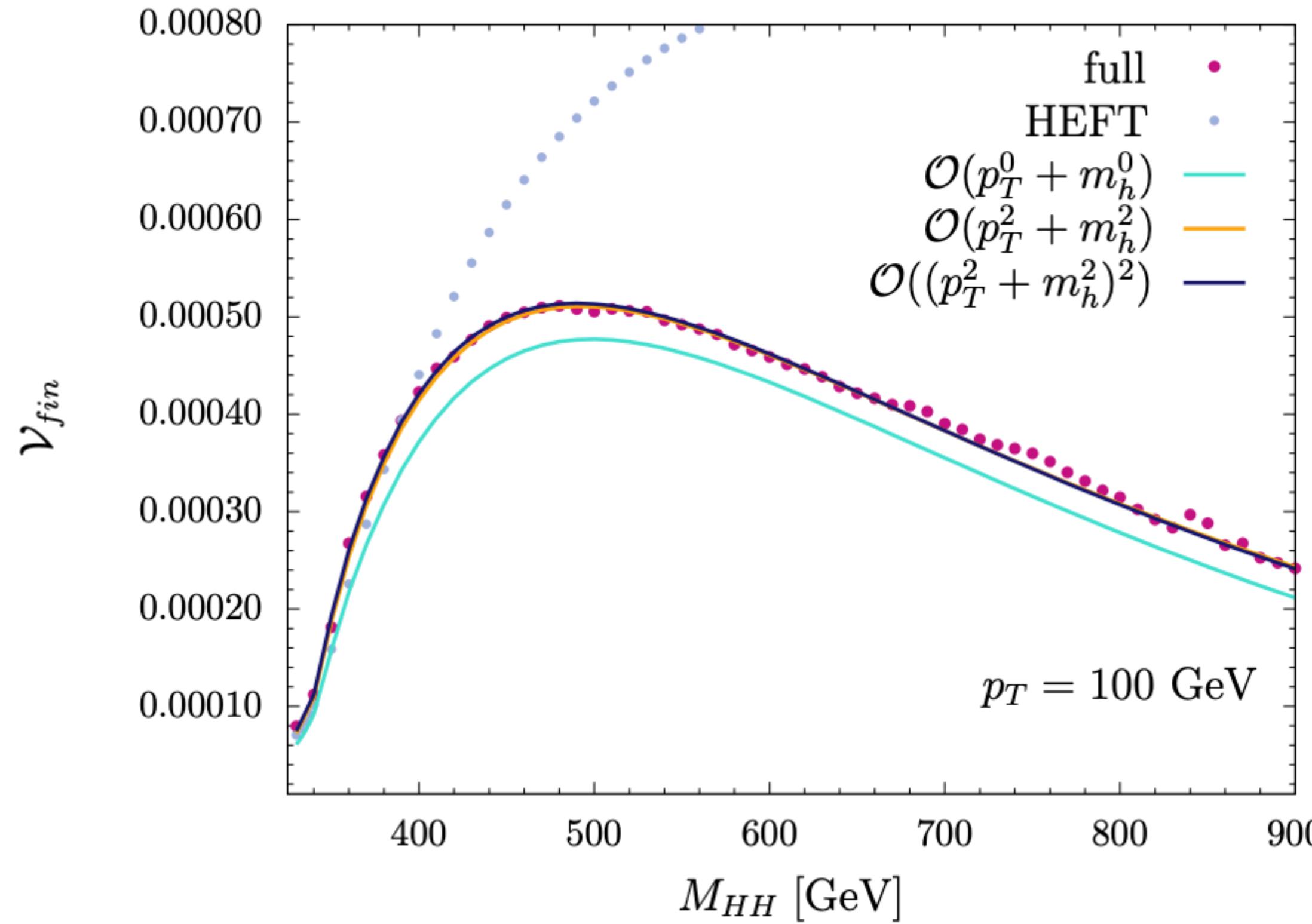
Baglio, Campanario, Glaus, Mühlleitner, Spira, Streicher '19



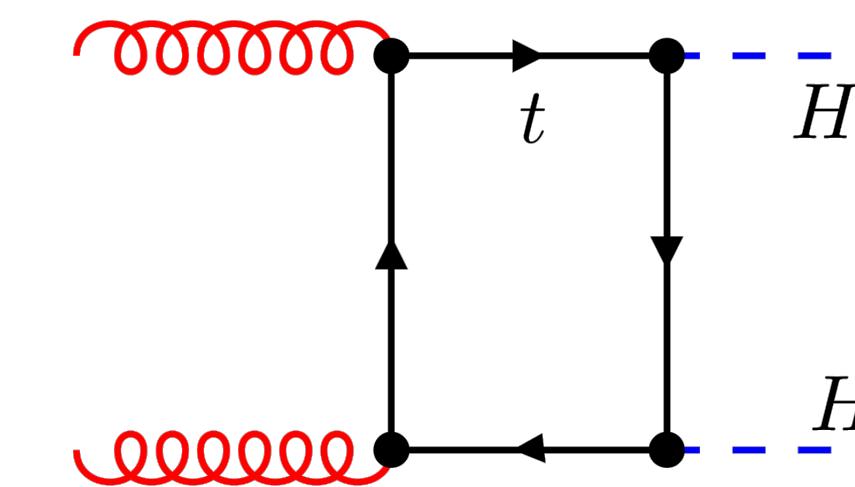
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 Baglio, Campanario, Glaus, Mühlleitner, Spira, Streicher '19

Baglio, Campanario, Glaus, Mühlleitner, Ronca, Spira '20

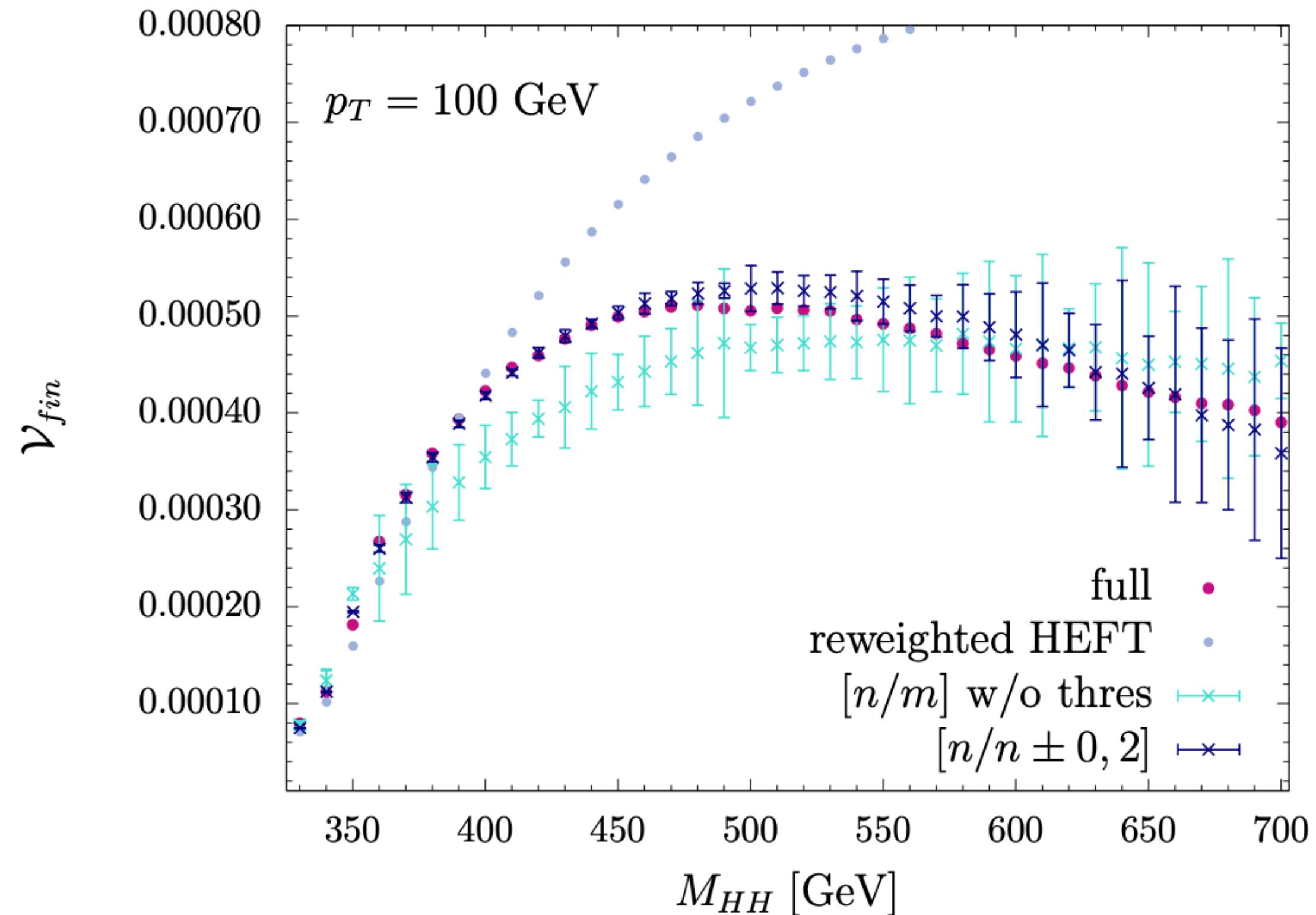
Improved/alternative approximations:



small p_T expansion
Bonciani, Degrassi, Giardino, Gröber '18



Improved/alternative approximations:



Padé approximation

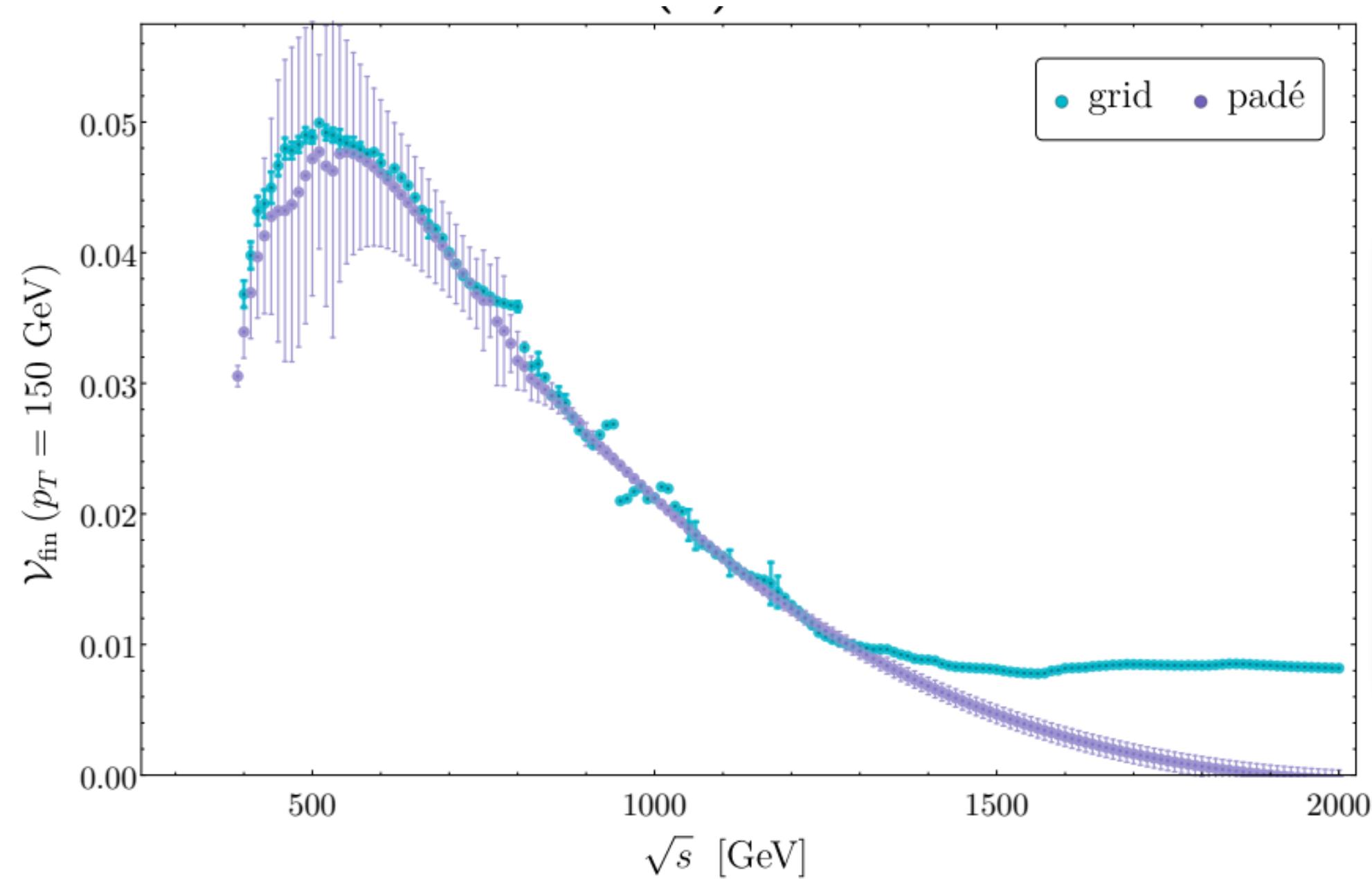
Gröber, Maier, Rauh '18

$1/m_t$ with threshold expansion:

$$F_{\Delta, C_F}^{2l} \underset{z \rightarrow 1}{\asymp} \pi^2(1-z) \ln(1-z) - \frac{\pi(40-3\pi^2)}{12}(1-z)^{3/2} + \frac{2\pi^2}{3}(1-z)^2 \ln(1-z) + \mathcal{O}((1-z)^{5/2}), \quad (10)$$

Combination with high-energy form factors ($\sqrt{\hat{s}} \gg m_t > m_H, m_W$):

Davies, Mishima, Steinhauser, Wellmann '18



Double Higgs boson production at NLO: combining
the exact numerical result and high-energy expansion

P3H-19-015

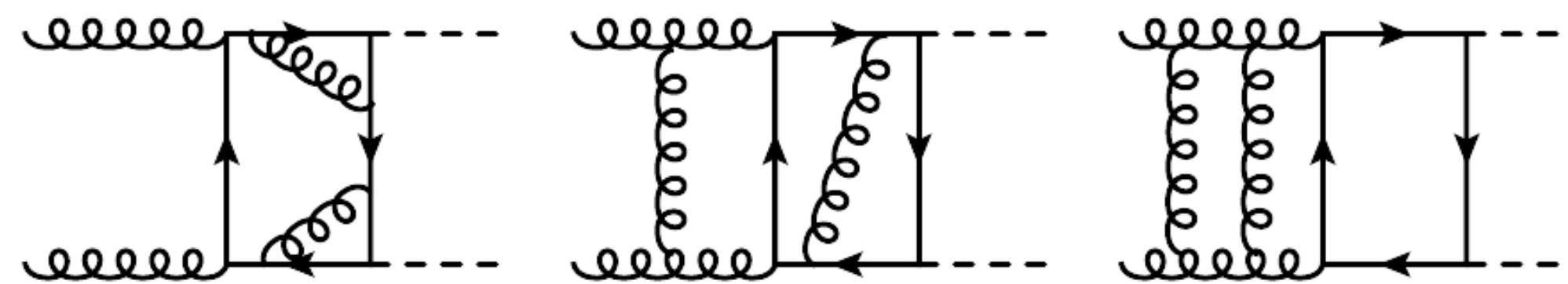
Joshua Davies^a, Gudrun Heinrich^b, Stephen P. Jones^c, Matthias Kerner^d,
Go Mishima^{a,e}, Matthias Steinhauser^a, David Wellmann^a

Combination with NNLO

Grazzini, Heinrich, Jones, Kallweit, Kerner, Lindert, Mazzitelli '18

N³LO in heavy-top limit

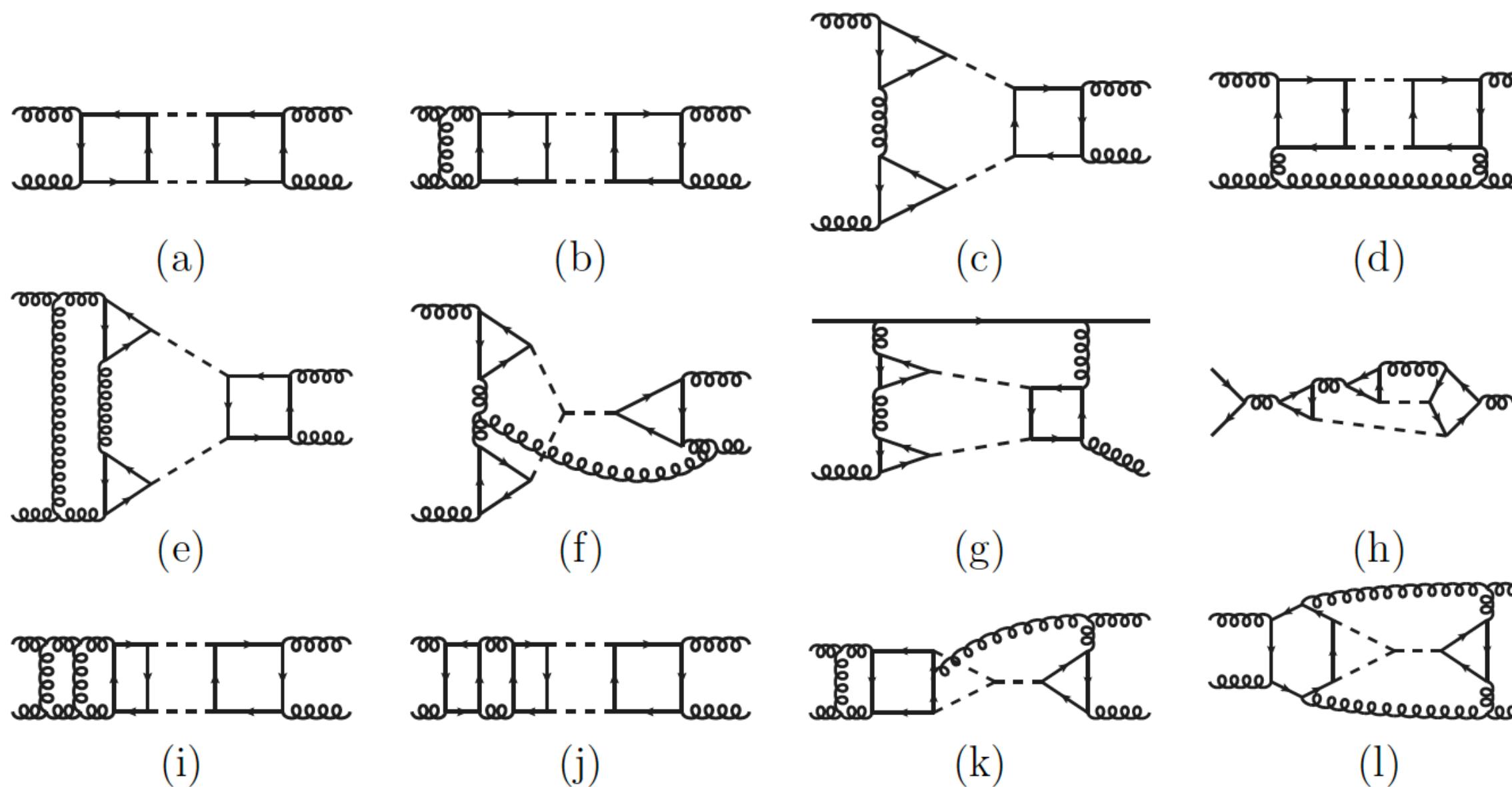
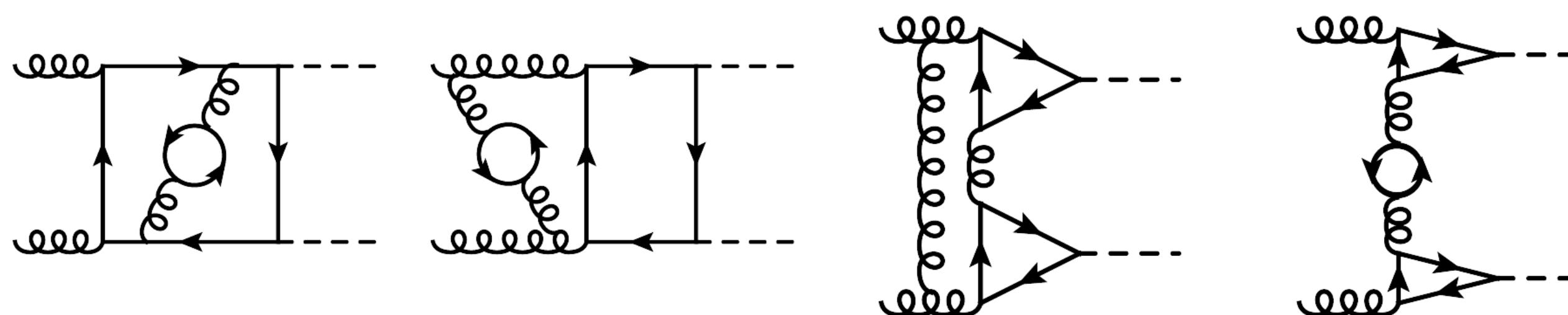
Chen, Li, Shao, Wang '20



Three-loop form factors for Higgs boson pair production in the large top mass limit

P3H-19-030

Joshua Davies and Matthias Steinhauser,



Real-virtual corrections to Higgs boson pair production at NNLO: three closed top quark loops

P3H-19-009

Joshua Davies^a, Florian Herren^a, Go Mishima^{a,b}, Matthias Steinhauser^a

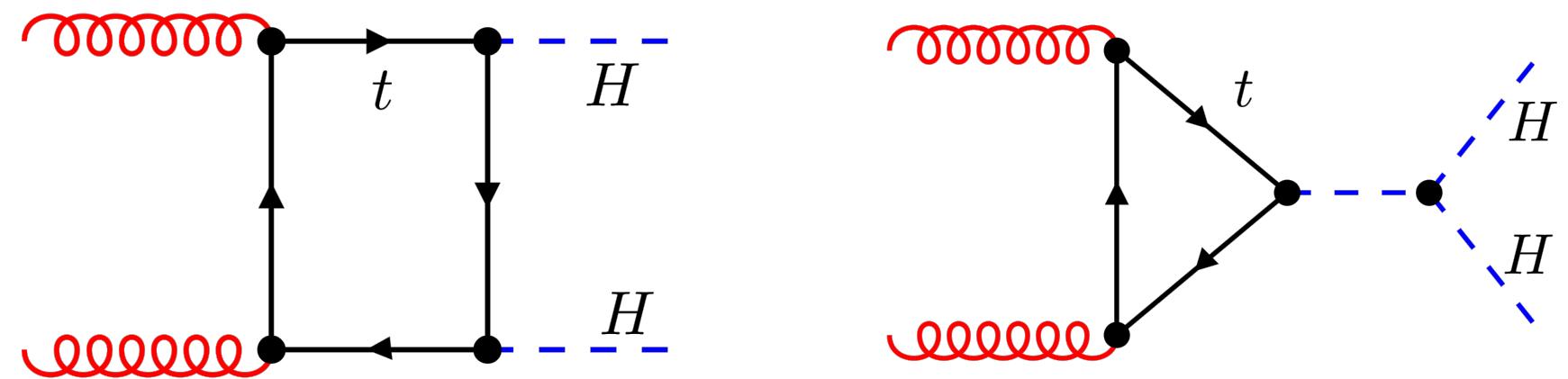
Application to BSM

$$\mathcal{L} \supset -m_t \left(c_t \frac{h}{v} + c_{tt} \frac{h^2}{v^2} \right) \bar{t} t - c_{hhh} \frac{m_h^2}{2v} h^3 + \frac{\alpha_s}{8\pi} \left(c_{ggh} \frac{h}{v} + c_{gggh} \frac{h^2}{v^2} \right) G_{\mu\nu}^a G^{a,\mu\nu}$$

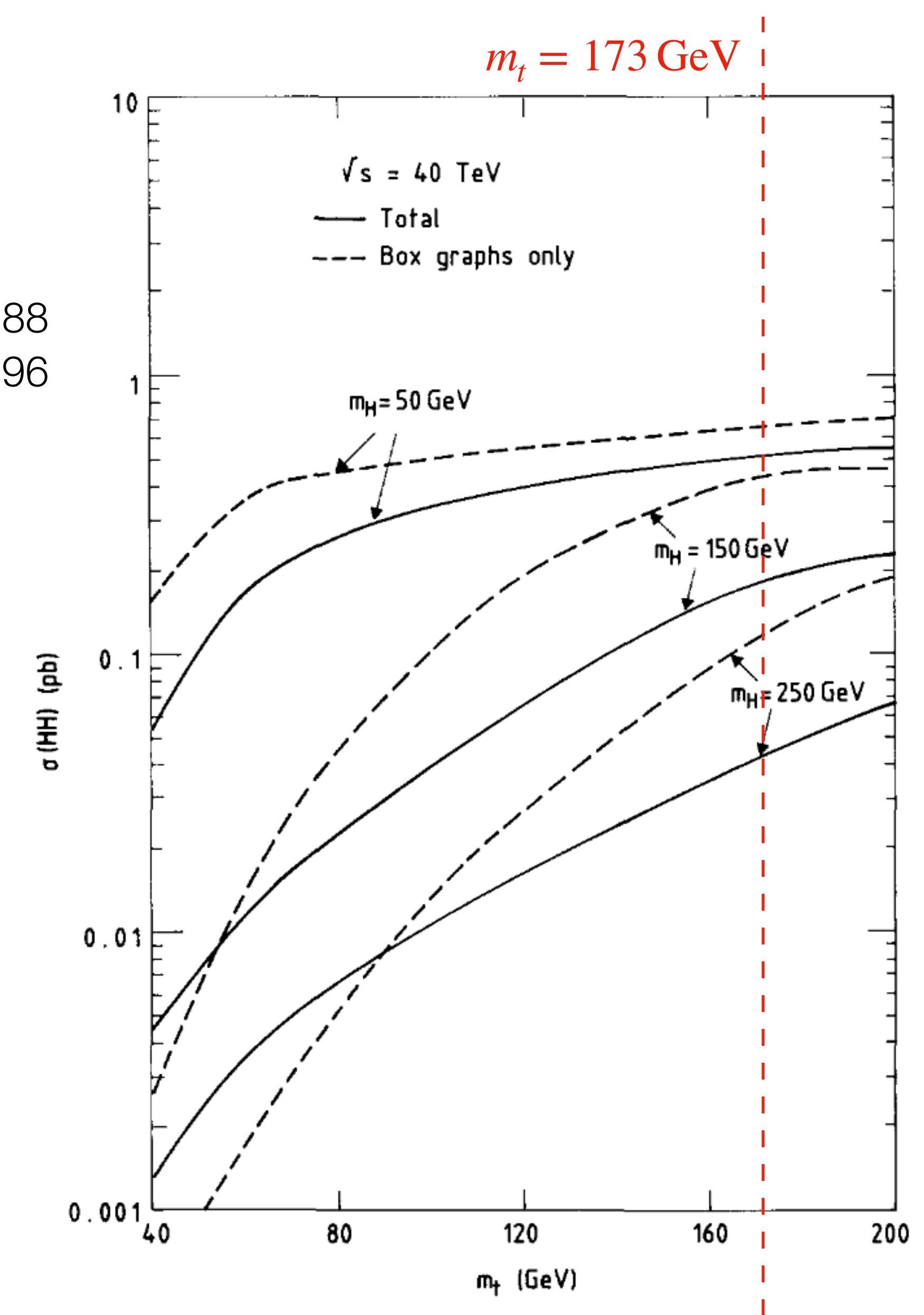
A non-linear EFT description of $gg \rightarrow HH$ at
NLO interfaced to POWHEG P3H-20-026

Gudrun Heinrich,^a Stephen P. Jones,^b Matthias Kerner,^c Ludovic Scyboz^d

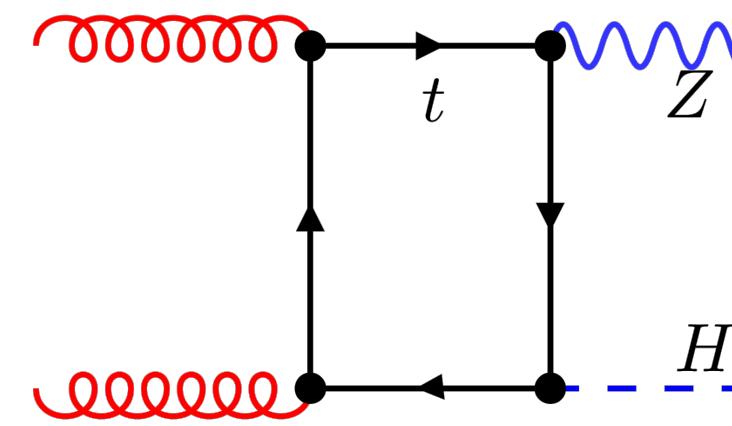
HH production



v.d.Bij, Glover '88
Plehn, Spira, Zerwas '96

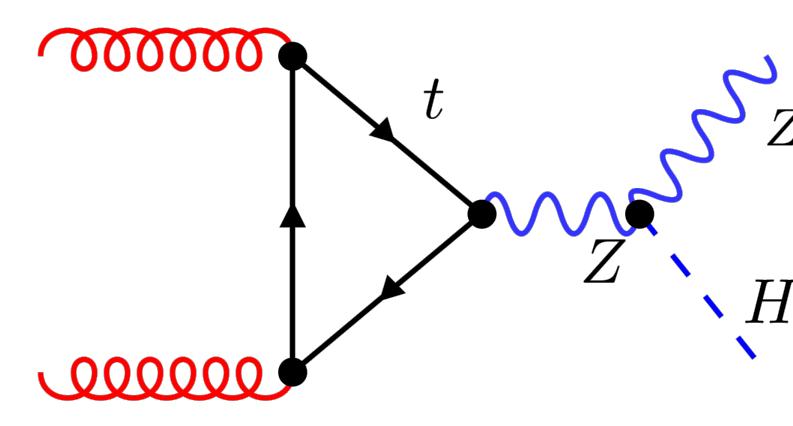


ZH production



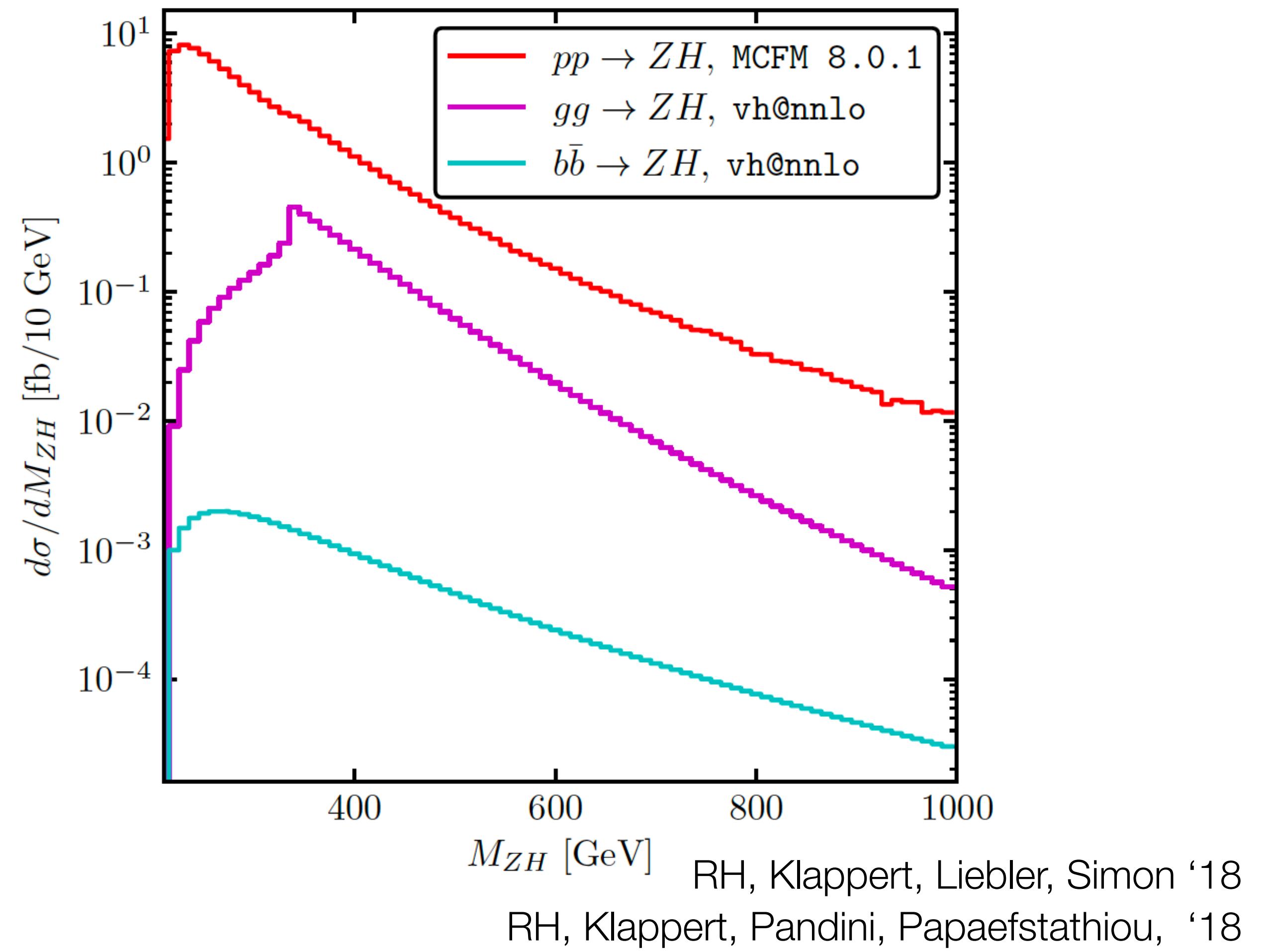
LO

Kniehl '90

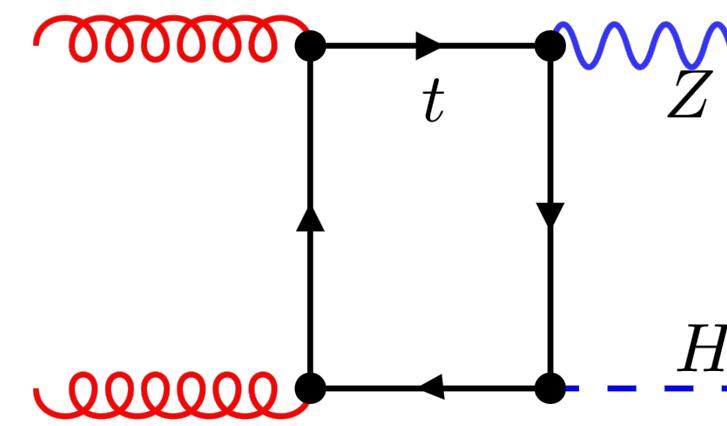
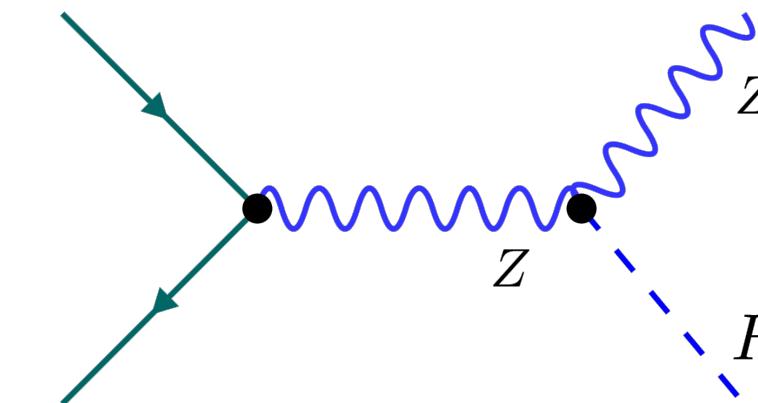


NLO 1/mt

Altenkamp, Dittmaier, RH, Rzehak, Zirke '13
Hasselhuhn, Luthe, Steinhauser '17



ZH production



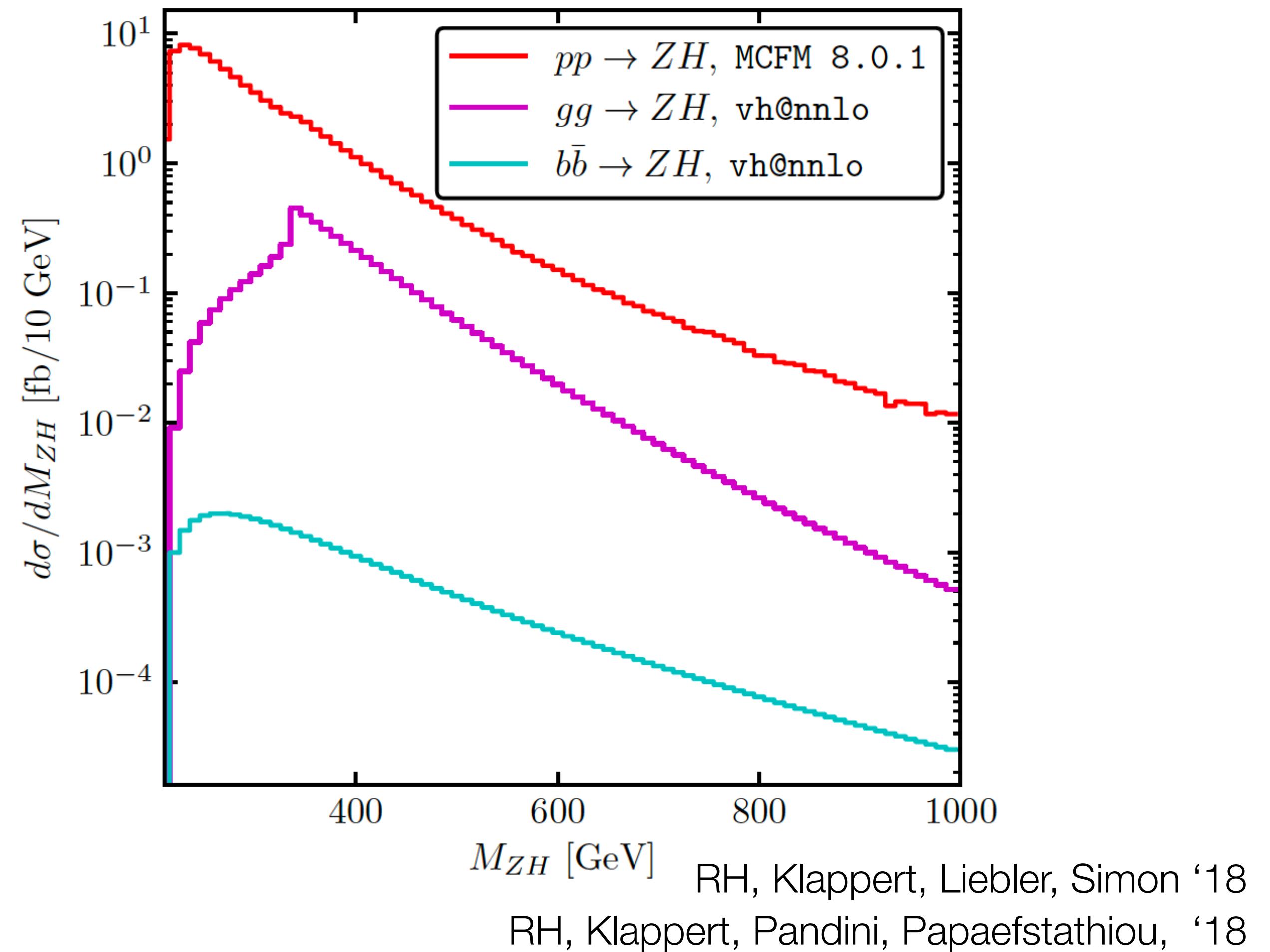
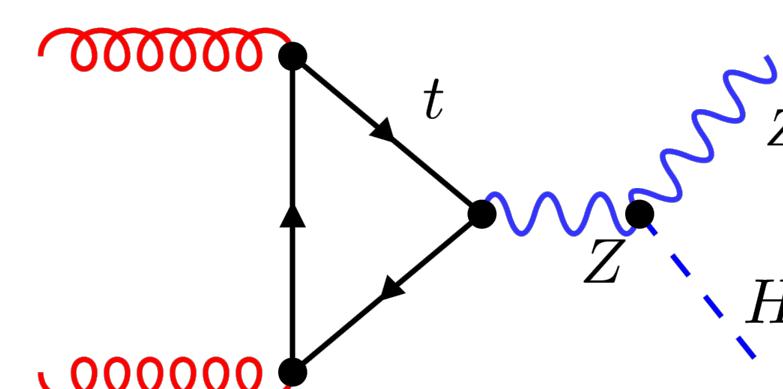
LO

Kniehl '90

NLO 1/mt

Altenkamp, Dittmaier, RH, Rzehak, Zirke '13

Hasselhuhn, Luthe, Steinhauser '17

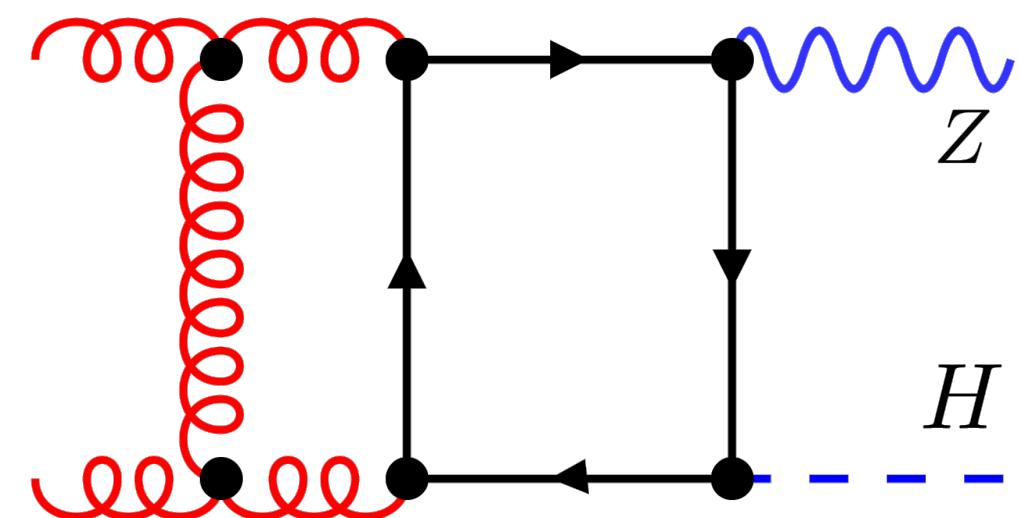


exact virtual:

ZH production in gluon fusion: two-loop amplitudes with full top quark mass dependence

P3H-20-076

Long Chen,^{a,f} Gudrun Heinrich,^b Stephen P. Jones,^{c,d} Matthias Kerner,^e Jonas Klappert,^f Johannes Schlenk^g

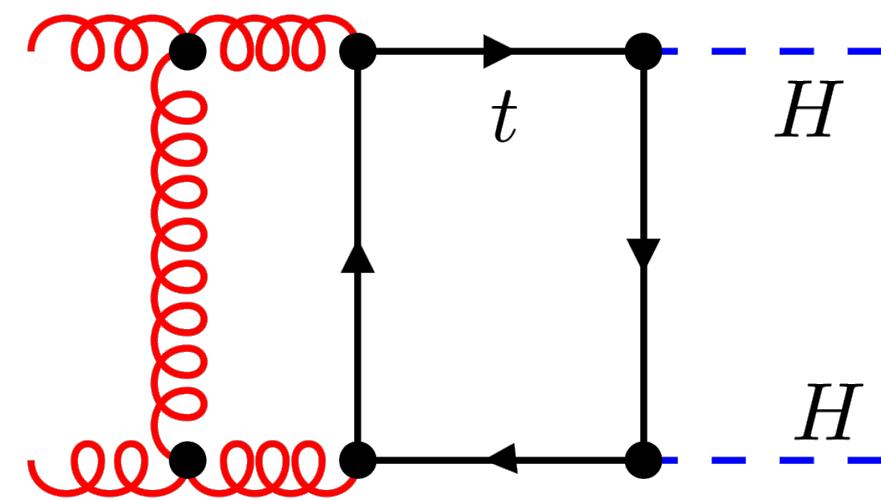


using Kira+FireFly:

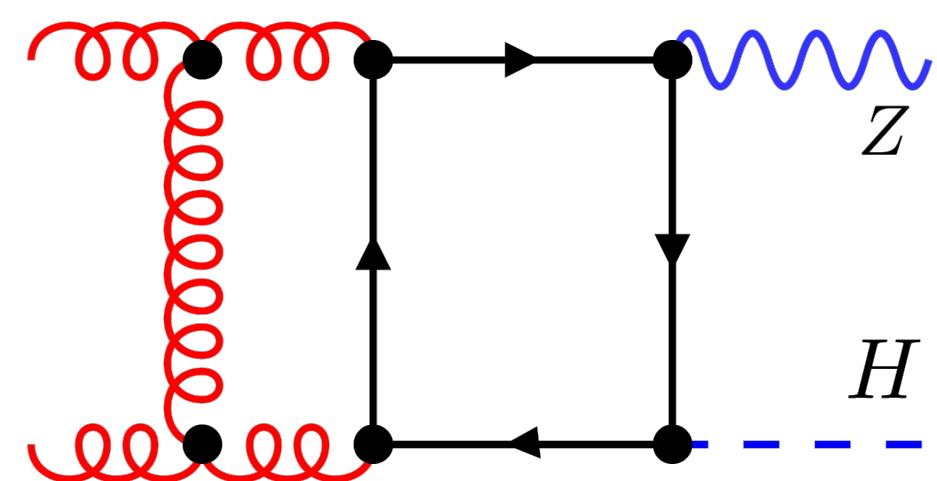
Integral Reduction with Kira 2.0 and Finite Field Methods

P3H-20-041

Jonas Klappert*,^a, Fabian Lange^{†,a}, Philipp Maierhöfer^{‡,b}, and Johann Usovitsch^{§,c}



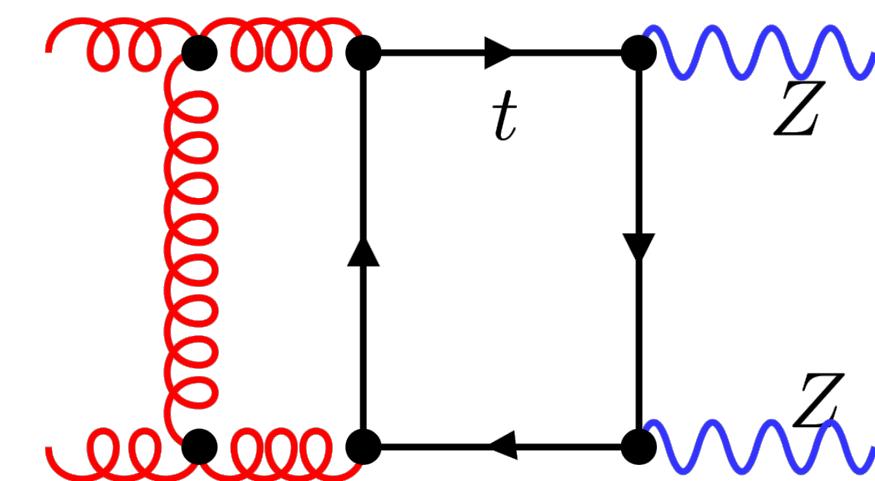
Borowka, Greiner, Heinrich, Jones, Kerner, Schlenk, Schubert, Zirke '16
 Baglio, Campanario, Glaus, Mühlleitner, Spira, Streicher '19



ZH production in gluon fusion: two-loop amplitudes with full top quark mass dependence

P3H-20-076

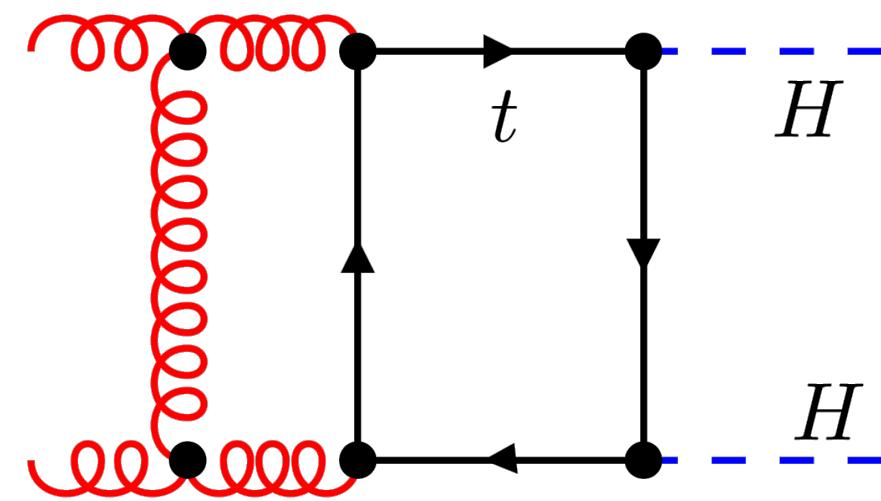
Long Chen,^{a,f} Gudrun Heinrich,^b Stephen P. Jones,^{c,d} Matthias Kerner,^e Jonas Klappert,^f Johannes Schlenk^g



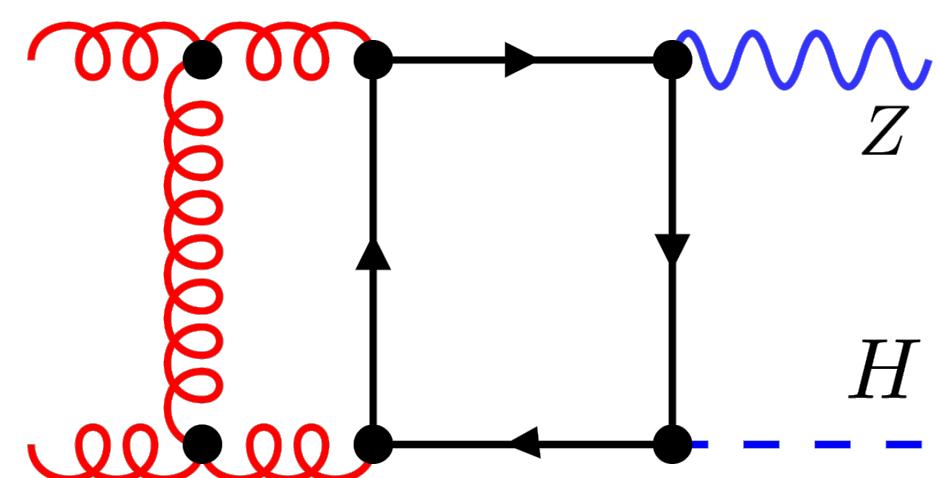
Top quark contribution to two-loop helicity amplitudes for Z boson pair production in gluon fusion

P3H-21-008

Christian Brønnum-Hansen and Chen-Yu Wang



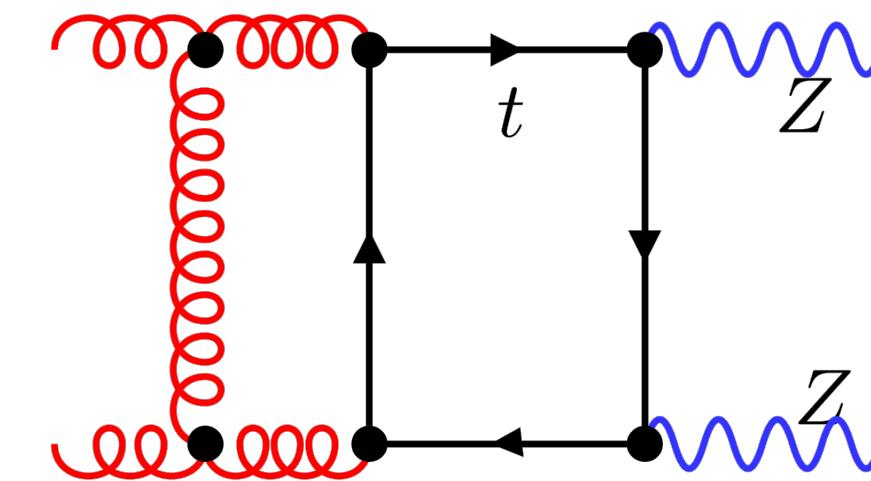
Borowka, Greiner, Heinrich, Jones, Kerner, Schlenk, Schubert, Zirke '16
 Baglio, Campanario, Glaus, Mühlleitner, Spira, Streicher '19



ZH production in gluon fusion: two-loop amplitudes with full top quark mass dependence

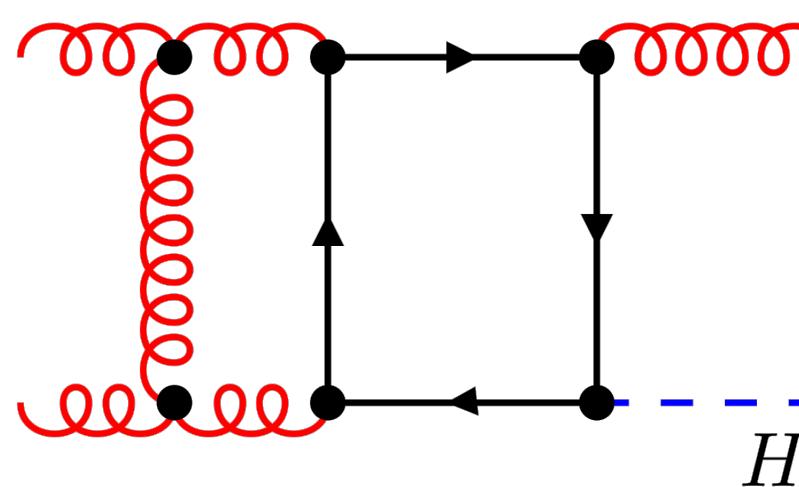
P3H-20-076

Long Chen,^{a,f} Gudrun Heinrich,^b Stephen P. Jones,^{c,d} Matthias Kerner,^e Jonas Klappert,^f Johannes Schlenk^g

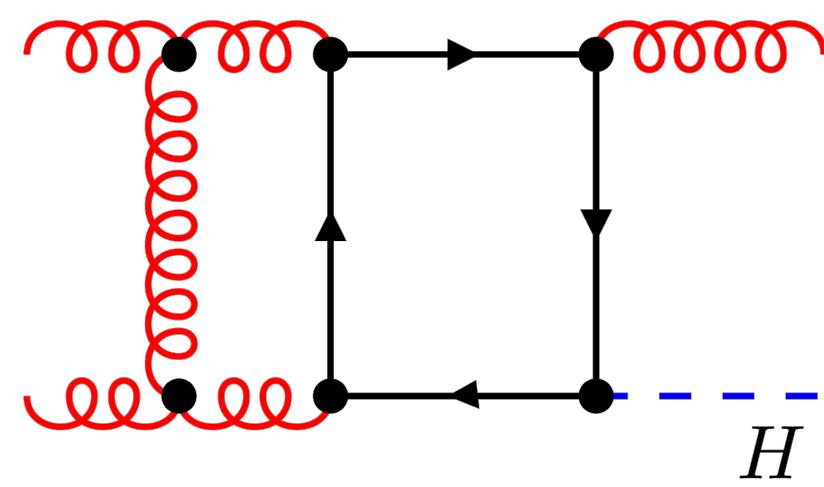


Top quark contribution to two-loop helicity amplitudes for Z boson pair production in gluon fusion

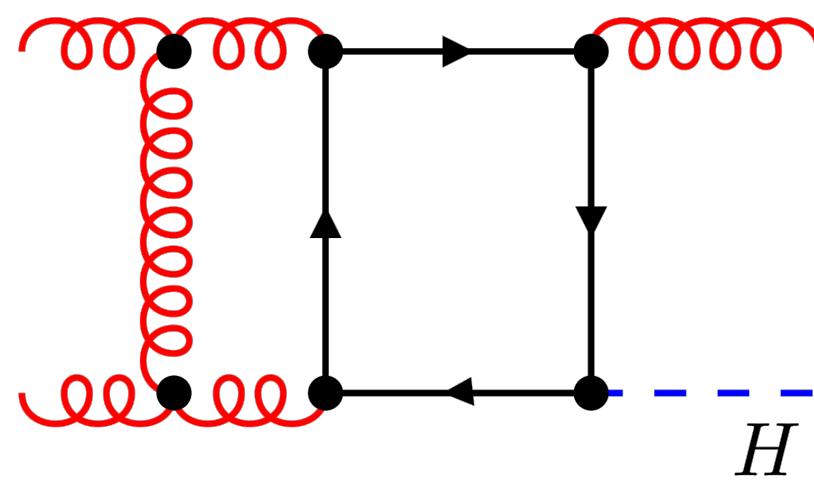
P3H-21-008



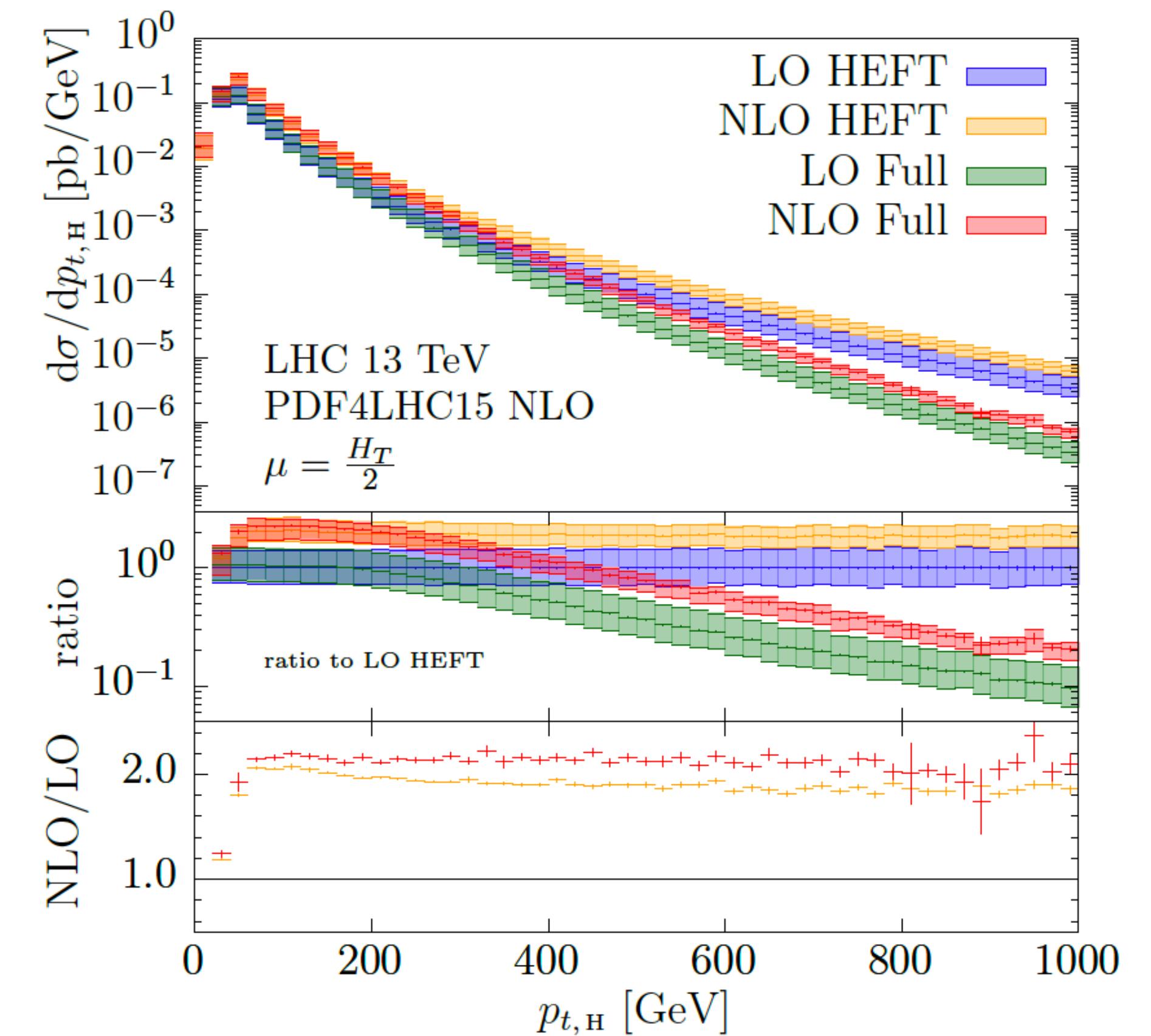
Christian Brønnum-Hansen and Chen-Yu Wang

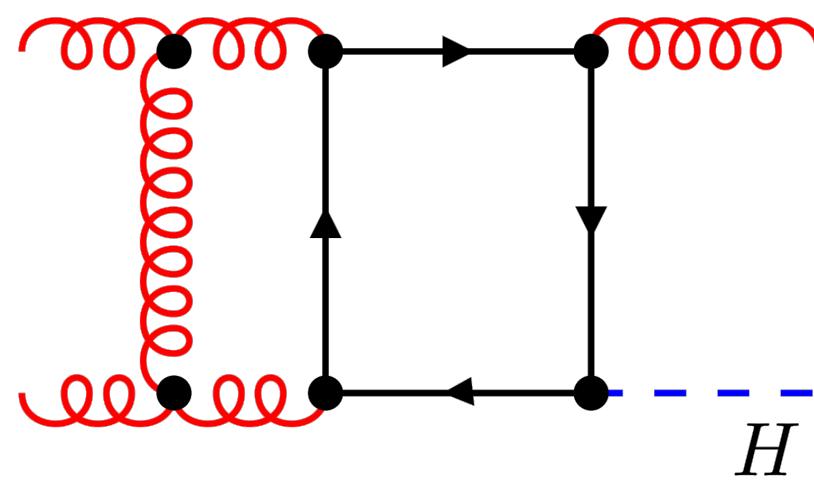


Lindert, Kudashkin, Melnikov, Wever '18 (large p_T)
Jones, Kerner, Luisoni '18

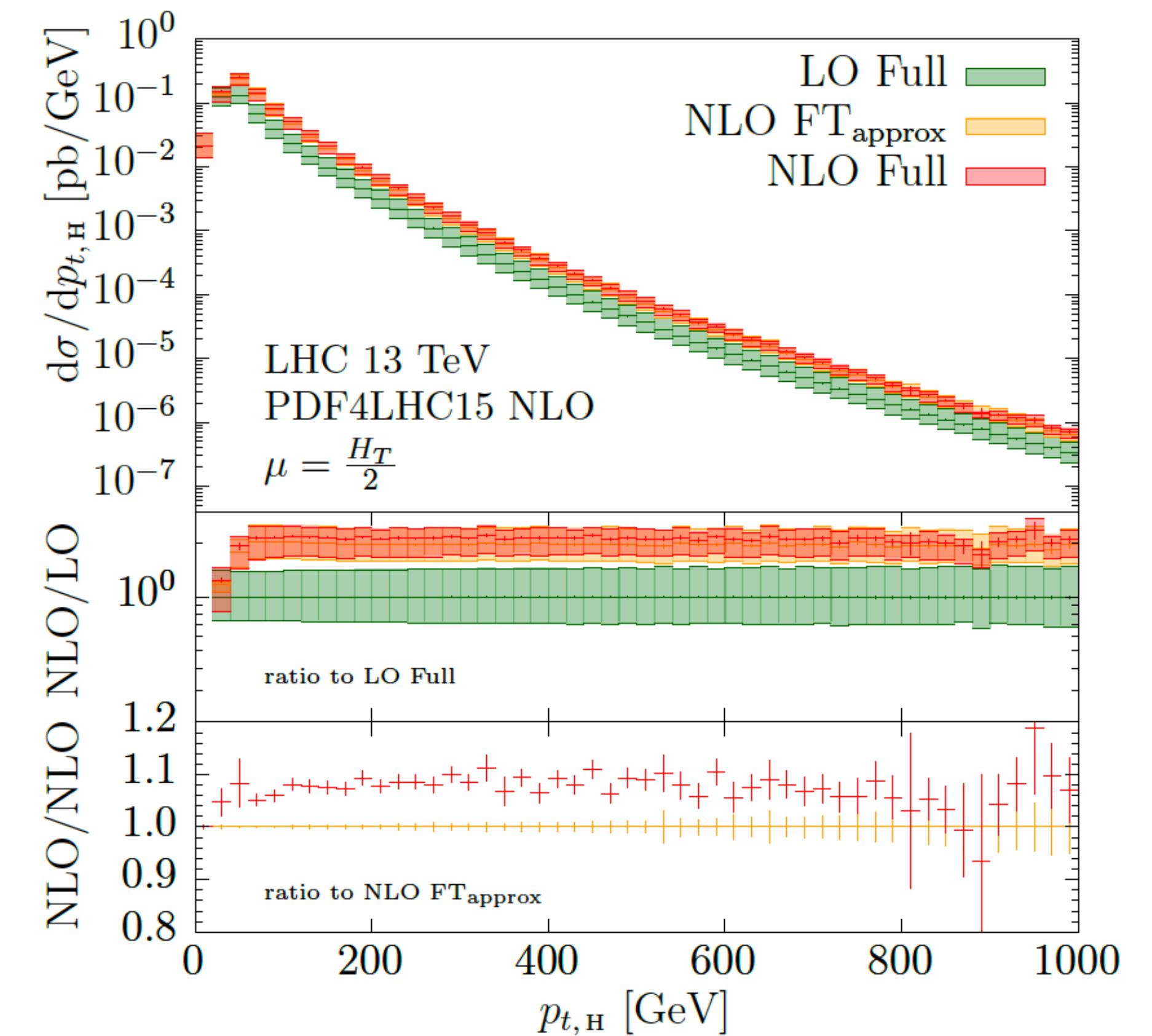


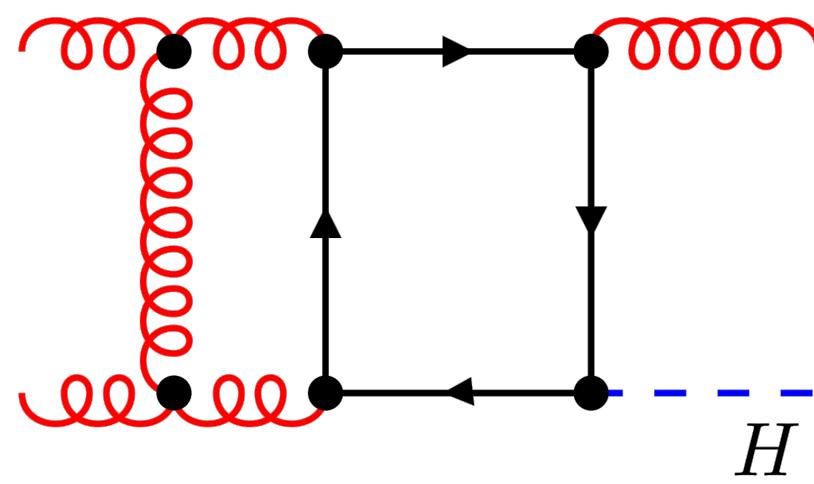
Lindert, Kudashkin, Melnikov, Wever '18 (large p_T)
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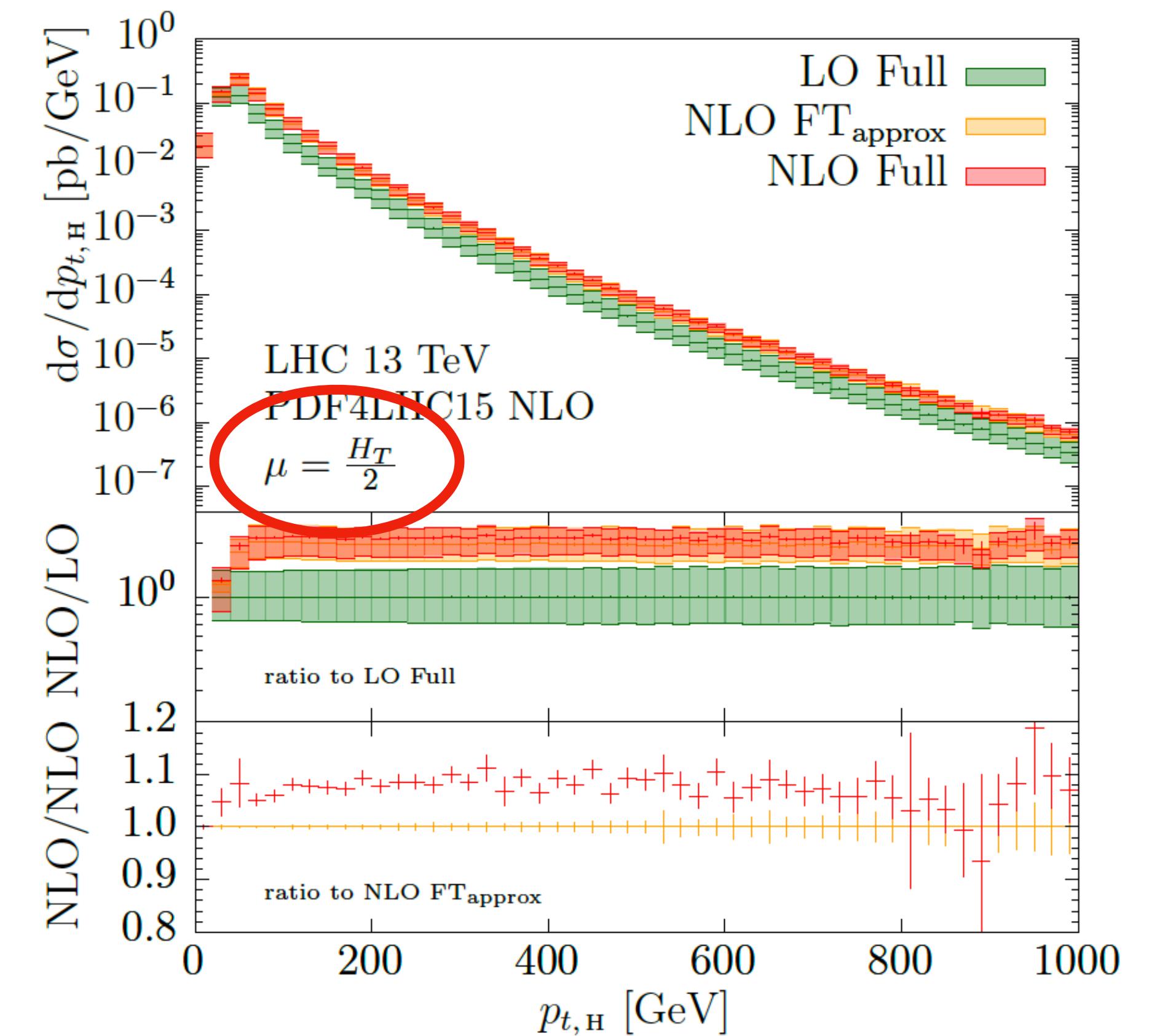


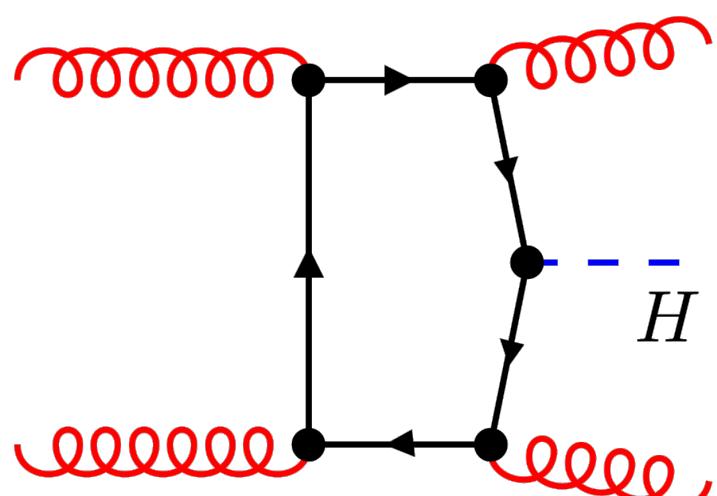
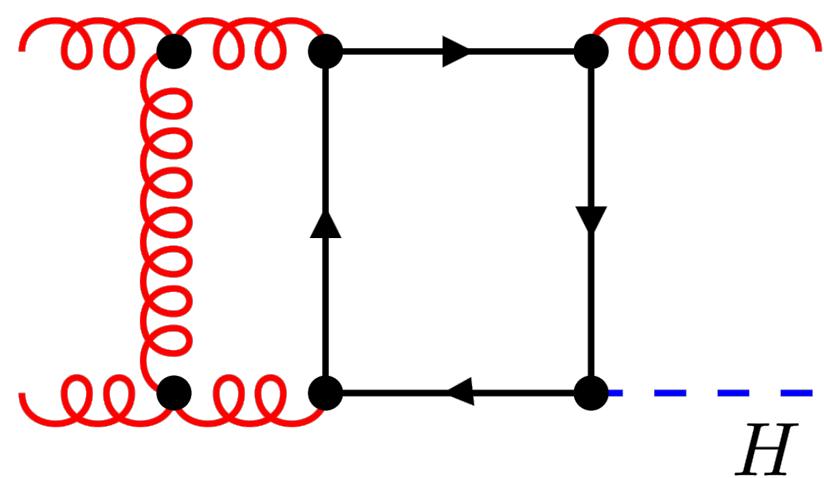
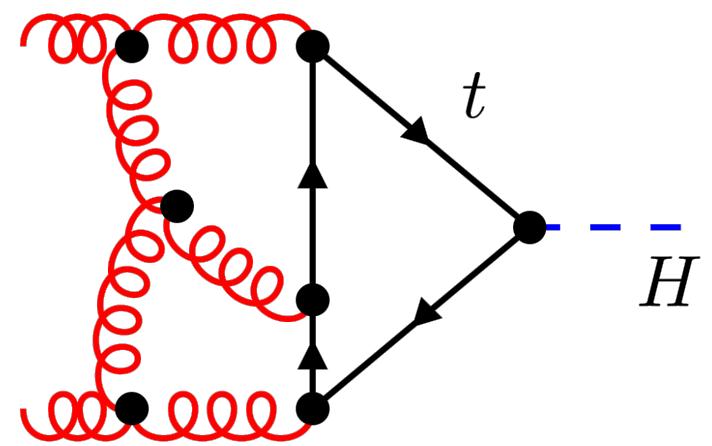
Lindert, Kudashkin, Melnikov, Wever '18 (large p_T)
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Lindert, Kudashkin, Melnikov, Wever '18 (large p_T)
 Jones, Kerner, Luisoni '18



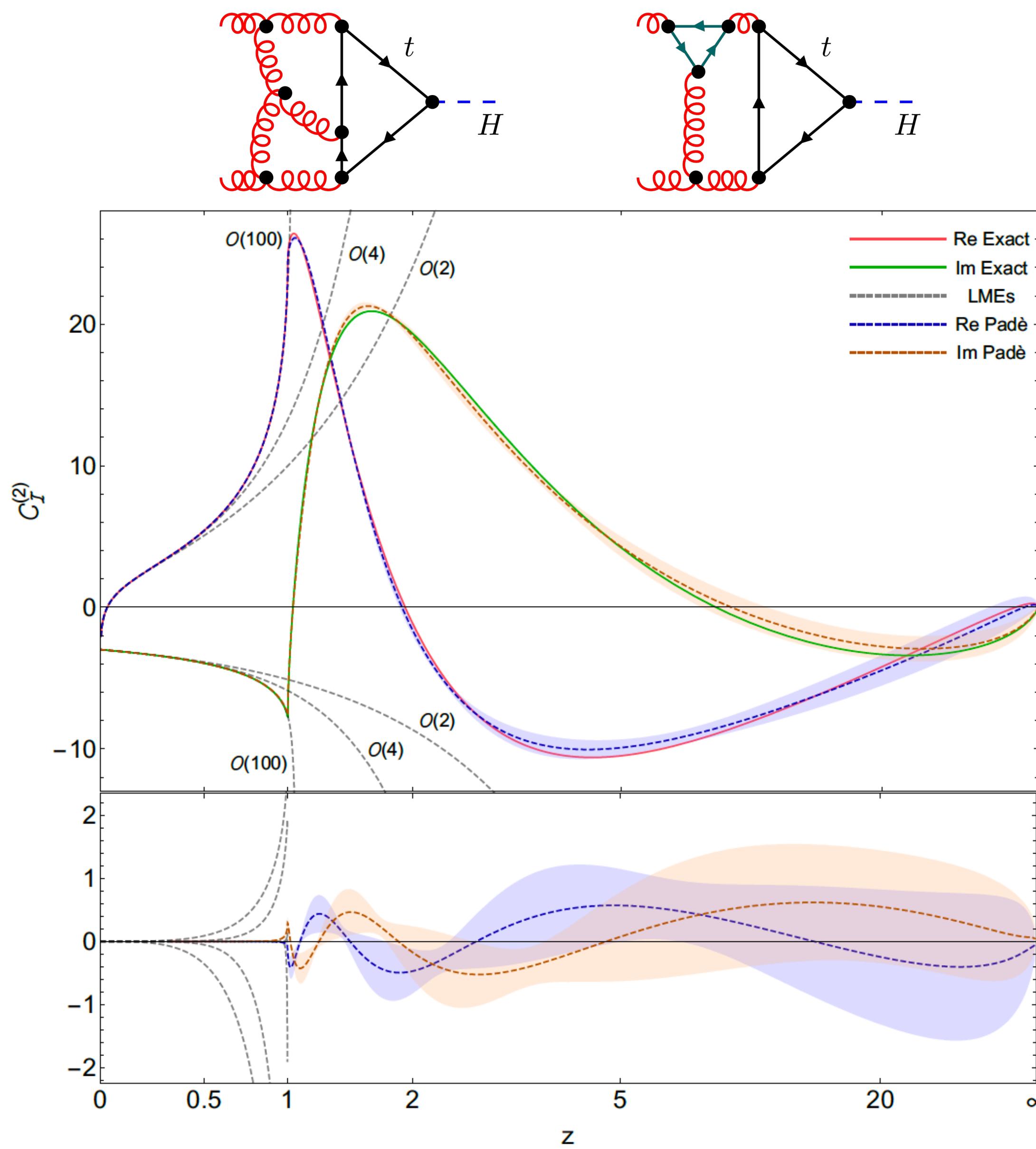


Lindert, Kudashkin, Melnikov, Wever '18 (large p_T)

Jones, Kerner, Luisoni '18

Del Duca, Kilgore, Schmidt, Oleari, Zeppenfeld '01

meanwhile: OpenLoops, MCFM, ...



Top quark mass dependence of the Higgs-gluon form factor at three loops

P3H-19-012

Joshua Davies^(a), Ramona Gröber^(b), Andreas Maier^(c),
Thomas Rauh^(d), Matthias Steinhauser^(a),

The light-fermion contribution to the exact Higgs-gluon form factor in QCD

P3H-19-021

Robert V. Harlander¹, Mario Prausa², and Johann Usovitsch³

Exact quark-mass dependence of the Higgs-gluon form factor at three loops in QCD

P3H-20-001

Michał Czakon and Marco Niggetiedt

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

	α_s^2	α_s^3	α_s^4		
gg	$16.30 + 19.64 + 8.76$			0.2%	2.8%
qg		$1.49 + 0.84$		24.8%	-1.4%
qq		$0.02 + 0.10$		161%	-50%
total	$16.30 + 21.15 + 9.79$			3.5%	$+1.9\%$

Exact top-quark mass dependence in hadronic Higgs production

NLO+NNLO: -0.26%

M. Czakon, R.V. Harlander, J. Klappert, M. Niggetiedt

*Institute for Theoretical Particle Physics and Cosmology,
RWTH Aachen University, 52056 Aachen, Germany*

(Dated: May 11, 2021)

P3H-21-031

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

	α_s^2	α_s^3	α_s^4		
gg	$16.30 + 19.64 + 8.76$			0.2%	2.8%
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P3H-21-031

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

	α_s^2	α_s^3	α_s^4			
gg	$16.30 + 19.64 + 8.76$			0.2%	2.8%	$\rightarrow + 0.5 \%$
qg		$1.49 + 0.84$		24.8%	-1.4%	$\rightarrow - 0.02 \%$
qq		$0.02 + 0.10$		161%	-50%	$\rightarrow - 0.1 \%$
total	$16.30 + 21.15 + 9.79$			3.5%	$+1.9 \%$	$\rightarrow + 0.38 \%$

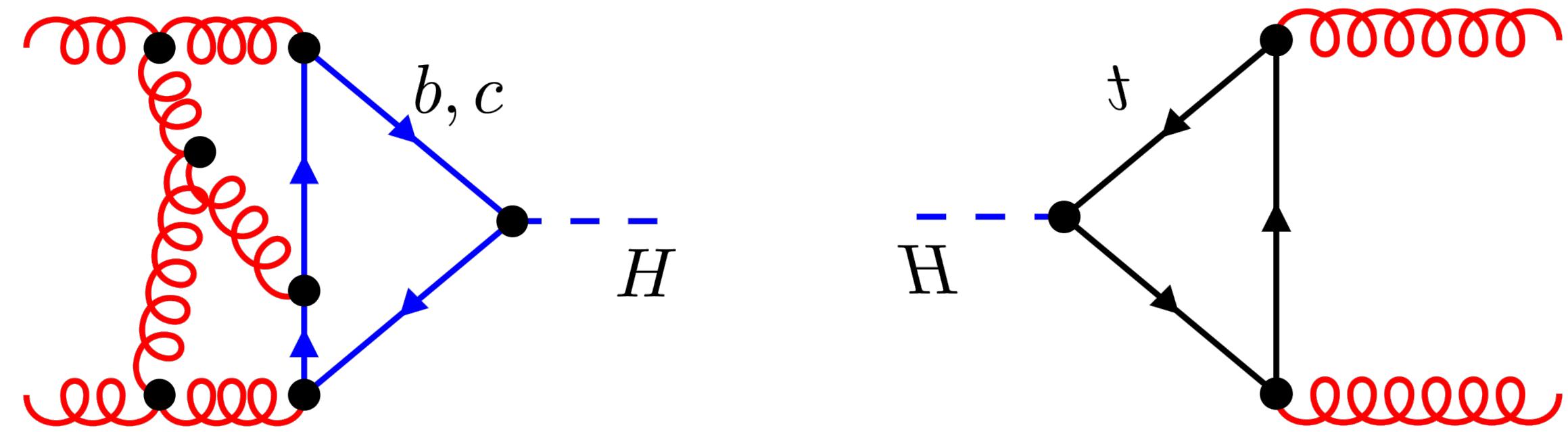
Exact top-quark mass dependence in hadronic Higgs production

M. Czakon, R.V. Harlander, J. Klappert, M. Niggetiedt **YSF**
*Institute for Theoretical Particle Physics and Cosmology,
RWTH Aachen University, 52056 Aachen, Germany*
(Dated: May 11, 2021)

NLO+NNLO: -0.26%

P3H-21-031

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
$+0.10 \text{ pb}$ -1.15 pb	$\pm 0.18 \text{ pb}$	$\pm 0.56 \text{ pb}$	$\pm 0.49 \text{ pb}$	$\pm 0.40 \text{ pb}$	$\pm 0.49 \text{ pb}$
$+0.21\%$ -2.37%	$\pm 0.27\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$



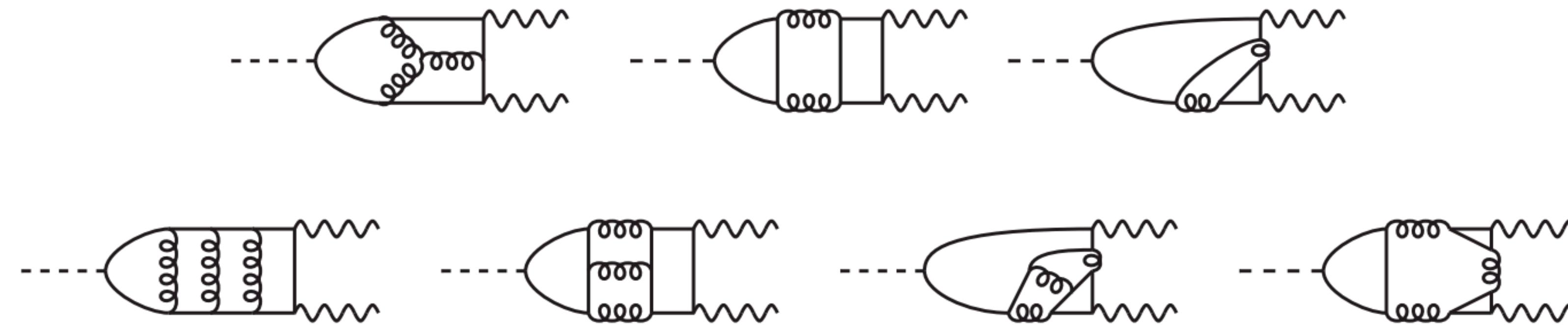
intricate structure of $\ln \frac{m_b}{m_H}$

Melnikov, Penin '16

Anastasiou, Penin '20

Liu, Mecaj, Neubert, Wang '20

What I could not talk about...



Exact quark-mass dependence of the Higgs-photon form factor at three loops in QCD

P3H-20-050

Marco Niggetiedt

3-loop exact

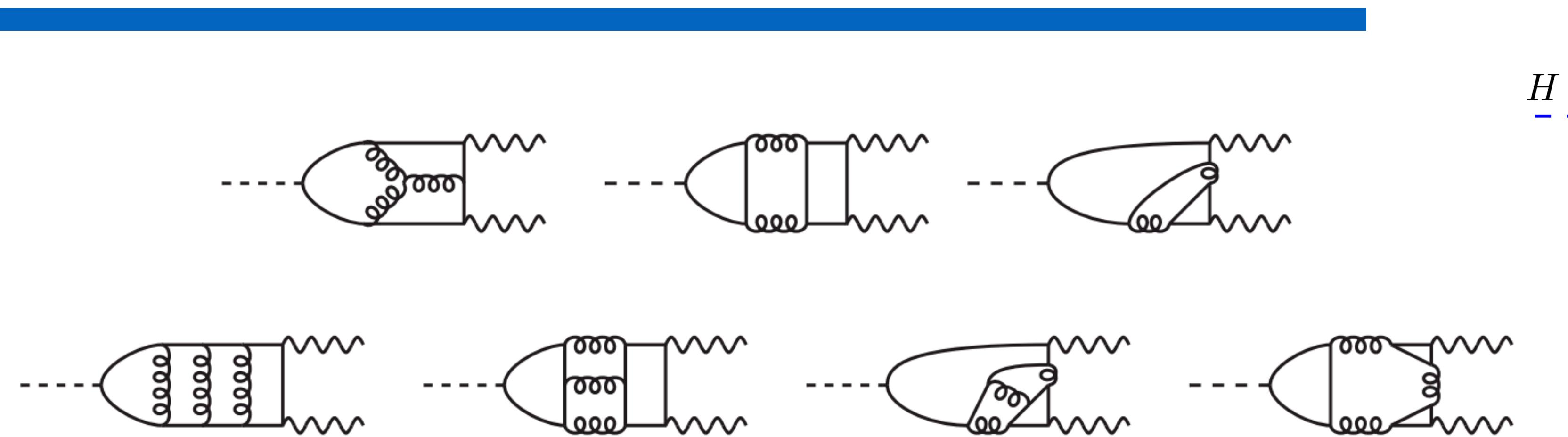
Higgs boson decay into photons at four loops

P3H-21-028

Joshua Davies¹ and Florian Herren^{2*}

4-loop, expansion in $1/m_t$

What I could not talk about...



Exact quark-mass dependence of the Higgs-photon
form factor at three loops in QCD

P3H-20-050

Marco Niggetiedt

3-loop exact

Higgs boson decay into photons at four loops

P3H-21-028

Joshua Davies¹ and Florian Herren^{2*}

4-loop, expansion in $1/m_t$

NNLO QCD corrections to three-photon production at the LHC

P3H-19-041

Herschel A. Chawdhry,^a Michał Czakon,^b Alexander Mitov,^a Rene Poncelet^a

Two-loop leading-colour QCD helicity amplitudes for two-photon plus jet production at the LHC

P3H-21-014

Herschel A. Chawdhry,^a Michał Czakon,^b Alexander Mitov,^c Rene Poncelet^c

NNLO QCD corrections to diphoton production with an additional jet at the LHC

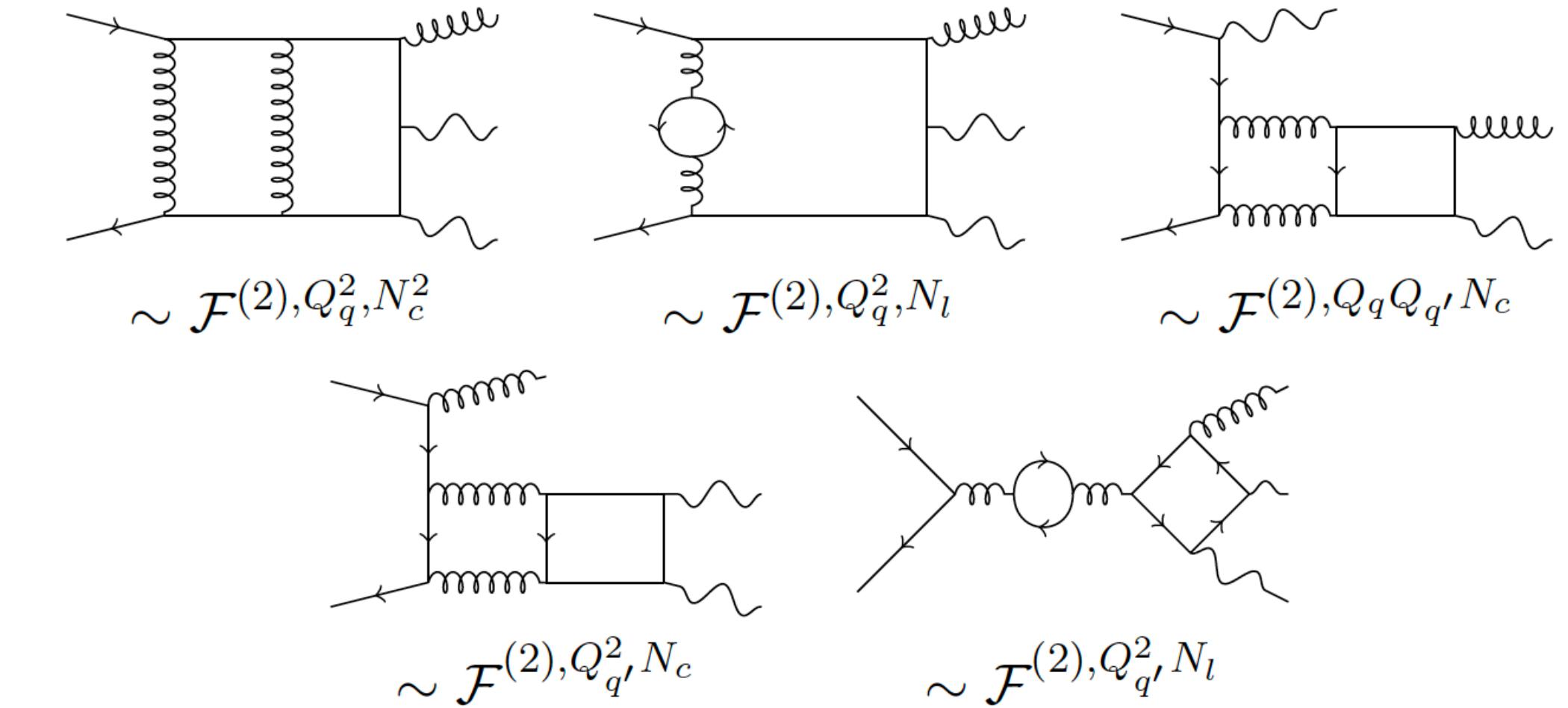
P3H-21-032

Herschel A. Chawdhry,^a Michał Czakon,^b Alexander Mitov,^c Rene Poncelet^c

Two-loop leading-color helicity amplitudes for three-photon production at the LHC

P3H-21-084

Herschel A. Chawdhry,^a Michał Czakon,^b Alexander Mitov,^c Rene Poncelet^c



also: Agarwal, Buccioni, von Manteuffel, Tancredi '21

What I could not talk about...

**Energy-energy correlation in hadronic Higgs decays:
analytic results and phenomenology at NLO**

P3H-20-060

Jun Gao,^{a,b} Vladyslav Shtabovenko,^{c,e} Tong-Zhi Yang,^{d,e}

Bottom quark mass effects in associated WH production

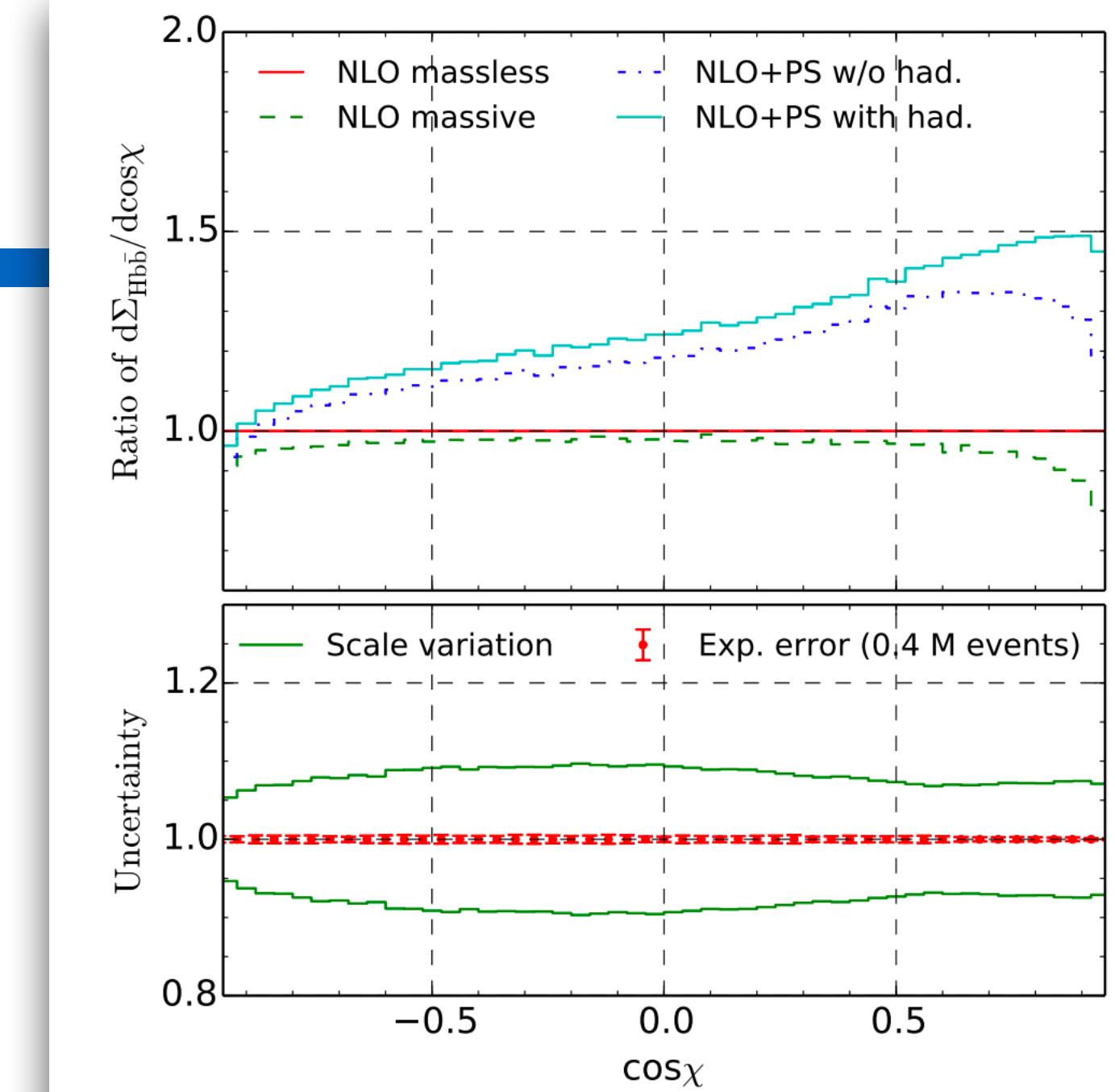
with $H \rightarrow b\bar{b}$ decay through NNLO QCD
P3H-20-009

Arnd Behring,^{1,*} Wojciech Bizoń,^{1,2,†} Fabrizio Caola,^{3,‡}
Kirill Melnikov,^{1,§} and Raoul Röntsch^{4,¶}

**Higgs decay into massive b -quarks at NNLO QCD
in the nested soft-collinear subtraction scheme**

P3H-19-047

Arnd Behring^a, Wojciech Bizoń^{a,b}.



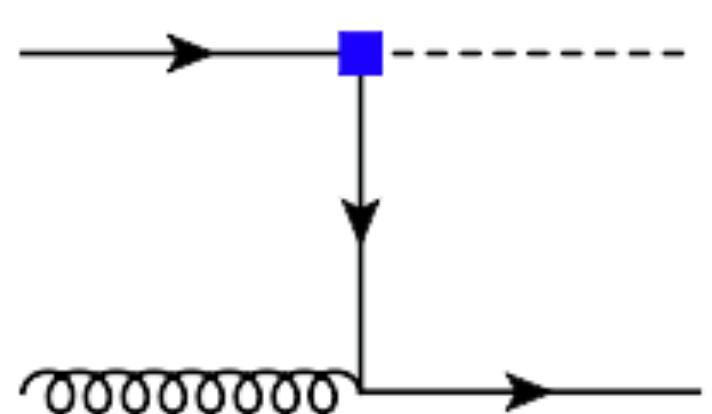
Order	b quarks	σ_{fid} [fb]	$\sigma_{\text{fid}}(\text{boosted})$ [fb]
LO	massive	$22.623^{+0.845}_{-1.047}$	$3.735^{+0.000}_{-0.016}$
	massless	$22.501^{+0.796}_{-1.007}$	$3.638^{+0.000}_{-0.009}$
NLO	massive	$25.364(1)^{+0.778}_{-0.756}$	$4.586(1)^{+0.158}_{-0.141}$
	massless	$24.421(1)^{+0.852}_{-0.879}$	$4.333(1)^{+0.165}_{-0.154}$
NNLO	massive	$24.225(4)^{+0.642}_{-0.742}$	$4.530(2)^{+0.071}_{-0.096}$
	massless	$22.781(3)^{+0.791}_{-0.898}$	$4.207(1)^{+0.097}_{-0.116}$

What I could not talk about...

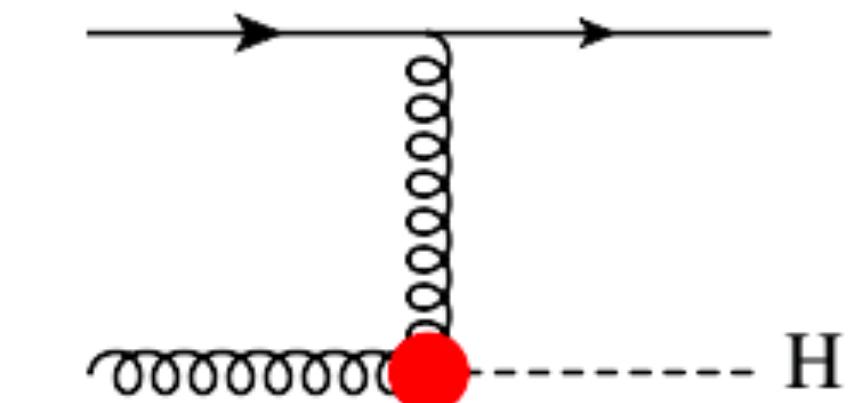
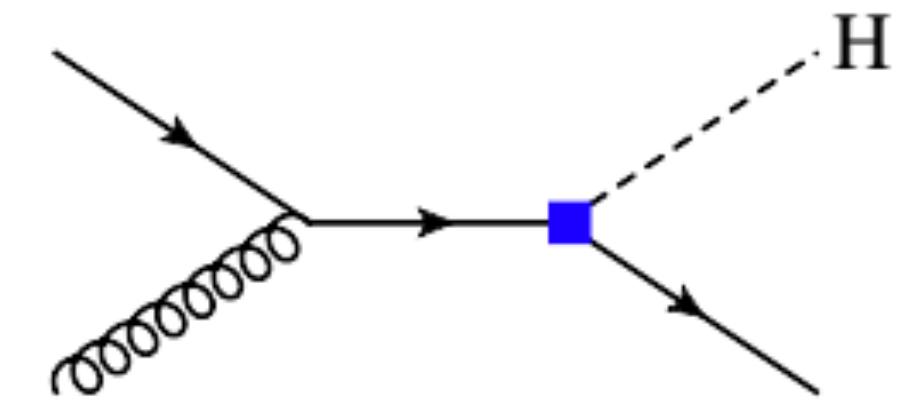
On the interference of ggH and $c\bar{c}H$ Higgs production mechanisms and the determination of charm Yukawa coupling at the LHC

P3H-21-010

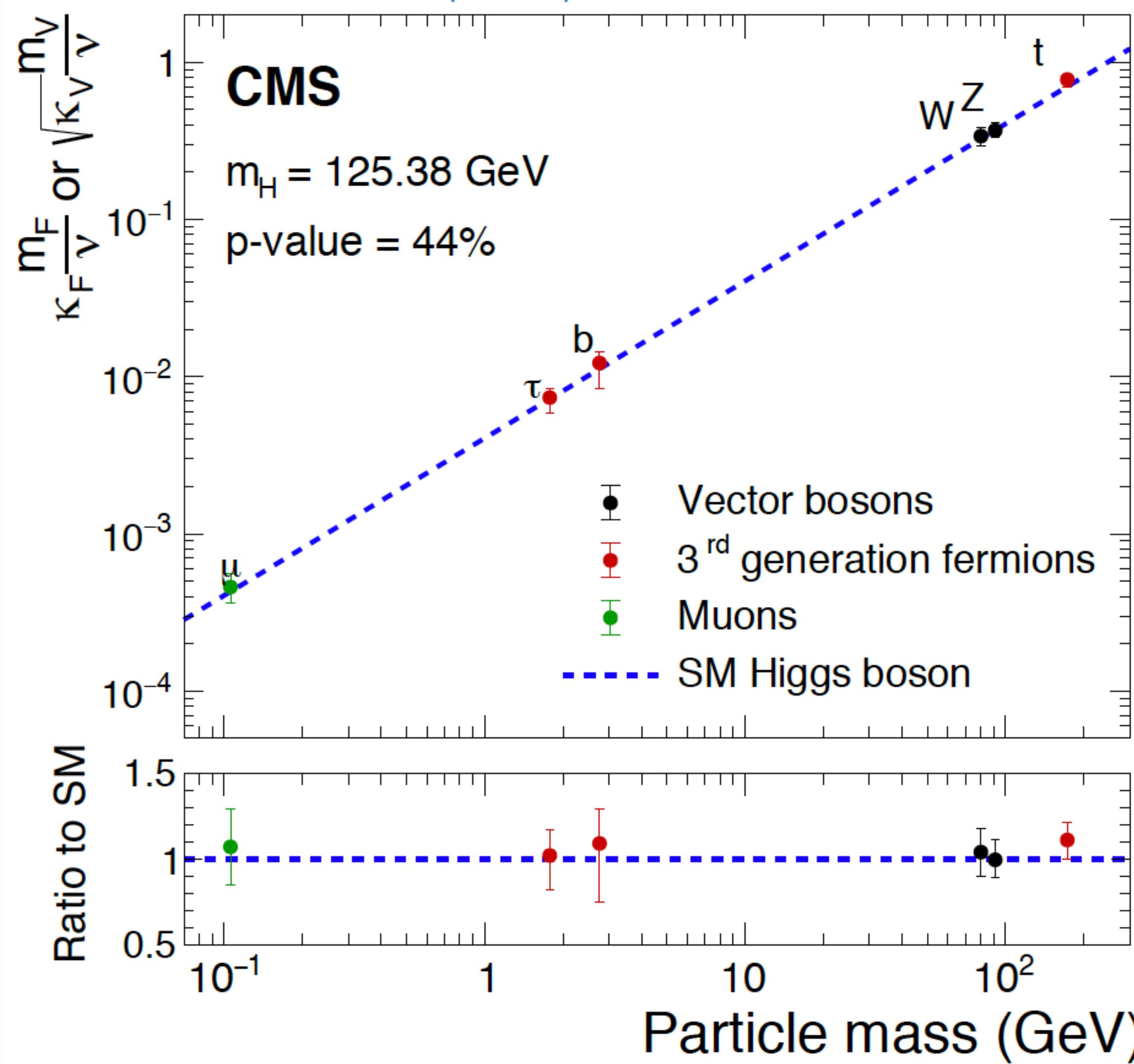
Wojciech Bizoń,^{a,b} Kirill Melnikov,^a Jérémie Quarroz^a



(a) Yukawa coupling

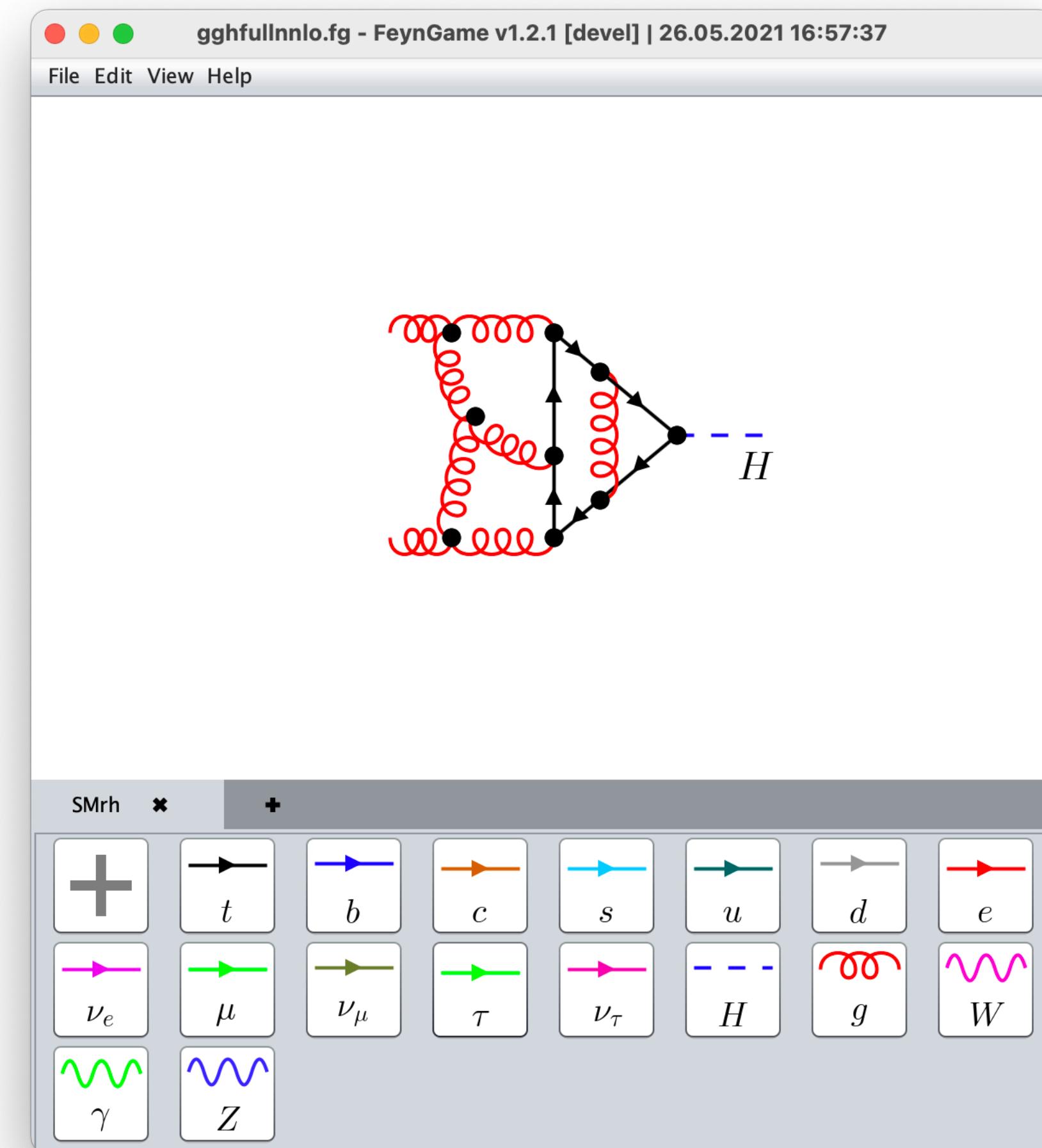
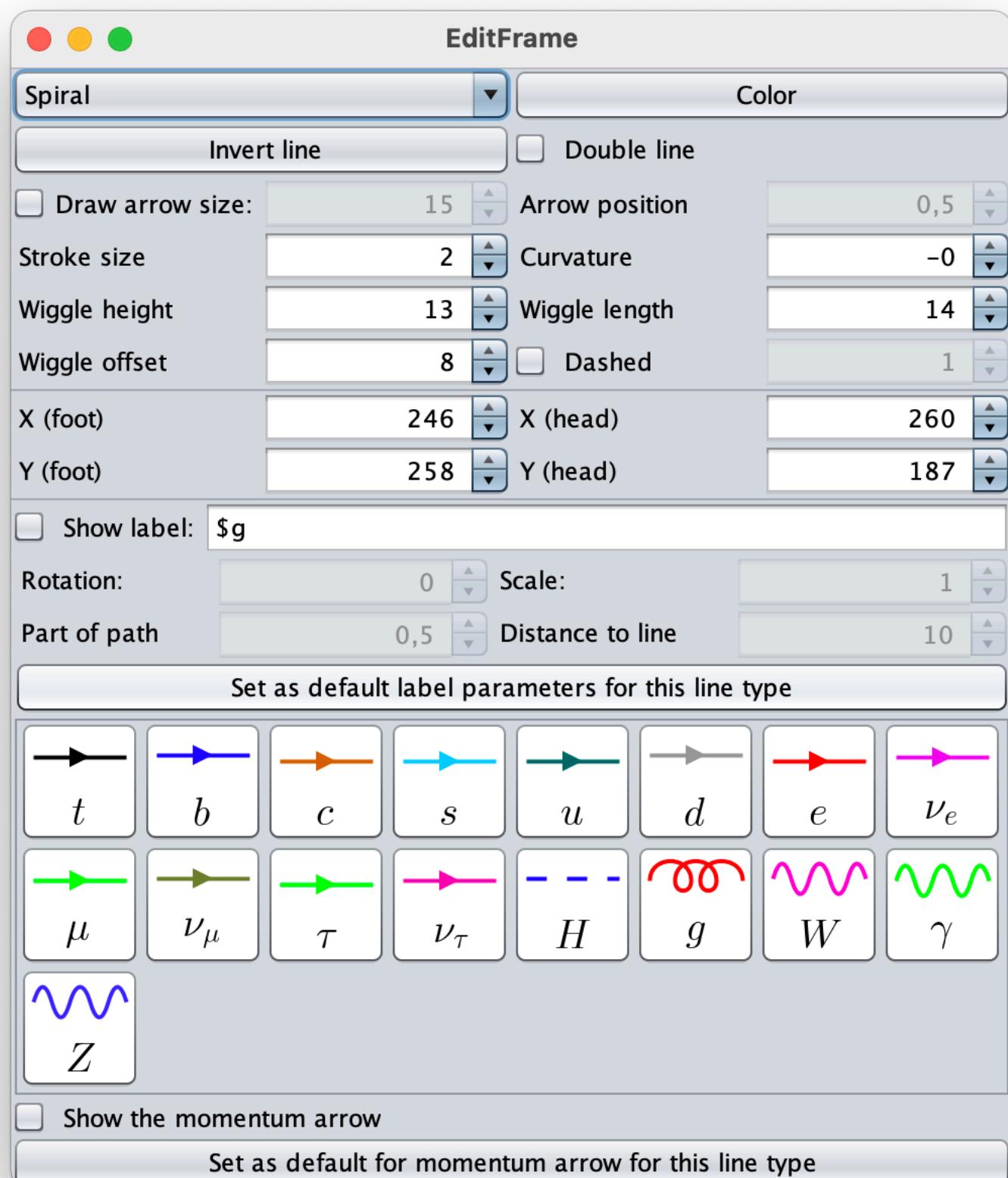


(b) Effective ggH vertex



interaction is multiplied by a factor $-\frac{1}{2}$ rather than $\frac{3}{2}$. Of course our model has too many arbitrary features for these predictions to be taken very seriously, but it is worth keeping in mind that the standard calculation⁸ of the

Weinberg '67



Thank you!