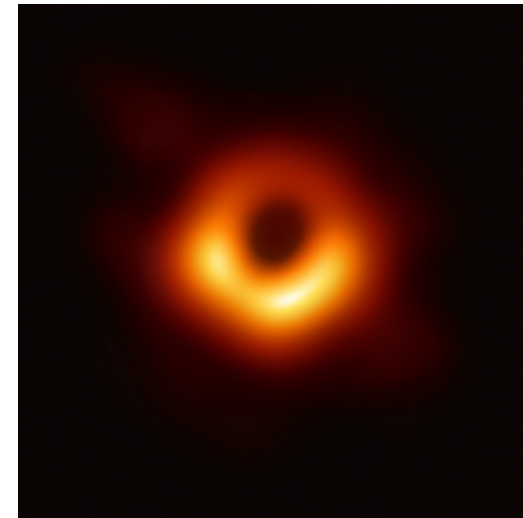
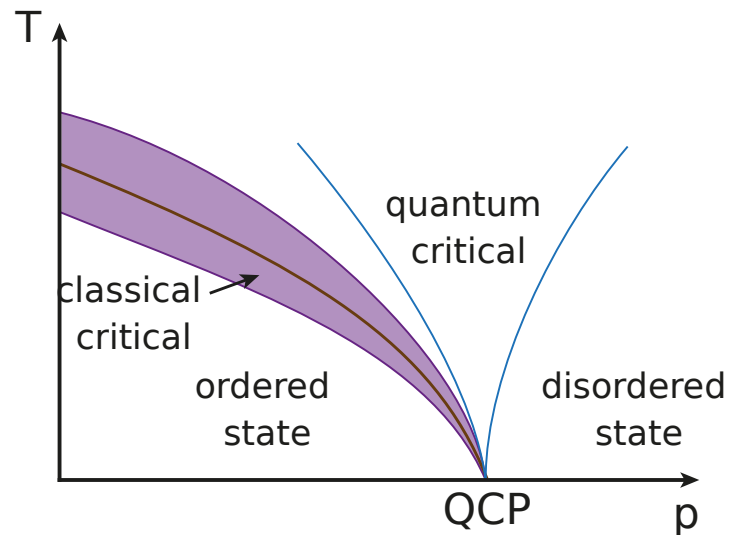


The Dirac-SYK Superconductor

Veronika Stangier, TKM

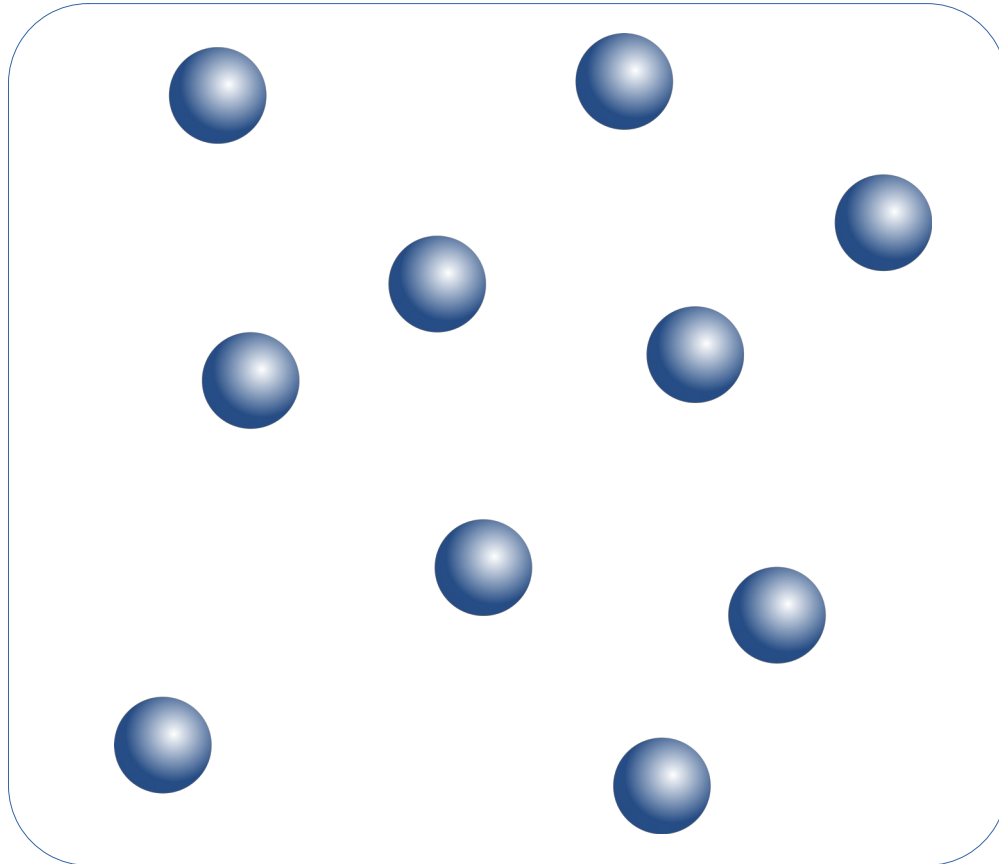


(Don't ask)

The TKM

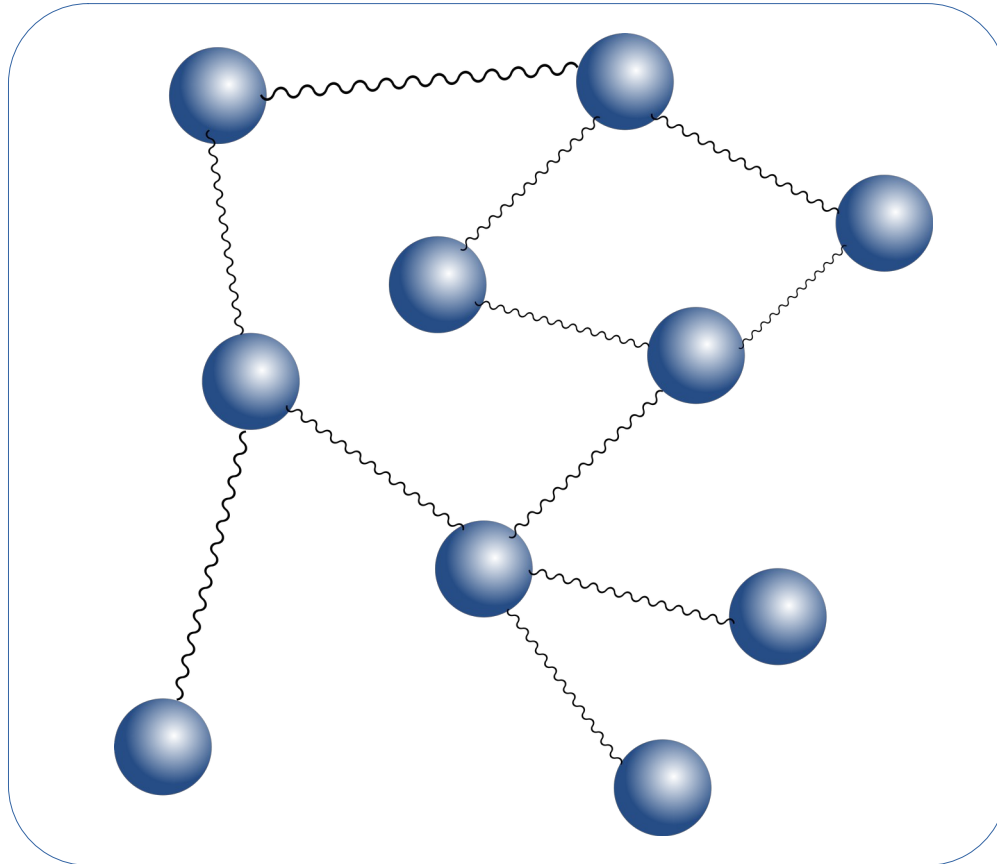


Fermi Liquid Theory



Spin : $\vec{\sigma}$, Charge: e , Momentum: k
Mass: m , Magnetic moment: μ

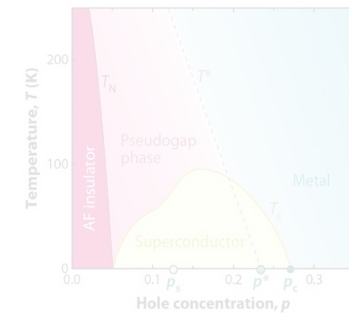
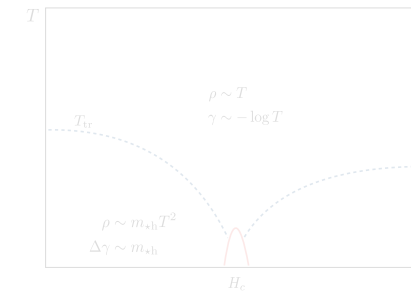
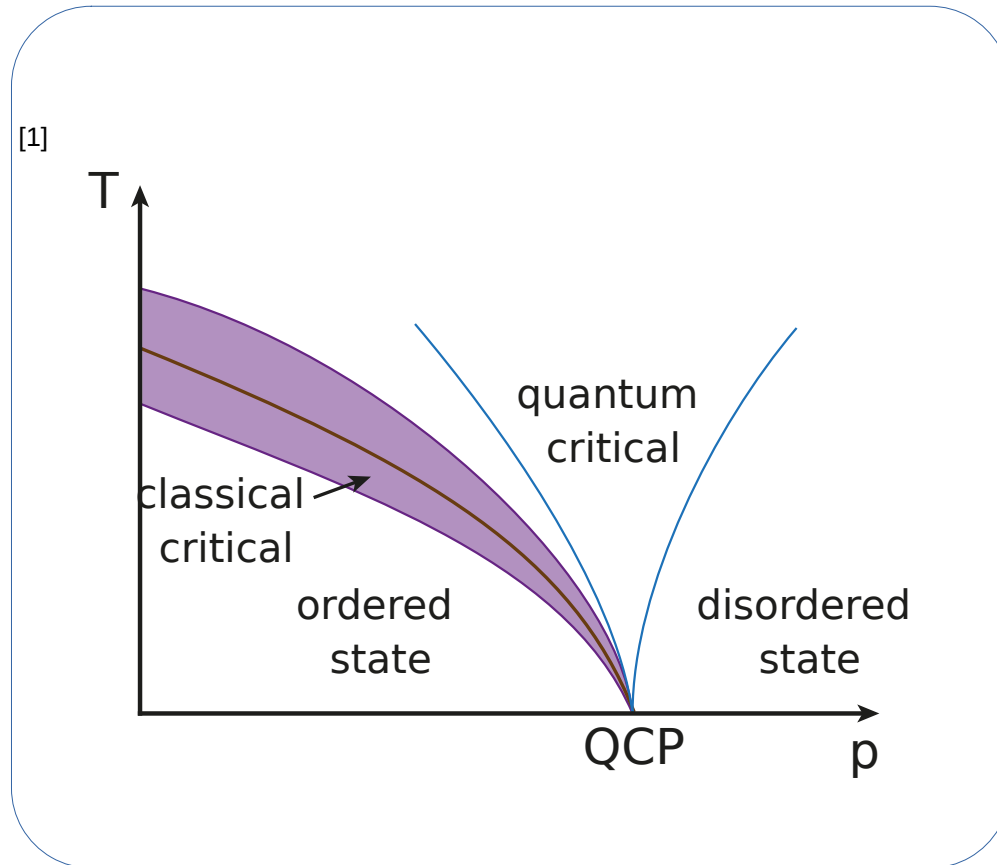
Fermi Liquid Theory



Spin : $\vec{\sigma}$, Charge: e , Momentum: k
Mass: m^* , Magnetic moment: μ^*

Resistivity : $\rho \sim T^2$
Specific heat: $c_V \sim T$

Background



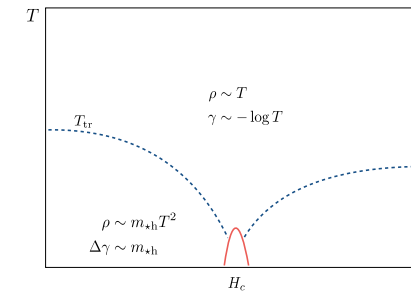
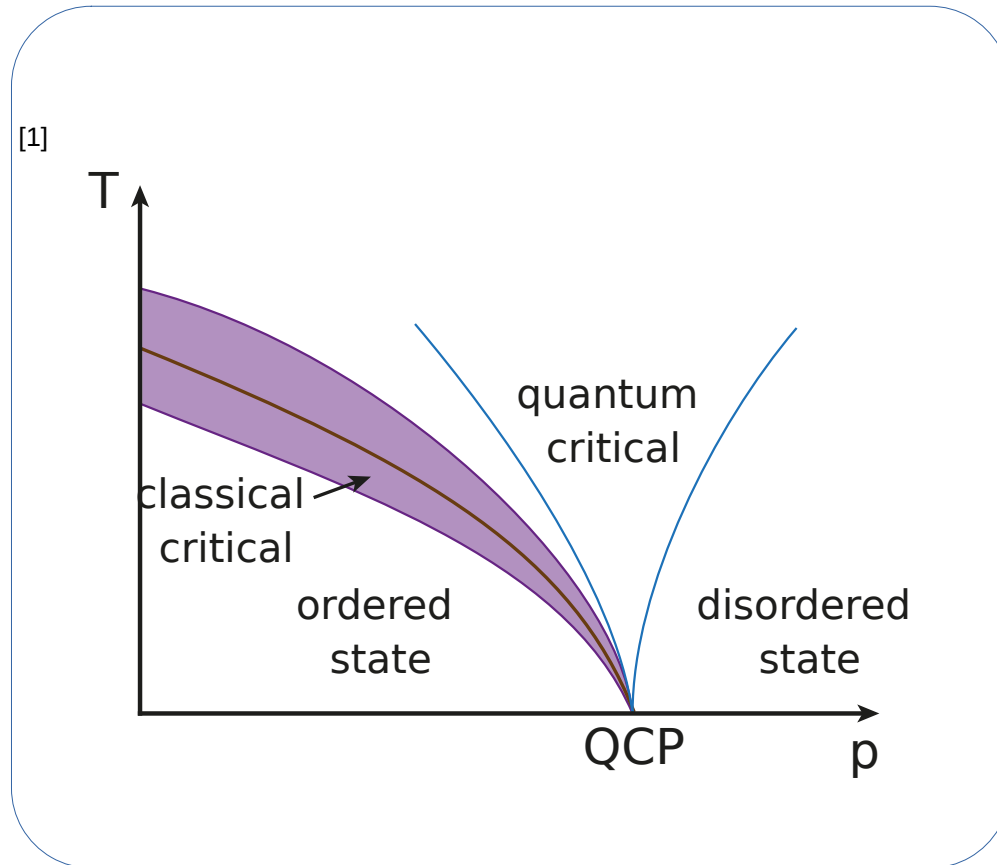
- Breakdown of the quasi-particle picture
- Anomalous transport properties
- Often: Superconductivity around QCP

[1] <https://commons.wikimedia.org/wiki/File:QuantumPhaseTransition.png>

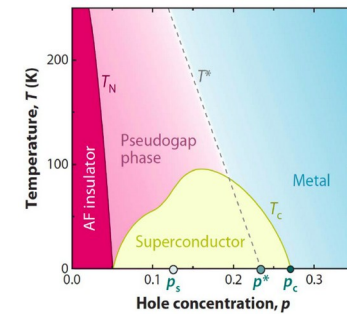
[2] C. Mousatov, S. Hartnoll, E. Berg, PNAS 117 (6) 2852-2857 (2020)

[3] L. Taillefer, arxiv:1003.2972

Background



Strange Metals ($\text{Sr}_3\text{Ru}_2\text{O}_7$) [2]



Cuprates [3]

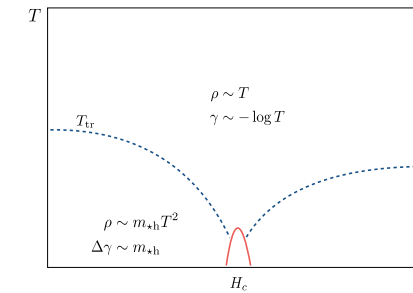
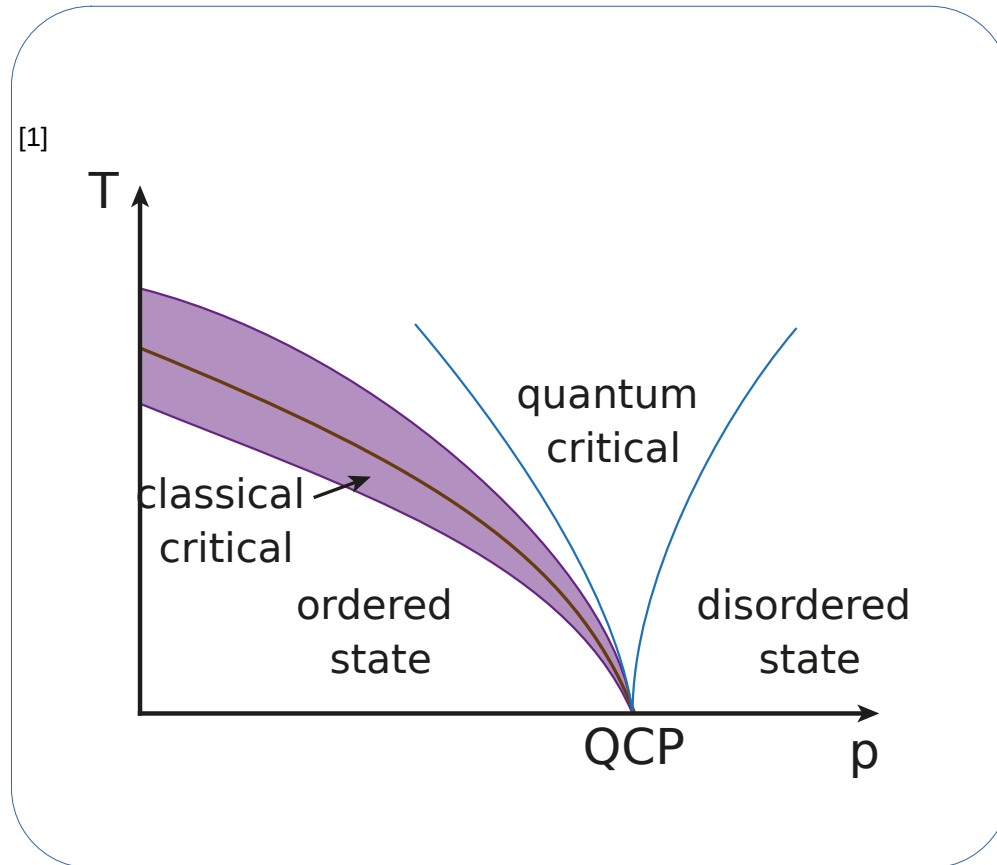
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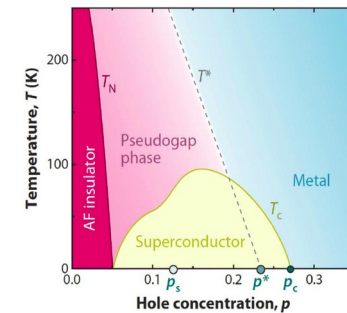
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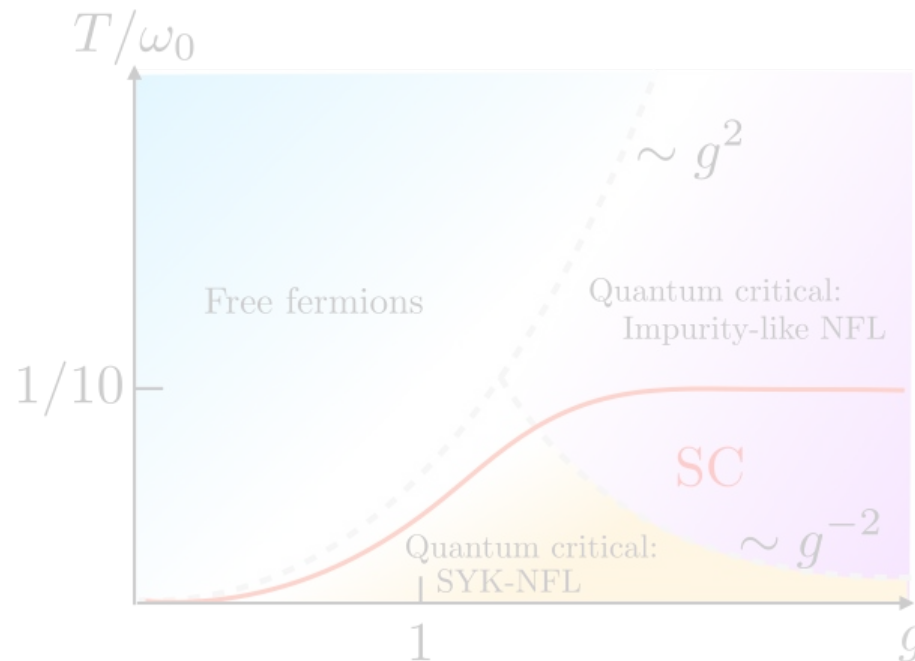
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The Yukawa-SYK model

$$H = - \sum_{i\sigma} \mu c_{i\sigma}^\dagger c_{i\sigma} + \frac{1}{2} \sum_k (\pi_k^2 + \omega_0^2 \phi_k^2) + \frac{\sqrt{2}}{N} \sum_{ijk,\sigma} g_{ij,k} c_{i\sigma}^\dagger c_{j\sigma} \phi_k$$

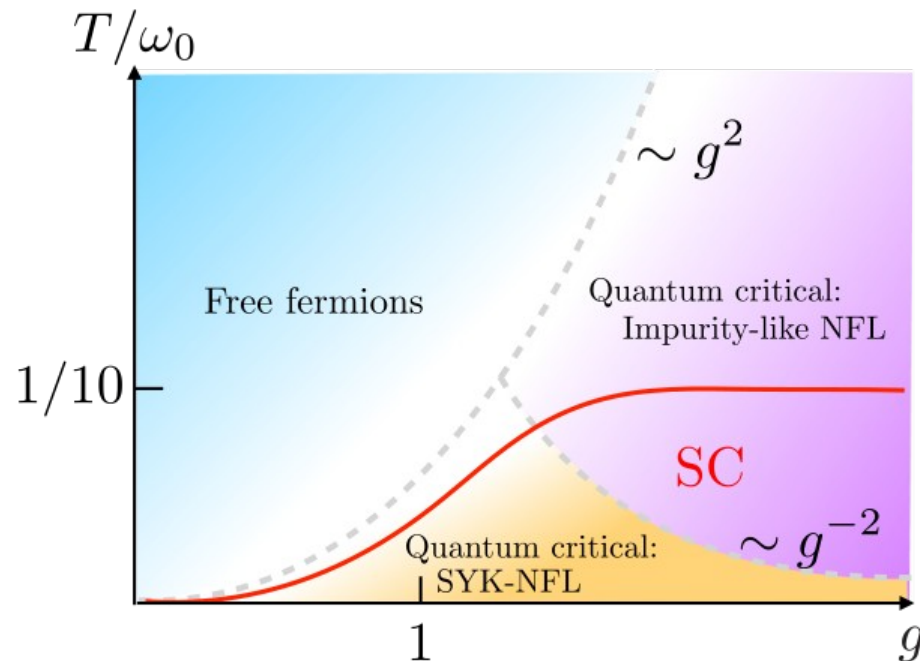


- ✓ Exactly solvable
- ✓ Critical solutions
- ✓ Superconducting
- ✗ Finite dimensions

I. Esterlis, J. Schmalian. Phys. Rev. B 100.11 (2019): 115132

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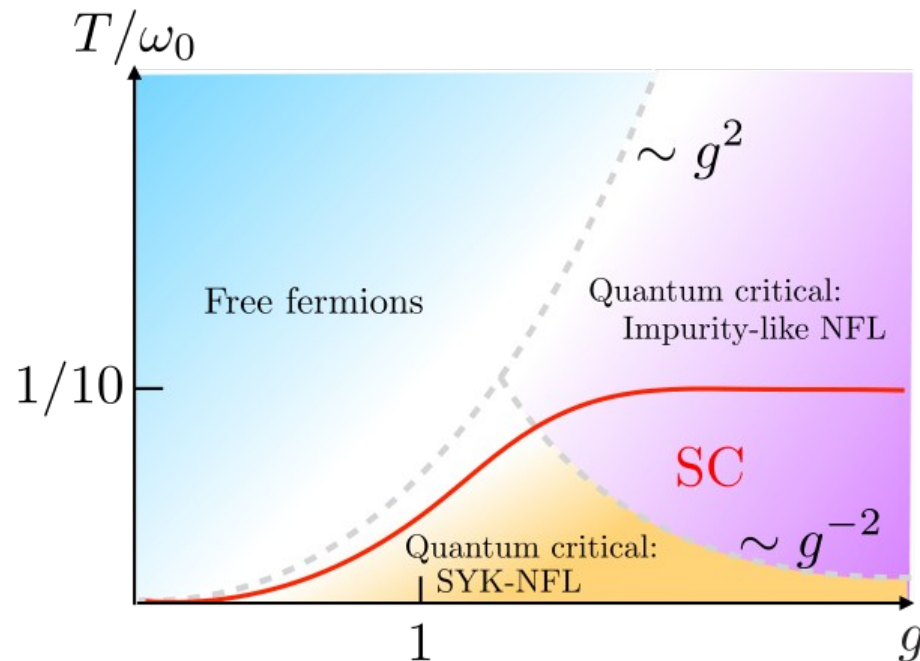


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I. Esterlis, J. Schmalian. Phys. Rev. B 100.11 (2019): 115132

The Dirac-SYK model – normal state

$$\mathcal{L} = \sum_{i=1}^N \bar{\psi}_{i\sigma} \not{\partial} \psi_{i\sigma} + \frac{1}{2} \sum_{k=1}^M \phi_k (m_0^2 - b \partial_\mu^2) \phi_k + \frac{1}{N} \sum_{ijk} g_{ijk} \bar{\psi}_{i\sigma} \hat{\Gamma}_{\sigma\sigma'} \psi_{j\sigma'} \phi_k$$

Disorder average

$$\overline{F} = -k_B T \overline{\log Z}$$

Bi-local fields

$$G(x, x') = \frac{1}{N} \sum_i \psi_i(x) \psi_i^\dagger(x')$$

$$D(x, x') = \frac{1}{M} \sum_k \phi_k(x) \phi_k(x')$$

Critical Solutions

$$(G^{-1})(x', x) = -((G_0^{-1})(x, x') - \Sigma(x, x'))$$

$$D(x', x) = \frac{1}{(D_0^{-1})(x, x') - \Pi(x, x')}$$

$$\hat{\Sigma}(x', x) = -g^2 n_s \xi [\hat{\Gamma} \hat{G}(x, x') \hat{\Gamma}'] D(x, x')$$

$$\Pi(x', x) = -g^2 \text{Tr} [\hat{\Gamma} \hat{G}(x', x) \hat{\Gamma} \hat{G}(x, x')]$$

J. Kim, E. Altman, X. Cao. Phys. Rev. B 103.8 (2021): L081113.

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Critical Solutions

(for large g)

$$G(k) \sim k_\mu \gamma^\mu k^{2\Delta+D-1}$$

$$D(q) \sim q^{-4\Delta+D}$$

$$\Sigma(k) \sim -k_\mu \gamma^\mu k^{-2\Delta+D-1}$$

$$\Pi(q) \sim q^{4\Delta-D}$$

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Disorder Average

$$S_{int} = \frac{1}{N} \sum_{ijk} g_{ijk} \bar{\psi}_{i\sigma} \hat{\Gamma}_{\sigma\sigma'} \psi_{j\sigma'} \phi_k$$

Gaussian Unitary Ensemble

$$g_{ijk}^* = g_{jik}$$

Gaussian Orthogonal Ensemble

$$g_{ijk} = g_{jik}$$

Disorder Average

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➔ No superconductivity

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➔ No superconductivity

Gaussian Orthogonal Ensemble

$$g_{ijk} = g_{jik}$$

➔ Maybe superconductivity?

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Additional bi-local fields:

$$F(x, x') = \frac{1}{N} \sum_i \psi_i(x) \psi_i(x')$$

$$\hat{\Phi}(x', x) = g^2 n_s \xi \left(\hat{\Gamma}^T \hat{F}^\dagger(x', x) \hat{\Gamma} \right) D(x, x').$$

Linearized Gap Equation

$$\hat{\Phi}(k) = C \int_q \left(\hat{\Gamma}^T \hat{G}_n^T(-q) \hat{\Phi}(q) \hat{G}_n(q) \hat{\Gamma} \right) D(q - k)$$

The Dirac-SYK model - superconducting state

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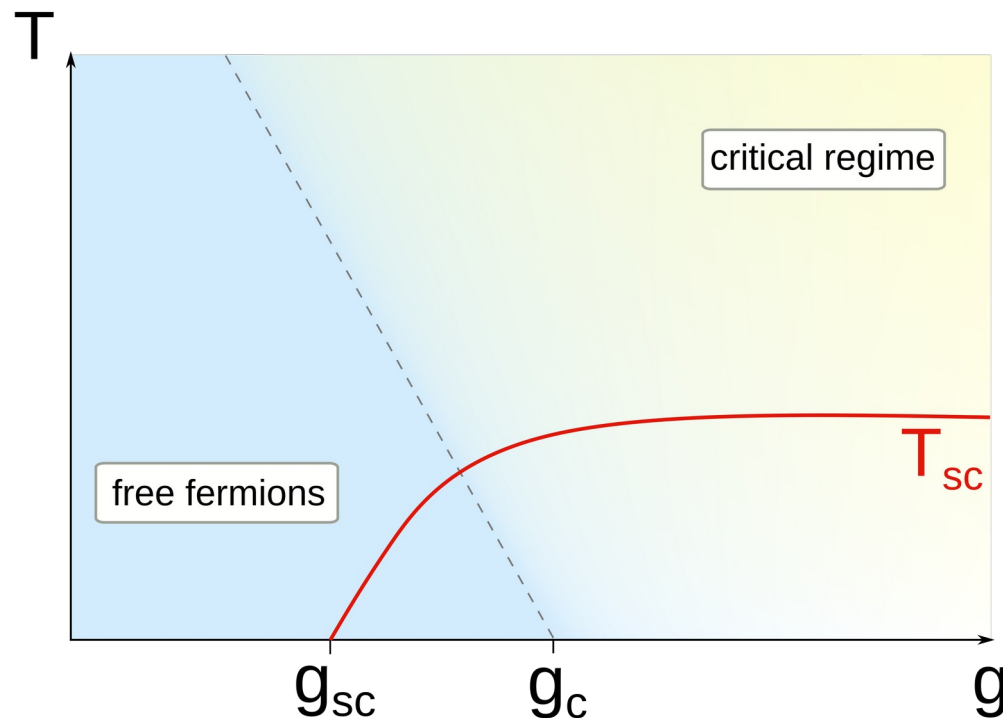
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→ Solvable by tuning fermion-to-boson ratio

The Dirac-SYK superconductor

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