

# Higgs pair production at the HL-LHC in the 2HDM: insight into trilinear Higgs couplings

Kateryna Radchenko

in collaboration with Francisco Arco, Sven Heinemeyer and Margarete Mühlleitner

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Instituto de  
Física  
Teórica  
UAM-CSIC



# 1. MOTIVATION

A Higgs boson was discovered in 2012!



What we know:

Scalar fields exist  
Higgs mechanism works

What we do **not** know:

Higgs content  
Shape of the potential  
(see Olallas' talk)

What we have to explain:

Nature of Dark Matter  
Baryon Asymmetry of the Universe  
(see Lisa's talk)

**OUR GOAL:** What can we learn about **triple Higgs couplings** (and ultimately about the Higgs potential) from measurements at the **HL-LHC**?



### 3. OUR STRATEGY

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1. Choose a versatile theoretical framework → the **2 Higgs doublet model**
2. Choose a collider → Large Hadron Collider (**LHC**)
3. Where to look for deviations? → choose **observables**:
  - Observable 1: **total di-Higgs production cross section**
    - Sensitivity to BSM triple Higgs couplings:  $\kappa_\lambda, \lambda_{hhH}$
  - Observable 2: **invariant mass distribution**
    - Information about  $\kappa_\lambda$
    - Information about resonant production:  $m_H, \Gamma_H$
    - Experimental challenges

# 4. THE THEORETICAL FRAMEWORK: 2HDM

[Santos, Barroso: [arXiv:hep-ph/9701257](https://arxiv.org/abs/hep-ph/9701257)]

CP conserving 2HDM with two complex doublets with a softly broken  $\mathbb{Z}_2$  symmetry ( $\Phi_1 \rightarrow \Phi_1; \Phi_2 \rightarrow \Phi_2$ )  
(only Type I discussed here)



Pontental: 
$$V_{2\text{HDM}} = m_1^2 |\phi_1|^2 + m_2^2 |\phi_2|^2 - m_3^2 (\phi_2^\dagger \phi_1 + \phi_1^\dagger \phi_2) + \frac{\lambda_1}{2} |\phi_1|^4 + \frac{\lambda_2}{2} |\phi_2|^4 + \lambda_3 |\phi_1|^2 |\phi_2|^2 + \lambda_4 |\phi_2^\dagger \phi_1|^2 + \frac{\lambda_5}{2} ((\phi_2^\dagger \phi_1)^2 + (\phi_1^\dagger \phi_2)^2)$$

Free parameters:  $m_h, m_A, m_H, m_{H^\pm}, m_{12}^2, v, \cos(\beta - \alpha), \tan\beta$

$$\tan\beta = v_2/v_1$$

$$v^2 = v_1^2 + v_2^2 \sim (246 \text{ GeV})^2$$

→ couplings to fermions and gauge bosons of the SM-like Higgs change.

(we need to be careful that all current data can be reproduced!) [Arco, Heinemeyer, Herrero: [arXiv2003.12684](https://arxiv.org/abs/2003.12684)]

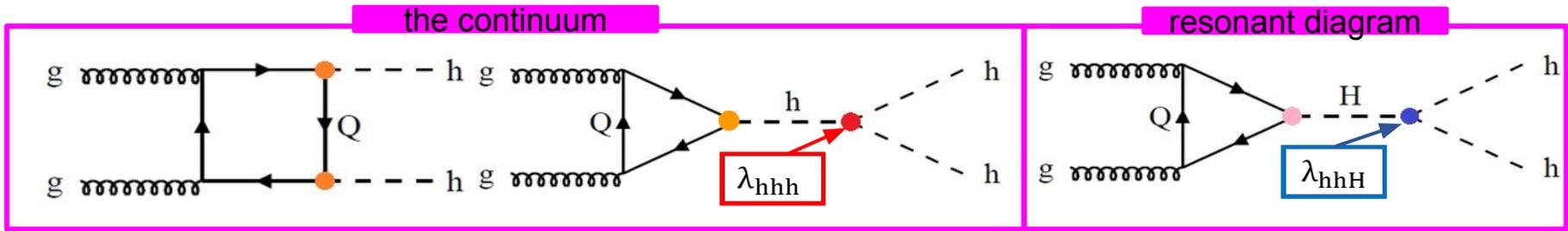
→ many more possibilities for “triple Higgs couplings” (we will look at two:  $\kappa_\lambda \in [-0.5, 1.7], \lambda_{\text{hhH}} \in [-1.7, 1.6]$ )

These lead to **different phenomenology** w.r.t the SM but also the contribution of the **heavy Higgses** in the loops.

# 5. DI HIGGS PRODUCTION

[Plehn, Spira, Zerwas : [arXiv:hep-ph/9603205](https://arxiv.org/abs/hep-ph/9603205)]

- Triple Higgs couplings can be accessed through Higgs pair production
- The dominant process at a hadron collider is gluon fusion involving a quark loop



$\sigma_{SM} \sim 38 \text{ fb at NLO}$

Diagrams that exist in the SM:  
They have a negative interference

Diagrams that are sensitive  
to triple Higgs couplings

~ 1 out of  $10^9$  events in the LHC is a Higgs  
~ 1 out of  $10^{13}$  events in the LHC is a Higgs pair

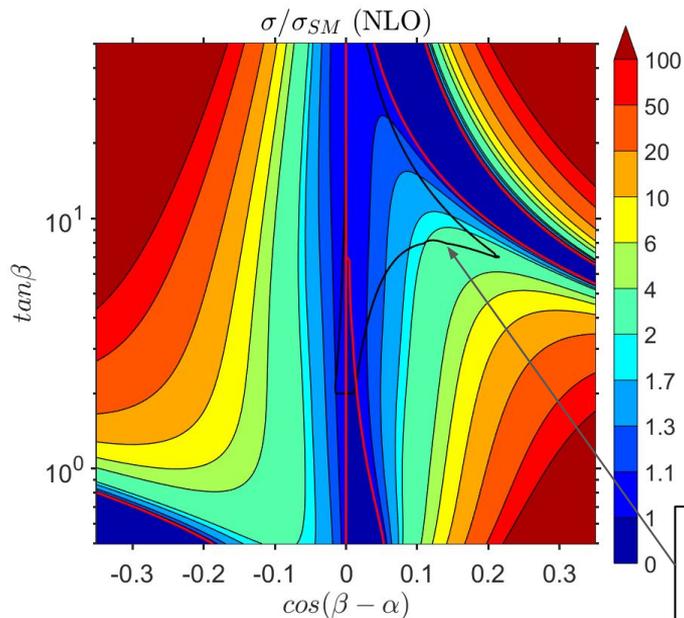
→ All calculations were done using a modified version of the code HPAIR

[Abouabid, Arhrib, Azevedo, El Falaki, Ferreira, Muhlleitner, Santos: [arXiv:hep-ph/2112.12515](https://arxiv.org/abs/hep-ph/2112.12515)]

# 6. TOTAL DI-HIGGS PRODUCTION CROSS SECTION

$$m_H = m_A = m_{H^\pm} = 1000 \text{ GeV}$$

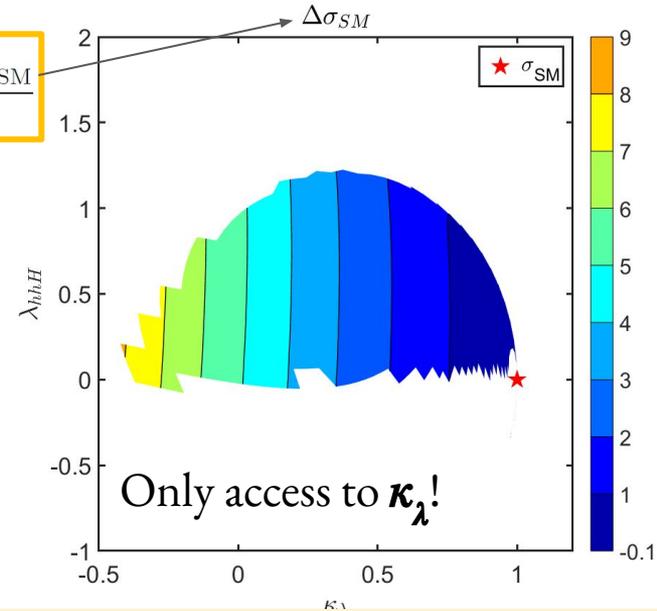
$$m_{12}^2 = (m_H^2 \cos^2 \alpha) / \tan \beta$$



$$\Delta\sigma_{SM} \equiv \frac{\overline{xs}_{2HDM} - \overline{xs}_{SM}}{\overline{\delta xs}}$$

$\overline{\delta xs} = \overline{xs}/4.5$   
 projected significance in standard deviation of the total di-Higgs production:  $4.5\sigma$

Allowed region inside the black contour!



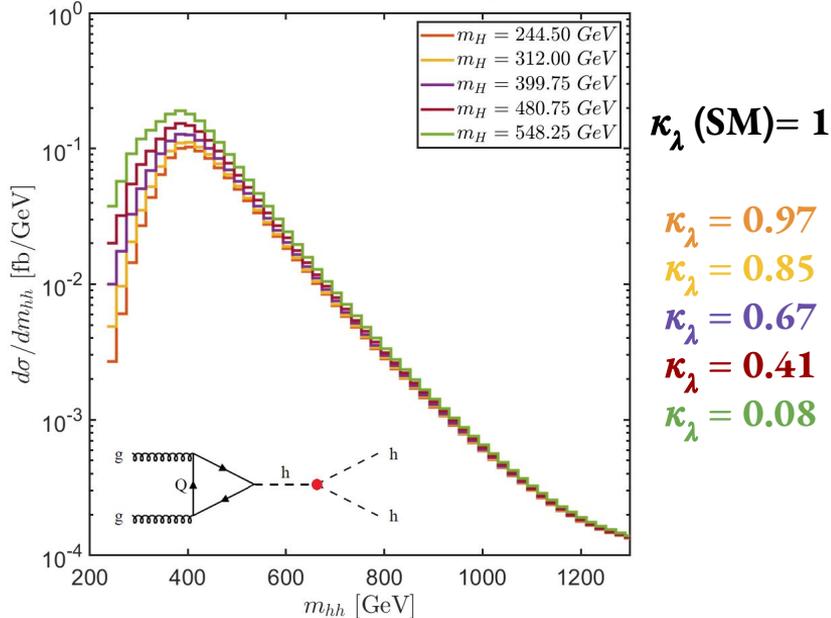
Only access to  $\kappa_\lambda$ !

- NLO QCD corrections implemented in HPAIR
- **Largest enhancements** inside the allowed region (black contour)  $\sim 3\sigma_{SM} \rightarrow$  due to deviations in  $\kappa_\lambda$
- **Expected sensitivity** to the deviation of the cross section: up to  $8\sigma$  away from the SM

# 7. INVARIANT MASS DISTRIBUTION

## 1) Non-resonant production:

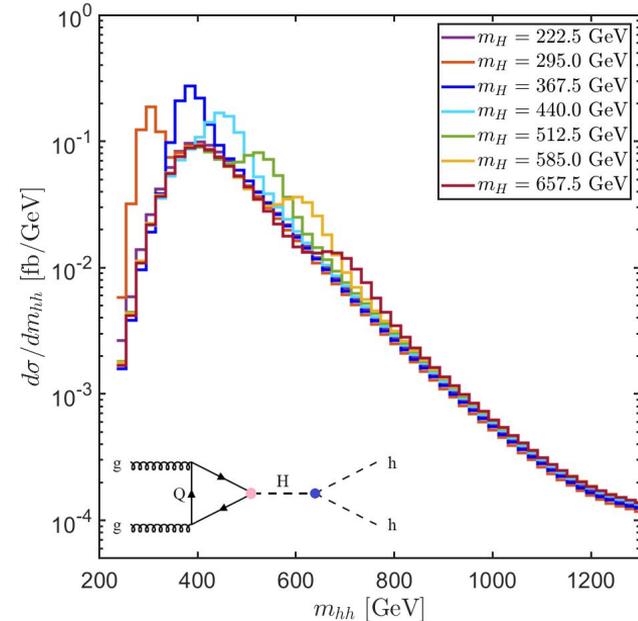
the contribution of the heavy Higgses is negligible:



- Larger sensitivity to deviations in  $\lambda_{hh}$  in the low  $m_{hh}$  region

## 2) Resonant production:

the contribution of the heavy Higgses is important:



- Sensitivity to the properties of the BSM resonance:  $m_H, \Gamma_H$

# CONCLUSIONS

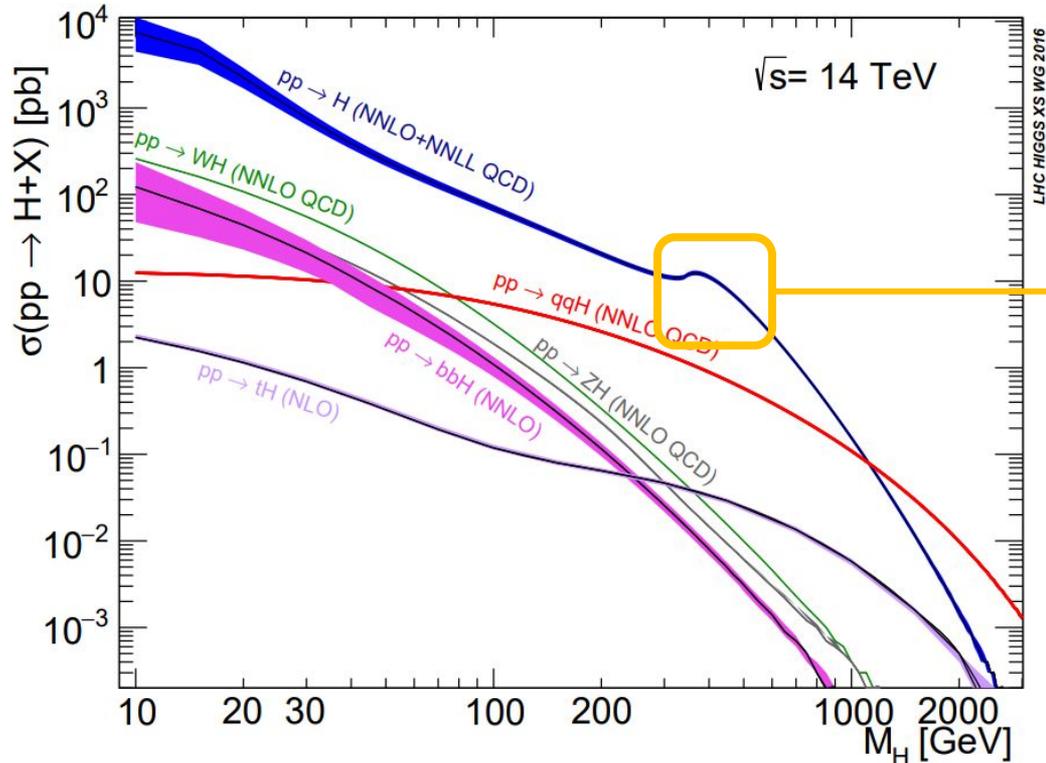
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- Many Beyond the Standard Model theories propose **extended Higgs sectors** and explain some open problems of the Standard Model.
- The next step to establish the Higgs potential is the measurement of **triple Higgs couplings**.
- Deviations of this parameter w.r.t. the Standard Model can affect the **Higgs pair production**.
- Measuring the **total production cross section** is not enough to disentangle the effects from deviations in the **Higgs triple self-interactions** and **contribution of additional particles** → **invariant mass distributions** are a complementary and promising avenue.

THANK YOU FOR YOUR ATTENTION!



# SINGLE HIGGS PRODUCTION



Top pair threshold  $\rightarrow$  gives a hint on the results for Higgs pair production

[LHC Higgs Working Group:  
[CERN Yellow Report 4](#)]

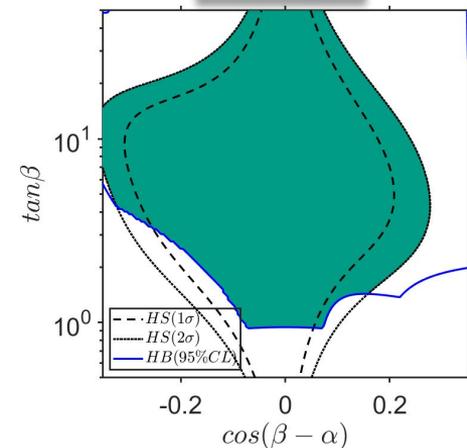
# BENCHMARK PLANES

[Arco, Heinemeyer, Herrero [arXiv2005.10576](#)]

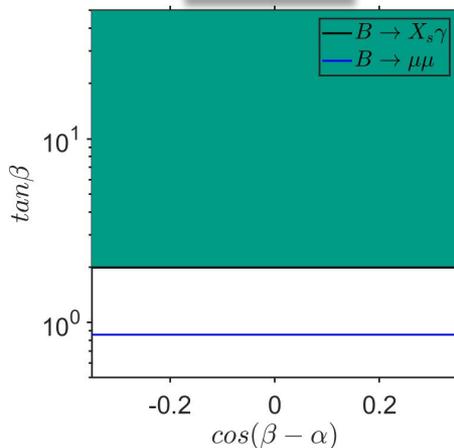
We scan the parameter space fixing all but two parameters and look for large deviations from the SM in the resulting benchmark planes.

Colored area is allowed

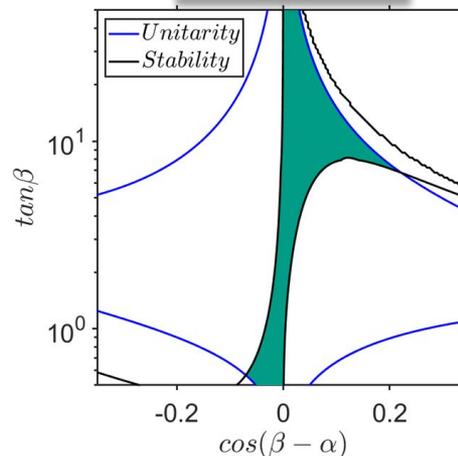
Collider



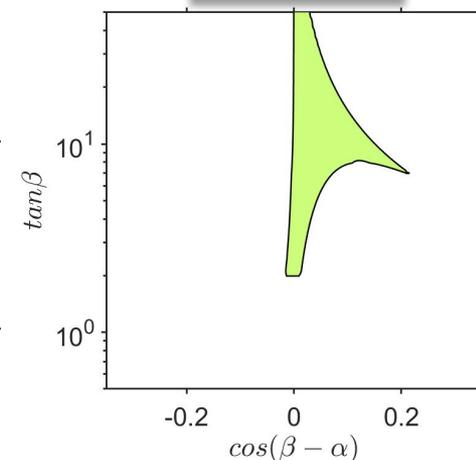
Flavour



Theoretical



Total

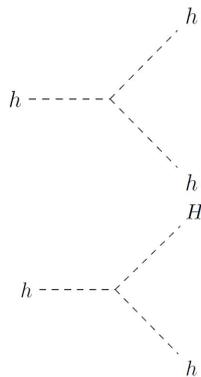


**PLANE 1:** Type I,  $m_H = m_A = m_{H^\pm} = 1000$  GeV,  $m_{12}^2 = (m_H^2 \cos^2 \alpha) / \tan \beta$   
Free parameters:  $\cos(\beta - \alpha)$ ,  $\tan(\beta)$

[Ren: [arXiv:1706.05980](#)]

→ Special equation for  $m_{12}$  that enlarges the area allowed by theoretical constraints.

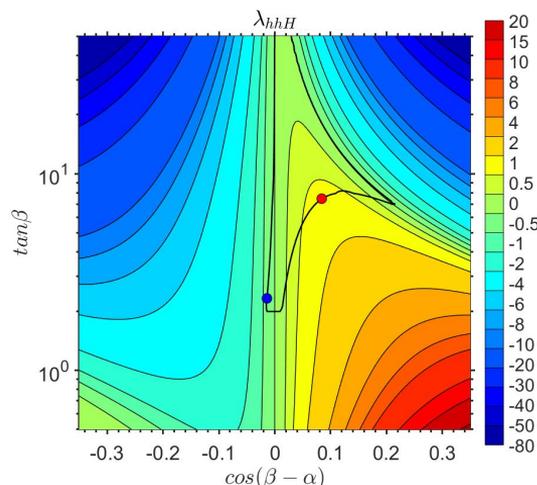
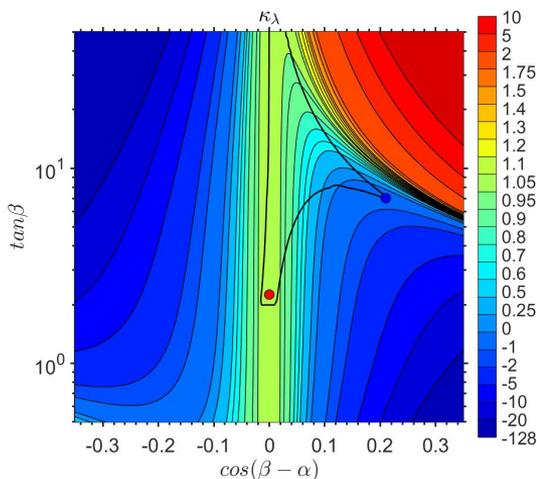
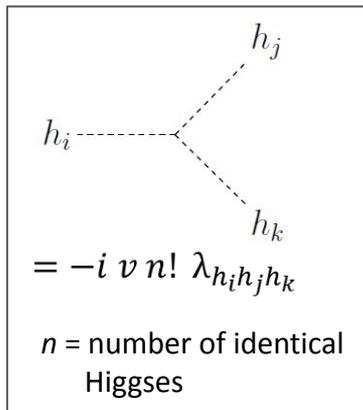
# FEYNMAN RULES FOR 2HDM TRIPLE HIGGS COUPLINGS



$$\lambda_{hhh} = \frac{1}{2v^2} \left\{ m_h^2 s_{\beta-\alpha}^3 + (3m_h^2 - 2\bar{m}^2) c_{\beta-\alpha}^2 s_{\beta-\alpha} + 2 \cot 2\beta (m_h^2 - \bar{m}^2) c_{\beta-\alpha}^3 \right\}$$

$$\lambda_{hhH} = \frac{-c_{\beta-\alpha}}{2v^2} \left\{ (2m_h^2 + m_H^2 - 4\bar{m}^2) s_{\beta-\alpha}^2 + 2 \cot 2\beta (2m_h^2 + m_H^2 - 3\bar{m}^2) s_{\beta-\alpha} c_{\beta-\alpha} - (2m_h^2 + m_H^2 - 2\bar{m}^2) c_{\beta-\alpha}^2 \right\}.$$

Notation:

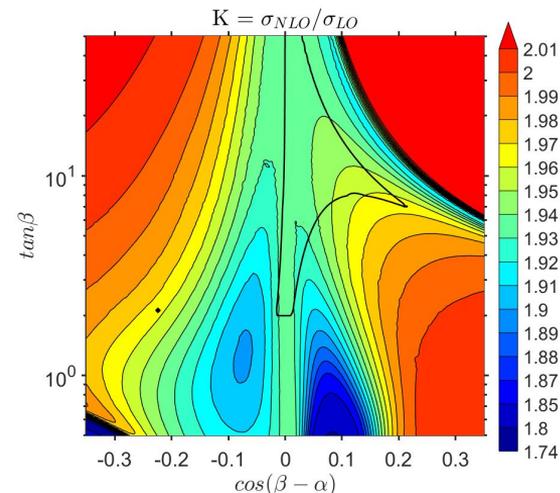
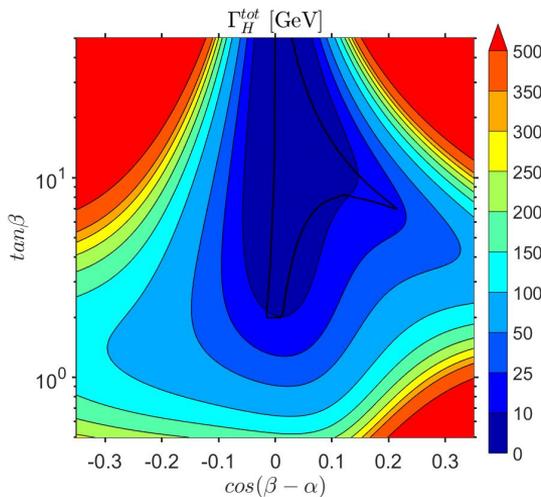
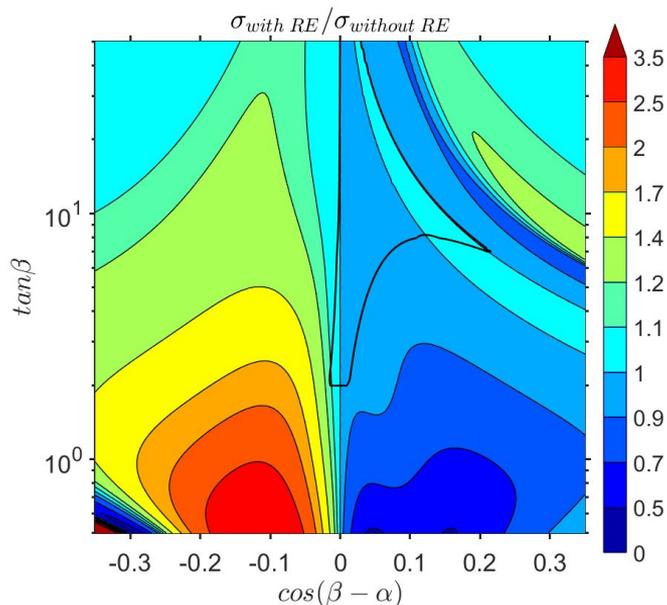


$$\bar{m}^2 = \frac{m_{12}^2}{\sin(\beta)\cos(\beta)}$$

Deviations from the alignment limit:

**MIN.**  $\kappa_\lambda \sim -0.4$   
 $\tan \beta \sim 7$   
 $\cos(\beta - \alpha) \sim 0.2$   
**MAX.**  $\lambda_{hhH} \sim 1.2$   
 $\tan \beta \sim 7$   
 $\cos(\beta - \alpha) \sim 0.1$

# MORE ABOUT TOTAL CROSS SECTION



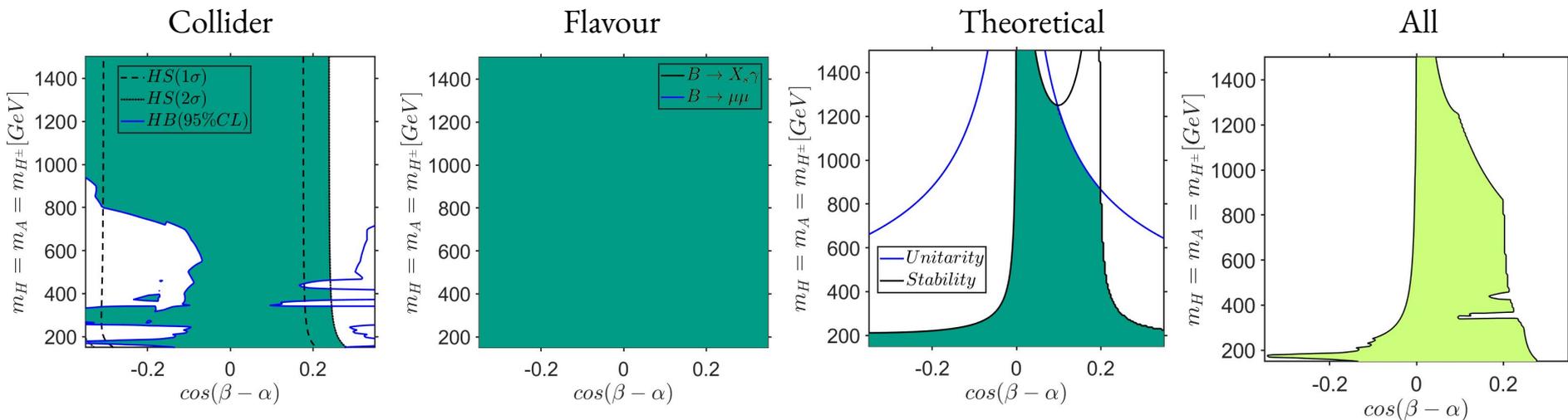
$$m_H = m_A = m_{H^\pm} = 1000 \text{ GeV}$$

$$m_{12}^2 = (m_H^2 \cos^2 \alpha) / \tan \beta$$

- Without resonant enhancement  $\rightarrow$  excluding the resonant diagram
- Resonant contribution is negligible in the allowed region
- Total decay width  $\sim 25$  GeV at the tip

# CONSTRAINTS P3

BP: Type I,  $\cos(\beta - \alpha) = 0.1$ ,  $\tan \beta = 10$ ,  $m_{12}^2 = m_H^2 \cos^2 \alpha / \tan \beta$ ,  $m_H = m_A = m_{H^\pm}$



# INVARIANT MASS DISTRIBUTION: EFFECTS OF DEVIATIONS IN $\kappa_\lambda$

BP: Type I,  $\cos(\beta - \alpha) = 0.1$ ,  $\tan \beta = 10$ ,  $m_{12}^2 = m_H^2 \cos^2 \alpha / \tan \beta$ ,  $m_H = m_A = m_{H^\pm}$

Prediction for BSM couplings:

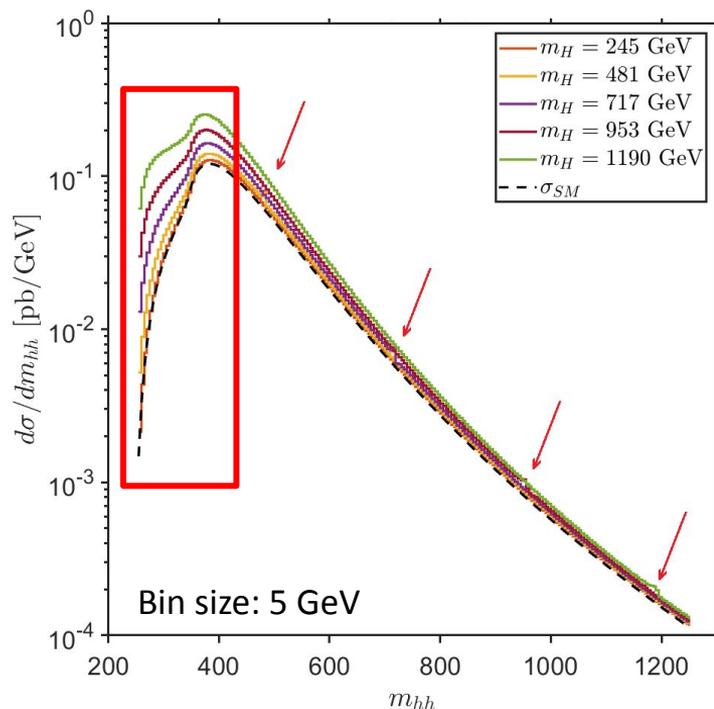
$$\kappa_\lambda = 0.97, \lambda_{hhH} = 0.05$$

$$\kappa_\lambda = 0.85, \lambda_{hhH} = 0.19$$

$$\kappa_\lambda = 0.67, \lambda_{hhH} = 0.42$$

$$\kappa_\lambda = 0.41, \lambda_{hhH} = 0.74$$

$$\kappa_\lambda = 0.08, \lambda_{hhH} = 1.15$$



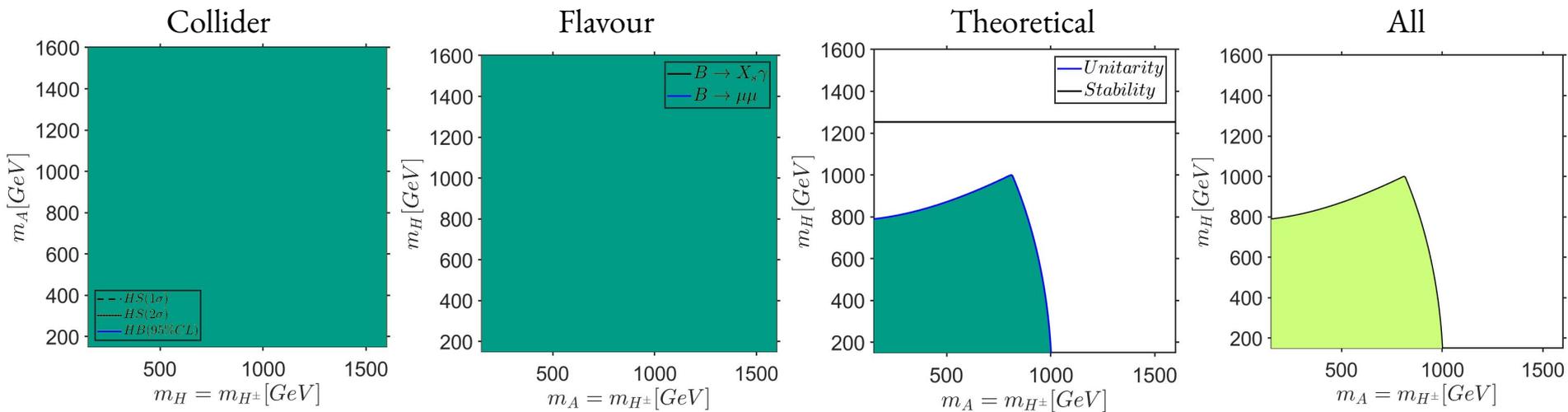
- Larger sensitivity to  $\kappa_\lambda$  in the low  $m_{hh}$  region.

- Resonant contribution very suppressed due to very small top Yukawa  $\xi_H^t \sim 10^{-4}$ .

$$\xi_H^t = \cos(\beta - \alpha) - \sin(\beta - \alpha) / \tan(\beta)$$

# CONSTRAINTS P8

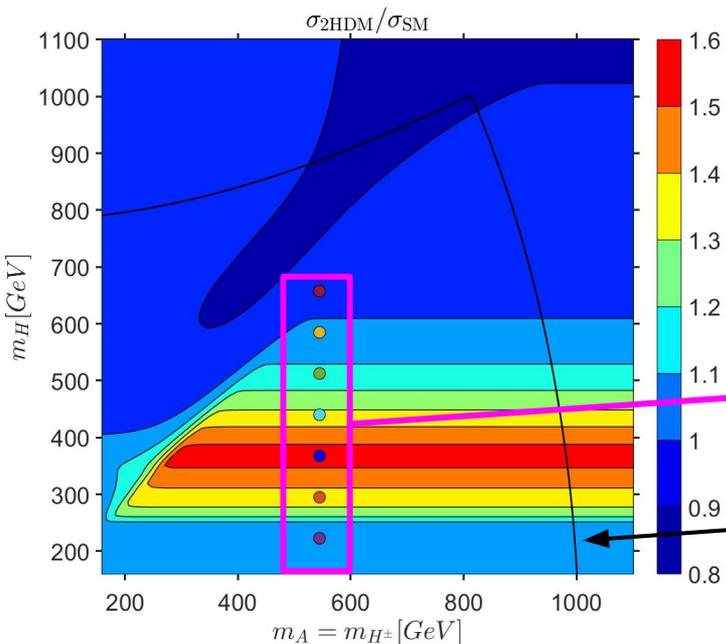
BP: Type I,  $\cos(\beta - \alpha) = 0.2$ ,  $\tan \beta = 10$ ,  $m_{12}^2 = m_H^2 \cos^2 \alpha / \tan \beta$



# EFFECT OF THE MASS OF THE HEAVY HIGGS

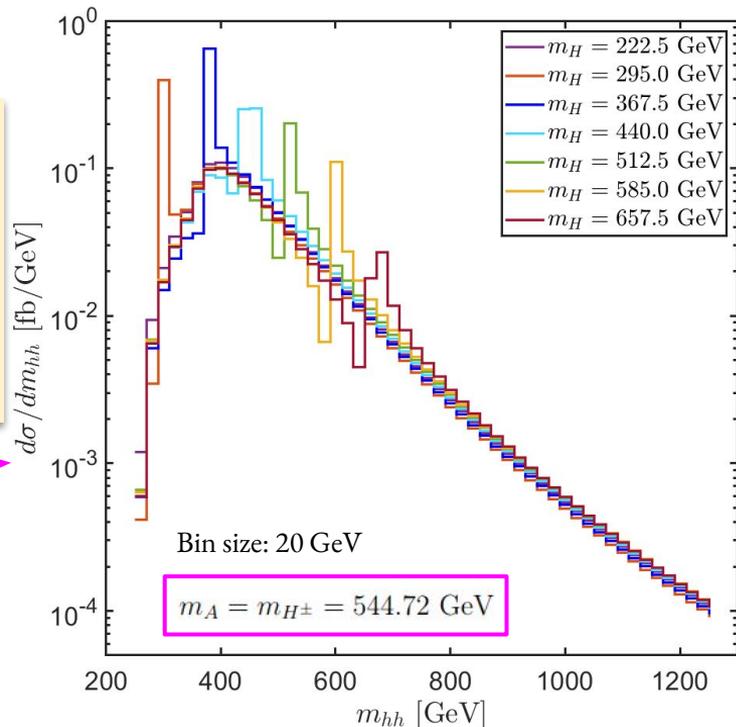
- We vary the mass of the heavy Higgs boson leaving the rest of the parameters of the model fixed.

BP: Type I,  $\cos(\beta - \alpha) = 0.2$ ,  $\tan \beta = 10$ ,  $m_{12}^2 = m_H^2 \cos^2 \alpha / \tan \beta$



Enhancement in the total cross section is resonance dominated. **Location** of the resonance is related to the mass of **H**.

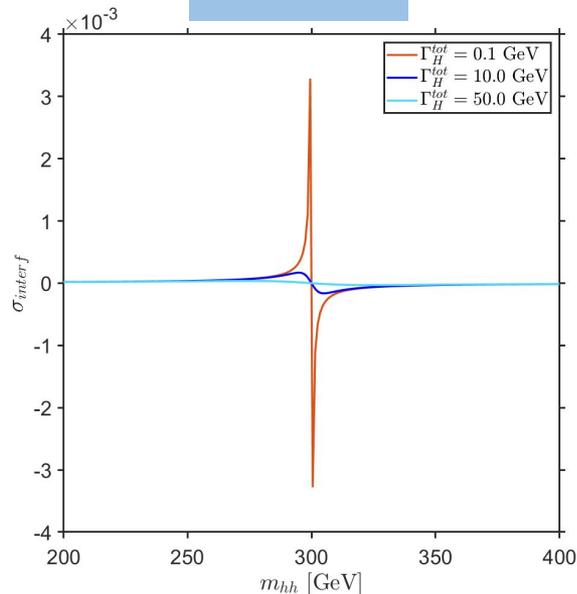
Allowed region inside the black contour.



# EFFECT OF THE TOTAL DECAY WIDTH

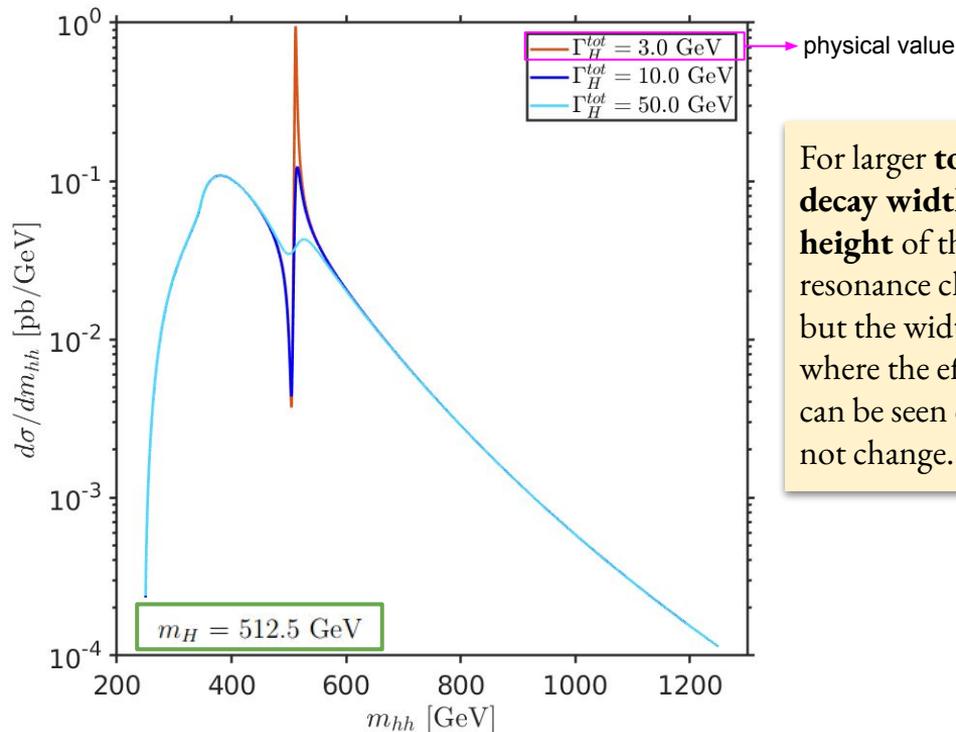
$$\frac{1}{Q^2 - M_{h/H}^2 + i\Gamma_{h/H}M_{h/H}}$$

## Toy model



$$\sigma_{\text{interf}} \propto \frac{Q^2 - m_H^2}{(Q^2 - m_H^2)^2 + m_H^2 \Gamma_H^2}$$

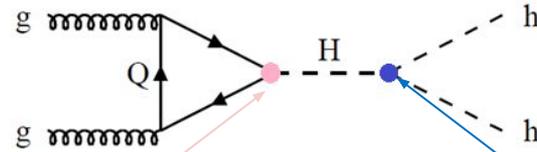
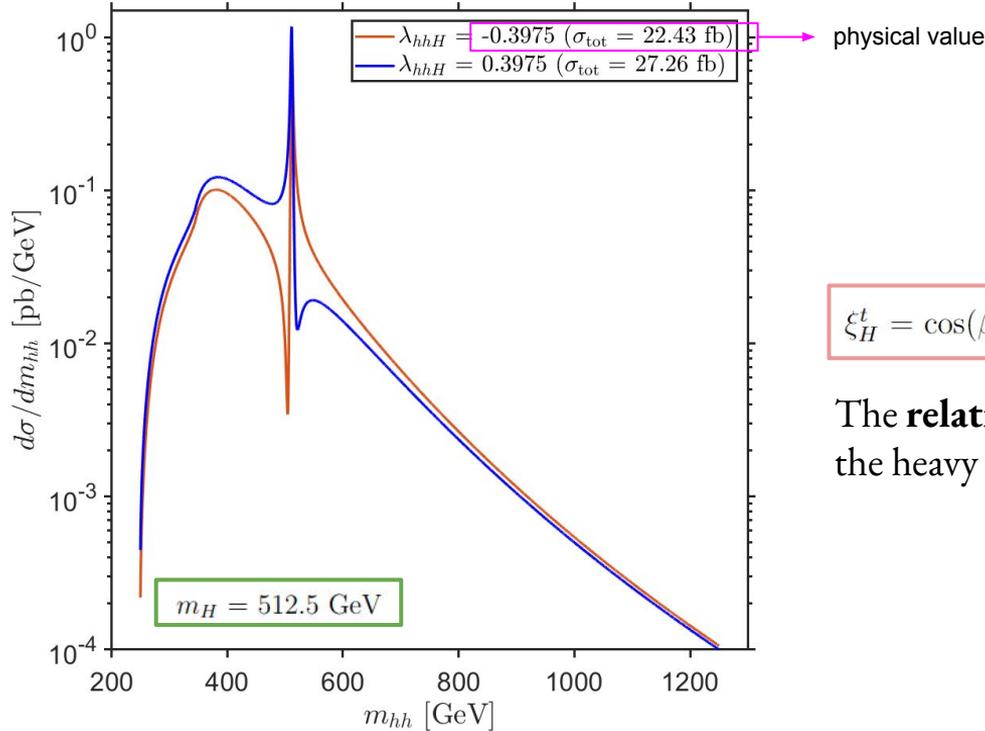
- For the green point of the previous benchmark plane we artificially change the total decay width of the heavy Higgs H:



For larger **total decay widths** the **height** of the resonance changes but the width where the effect can be seen does not change.

# EFFECT OF THE COUPLINGS

- What is the effect of the couplings involved in the resonant diagram on the invariant mass distributions ?



$$\xi_H^t = \cos(\beta - \alpha) - \sin(\beta - \alpha) / \tan(\beta) = 0.104$$

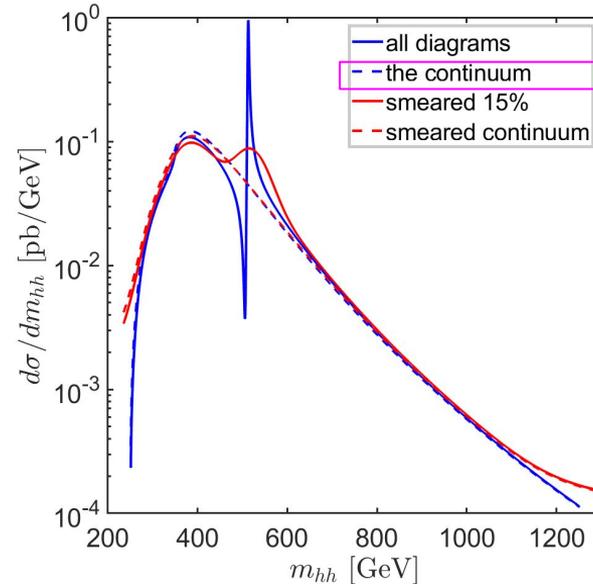
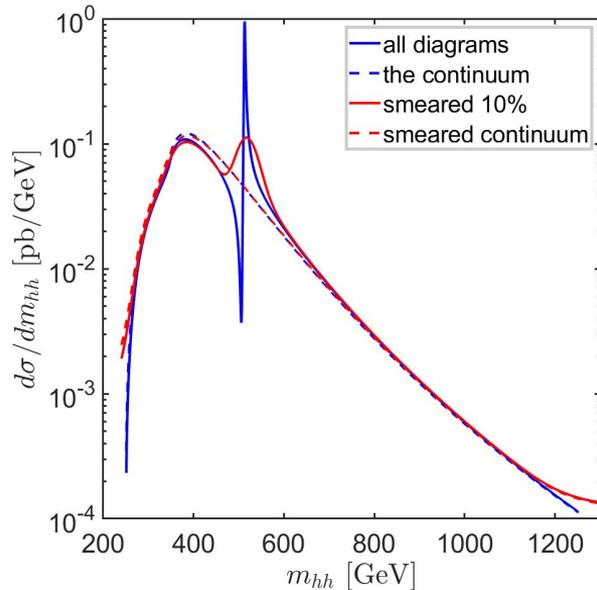
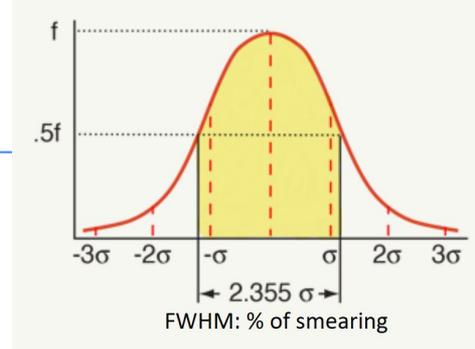
$$\lambda_{hhH}$$

The **relative sign** of the top Yukawa and the BSM coupling to the heavy Higgs gives a **structure** to the resonance:

sign ( $\lambda_{hhH} \cdot \xi_H^t$ )	structure
+	peak-dip
-	dip-peak

# EXPERIMENTAL CHALLENGES: SMEARING

- Differential cross section measurements are affected by the finite resolution of particle detectors → observed spectrum is “**smeared**”.
- We try to mimic this effect by artificially smearing the theoretical prediction introducing **Gaussian uncertainties** in the invariant mass.



→ box diagram + SM-like Higgs exchange