## Status and Prospects of Nonleptonic *B* Decays

#### K. Keri Vos

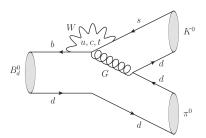
Maastricht University & Nikhef

## Puzzles in nonleptonic decays



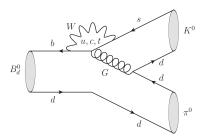
## The challenge of nonleptonic *B* decays

- Nonleptonic decays are important probes of CP violation
  - Direct CP violation due to different strong and weak phases
  - Mixing-induced CP violation in neutral decays probe mixing phase  $\phi_{d,s}$
  - Sensitivity to NP in loops (penguins)
- CP violation in the SM is too small and peculiar!
  - CKM CP violating effects only from flavour changing currents
  - Flavour diagonal CP violation tiny in SM (EDMs)
  - Large CP asymmetries with processes with tiny BRs and vice versa



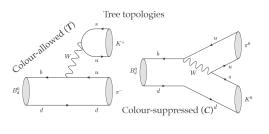
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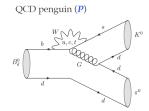
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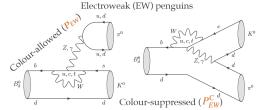
Challenge: Calculation of Hadronic matrix elements

# Why $B \rightarrow \pi K$ decays?



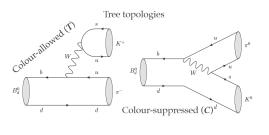


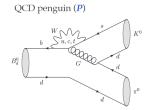
- Tree topologies suppressed by  $V_{ub}$
- QCD penguins dominant
- EW penguins at same level as tree



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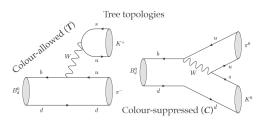
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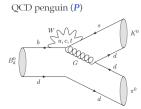




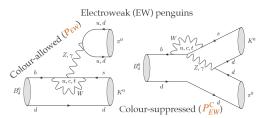
- Electroweak (EW) penguins
- Colour-allowed  $P_{EW}^{(a)}$   $Q_{EW}^{(a)}$   $Q_{EW}^{(a)}$
- ullet Tree topologies suppressed by  $V_{ub}$
- QCD penguins dominant
- EW penguins at same level as tree
  - Interesting probes of New Physics

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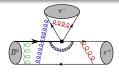




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- QCD penguins dominant
- EW penguins at same level as tree
- Interesting probes of New Physics
  - Search for tiny deviations of SM predictions



## How to handle nonleptonic B decays?



#### QCD Factorization Beneke, Buchalla, Neubert, Sachrajda

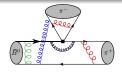
- Disentangle perturbative (calculable) and non-perturbative dynamics using HQE
- Systematic expansion in  $\alpha_s$  and  $1/m_b$  (studied up to  $\alpha_s^2$ ) Bell, Beneke, Huber, Li

$$\langle \pi^+ \pi^- | \mathcal{Q}_i | B \rangle = T_i^I \otimes F^{B \to \pi^+} \otimes \Phi_{\pi^-} + T_i^{II} \otimes \Phi_{\pi^-} \otimes \Phi_{\pi^+} \otimes \Phi_B$$

- Non-perturbative form factors and LCDAs
  - from data, lattice or Light-Cone Sum Rules
- No systematic framework to compute power corrections (yet?)
- · Strong phases suffer from large uncertainties
- Theoretical challenge: reliable computations of observables
- Recent progress: include QED corrections Talks by Martin and GaelBeneke, Boer, Toelstede, KKV [2020]

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## How to handle nonleptonic B decays?



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### Flavour symmetries (Isospin or SU(3))

- Many studies e.g. Fleischer, Jaarsma, KKV, Malami [2017,2018]
- Recent progress: Global SU(3) fit to B o PP decays Huber, Tetlalmatzi-Xolocotzi [2111.06418]

#### Light-cone sumrules

• Work in progress by Jung, Melic, Khodjamirian

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$$A_{M_1M_2} \equiv i \frac{G_F}{\sqrt{2}} m_B^2 F_0^{BM_1} f_{M_2}$$

Amplitude parametrization a la QCDF

[Beneke, Neubert [2003]]

$$\mathcal{A}_{B^- \to \pi^- \bar{K}^0} = A_{\pi K} \hat{\alpha}_4^{\rho} ,$$

$$\sqrt{2} \mathcal{A}_{B^- \to \pi^0 K^-} = A_{\pi K} \left[ \delta_{\rho u} \alpha_1 + \hat{\alpha}_4^{\rho} \right] + A_{K\pi} \left[ \delta_{\rho u} \alpha_2 + \delta_{\rho c} \frac{3}{2} \alpha_{3, \text{EW}}^{c} \right] ,$$

$$\mathcal{A}_{\bar{B}^0 \to \pi^+ K^-} = A_{\pi K} \left[ \delta_{\rho u} \alpha_1 + \hat{\alpha}_4^{\rho} \right] ,$$

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- $\alpha_1$  and  $\alpha_2$  color-allowed and color-suppressed tree coefficients
- $\alpha_4$  and  $\alpha_{3,EW}$  penguin and electromagnetic penguin coefficients
- ullet contain all perturbative effects up to NNLO  $(lpha_{
  m s}^2)$

e.g. [Bell, Beneke, Huber, Li]

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Amplitude parametrization a la QCDF

$$A_{P^-} = \bar{\kappa}_0 = A_{\pi K} \hat{\alpha}_A^P ,$$

$$\sqrt{2}\mathcal{A}_{B^-\to\pi^0K^-} = \mathcal{A}_{\pi K} \left[ \delta_{pu} \alpha_1 + \hat{\alpha}^\rho_4 \right] + \mathcal{A}_{K\pi} \left[ \delta_{pu} \alpha_2 + \delta_{pc} \frac{3}{2} \alpha^c_{3,\mathrm{EW}} \right] \; , \label{eq:delta_B}$$

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QED can be included! Beneke, Boer, Toelstede, KKV, JHEP 11 (2020) 081 [2008.10615]

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$$\mathcal{A}(M_{1}M_{2}) \equiv i \frac{G_{F}}{\sqrt{2}} m_{B}^{2} \mathcal{F}_{Q_{2}}^{BM_{1}}(0) F_{M_{2}}$$

$$\langle M_1 M_2 | Q_i | B \rangle = \mathcal{A}(M_1 M_2) \alpha_i(M_1 M_2) = A_{M_1 M_2} \Big( \alpha_i^{\text{QCD}}(M_1 M_2) + \delta \alpha_i(M_1 M_2) \Big)$$

- Electroweak scale to  $m_B$ : QED corrections to the Wilson coefficients
- $m_B$  to  $\mu_c$ : QED corrections to the hard-scattering kernels, form factors and decay constants
- below  $\Lambda_{\rm QCD}$ : Ultrasoft QED effects (for the rate!)

$$\delta\alpha_i(\textit{M}_1\textit{M}_2) \equiv \delta\alpha_i^{\rm WC}(\textit{M}_1\textit{M}_2) + \delta\alpha_i^{\rm K}(\textit{M}_1\textit{M}_2) + \delta\alpha_i^{\rm F,V}(\textit{M}_1\textit{M}_2) + \delta\alpha_i^{\rm F,sp}(\textit{M}_1\textit{M}_2) \,.$$

$$\rightarrow \delta \alpha_i^{\rm WC} = \mathcal{O}(10^{-3})$$

[Huber, Lunghi, Misiak, Wyler [2006]]

$$\rightarrow \delta \alpha_i^{\rm K} = \mathcal{O}(10^{-3})$$

$$\rightarrow \delta \alpha_i^{\mathrm{F,V}} = ??$$

[Beneke, Boer, Toelstede, KKV [2021]]

$$\rightarrow \delta \alpha_i^{\mathrm{F,sp}} = ?? \text{ but } \mathcal{O}(\alpha_{\mathrm{em}} \alpha_s)$$

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- Ultrasoft effects dress braching ratio
- Key point: scale dependence cancels!!

$$U(M_1 M_2) = \left(\frac{2\Delta E}{m_B}\right)^{-\frac{\alpha_{\rm em}}{\pi}} \left(Q_B^2 + Q_{M_1}^2 \left[1 + \ln \frac{m_{M_1}^2}{m_{B_q}^2}\right] + Q_{M_2}^2 \left[1 + \ln \frac{m_{M_2}^2}{m_B^2}\right]\right)$$

- Recover the standard QED factor
- $\Delta E$  is the window of the  $\pi K$  invariant mass around  $m_B$
- Theory requires  $\Delta E \ll \Lambda_{\rm QCD} = 60 \, \text{MeV}$

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### **Ultrasoft Contribution**

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$$\rightarrow~U(\pi^+K^-)=0.914, U(\pi^0K^-)=U(K^-\pi^0)=0.976$$
 and  $U(\pi^-\bar{K}^0)=0.954$ 

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- Experimentally usoft effects included using PHOTOS
- Challenging to compare theory with experiment! Work in progress!

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## Ratios and isospin sumrules

Beneke, Boer, Toelstede, KKV, JHEP 11 (2020) 081 [2008.10615]

QED gives sub-percent corrections to Branching ratios

[Talk by Gael F.]

Beneke, Boer, Toelstede, KKV, JHEP 11 (2020) 081 [2008.10615]

• Beneficial to consider ratios in which QCD is suppressed

$$R_L = \frac{2\mathrm{Br}(\pi^0 K^0) + 2\mathrm{Br}(\pi^0 K^-)}{\mathrm{Br}(\pi^- K^0) + \mathrm{Br}(\pi^+ K^-)} = R_L^{\mathrm{QCD}} + \cos \gamma \mathrm{Re} \ \delta_{\mathrm{E}} + \delta_U$$

new structure dependent QED corrections enter linearly, QCD only quadratically

$$\delta_E = (-1.12 + 0.16i) \cdot 10^{-3}$$

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$$\delta_U \equiv \frac{1 + U(\pi^0 K^-)}{U(\pi^- \bar{K}^0) + U(\pi^+ K^-)} - 1 = 5.8\%$$

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• Combined QED effect larger than QCD uncertainty!

## $B \to \pi K$ puzzle



e.g. Buras, Fleischer, Recksiegel, Schwab [2004, 2007];Fleischer, Jaeger, Pirjol, Zupan [2008] Neubert, Rosner [1998]; Beaudry, Datta, London, Rashed, Roux [2018]; Fleischer, Jaarsma, KKV [2018]

### (Longstanding) Puzzling patterns in $B \to \pi K$ data

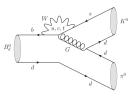
First Example:

$$\delta(\pi K) \equiv A_{\rm CP}(\pi^0 K^-) - A_{\rm CP}(\pi^+ K^-)$$

- Recent LHCb measurement for  $A_{\rm CP}(K^-\pi^0)$  LHCb Collaboration, PRL 126, 091802 [2021]
- Confirms and enhances the observed difference

- 
$$\delta(\pi K)^{\text{exp}} = (11.5 \pm 1.4)\%$$

-  $8\sigma$  from 0



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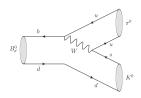
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  - $\delta(\pi K)^{
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  - or via SU(3) [Fleischer, Jaarsma, Malami, KKV [2018]]



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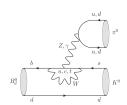
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- Hint for NP in the FWP sector?



Work in progress Jung, Melic, Khodjamirian (see MITP workshop 2019)

#### Preliminary!

Decay mode	BR-exp (in $10^{-6}$ )	$A_{CP} = -C_{CP}$	BR-th	$A_{CP}$ -th
$\Delta S = -1$				
$B^-  o \pi^0 K^-$	$12.7 \pm 0.6$	$0.040 \pm 0.021$	13.74	0.050
$B^-  o \pi^- ar K^0$	$23.3 \pm 0.8$	$-0.017 \pm 0.016$	24.56	-0.012
$\bar{B}^0  o \pi^+ K^-$	$20.0 \pm 0.6$	$-0.082 \pm 0.006$	20.10	0.057
$ar{B}^0  ightarrow \pi^0 ar{K}^0$	$10.1\pm0.5$	$-0.01 \pm 0.10$	8.87	-0.021

- LCSR calculations (+some QCDF input)
- More reliable than for  $B o \pi\pi$
- Different sign for  $B o K^+\pi^-$  (as in QCDF)

e.g. Gronau [2005]: Gronau, Rosner [2006]

$$\begin{split} \Delta(\pi K) &\equiv A_{\mathrm{CP}}(\pi^+ K^-) + \frac{\Gamma(\pi^- \bar{K}^0)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^- \bar{K}^0) - \frac{2\Gamma(\pi^0 K^-)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^0 K^-) \\ &- \frac{2\Gamma(\pi^0 \bar{K}^0)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^0 \bar{K}^0) \equiv \Delta(\pi K)^{\mathrm{QCD}} + \delta\Delta(\pi K) \end{split}$$

- Sensitive to new physics effects:  $\Delta(\pi K)^{\rm QCD}=(0.5\pm 1.1)\%$  [Bell, Beneke, Huber, Li]
- Recent progress: QED effects:  $\delta\Delta(\pi K) = -0.42\%$  [Beneke, Boer, Toelstede, KKV [2020]]
- Isospin sumrule also robust against QED effects!

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- Updates of modes with neutral pions necessary → Belle II

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- Isospin sumrule also robust against QED effects!
- $\bullet$  Updates of modes with neutral pions necessary  $\to$  Belle II
- ullet Or can be used to predict the direct CP in  $B o\pi^0 K^0$
- Mixing-induced CP asymmetry in  $B \to \pi^0 K^0$  provides additional test Fleischer, Jaarsma, Malami, KKV [2016,2018]

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Nir. Quin [1991]: Gronau. Hernandez. London. Rosner [1995]

Fleischer, Jaeger, Pirjol, Zupan [2008]; Fleischer, Jaarsma, KKV [2018]

$$\begin{split} \overline{\sqrt{2}A(B^0 \to \pi^0 K^0) + A(B^0 \to \pi^- K^+)} \\ &= \sqrt{2}A(B^+ \to \pi^0 K^+) + A(B^+ \to \pi^+ K^0) \\ &= -(\hat{T} + \hat{C}) \left( e^{i\gamma} - q e^{i\phi} e^{i\omega} \right) \equiv 3A_{3/2} = 3|A_{3/2}|e^{i\phi_{3/2}} \;, \end{split}$$

- QCD penguin and colour-suppressed EWPs cancel
- Minimal SU(3) input

$$|\hat{T} + \hat{C}| = R_{T+C} |V_{us}/V_{ud}| \sqrt{2} |A(B^+ \to \pi^+ \pi^0)|$$

$$R_{T+C}|_{\mathsf{fact}} = f_{\mathcal{K}}/f_{\pi} = 1.2 \pm 0.2$$

Uncertainty accounts for non-factorizable SU(3) breaking

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Nir, Quin [1991]; Gronau, Hernandez, London, Rosner [1995]

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- QCD penguin and colour-suppressed EWPs cancel
- EWPs described by  $(\phi(\omega)$  CP-violating (conserving) phases)

$$qe^{i\phi}e^{i\omega} \equiv -\left(\frac{\hat{P}_{EW}+\hat{P}_{EW}^{C}}{\hat{T}+\hat{C}}\right) \stackrel{\text{SM}}{=} \frac{-3}{2\lambda^{2}R_{b}}R_{q}\frac{C_{9}+C_{10}}{C_{1}+C_{2}} = 0.68R_{q}$$

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Nir. Quin [1991]: Gronau. Hernandez. London. Rosner [1995]

Fleischer, Jaeger, Pirjol, Zupan [2008]; Fleischer, Jaarsma, KKV [2018]

$$\overline{\sqrt{2}A(B^0 \to \pi^0 K^0) + A(B^0 \to \pi^- K^+)} 
= \sqrt{2}A(B^+ \to \pi^0 K^+) + A(B^+ \to \pi^+ K^0) 
= -(\hat{T} + \hat{C}) (e^{i\gamma} - qe^{i\phi}e^{i\omega}) \equiv 3A_{3/2} = 3|A_{3/2}|e^{i\phi_{3/2}},$$

QCD penguin and colour-suppressed EWPs cancel

$$qe^{i\phi}e^{i\omega} \equiv -\left(\frac{\hat{P}_{EW}+\hat{P}_{EW}^{C}}{\hat{T}+\hat{C}}\right) \stackrel{\text{SM}}{=} \frac{-3}{2\lambda^{2}R_{b}}R_{q}\frac{C_{0}+C_{10}}{C_{1}+C_{2}} = 0.68R_{q}$$

New CP violating physics might enter with a large phase  $\phi$ 

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## **Isospin Amplitude Triangles**

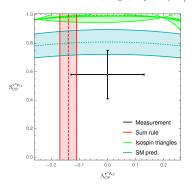
Fleischer, Jaeger, Pirjol, Zupan [2008]; Fleischer, Jaarsma, KKV [2018]

$$A_{
m mix}^{\pi^0 K_S} = \sin(\phi_d - \phi_{00}) \sqrt{1 - (A_{
m dir}^{\pi^0 K_S})^2}$$

- $\sqrt{2} \, \bar{A}_{00}$   $\sqrt{2} \, A_{00}$   $\sqrt{2} \, A_{00}$   $3 \, \bar{A}_{3/2}$   $3 \, A_{3/2}$
- ullet Only  $\pi K$  mode with mixing-induced CP asymmetry
- $\phi_d$   $B_d$  mixing phase
- Amplitude triangles give  $\phi_{00}$
- ullet Gives a clean correlation between the CP asymmetries in  $B_d o \pi^0 K_S$

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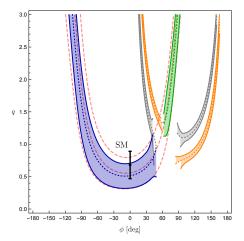
Nir, Quin [1991]; Gronau, Hernandez, London, Rosner [1995] Fleischer, Jaeger, Pirjol, Zupan [2008]; Fleischer, Jaarsma, KKV [2018]



Hints at New Physics in the EWP sector?

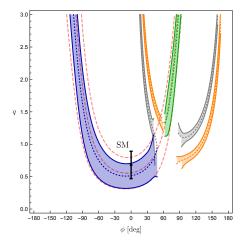
## New Physics in the EWP sector





- Allow for New Physics in EWPs  $qe^{i\phi}e^{i\omega}\equiv -\left(\frac{\hat{P}_{EW}+\hat{P}_{EW}^{C}}{\hat{T}+\hat{C}}\right)$
- Isospin relation holds for neutral and charged decays
- Current uncertainties in neutral decays still large → charged decays
- Only minimal SU(3)

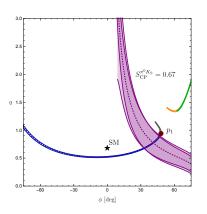




- Allow for New Physics in EWPs  $qe^{i\phi}e^{i\omega}\equiv -\left(\frac{\hat{P}_{EW}+\hat{P}_{EW}^{C}}{\hat{T}+\hat{C}}\right)$
- Isospin relation holds for neutral and charged decays
- $\hbox{ Current uncertainties in neutral decays} \\ \hbox{ still large} \rightarrow \hbox{ charged decays} \\$
- Only minimal SU(3)
- Further input necessary

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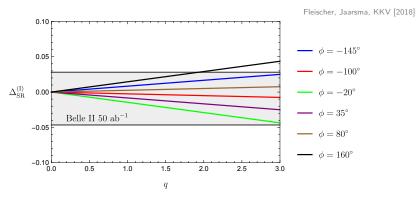




Perpendicular constraint from mixing-induced CP asymmetry  $B_d o \pi^0 K_S$ 

- Theory uncertainties (wide band) match experimental (smaller band)
- Include hadronic uncertainties in data-driven way

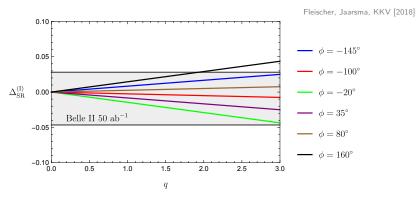
#### Constraints on new physics from the sum rule



- Projected sensitivity at Belle II for  $\pi^0 K_S$
- $\bullet$  Limited sensitivity to q and  $\phi$  for q<3

$$\begin{split} \Delta(\pi K) &\equiv A_{\mathrm{CP}}(\pi^+ K^-) + \frac{\Gamma(\pi^- \bar{K}^0)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^- \bar{K}^0) - \frac{2\Gamma(\pi^0 K^-)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^0 K^-) \\ &- \frac{2\Gamma(\pi^0 \bar{K}^0)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^0 \bar{K}^0) \equiv \Delta(\pi K)^{\mathrm{QCD}} + \delta\Delta(\pi K) \end{split}$$

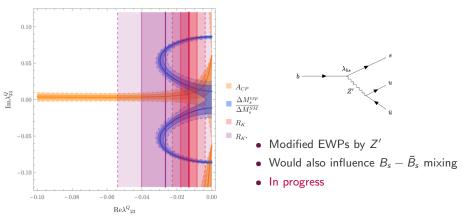
#### Constraints on new physics from the sum rule



- Projected sensitivity at Belle II for  $\pi^0 K_S$
- $\bullet$  Limited sensitivity to q and  $\phi$  for q<3
- Use mixing-induced CP asymmetry to pin down NP

$$\begin{split} \Delta(\pi K) &\equiv A_{\mathrm{CP}}(\pi^+ K^-) + \frac{\Gamma(\pi^- \bar{K}^0)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^- \bar{K}^0) - \frac{2\Gamma(\pi^0 K^-)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^0 K^-) \\ &- \frac{2\Gamma(\pi^0 \bar{K}^0)}{\Gamma(\pi^+ K^-)} A_{\mathrm{CP}}(\pi^0 \bar{K}^0) \equiv \Delta(\pi K)^{\mathrm{QCD}} + \delta\Delta(\pi K) \end{split}$$





# Extracting $\gamma$ from QCD penguin decays

Fleischer [1999, 2007]; Fleischer, Knegjens [2011]

$$A(B_s \to K^+ K^-) = \sqrt{\epsilon} e^{i\gamma} \mathcal{C}' \left[ 1 + \frac{1}{\epsilon} d' e^{i\theta'} e^{-i\gamma} \right]$$

$$A(B_d \to \pi^+ \pi^-) = e^{i\gamma} \mathcal{C} \left[ 1 - d e^{i\theta} e^{-i\gamma} \right]$$

$$\mathcal{C}' \propto T' + P^{(ut)'} + E' + PA^{(ut)'} \text{ and } d' e^{i\theta'} \propto \frac{P^{(ct)'} + PA^{(ct)'}}{T' + P^{(ut)'} + E' + PA^{(ut)'}}$$

*U*-spin symmetry

$$de^{i\theta} = d'e^{i\theta'}$$

• Extract hadronic parameters from direct and mixing-induced CP asymmetries

Fleischer [1999, 2007]; Fleischer, Knegjens [2011]

$$A(B_s \to K^+ K^-) = \sqrt{\epsilon} e^{i\gamma} \mathcal{C}' \left[ 1 + \frac{1}{\epsilon} d' e^{i\theta'} e^{-i\gamma} \right]$$

$$A(B_d \to \pi^+ \pi^-) = e^{i\gamma} \mathcal{C} \left[ 1 - d e^{i\theta} e^{-i\gamma} \right]$$

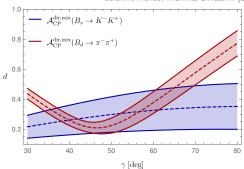
$$\mathcal{C}' \propto T' + P^{(ut)'} + E' + PA^{(ut)'} \text{ and } d' e^{i\theta'} \propto \frac{P^{(ct)'} + PA^{(ct)'}}{T' + P^{(ut)'} + E' + PA^{(ut)'}}$$

*U*-spin symmetry

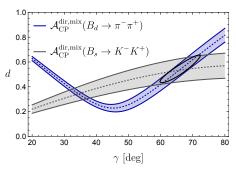
$$de^{i\theta} = d'e^{i\theta'}$$

- Extract hadronic parameters from direct and mixing-induced CP asymmetries
- Or assume d = d' and extract  $\gamma$
- Limited by *U*-spin breaking corrections

Fleischer [1999,2007]; Fleischer, Knegjens [2011]; Fleischer, Malami, Jaarsma, KKV [2016,in progress]
Cuichini, Franco, Mishima, Silvestrini [2012], Data from LHCb [2022]



Fleischer [1999,2007]; Fleischer, Knegjens [2011]; Fleischer, Malami, Jaarsma, KKV [2016,in progress] Cuichini, Franco, Mishima, Silvestrini [2012], Data from LHCb [2022]



- $\gamma = (65^{+7}_{-5})^\circ$  Jaarsma, Fleischer, KKV [in progress]
- Agrees with tree determinations
- $\circ \ \gamma = (64.9 \pm 4.5)^\circ \ {
  m LHCb} \ {
  m [2021]} \ {
  m without} \ {
  m \textit{B}_{\it s}} \ {
  m modes}$

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# Challenges in (nonleptonic) B decays

We are in the High-precision Era in Flavour Physics!

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- Stay tuned for new data and updated theory predictions
  - Neutral pion modes!

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  - Neutral pion modes!

Close collaboration between theory and experiment necessary!

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#### We are in the High-precision Era in Flavour Physics!

- QED is important!
- Rethink our previous assumptions
- Stay tuned for new data and updated theory predictions
  - Neutral pion modes!

Next time in Maastricht?

# Backup

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# The Challenge of QED Corrections

Beneke, Boer, Toelstede, KKV, JHEP 11 (2020) 081 [2008.10615]

$$\Gamma[\bar{B} \to M_1 M_2](\Delta E) \equiv \Gamma[\bar{B} \to M_1 M_2 + X_s] \big|_{E_{X_s} \leq \Delta E} \,,$$

- IR finite observable (width) must include ultra-soft photon radiation
- $X_s$  are soft photons with total energy less than ultrasoft scale  $\Delta E$
- Factorizes in non-radiative amplitude and ultrasoft function

$$\Gamma[\bar{B} \to M_1 M_2](\Delta E) = |\mathcal{A}(\bar{B} \to M_1 M_2)|^2 \sum_{X_s} |\langle X_s | (\bar{S}_v^{(Q_B)} S_{v_1}^{\dagger(Q_{M_1})} S_{v_2}^{\dagger(Q_{M_2})}) |0\rangle|^2 \theta(\Delta E - E_{X_s})$$

#### Simple classification:

• Ultra-soft photons: eikonal approximation, well understood

$$\Delta E \ll \Lambda_{\rm QCD}$$

- NEW: Non-universal, structure dependent corrections Beneke, Boer, Toelstede, KKV [2020]
- Both effects important: virtual photons can resolve the structure of the meson!

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Beneke, Boer, Toelstede, KKV, JHEP 11 (2020) 081 [2008.10615]

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#### Simple classification:

• Ultra-soft photons: eikonal approximation, well understood

$$\Delta E \ll \Lambda_{\rm QCD}$$

- Often done: Assume pointlike approximation up to the scale  $m_B$  [Baracchini, Isidori]
  - → fails to account for all large logarithms (and scales)!
  - $\rightarrow$  photons with energy  $\gtrsim \Lambda_{\rm QCD}$  probe the partonic structure of the mesons

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### **Ultrasoft Contribution**

• Ultrasoft effects dress braching ratio

$$\textit{U(M$_{1}$M$_{2}$)} = \left(\frac{2\Delta \textit{E}}{\textit{m}_{\textit{B}}}\right)^{-\frac{\alpha_{\rm em}}{\pi}\left(\textit{Q}_{\textit{B}}^{2} + \textit{Q}_{\textit{M}_{1}}^{2}\left[1 + \ln\frac{\textit{m}_{\textit{M}_{1}}^{2}}{\textit{m}_{\textit{B}_{q}}^{2}}\right] + \textit{Q}_{\textit{M}_{2}}^{2}\left[1 + \ln\frac{\textit{m}_{\textit{M}_{2}}^{2}}{\textit{m}_{\textit{B}}^{2}}\right]\right)}$$

- Recover the standard eikonal/QED factor Beneke, Boer, Toelstede, KKV [2020]
- $\Delta E$  is the window of the  $\pi K$  invariant mass around  $m_B$
- Theory requires  $\Delta E \ll \Lambda_{\rm QCD} = 60 \, \text{MeV}$

• Ultrasoft effects dress braching ratio

$$U(M_1M_2) = \left(\frac{2\Delta E}{m_B}\right)^{-\frac{\alpha_{\rm em}}{\pi}\left(Q_B^2 + Q_{M_1}^2\left[1 + \ln\frac{m_{M_1}^2}{m_{B_q}^2}\right] + Q_{M_2}^2\left[1 + \ln\frac{m_{M_2}^2}{m_B^2}\right]\right)}$$

- Recover the standard eikonal/QED factor Beneke, Boer, Toelstede, KKV [2020]
- $\Delta E$  is the window of the  $\pi K$  invariant mass around  $m_B$
- Theory requires  $\Delta E \ll \Lambda_{\rm QCD} = 60 \, \text{MeV}$
- Large effects:

$$\rightarrow \ U(\pi^+ K^-) = 0.914, U(\pi^0 K^-) = U(K^- \pi^0) = 0.976 \ \text{and} \ U(\pi^- \bar{K}^0) = 0.954$$

Puzzle 02.05.2022 29 / 25 • Ultrasoft effects dress braching ratio

$$U(M_1M_2) = \left(\frac{2\Delta E}{m_B}\right)^{-\frac{\alpha_{\rm em}}{\pi}\left(Q_B^2 + Q_{M_1}^2\left[1 + \ln\frac{m_{M_1}^2}{m_{B_q}^2}\right] + Q_{M_2}^2\left[1 + \ln\frac{m_{M_2}^2}{m_B^2}\right]\right)}$$

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- Large effects:

$$\rightarrow U(\pi^+K^-) = 0.914, U(\pi^0K^-) = U(K^-\pi^0) = 0.976 \text{ and } U(\pi^-\bar{K}^0) = 0.954$$

- Experimentally usoft effects included using PHOTOS
- Challenging to compare theory with experiment! In progress... See e.g. Isidori, Zwicky, Nabeebaccus, Bordone, Pattori

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### **Hadronic parameters**

Isospin parametrization Buras, Fleischer, Recksiege, Schwab [2004]

$$\begin{split} A(B^+ \to \pi^+ K^0) &= -P' \left[ 1 + \rho_c e^{i\theta_c} e^{i\gamma} \right] \\ \sqrt{2} A(B^+ \to \pi^0 K^+) &= P' \left[ 1 + \rho_c e^{i\theta_c} e^{i\gamma} - \left( e^{i\gamma} - q e^{i\phi} e^{i\omega} \right) r_c e^{i\delta_c} \right] \\ A(B_d^0 \to \pi^- K^+) &= P' \left[ 1 - r e^{i\delta} e^{i\gamma} \right] \\ \sqrt{2} A(B_d^0 \to \pi^0 K^0) &= -P' \left[ 1 - r e^{i\delta} e^{i\gamma} + \left( e^{i\gamma} - q e^{i\phi} e^{i\omega} \right) r_c e^{i\delta_c} \right] \; . \end{split}$$

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Isospin parametrization Buras, Fleischer, Recksiege, Schwab [2004]

$$\begin{split} &A(B^+\to\pi^+K^0)=-P'\left[1+\rho_ce^{i\theta_c}e^{i\gamma}\right]\\ &\sqrt{2}A(B^+\to\pi^0K^+)=P'\left[1+\rho_ce^{i\theta_c}e^{i\gamma}-\left(e^{i\gamma}-qe^{i\phi}e^{i\omega}\right)r_ce^{i\delta_c}\right]\\ &A(B_d^0\to\pi^-K^+)=P'\left[1-re^{i\delta}e^{i\gamma}\right]\\ &\sqrt{2}A(B_d^0\to\pi^0K^0)=-P'\left[1-re^{i\delta}e^{i\gamma}+\left(e^{i\gamma}-qe^{i\phi}e^{i\omega}\right)r_ce^{i\delta_c}\right]\;. \end{split}$$

ullet Hadronic parameters determined from  $B o \pi\pi$ 

$$r_c e^{i\delta_c} \equiv (\hat{T} + \hat{C})/P' = (0.17 \pm 0.06)e^{i(1.9 \pm 23.9)^{\circ}} ,$$

$$re^{i\delta} \equiv (\hat{T} - \hat{P}_u + \hat{P}_t)/P' = (0.09 \pm 0.03)e^{i(28.6 \pm 21.4)^{\circ}} ,$$

$$\rho_c e^{i\theta_c} \equiv (\hat{P}_t - \hat{P}_u)/(\hat{P}_t - \hat{P}_c) \sim 0$$

- Hadronic parameters of  $\mathcal{O}(\lambda)$  and include 20% SU(3) corrections
- Can be compared to the QCDF calculations

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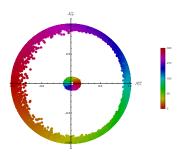
ullet Complement the isospin analysis with  $S_{\mathsf{CP}}^{\pi^0 K_{\mathcal{S}}}$ 

$$\tan \phi_{00} = 2(r\cos \delta - r_c\cos \delta_c)\sin \gamma + 2r_c(\cos \delta_c - 2\tilde{a}_C/3)q\sin \phi + \mathcal{O}(\lambda^2)$$

- $r, \delta, r_c$  and  $\delta_c$  hadronic parameters determined from  $B \to \pi\pi$
- Only cosines of small phases, low sensitivity to variations
- Includes color-suppressed EWPs  $\tilde{a}_{\mathsf{C}} = a_{\mathsf{C}} \cos(\Delta_{\mathsf{C}} + \delta_{c})$
- Effects included in a data-driven way

$$R \equiv \frac{\mathsf{Br}(\pi^- K^+)}{\mathsf{Br}(\pi^+ K^0)} = 0.89 \pm 0.04 = 1 - 2r\cos\delta\cos\gamma + 2r_c\tilde{a}_C q\cos\phi + \mathcal{O}(\lambda^2)$$

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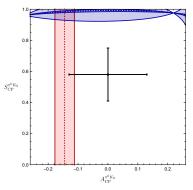


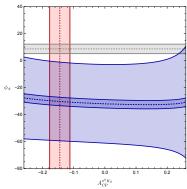
- Correlation between CP asymmetries in  $B_s \to \pi\pi$  (small circle) and  $B_d \to KK$ (wide) for different strong phases
- Important also to improve QCDF

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New element: constraint on angle  $\phi_+ = \arg(\bar{A}_+ A_+^*)$ 

$$\phi_{\pm}|_{\mathrm{SM},\phi=0} = 2r\cos\delta\sin\gamma + \mathcal{O}(\lambda^2) = (8.7 \pm 3.5)^{\circ}$$

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