

LHCb outlook for non-leptonic decays

(including a_{fs} and the determination of γ)

Status and prospects of non-leptonic B meson decays

"The mascots of Siegen, miner Henner and iron industry worker Frieder, meet the Higgs particle"

Collaborative Research Center TRR 257

MANCHESTER
1824

The University of Manchester

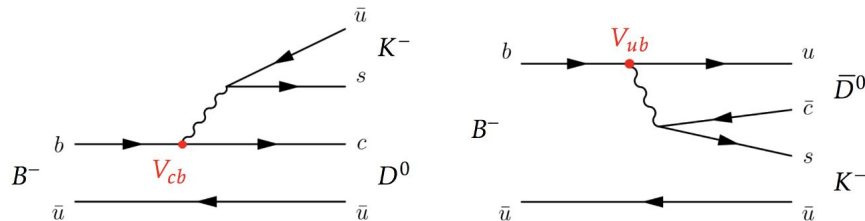
N. Skidmore on behalf of the LHCb
collaboration
June 2022

The many measurements of γ ...

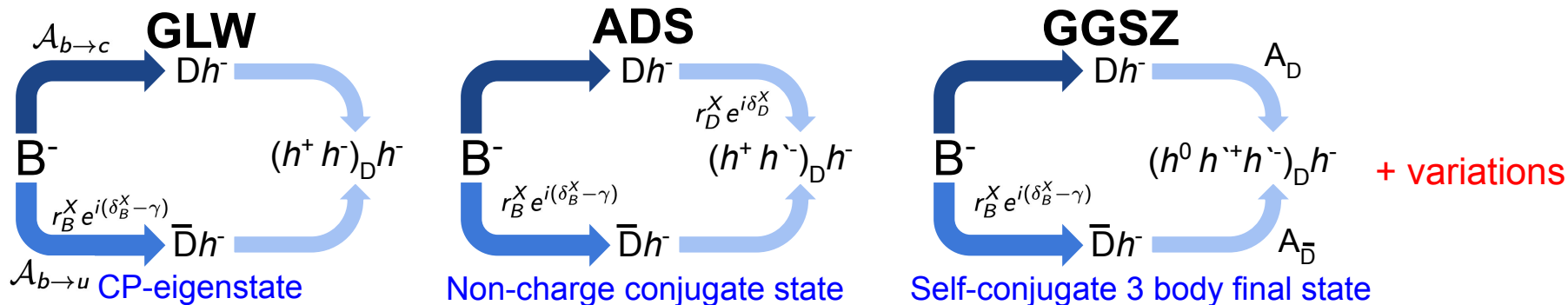
Here we discuss TI measurements, for TD see [Stephie's talk](#)

CKM phase γ can be measured solely with **tree-level $B \rightarrow Dh$ decays**

In the absence of NP effects at tree-level - SM benchmark- negligible theoretical uncertainties when interpreting observables in terms of γ



Compare to γ interpretations from decays susceptible to NP contributions $\rightarrow$ sensitive to **BSM effects at mass scales greater than those accessible directly**

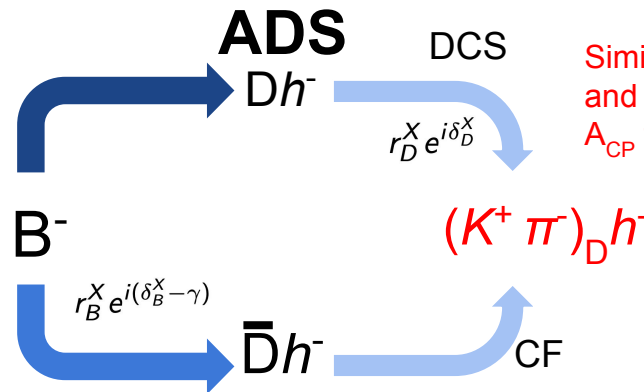
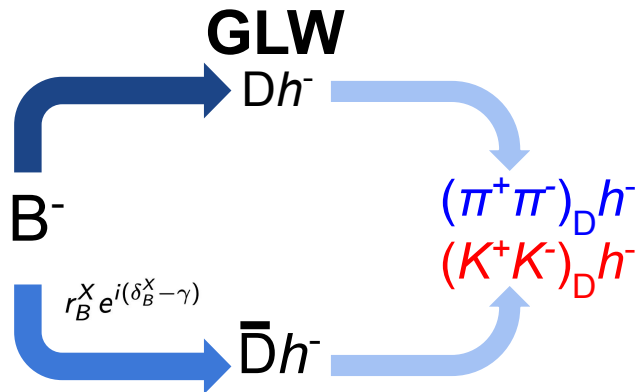


Choice of measured observables is dependent on the method that is most appropriate for the D decay analysed - all sensitive to the same B decay parameters, $r_B^X, \delta_B^X, \gamma \rightarrow$ **combination fits**

CP observables in $B^\pm \rightarrow D^{(*)} h^\pm$ decays

8.7 fb⁻¹

Use $B^\pm \rightarrow D^{(*)} h^\pm$ decays to measure **experimentally robust** CP observables in GLW and ADS modes



Similar magnitude of $r_D^{K\pi}$ and r_B^{DK} leads to large A_{CP} for ADS modes

$$\Gamma(B^\mp \rightarrow [[h^+ h^-]_D h^\mp]_X) \propto 1 + (r_B^X)^2 + 2r_B^X \cos(\delta_B^X \mp \gamma) - \alpha y (1 + (r_B^X)^2) - \alpha y 2r_B^X \cos(\delta_B^X \mp \gamma)$$

$$\Gamma(B^\mp \rightarrow [[K^\pm \pi^\mp]_D h^\mp]_X) \propto (r_D^{K\pi})^2 + (r_B^X)^2 + 2r_D^{K\pi} r_B^X \cos(\delta_B^X + \delta_D^{K\pi} \mp \gamma) - \alpha y (1 + (r_B^X)^2) r_D^{K\pi} \cos \delta_D^{K\pi} - \alpha y (1 + (r_D^{K\pi})^2) r_B^X \cos(\delta_B^X \mp \gamma) + \alpha x (1 - (r_B^X)^2) r_D^{K\pi} \sin \delta_D^{K\pi} - \alpha x (1 - (r_D^{K\pi})^2) r_B^X \sin(\delta_B^X \mp \gamma)$$

- 28 CP observables are measured - partial width ratios, double ratios and A_{CP} s
- Can be interpreted in terms of $\gamma, r_B^{DK}, \delta_B^{DK}, r_B^{D\pi}, \delta_B^{D\pi}, r_B^{D^*K}, \delta_B^{D^*K}, r_B^{D^*\pi}, \delta_B^{D^*\pi}$

$r_D^{K\pi}, \delta_D^{K\pi}$ and **charm mixing parameters x and y** taken from HFLAV average Eur. Phys. J. C (2021) 81: 226

CP observables in $B^\pm \rightarrow D^{(*)} h^\pm$ decays

8.7 fb⁻¹

First measurements of CP observables using $B^\pm \rightarrow D^* h^\pm$ in ADS modes

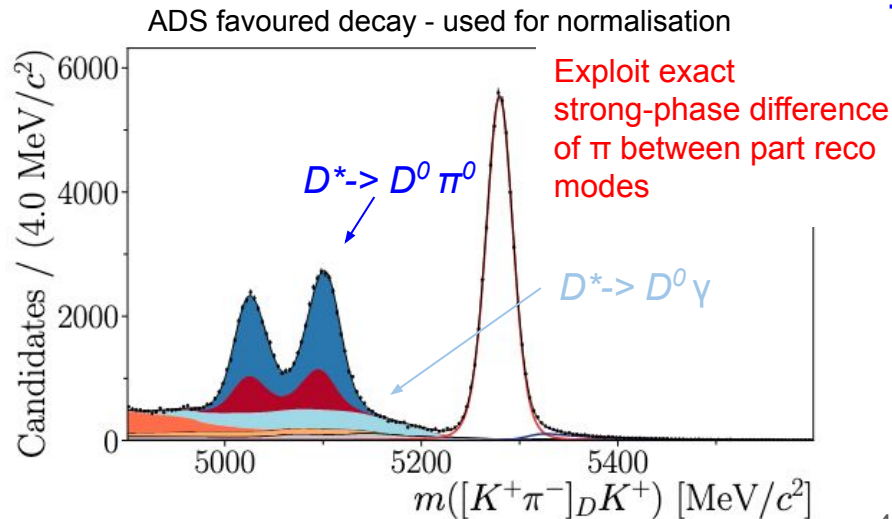
$D^* \rightarrow D^0 \pi^0 / \gamma$ is partially reconstructed - missing neutral

Reconstruction efficiency for low momentum neutral pions and photons is low - partial reconstruction provides significantly larger yields

Small $m_{D^*} - m_D$ and angular structure of $D^* \rightarrow D^0 \pi^0 / \gamma$ decays result in distinctive line shapes in $m(Dh^\pm)$ - yields can be obtained

Measurements of $\mathcal{B}(B^- \rightarrow D^{*0} \pi^-)$ and $\mathcal{B}(D^{*0} \rightarrow D^0 \pi^0)$ are made and agree with current world averages

Demonstrates method of partial reconstruction accurately measures $B^\pm \rightarrow D^* h^\pm$ modes despite presence of multiple partially reconstructed background sources



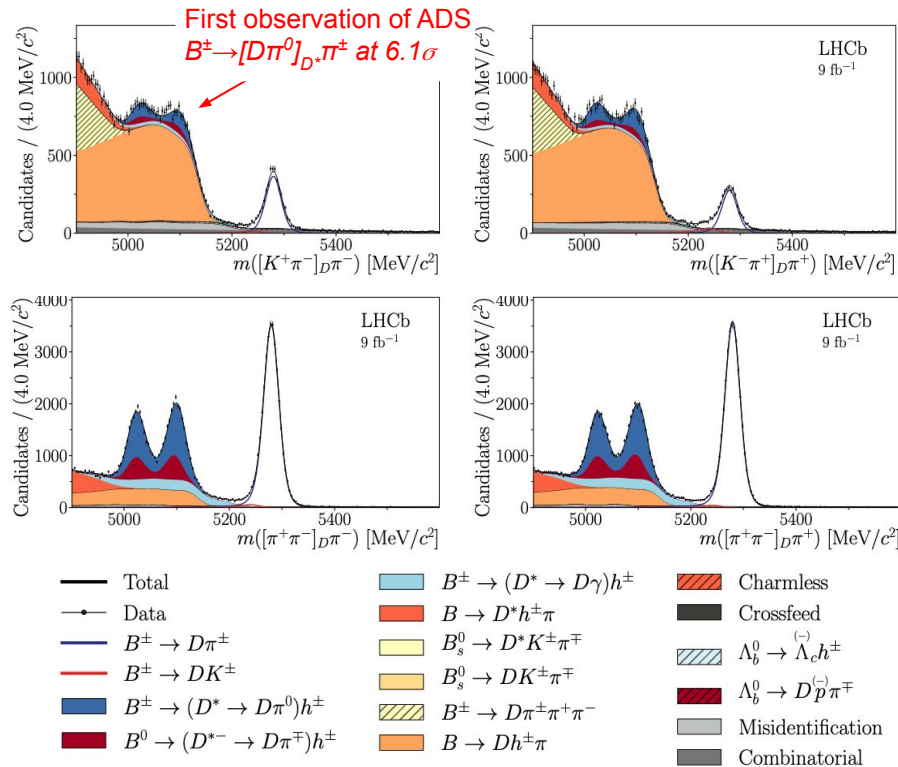
CP observables in $B^\pm \rightarrow D^{(*)} h^\pm$ decays

8.7 fb⁻¹

JHEP 04 (2021) 081

ADS

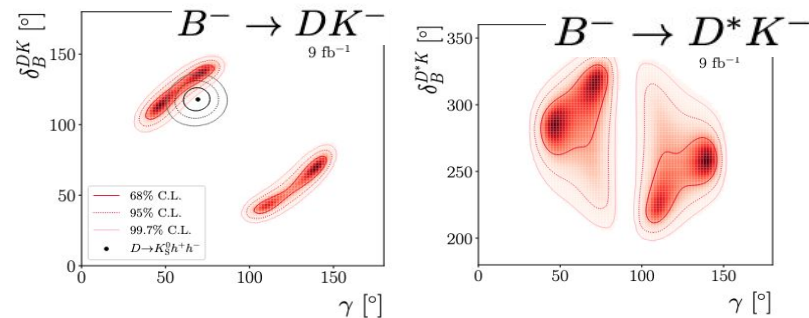
GLW



All CP observables measured to **worlds best precision**

Profile likelihood contours dominated by $B^\pm \rightarrow DK^\pm$ - good agreement with LHCb/B factory results

Constructed using charm parameters taken as external constraints from HFLAV average but these modes provide significant sensitivity to these parameters → **motivates combination of gamma and charm results**

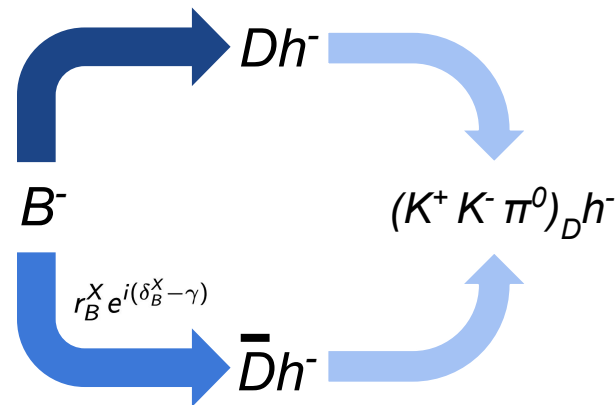
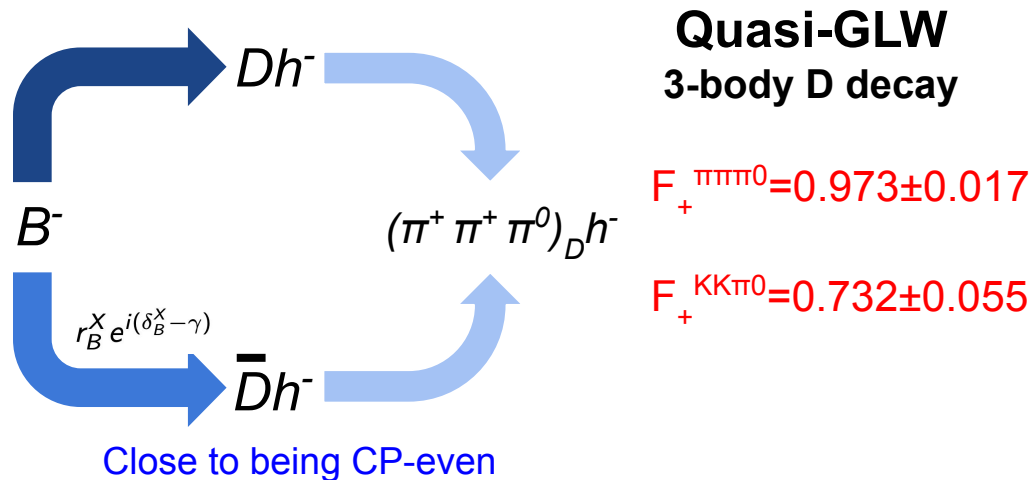


16 independent samples fit simultaneously

Constraints on γ from $B^\pm \rightarrow Dh^\mp$ with $D \rightarrow h^\pm h^\mp \pi^0$

9fb⁻¹

arXiv:2112.10617, submitted to JHEP



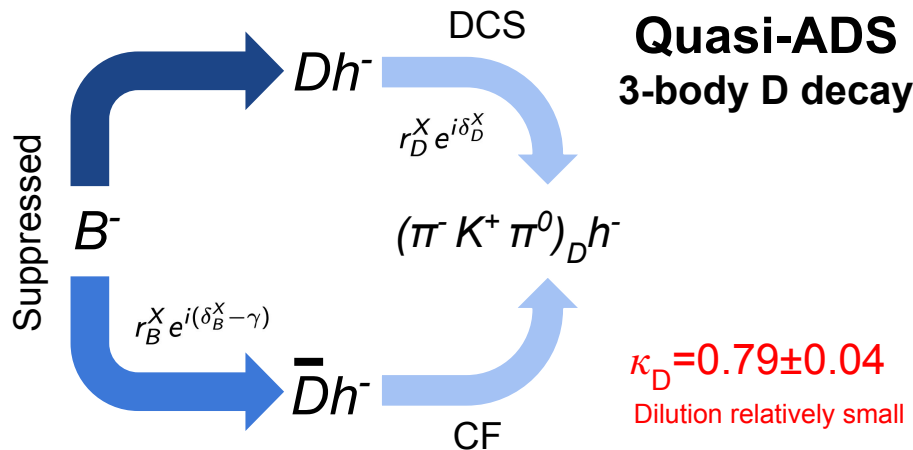
$$\Gamma(B^\mp \rightarrow [\pi^+ \pi^- \pi^0]_D h^\mp) \propto 1 + r_B^2 + 2r_B(2F_+^{\pi\pi\pi^0} - 1) \cos(\delta_B \mp \gamma) - \alpha \gamma (1 + r_B^2)(2F_+^{\pi\pi\pi^0} - 1) - 2\alpha \gamma r_B \cos(\delta_B \mp \gamma)$$

$$\Gamma(B^\mp \rightarrow [K^+ K^- \pi^0]_D h^\mp) \propto 1 + r_B^2 + 2r_B(2F_+^{KK\pi^0} - 1) \cos(\delta_B \mp \gamma) - \alpha \gamma (1 + r_B^2)(2F_+^{KK\pi^0} - 1) - 2\alpha \gamma r_B \cos(\delta_B \mp \gamma)$$

- Admixture of CP-odd and CP-even eigenstates
- Integration over 3-body D decay phasespace dilutes A_{CP}
- CP-even fractions, $F_+^{hh\pi^0}$, measured at CLEO-c

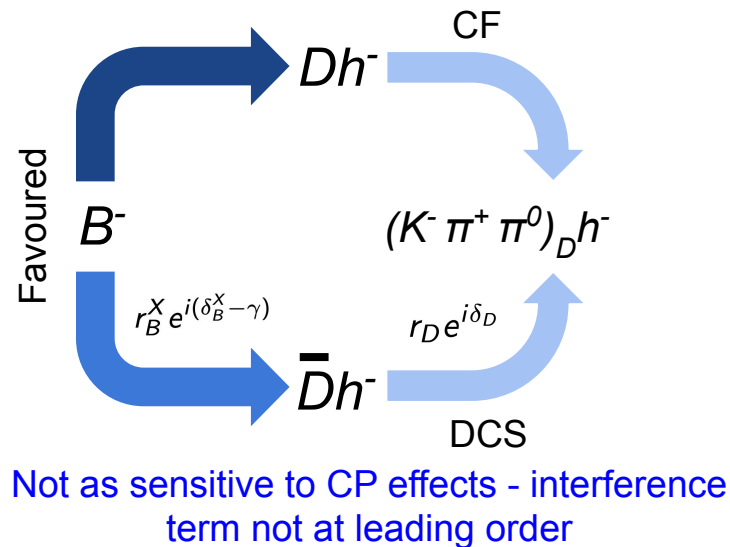
Constraints on γ from $B^\pm \rightarrow Dh^\mp$ with $D \rightarrow h^\pm h^\mp \pi^0$

9fb⁻¹



Balanced competing amplitudes maximises interference

$$\Gamma(B^\mp \rightarrow [\pi^\mp K^\pm \pi^0]_D h^\mp) \propto r_D^2 + r_B^2 + 2r_D r_B \kappa_D \cos(\delta_B + \delta_D \mp \gamma) \\ - \alpha y (1 + r_B^2) r_D \kappa_D \cos \delta_D - \alpha y (1 + r_D^2) r_B \cos(\delta_B \mp \gamma) \\ + \alpha x (1 - r_B^2) r_D \kappa_D \sin \delta_D - \alpha x (1 - r_D^2) r_B \sin(\delta_B \mp \gamma)$$



$$\Gamma(B^\mp \rightarrow [\pi^\pm K^\mp \pi^0]_D h^\mp) \propto 1 + r_D^2 r_B^2 + 2r_D r_B \kappa_D \cos(\delta_B - \delta_D \mp \gamma) \\ - \alpha y (1 + r_B^2) r_D \kappa_D \cos \delta_D - \alpha y (1 + r_D^2) r_B \cos(\delta_B \mp \gamma) \\ - \alpha x (1 - r_B^2) r_D \kappa_D \sin \delta_D + \alpha x (1 - r_D^2) r_B \sin(\delta_B \mp \gamma)$$

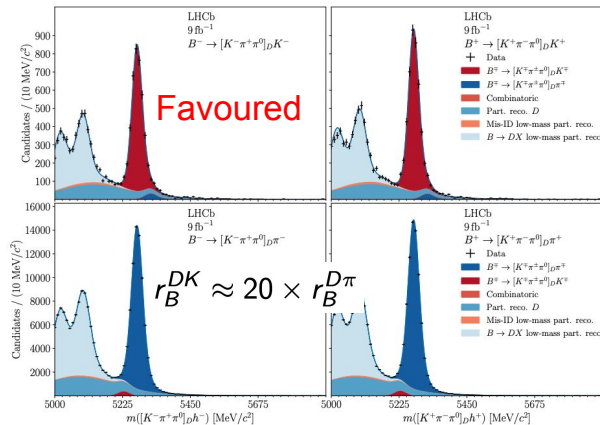
- Integration over 3-body D decay phasespace **dilutes sensitivity**
- Dilution factor, κ_D , and average strong phase, δ_D taken from BESIII/ CLEO-c/ LHCb

9fb⁻¹

Constraints on γ from $B^\pm \rightarrow Dh^\mp$ with $D \rightarrow h^\pm h^\mp \pi^0$

arXiv:2112.10617, submitted to JHEP

Quasi-ADS

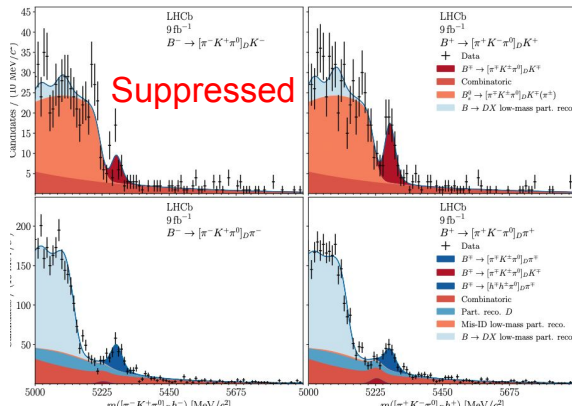


Interpreted as γ , r_B , δ_B using inputs from CLEO-c/BES III
Due to trigonometric ambiguities 4 solutions for γ

$$r_B = (9.3_{-0.9}^{+1.0}) \times 10^{-2}$$

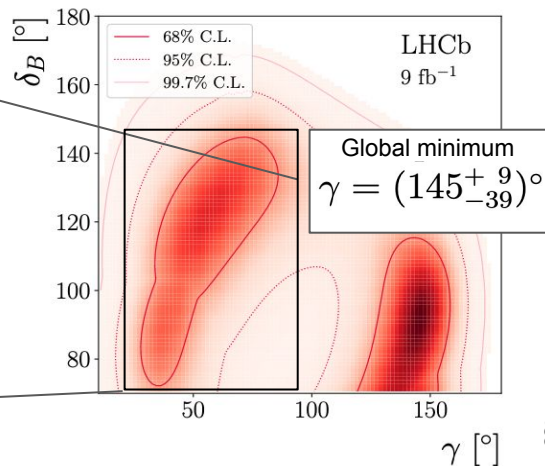
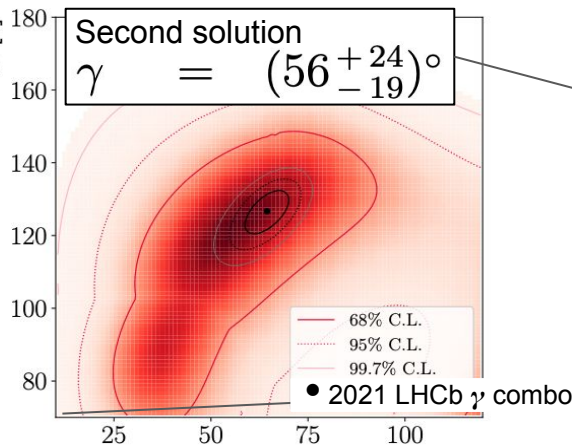
$$\delta_B = (122_{-23}^{+19})^\circ$$

Consistent with LHCb γ combo



World's best precision on observables

First observation
of $B^\mp \rightarrow [\pi^\mp K^\pm \pi^0]_D K^\mp$ at $>7\sigma$!

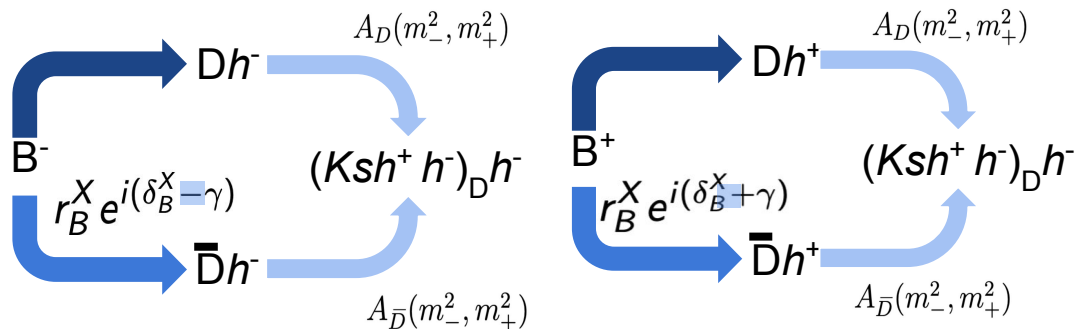


Measurement of gamma in $B^\pm \rightarrow D[Kshh]h^\pm$

GGSZ

9fb⁻¹

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Analyse $B^\pm \rightarrow D[Kshh]h^\pm$ decays in bins of D -decay phasespace

To access gamma need knowledge of how strong phase of D decay varies over phasespace

$$A_D(m_-^2, m_+^2) = |A_D(m_-^2, m_+^2)| e^{i\delta_D(m_-^2, m_+^2)}$$

Can construct yield observables for each bin i

$$N_{+i}^- = h_{B^-} \left[F_{+i} + \left((x_-^{DK})^2 + (y_-^{DK})^2 \right) F_{-i} + 2\sqrt{F_i F_{-i}} (x_-^{DK} c_{+i} + y_-^{DK} s_{+i}) \right]$$

Normalisation absorbs production and detection asymmetries

Efficiency corrected fraction of events that fall in bin i

Cosine/sine of strong phase difference between D^0 and $\bar{D}^0$ averaged over bin i

c_i and s_i inputs taken from combination of BESIII and CLEO results
PRD101 (2020) 112002, PRD102 (2020) 052008

F_i normally determined through semileptonic decays (no CPV) however to increase sensitivity $B \rightarrow D\pi$ channel is fit simultaneously with $B \rightarrow DK$

Negligible contribution to gamma precision from $B \rightarrow D\pi$ as $r_B^{DK} \approx 20 \times r_B^{D\pi}$

Use of $D\pi$ channel controls selection and reconstruction effects

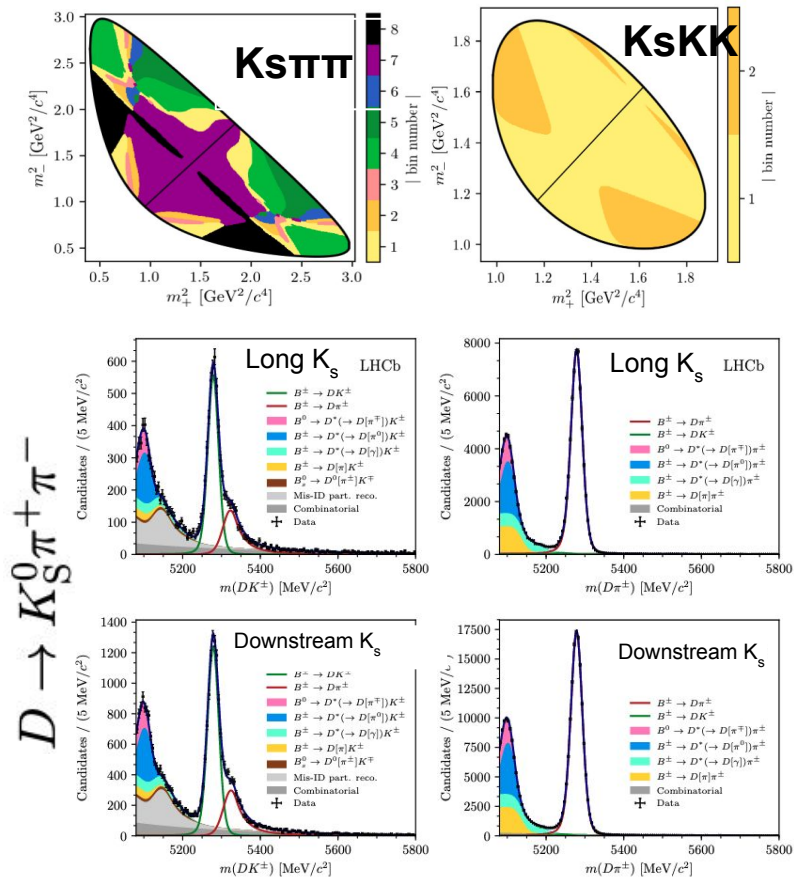
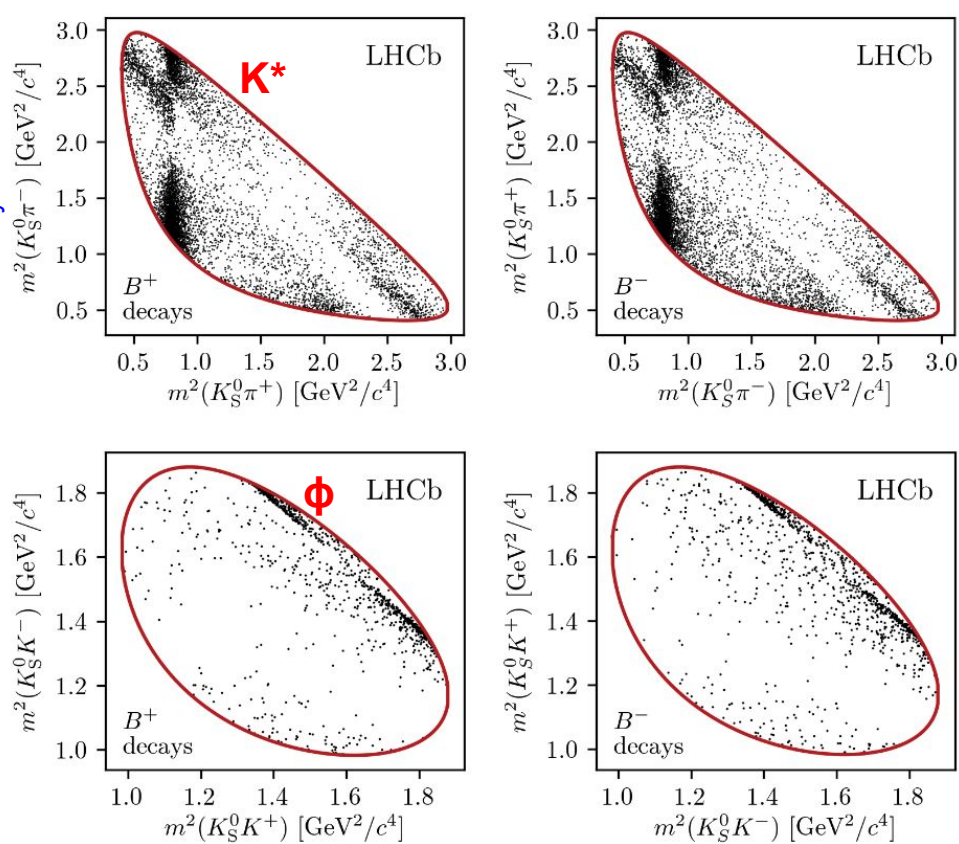
See also Martin's [Belle \(II\) talk](#)

Measurement of gamma in $B^\pm \rightarrow D[Ks0hh]h^\pm$

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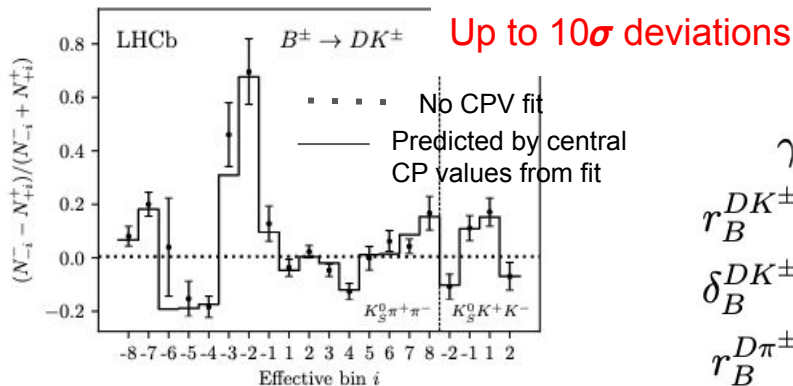
D decay has rich resonance structure - increases sensitivity



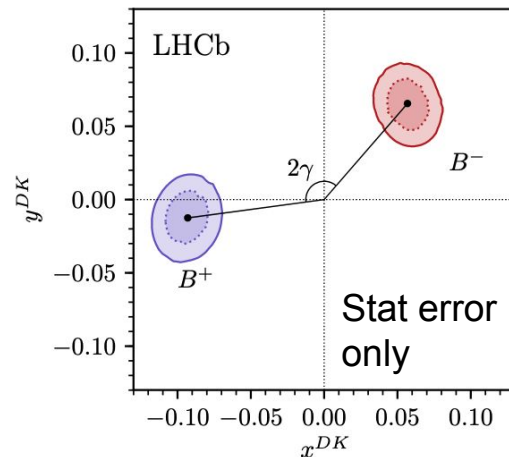
Measurement of gamma in $B^\pm \rightarrow D[Kshh]h^\pm$

9fb⁻¹

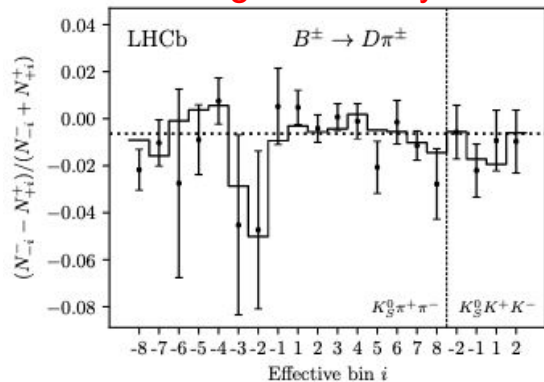
Binned asymmetries



$$\begin{aligned}
 \gamma &= (68.7^{+5.2}_{-5.1})^\circ, \\
 r_B^{DK^\pm} &= 0.0904^{+0.0077}_{-0.0075}, \\
 \delta_B^{DK^\pm} &= (118.3^{+5.5}_{-5.6})^\circ, \\
 r_B^{D\pi^\pm} &= 0.0050 \pm 0.0017, \\
 \delta_B^{D\pi^\pm} &= (291^{+24}_{-26})^\circ.
 \end{aligned}$$

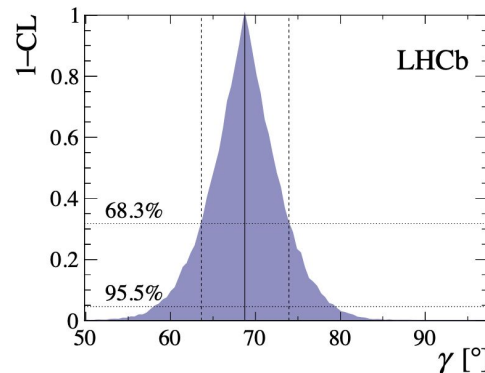


No significant asymmetries



Most precise measurement of gamma from a single analysis

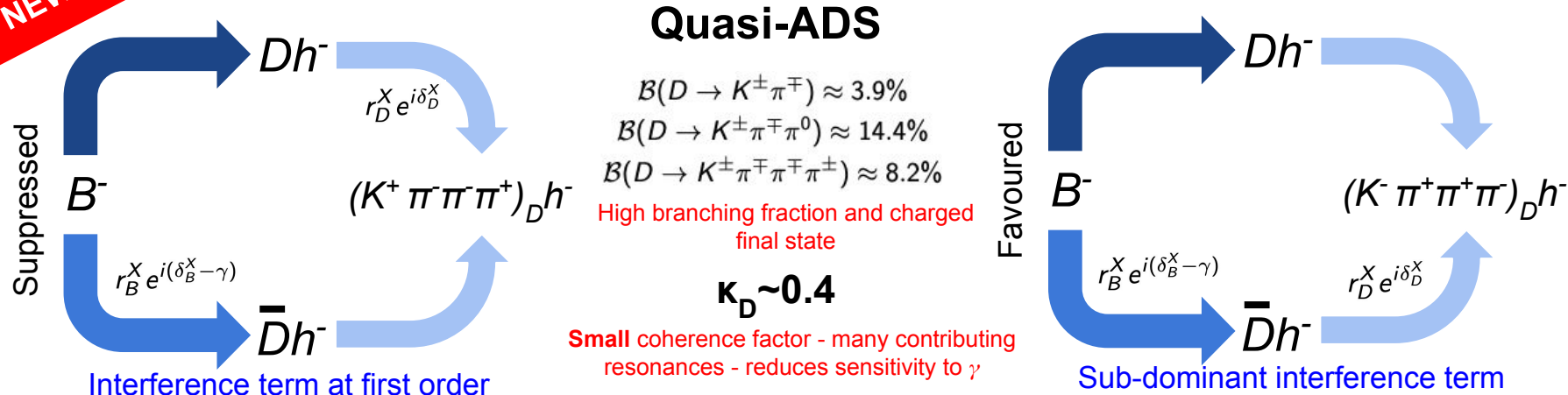
Systematic due to precision of strong phase parameters is of similar size to total LHCb-related systematic ~ 1 degree



Measurement of γ using $B^\pm \rightarrow D[K\pi\pi\pi]h^\pm$ - binned

9fb⁻¹

NEW!



Determine **CP asymmetry** of suppressed modes using favoured mode as a normalisation channel

Small coherence factor - bin phasespace into 4 regions such that local coherence factor, κ_D , in each bin is higher than phase-space integrated value. Binning scheme uses LHCb model but κ_D and average strong phase taken from CLEO-c/BESIII - **Analysis remains model-independent**

Simultaneously measure flavor tagged D^0 decays from inclusive SL modes to provide additional constraints to D parameters

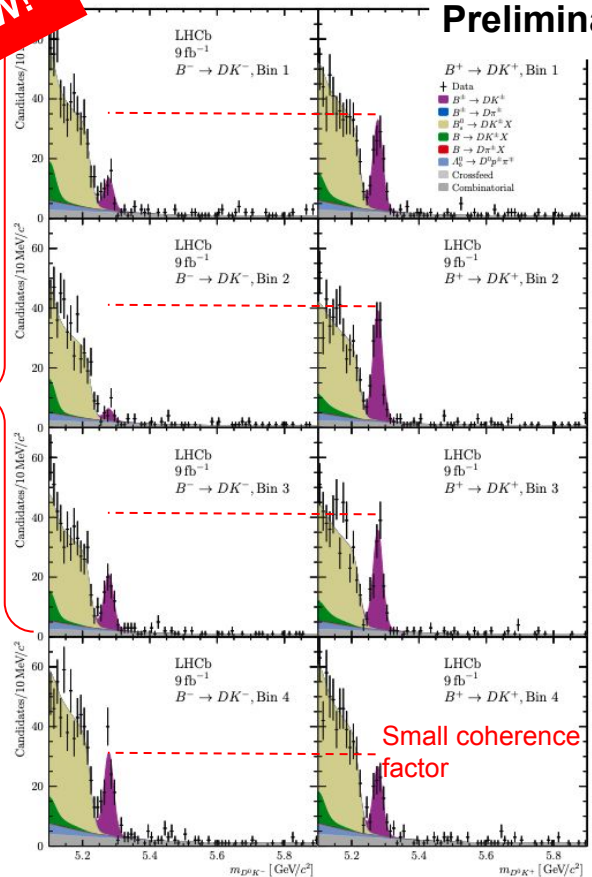
Measurement of γ using $B^\pm \rightarrow D[K\pi\pi\pi]h^\pm$ - binned

9fb⁻¹

NEW!

Preliminary

Large CP asymmetry



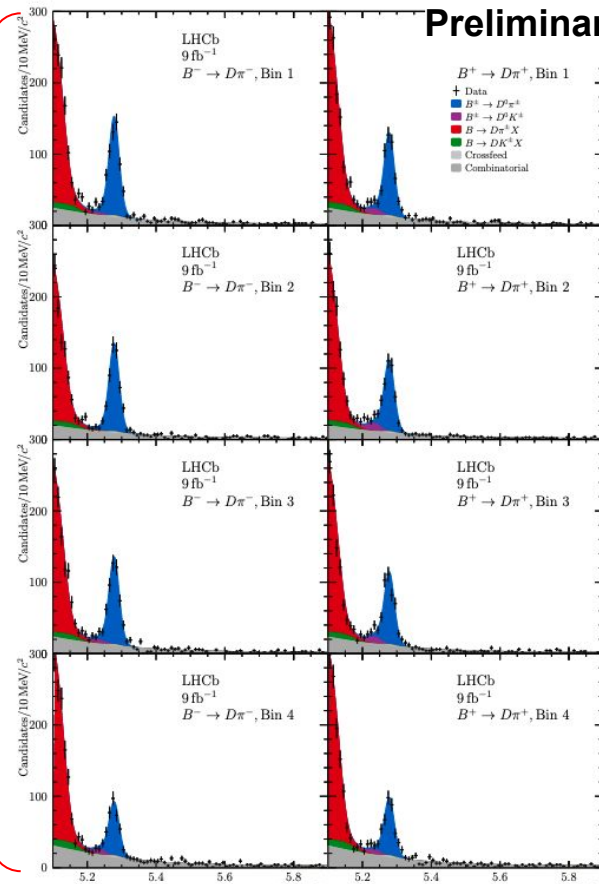
More abundant but lower interference effects as $r_B^{\pi^-} \sim 0.005$ compared with $r_B^{K^-} \sim 0.1$

CPV at 10 σ - largest CP asymmetry observed so far

Opposite sign asymmetries - “strong phase” difference of π



Preliminary



Measurement of γ using $B^\pm \rightarrow D[K\pi\pi\pi]h^\pm$ - binned

9fb⁻¹

NEW!

A_{CP} results of phase space integrated analysis consistent with Run 1 result

$$\begin{aligned} A_K &= -0.321 \pm 0.039 \pm 0.005 \\ A_\pi &= 0.070 \pm 0.019 \pm 0.006 \end{aligned}$$

This analysis

$$\begin{aligned} A_K &= -0.313 \pm 0.102 \pm 0.038 \\ A_\pi &= 0.023 \pm 0.048 \pm 0.005 \end{aligned}$$

Run 1 result

Uncertainties factor
~2.5 smaller

Using binned phasespace results with hadronic parameters of the D decay constrained from BESIII, CLEO-c and LHCb charm mixing results

$$\begin{aligned} \delta_B^K &= (134.6^{+6.0+0.7+8.6}_{-6.0-0.7+8.7})^\circ \\ \delta_B^\pi &= (311.8^{+14.7+3.0+14.7}_{-15.0-2.3+15.0})^\circ \\ r_B^K &= (94.6^{+3.1+0.5+3.0}_{-3.1-0.5+2.3}) \times 10^{-3} \\ r_B^\pi &= (4.5^{+1.1+0.3+0.4}_{-1.0-0.3+0.3}) \times 10^{-3}, \end{aligned}$$

$$\gamma = (54.8^{+6.0+0.6+6.7}_{-5.8-0.6+4.3})^\circ$$

Key input to next gamma combination!

Second most precise determination of gamma using any single D decay mode - systematic uncertainty order of magnitude smaller than statistical

First analysis of this kind!

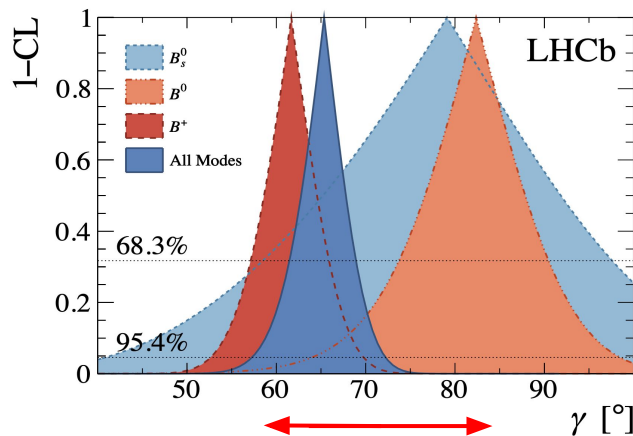
Beauty and Charm LHCb gamma combination

Measurements of gamma in B decays include **charm mixing parameters** and **strong phases** of subsequent D decays

151 observables to determine
52 parameters using
[GammaCombo](#) package

Decay	Parameters	Source
$B^\pm \rightarrow DK^{*\pm}$	$\kappa_{B^\pm}^{DK^{*\pm}}$	LHCb
$B^0 \rightarrow DK^{*0}$	$\kappa_{B^0}^{DK^{*0}}$	LHCb
$B^0 \rightarrow D^\mp \pi^\pm$	β	HFLAV
$B_s^0 \rightarrow D_s^\mp K^\pm (\pi\pi)$	ϕ_s	HFLAV
$D \rightarrow h^+ h^- \pi^0$	$F_{\pi\pi\pi^0}^+, F_{K\pi\pi^0}^+$	CLEO-c
$D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$	$F_{4\pi}^+$	CLEO-c
$D \rightarrow K^+ \pi^- \pi^0$	$r_D^{K\pi\pi^0}, \delta_D^{K\pi\pi^0}, \kappa_D^{K\pi\pi^0}$	CLEO-c+LHCb+BESIII
$D \rightarrow K^\pm \pi^\mp \pi^+ \pi^-$	$r_D^{K3\pi}, \delta_D^{K3\pi}, \kappa_D^{K3\pi}$	CLEO-c+LHCb+BESIII
$D \rightarrow K_S^0 K^\pm \pi^\mp$	$r_D^{K_S^0 K\pi}, \delta_D^{K_S^0 K\pi}, \kappa_D^{K_S^0 K\pi}$	CLEO
$D \rightarrow K_S^0 K^\pm \pi^\mp$	$r_D^{K_S^0 K\pi}$	LHCb

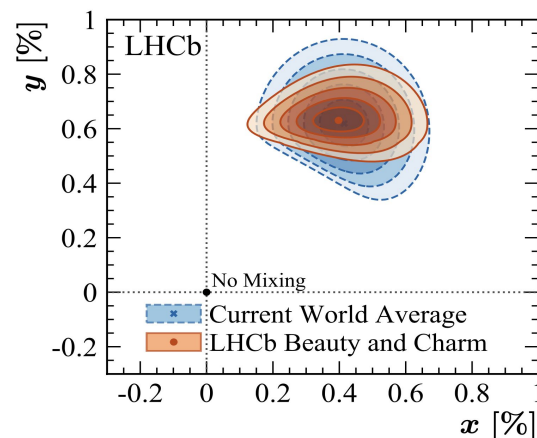
Rely on inputs from independent
measurements



2σ tension
between B^+ and B^0

Wait for full Run 1+2 B^0 results

$$\gamma = (65.4_{-4.2}^{+3.8})^\circ$$



Factor 2 more precise in y !

$$x = (0.400_{-0.053}^{+0.052})\%$$

$$y = (0.630_{-0.030}^{+0.033})\%$$

Dedicated [CERN seminar](#)

New inputs to γ combination

Measurement of the charm mixing parameter $y_{CP} - y_{CP}^{K\pi}$ using two-body D^0 meson decays

LHCb collaboration[†]

Abstract

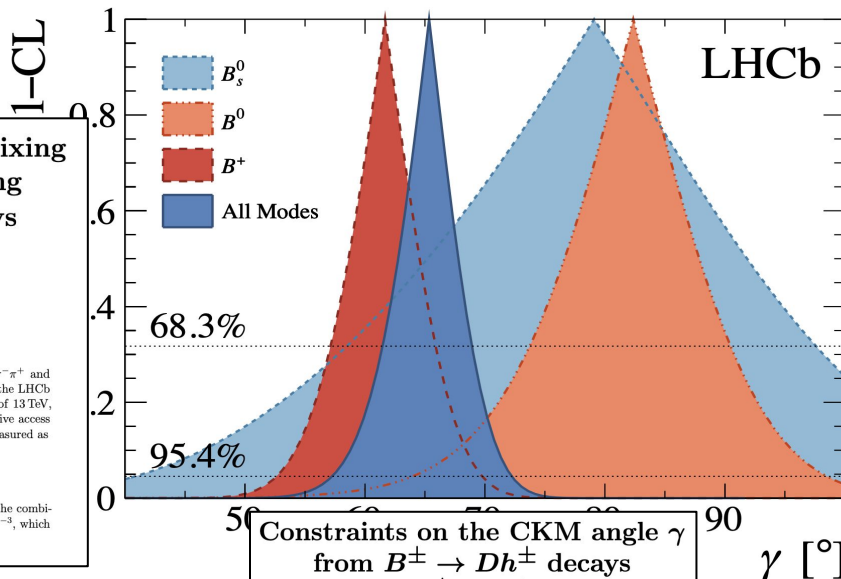
A measurement of the ratios of the effective decay widths of $D^0 \rightarrow \pi^+ \pi^-$ and $D^0 \rightarrow K^+ K^-$ decays over that of $D^0 \rightarrow K^- \pi^+$ decays is performed with the LHCb experiment using proton-proton collisions at a centre-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 6 fb^{-1} . These observables give access to the charm mixing parameters $y_{CP}^{\pi\pi} - y_{CP}^{K\pi}$ and $y_{CP}^{KK} - y_{CP}^{K\pi}$, and are measured as

$$y_{CP}^{\pi\pi} - y_{CP}^{K\pi} = (6.57 \pm 0.53 \pm 0.16) \times 10^{-3},$$

$$y_{CP}^{KK} - y_{CP}^{K\pi} = (7.08 \pm 0.30 \pm 0.14) \times 10^{-3},$$

where the first uncertainties are statistical and the second systematic. The combination of the two measurements is $y_{CP} - y_{CP}^{K\pi} = (6.96 \pm 0.26 \pm 0.13) \times 10^{-3}$, which is four times more precise than the previous world average.

Dedicated [CERN seminar](#)



Constraints on the CKM angle γ from $B^\pm \rightarrow D h^\pm$ decays using $D \rightarrow h^\pm h'^\mp \pi^0$ final states

LHCb collaboration[†]

This paper is dedicated to our friend and colleague Bernhard Spaan.

Abstract

A data sample collected with the LHCb detector, corresponding to an integrated luminosity of 9 fb^{-1} , is used to measure CP observables in $B^\pm \rightarrow D h^\pm$ decays, where $h^{(0)}$ is either a kaon or a pion, and the neutral D meson decay is reconstructed in the three-body final states $K^\pm \pi^\mp \pi^0$, $\pi^\pm \pi^\mp \pi^0$, and $K^\pm K'^\mp \pi^0$. The most suppressed of these modes, $B^\pm \rightarrow [\pi^\pm K'^\mp \pi^0]_D K^\pm$, is observed with a significance greater than seven standard deviations and constraints on the CKM angle γ are calculated from the combination of the measurements.

Measurement of the CKM angle γ with $B^\pm \rightarrow D[K^\mp \pi^\pm \pi^\pm \pi^\mp] h^\pm$ decays using a binned phase-space approach

LHCb collaboration[†]

Abstract

The CKM angle γ is determined from CP -violating observables measured in $B^\pm \rightarrow D[K^\mp \pi^\pm \pi^\pm \pi^\mp] h^\pm$, ($h = K, \pi$) decays, where the measurements are performed in bins of the decay phase space of the D meson. Using the full Run 1 and Run 2 LHCb data sets, γ is determined to be

$$\gamma = (54.8^{+6.0+0.6+6.7}_{-5.8-0.6+4.3})^\circ$$

where the first uncertainty is statistical, the second systematic and the third from the external inputs on the coherence factors and strong phases of the D -meson decay.

CP violation in $B^+ \rightarrow K^+ \pi^0$

5.4fb⁻¹

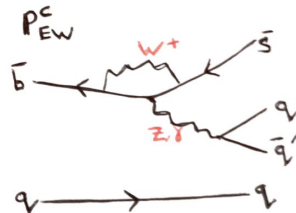
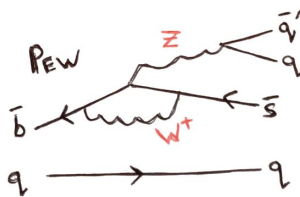
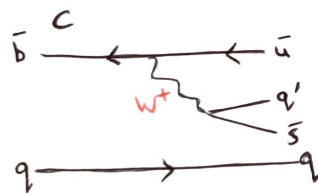
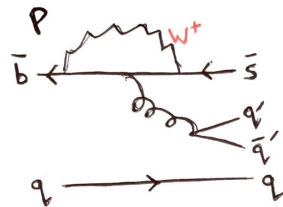
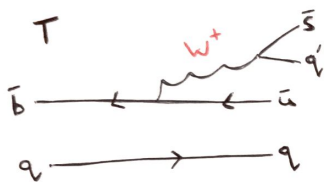
Family of $B \rightarrow K \pi$ decays proceed via several processes - CPV arises from interference - great probe for new physics contributions

Mode	Quark-level amplitude
$B^0 \rightarrow K^+ \pi^-$	$T + P + P_{EW}^C$
$B^+ \rightarrow K^+ \pi^0$	$T + P + C + P_{EW} + P_{EW}^C + A$
$B^+ \rightarrow K^0 \pi^+$	$P + P_{EW}^C + A$
$B^0 \rightarrow K^0 \pi^0$	$P + C + P_{EW} + P_{EW}^C$

Tree and penguin contributions dominate, expect sizable A_{CP} with ~ same magnitude and sign

Almost pure penguin - no sizable A_{CP} expected in SM

arxiv:1101.0675



CP violation in $B^+ \rightarrow K^+ \pi^0$

5.4 fb⁻¹

Within SM amplitudes should obey relations of Isospin symmetry - can predict branching fractions, A_{CP} and lifetime relations amongst the four $B \rightarrow K \pi$ channels

Full SM Isospin sum rule [PLB627 (2005) 82-88]

$$\left[A_{CP}^{K^0 \pi^+} \frac{\mathcal{B}^{K^0 \pi^+}}{\mathcal{B}^{K^+ \pi^-}} - A_{CP}^{K^+ \pi^0} \frac{2\mathcal{B}^{K^+ \pi^0}}{\mathcal{B}^{K^+ \pi^-}} \right] \frac{\tau_{B^0}}{\tau_{B^\pm}} + \left[A_{CP}^{K^+ \pi^-} - A_{CP}^{K^0 \pi^0} \frac{2\mathcal{B}^{K^0 \pi^0}}{\mathcal{B}^{K^+ \pi^-}} \right] = 0$$

K π puzzle: (long standing b anomaly)

$$A_{CP} = \frac{\Gamma(B^- \rightarrow K^- \pi^0) - \Gamma(B^+ \rightarrow K^+ \pi^0)}{\Gamma(B^- \rightarrow K^- \pi^0) + \Gamma(B^+ \rightarrow K^+ \pi^0)}$$

Expect

$$A_{CP}^{K^+ \pi^-} \sim A_{CP}^{K^+ \pi^0}$$



Observe

5.5 σ deviation from equity
(BaBar/Belle/LHCb)

$$A_{CP}(K_S^0 \pi^0) = -0.150 \pm 0.032$$

(non-zero)



$$A_{CP}(K_S^0 \pi^0) = 0.01 \pm 0.10$$

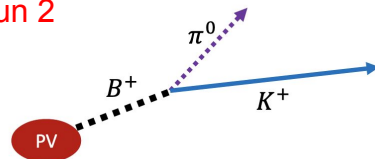
CP violation in $B^+ \rightarrow K^+ \pi^0$

5.4 fb⁻¹

$B^+ \rightarrow h^+ \pi^0$ poses unique challenges for p-p collision environment

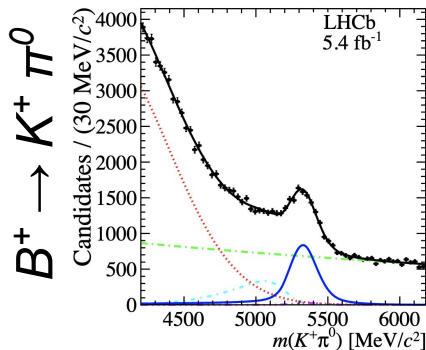
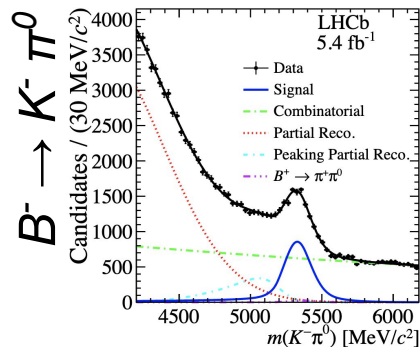
- One track + one ECAL cluster $\rightarrow$ First analysis of one track b hadron decay at hadron collider
- No reconstructible displaced vertex - key signature of b hadron decays

Introduced dedicated trigger for Run 2



$$A_{CP}(B^+ \rightarrow K^+ \pi^0) = A_{\text{raw}}(B^+ \rightarrow K^+ \pi^0) - A_{\text{prod.}}^B - A_{\text{det.}}^K$$

Combined effect of **nuisance asymmetries** measured with high statistics $B \rightarrow J/\psi K$ control channel



$$A_{CP}(B^+ \rightarrow K^+ \pi^0) = 0.025 \pm 0.015 \pm 0.006 \pm 0.003$$

Result is consistent with world average but more precise

$$A_{CP}^{K^+ \pi^-} \sim A_{CP}^{K^+ \pi^0} \text{ does not hold at } > 8\sigma$$

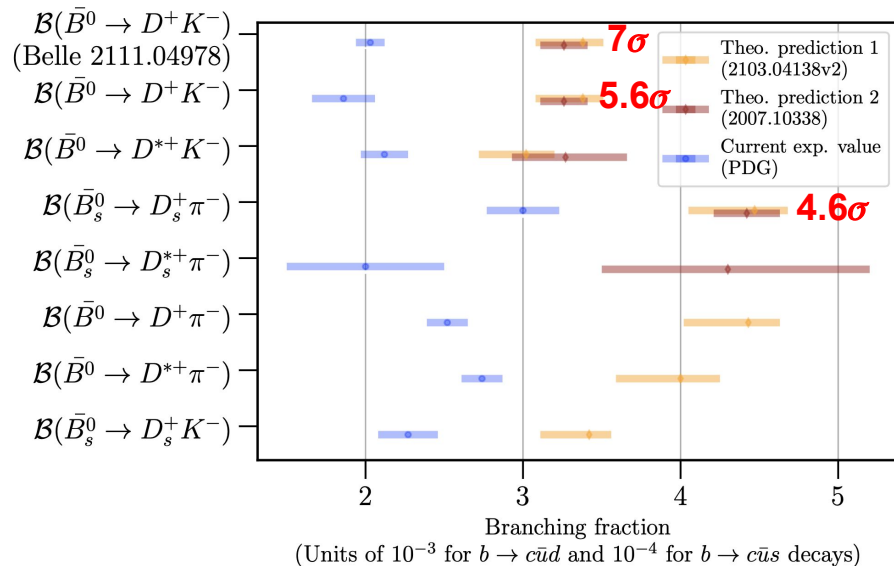
Using result in sum rule $A_{CP}(B^0 \rightarrow K^0 \pi^0) = -0.138 \pm 0.025$ - inconsistent with zero at **5.5 σ**

With Run 3 data projected statistical uncertainty 0.005
Proof of concept for $B^0 \rightarrow K^0 \pi^0$ - run 3 trigger in place!

Confirms and strengthens $K\pi$ puzzle!

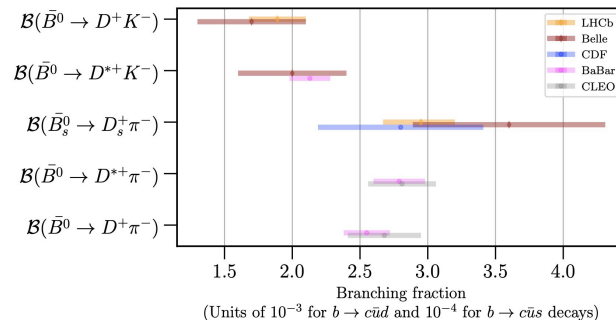
$b \rightarrow c \bar{u} q$ anomaly

Rates of colour allowed non-leptonic decays are not in agreement with SM expectations based on QCDf - tensions pull in the same direction



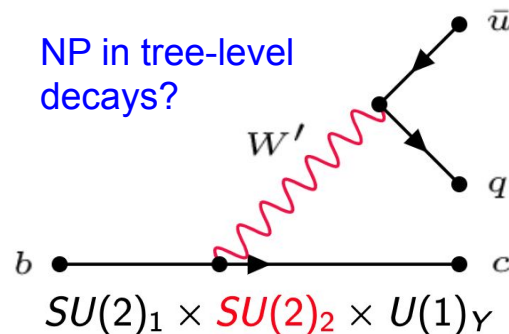
Potential causes?

Experimental inaccuracies?



Unaccounted for QCD effects?

NP in tree-level decays?



Non-leptonic tree-level decays have been considered insensitive to NP - how does this affect our interpretations of gamma? JHEP07(2020)177

$b \rightarrow c \bar{u} q$ anomaly - A_{fs}^s

In SM $\bar{B}_s \rightarrow D_s^+ \pi^-$ is flavour specific and CP-conserving

$$C1 : \mathcal{A}_{\bar{f}} = 0 = \bar{\mathcal{A}}_f \quad C2 : \bar{\mathcal{A}}_{\bar{f}} = \mathcal{A}_f$$

Standard Model

$$\begin{aligned} A_{fs}^s &= \frac{\Gamma(\bar{B}_s(t) \rightarrow f) - \Gamma(B_s(t) \rightarrow \bar{f})}{\Gamma(\bar{B}_s(t) \rightarrow f) + \Gamma(B_s(t) \rightarrow \bar{f})} \\ &= a_{fs}^s \quad \text{---} \quad \text{CPV in mixing} \\ &= a_{sl}^s = (-60 \pm 280) \cdot 10^{-5} \\ &\quad \text{Dominated by SL decays} \end{aligned}$$

General BSM model

Challenge C2 with general BSM contribution

$$\begin{aligned} \mathcal{A}_f &= |\mathcal{A}_f^{\text{SM}}| e^{i\phi^{\text{SM}}} e^{i\varphi^{\text{SM}}} + |\mathcal{A}_f^{\text{BSM}}| e^{i\phi^{\text{BSM}}} e^{i\varphi^{\text{BSM}}} \\ &=: |\mathcal{A}_f^{\text{SM}}| e^{i\phi^{\text{SM}}} e^{i\varphi^{\text{SM}}} (1 + r e^{i\phi} e^{i\varphi}), \\ \phi &= \phi^{\text{BSM}} - \phi^{\text{SM}} \quad \varphi = \varphi^{\text{BSM}} - \varphi^{\text{SM}} \quad r = |\mathcal{A}_f^{\text{BSM}}| / |\mathcal{A}_f^{\text{SM}}| \end{aligned}$$

$$\begin{aligned} A_{fs}^s &\approx a_{fs}^s - 2r \sin \phi \sin \varphi \\ &= a_{fs}^s - A_{\text{dir}}^s \end{aligned}$$

New physics effects may modify value of A_{fs}^s in $\bar{B}_s \rightarrow D_s^+ \pi^-$ to 10^{-2} or above

Neither $A_{fs}^s(D_s^+ \pi^-)$ or its friend $A_{fs}^s(D^+ K^-)$ has been explicitly measured.

$b \rightarrow c \bar{u} q$ anomaly - A_{fs}^s

Experimentally favourable to measure **untagged, time-integrated** asymmetry

→ Significant gain in statistics as no longer need flavour tagging (tagging efficiency $\sim 6\%$ at LHCb)

$$\langle A_{\text{untagged}}^q \rangle \approx A_{\text{dir}}^q - \frac{a_{\text{fs}}^q}{2} (1 - \rho_q) \quad \text{Lose no sensitivity to } A_{\text{dir}}$$

$(\rho_d \approx 0.63 \text{ and } \rho_s \approx 0.001)$

$$A_{\text{raw}} = \frac{N(D_s^+ \pi^-) - N(D_s^- \pi^+)}{N(D_s^+ \pi^-) + N(D_s^- \pi^+)} \quad \text{red arrow}$$

Symmetric $KK\pi\pi$ or $\pi\pi\pi\pi$ final state

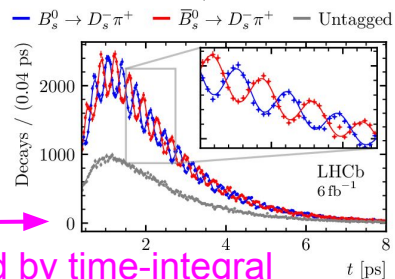
$$\langle A_{\text{untagged}}^s \rangle = A_{\text{raw}} - A_{\text{det}} \quad \text{blue arrow}$$

$$- A_{\text{prod}} \frac{\int_{t=0}^{\infty} e^{-\Gamma_s t} \cos(\Delta M_s t) \epsilon(t) dt}{\int_{t=0}^{\infty} e^{-\Gamma_s t} \cosh(\frac{\Delta \Gamma_s t}{2}) \epsilon(t) dt}$$

Backgrounds well understood through control samples - not a limiting systematic →

$$\sum_i f_{\text{bkg}}^i \cdot A_{\text{bkg}}^i$$

Small A_{prod} is diluted by time-integral



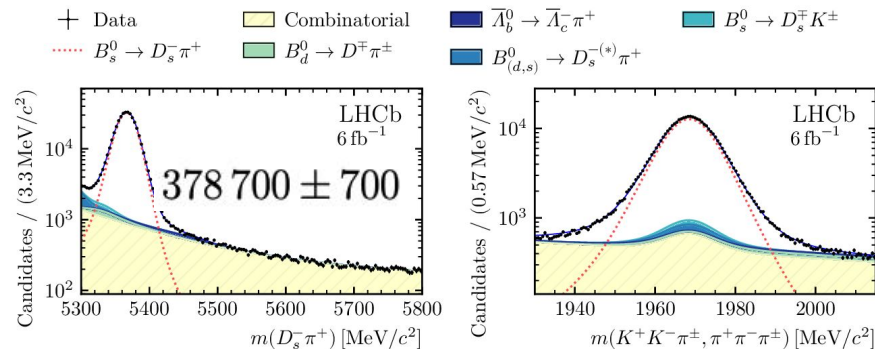
$b \rightarrow c \bar{u} q$ anomaly - A_{fs}^s

LHCb has best prospects for such a measurement

Run 2 sensitivity of 2×10^{-3}

Run 3 sensitivity of 6×10^{-4}

- 15 fb^{-1}
- New software trigger enhancing efficiency of hadronic modes in particular



Existing limits on A_{fs}^s from SL decays are consistent with tiny SM value - if $\langle A_{\text{untagged}}^s \rangle$ non-zero at 10^{-2} level would be clear evidence of NP causing direct CPV independent of any theoretical uncertainties

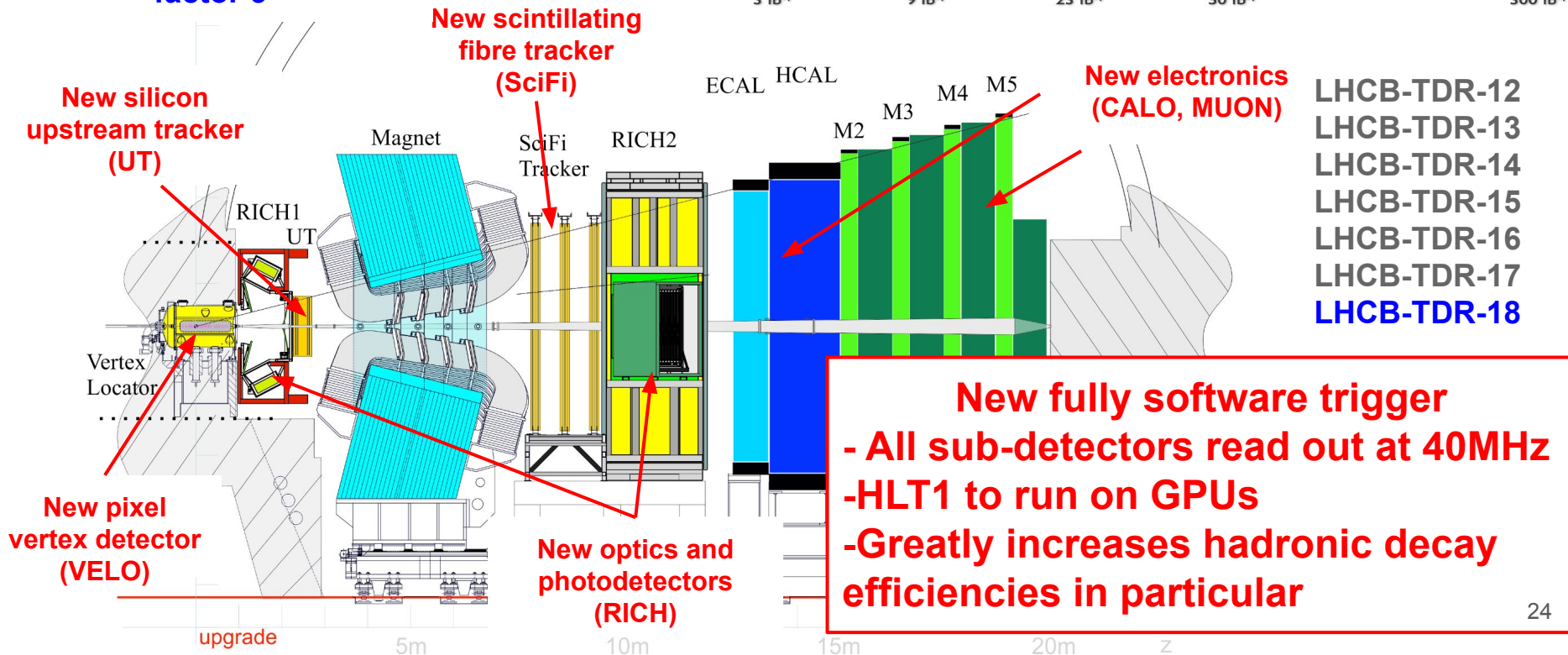
With much improved precision can consider

$$A_{fs}^s(D_s^+ \pi^-) - A_{fs}^s(D_s^+ \ell^- \bar{\nu}_\ell) = A_{dir}^s(D_s^+ \pi^-)$$

Any significant non-zero value would be **unambiguous** signal of NP **independent of any theoretical assumptions**

LHCb Upgrade I

Inst. Lumi. increasing by
factor 5

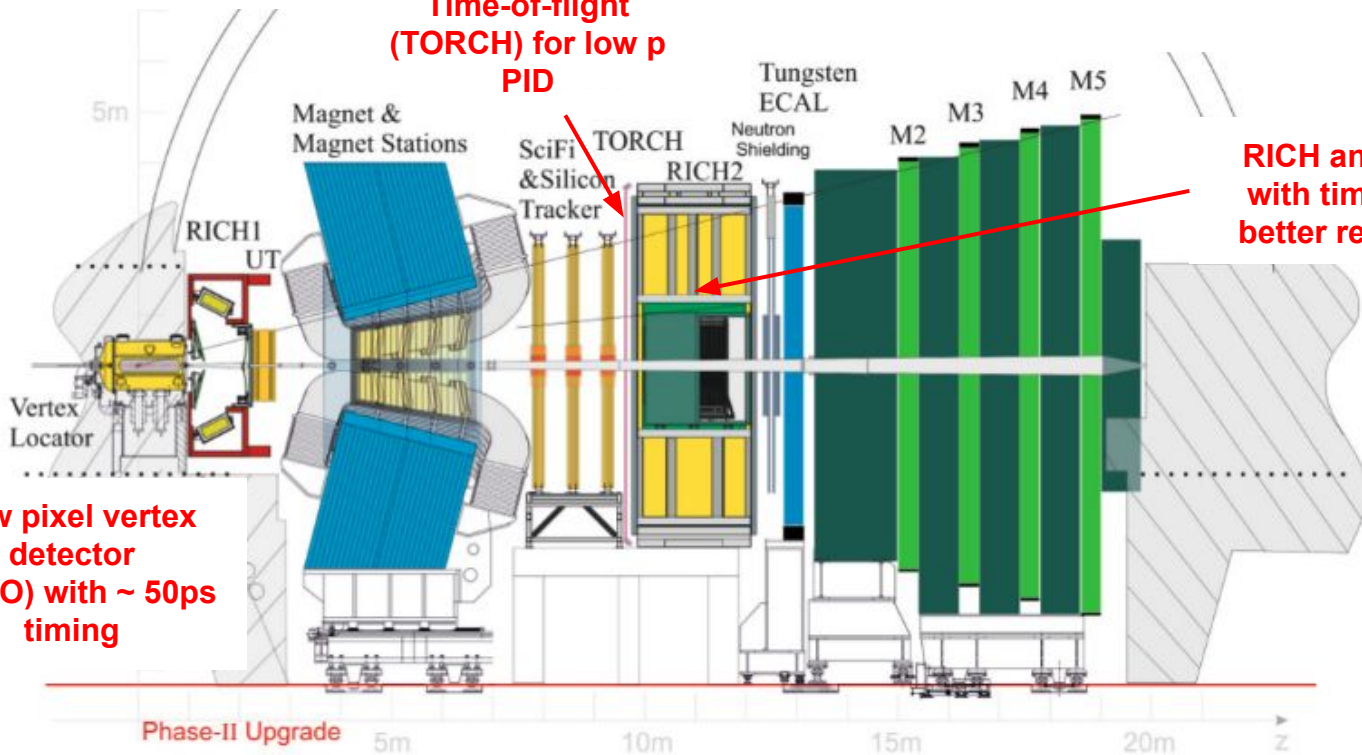


LHCb Upgrade II

Inst. Lumi. increasing by
factor 10

Time-of-flight
(TORCH) for low p
PID

New pixel vertex
detector
(VELO) with ~ 50ps
timing



1 visible interaction

5 visible interactions

50 visible interactions

$$\mathcal{L} = 2 \times 10^{33} \text{cm}^{-2}\text{s}^{-1}$$

$$\mathcal{L} = 1 - 2 \times 10^{34} \text{cm}^{-2}\text{s}^{-1}$$

You are here

HL-LHC



RICH and CALO
with timing and
better resolution

New technologies for
high occupancy
regions

Upgrade II
Physics case

Prospects for γ measurements

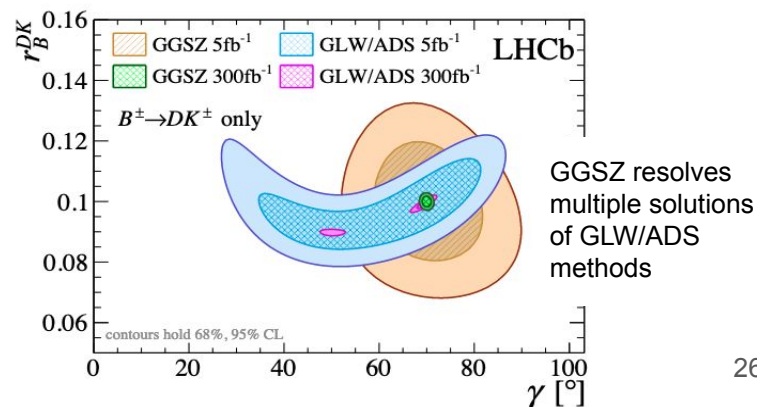
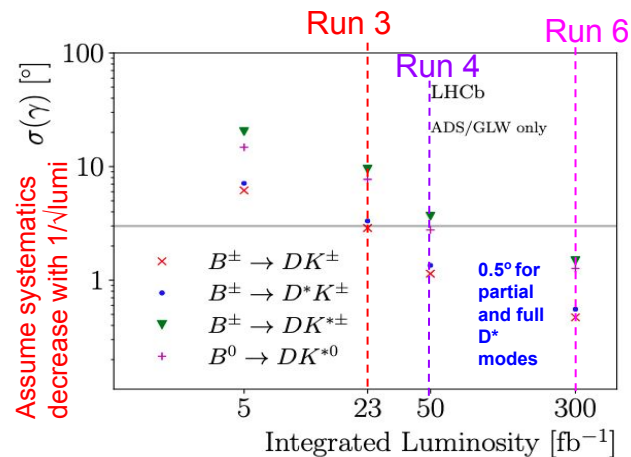
All **ADS/GLW** observables are statistically limited. Will the systematic uncertainty scale with statistics?

- Increased calibration samples for production and reconstruction asymmetries
- Larger MC samples for efficiencies - fast simulation methods
- Background uncertainties that will improve with higher statistics studies

GGSZ golden-mode model-independent analysis gives the most precise single measurement of gamma to date - also statistically limited

- Largest systematic due to external inputs from CLEO-c/BESIII(3fb⁻¹) currently 1 degree
- With 20fb⁻¹ planned from BESIII will reduce to 0.5 degrees [arxiv:1912.05983](https://arxiv.org/abs/1912.05983)

Can extend methods to $B^0 \rightarrow D[-\rightarrow K_s \pi^+ \pi^-] K^+ \pi^-$ “Double Dalitz” model-independent measurements, $r_B \sim 0.25$ - high sensitivity!



Prospects for γ measurements

Sensitivity to π^0 modes will increase

$D \rightarrow K\pi\pi^0$, $D \rightarrow KK\pi^0$, $D \rightarrow \pi\pi\pi^0$

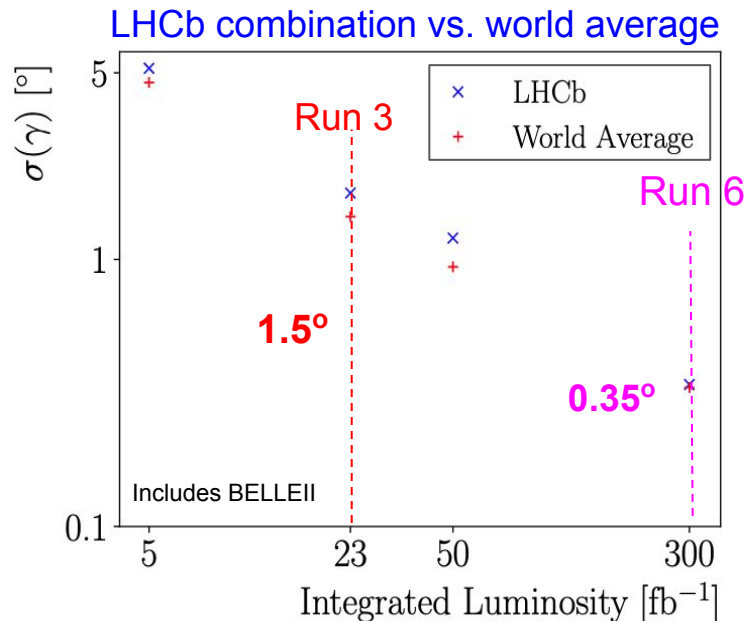
Branching fractions 3-10x larger than 2-body decays

Reconstruction and selection efficiency of π^0 3% currently

In Upgrade II can fully exploit ADS/GLW and GGSZ modes with additional π^0 due to CALO granularity and energy resolution improvements

$B \rightarrow D[K\pi\pi\pi^0]K$ will have sensitivities approaching that of golden $B \rightarrow D[K\pi\pi]K$ channel

CALO improvements also will drive $B \rightarrow K\pi^0$ and $B \rightarrow K^0\pi^0$ measurement sensitivity



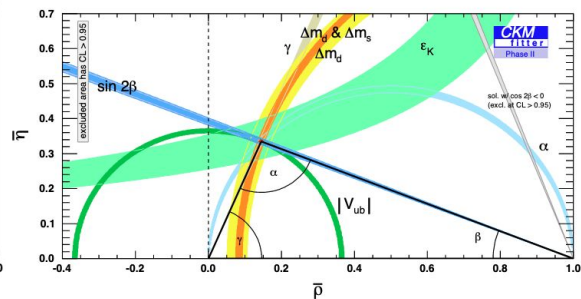
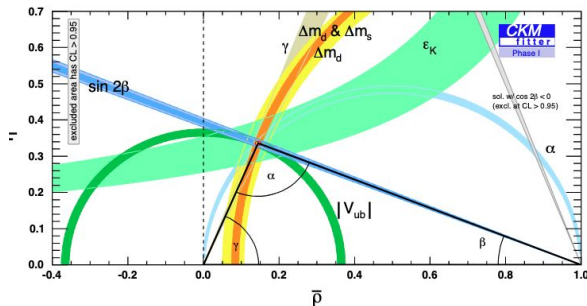
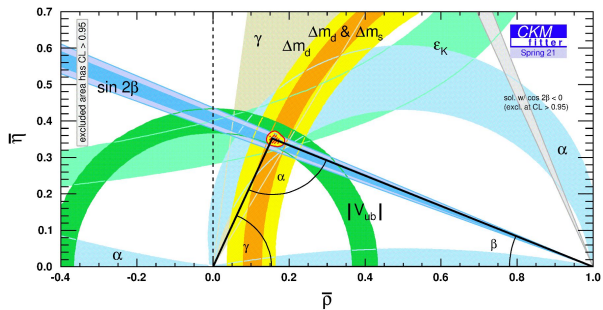
Prospects for γ measurements

Now

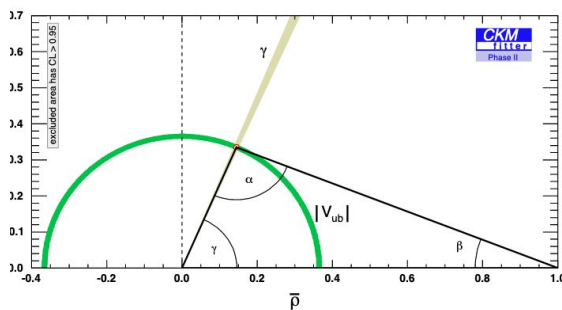
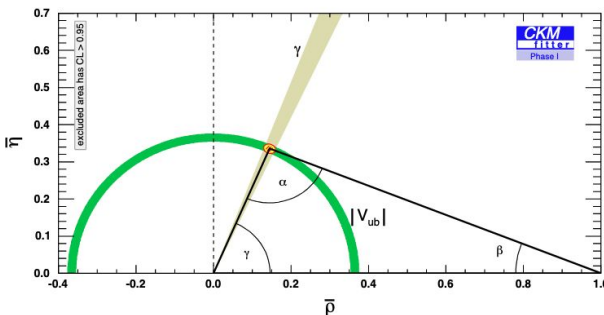
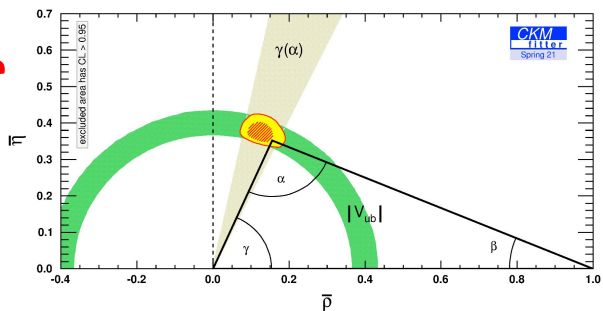
End of Run 3

End of HL-LHC

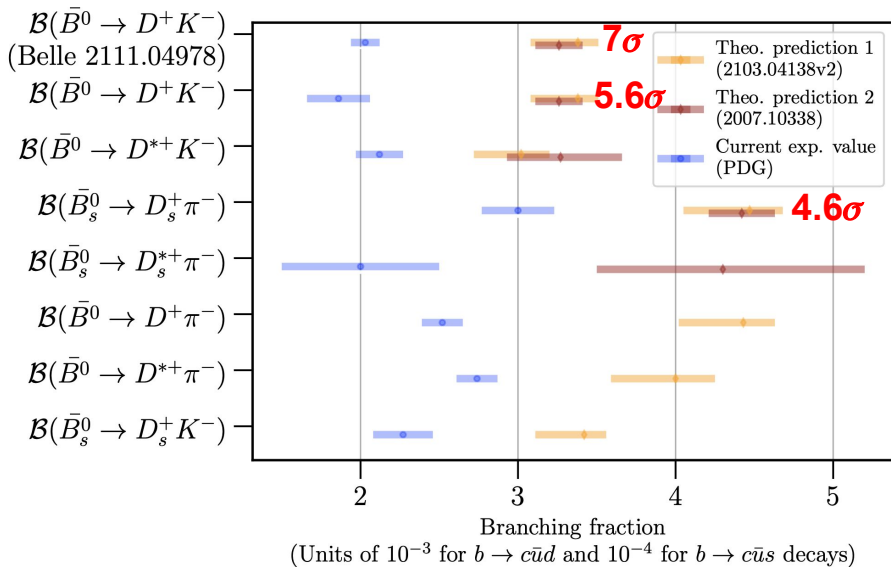
Global



Tree only



Prospects for $b \rightarrow c \bar{u} q$ anomaly



With 300fb^{-1} LHCb will be able to bring statistical and systematic uncertainties down by almost **an order of magnitude**.

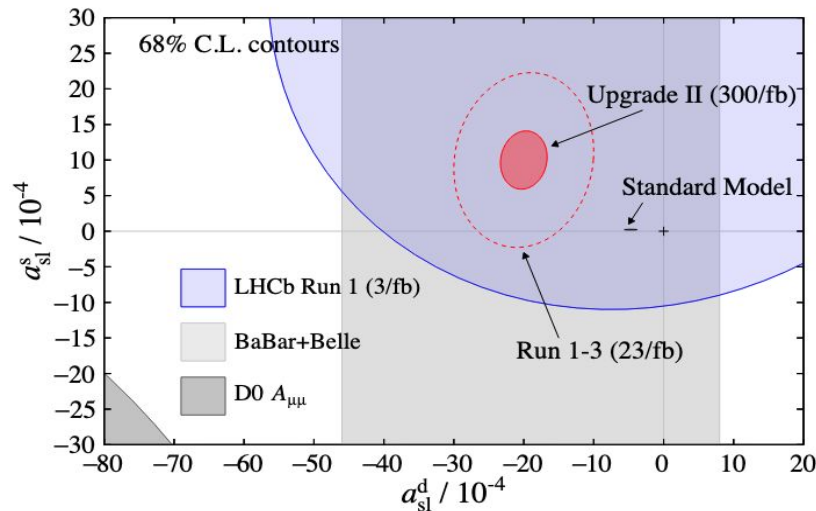
Normalisation channel is limiting systematic- need improvements on $B^0 \rightarrow D \pi$ and B_s mode or fs/fd

With 300fb^{-1} expect uncertainties on A_{sl}^s and A_{fs}^s at 10^{-4} *

Can study

$$A_{fs}^s(D_s^+ \pi^-) - A_{fs}^s(D_s^+ \ell^- \bar{\nu}_\ell) = A_{dir}^s(D_s^+ \pi^-)$$

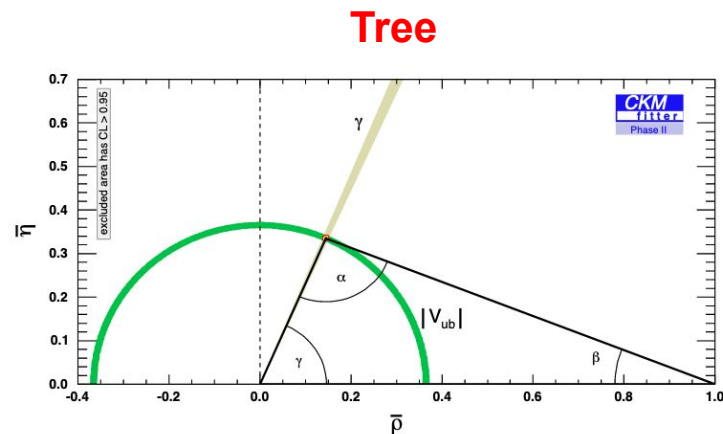
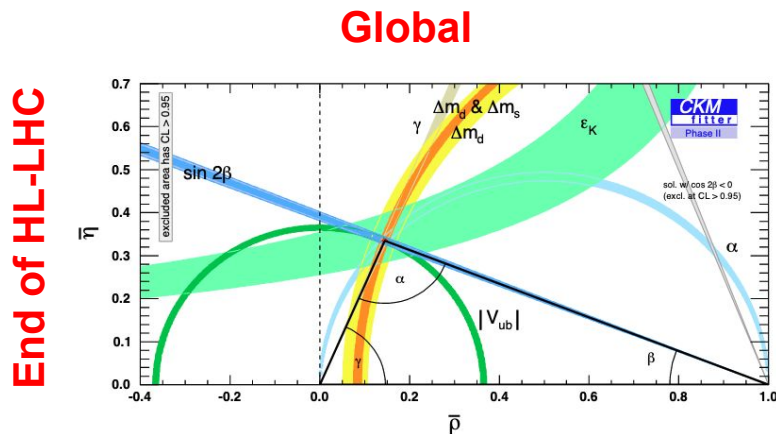
Any significant non-zero value would be unambiguous signal of NP independent of any theoretical assumptions



Summary



- LHCb continues to perform a uniquely wide range of CP violation studies in non-leptonic B meson decays with Run 1 and Run 2 data
- With upgrades I and II LHCb will be sensitive to new modes
- Extremely exciting prospects for anomalies such as K - π puzzle and $b \rightarrow c \bar{u} q$ anomaly



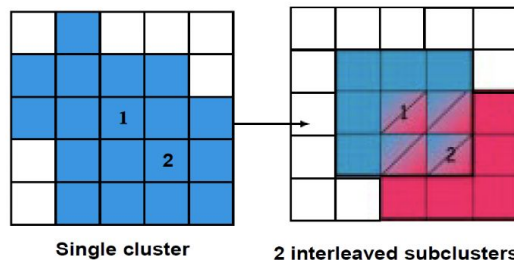
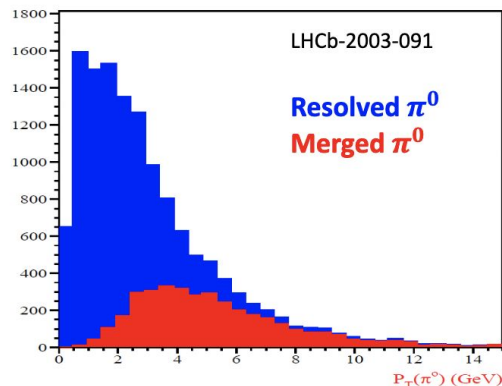
Backup - π^0 reconstruction



π^0 Reconstruction

- Neutral pions identified by decay to two photons
- Below $p_T = 3$ GeV photons can be **resolved** in two separate clusters, at higher energies clusters **merge**
- Cluster separated into two subclusters centered on highest energy deposits according to expected transverse profile
- Photon separation and invariant mass required to be consistent with π^0
- Merged π^0 :
 - + Higher p_T
 - + Reduced combinatorial
 - Wider mass resolution
- For $B^+ \rightarrow K^+ \pi^0$, keep only **merged π^0** to preserve trigger bandwidth

π^0 from $B^0 \rightarrow \pi^+ \pi^- \pi^0$



Taken from [Will Parker](#)

Backup - Beauty and Charm LHCb gamma combination

<i>B</i> decay	<i>D</i> decay	Dataset	Status since Ref. 17
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow h^+h^-$	Run 1&2	Updated
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow h^+\pi^-\pi^+\pi^-$	Run 1	As before
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow h^+h^-\pi^0$	Run 1	As before
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow K_S^0 h^+h^-$	Run 1&2	Updated
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow K_S^0 K^\pm \pi^\mp$	Run 1&2	Updated
$B^\pm \rightarrow D^* h^\pm$	$D \rightarrow h^+h^-$	Run 1&2	Updated
$B^\pm \rightarrow DK^{*\pm}$	$D \rightarrow h^+h^-$	Run 1&2(*)	As before
$B^\pm \rightarrow DK^{*\pm}$	$D \rightarrow h^+\pi^-\pi^+\pi^-$	Run 1&2(*)	As before
$B^\pm \rightarrow Dh^\pm \pi^+ \pi^-$	$D \rightarrow h^+h^-$	Run 1	As before
$B^0 \rightarrow DK^{*0}$	$D \rightarrow h^+h^-$	Run 1&2(*)	Updated
$B^0 \rightarrow DK^{*0}$	$D \rightarrow h^+\pi^-\pi^+\pi^-$	Run 1&2(*)	New
$B^0 \rightarrow DK^{*0}$	$D \rightarrow K_S^0 \pi^+ \pi^-$	Run 1	As before
$B^0 \rightarrow D^\mp \pi^\pm$	$D^+ \rightarrow K^- \pi^+ \pi^+$	Run 1	As before
$B_s^0 \rightarrow D_s^\mp K^\pm$	$D_s^+ \rightarrow h^+ h^- \pi^+$	Run 1	As before
$B_s^0 \rightarrow D_s^\mp K^\pm \pi^+ \pi^-$	$D_s^+ \rightarrow h^+ h^- \pi^+$	Run 1&2	New

151 observables to determine 52 parameters using [GammaCombo](#) package

Backup - Beauty and Charm LHCb gamma combination

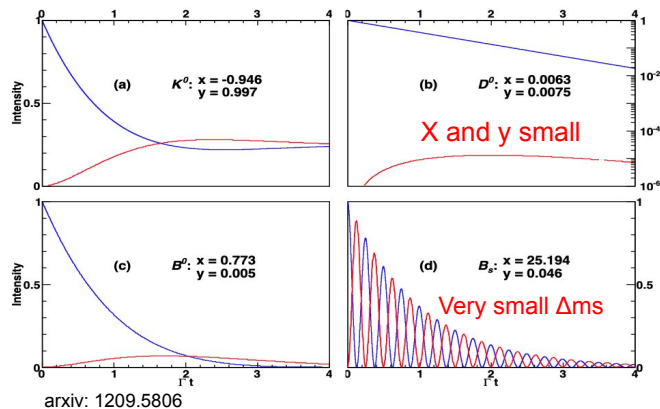
D decay	Observable(s)	Dataset	Status since Ref. 17
$D^0 \rightarrow h^+ h^-$	ΔA_{CP}	Run 1&2	New
$D^0 \rightarrow h^+ h^-$	y_{CP}	Run 1	New
$D^0 \rightarrow h^+ h^-$	ΔY	Run 1&2	New
$D^0 \rightarrow K^+ \pi^-$ (Single Tag)	$R^\pm, (x'^\pm)^2, y'^\pm$	Run 1	New
$D^0 \rightarrow K^+ \pi^-$ (Double Tag)	$R^\pm, (x'^\pm)^2, y'^\pm$	Run 1&2(*)	New
$D^0 \rightarrow K^\pm \pi^\mp \pi^+ \pi^-$	$(x^2 + y^2)/4$	Run 1	New
$D^0 \rightarrow K_S^0 \pi^+ \pi^-$	x, y	Run 1	New
$D^0 \rightarrow K_S^0 \pi^+ \pi^-$	$x_{CP}, y_{CP}, \Delta x, \Delta y$	Run 1	New
$D^0 \rightarrow K_S^0 \pi^+ \pi^-$	$x_{CP}, y_{CP}, \Delta x, \Delta y$	Run 2	New

151 observables to determine 52 parameters using [GammaCombo](#) package

Backup - Charm mixing parameter $y_{CP} - y_{CP}^{K\pi}$

6fb⁻¹

$$\text{Prob}(D^0 \rightarrow \bar{D}^0, t) = \left| \frac{q}{p} \right|^2 \frac{e^{-\Gamma t}}{2} (\cosh(y\Gamma t) - \cos(x\Gamma t))$$



Result compatible with the present world average and more precise by a factor of four!

$$y_{CP} - y_{CP}^{K\pi} = (6.96 \pm 0.26 \pm 0.13) \times 10^{-3}$$

Result will be added to LHCb Beauty and Charm combination soon!

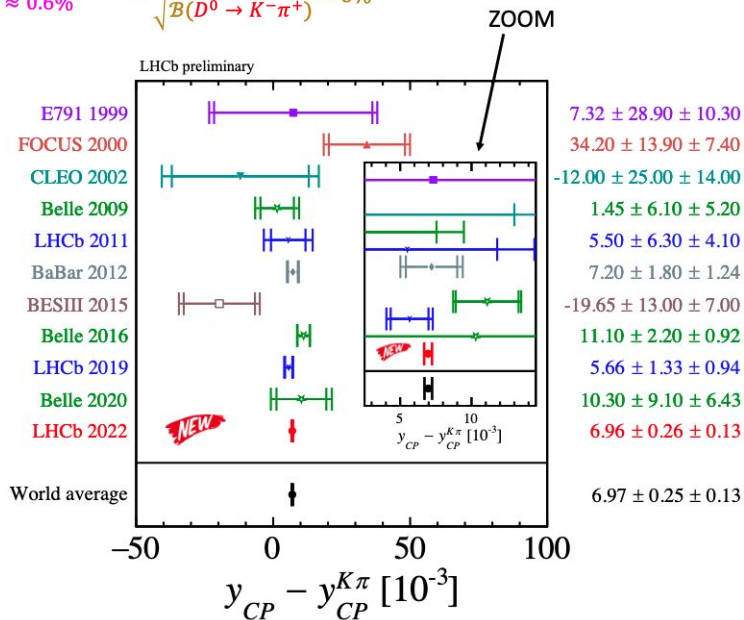
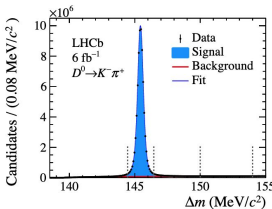
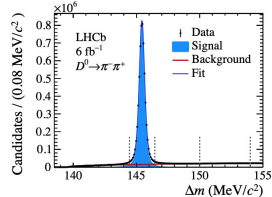
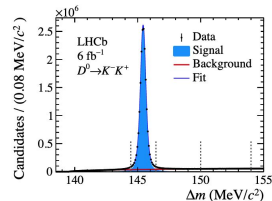
$$\frac{\tau(D^0 \rightarrow K^- \pi^+)}{\tau(D^0 \rightarrow f)} - 1 = y_{CP}^f - y_{CP}^{K\pi} \approx y(1 + \sqrt{R_D})$$

$f = K^- K^+$ and $\pi^- \pi^+$

$$= \frac{\Gamma_1 - \Gamma_2}{2\Gamma} \approx 0.6\%$$

$$= \sqrt{\frac{\mathcal{B}(D^0 \rightarrow K^- \pi^+)}{\mathcal{B}(D^0 \rightarrow K^- \pi^+)}} \approx 6\%$$

Extremely high yields and signal purity



Dedicated [CERN seminar](#)

Backup - Measurement of gamma in $B^\pm \rightarrow D[Kshh]h^\pm$

9fb⁻¹

Can construct observables for each bin i

$$N_{+i}^- = h_{B^-} \left[F_{+i} + \left((x_-^{DK})^2 + (y_-^{DK})^2 \right) F_{-i} + 2\sqrt{F_i F_{-i}} \left(x_-^{DK} c_{+i} + y_-^{DK} s_{+i} \right) \right]$$

Normalisation absorbs
production and detection
asymmetries

Efficiency corrected fraction of events
that fall in bin i

Cosine/sine of strong phase
difference between D^0 and $\bar{D}^0$
averaged over bin i

c_i and s_i inputs taken from combination of BESIII and CLEO results
PRD101 (2020) 112002, PRD102 (2020) 052008

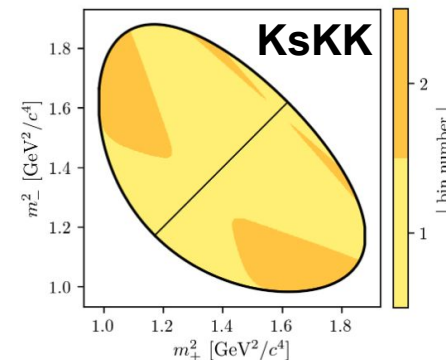
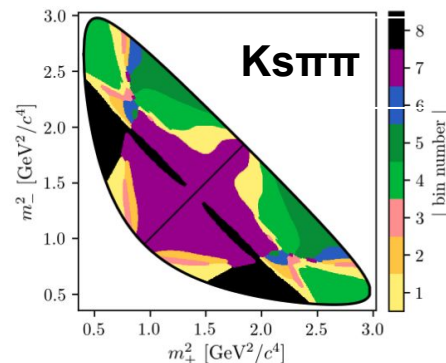
F_i normally determined through semileptonic decays (no CPV) however to
increase sensitivity B $\rightarrow$ D π channel is fit simultaneously with B $\rightarrow$ DK to
determine

$$F_i, x_{\pm}^{DK}, y_{\pm}^{DK}, x_{\pm}^{D\pi}, y_{\pm}^{D\pi}$$

In fact a slightly different parameterisation is used for fit stability

Use of D π channel controls
selections and reconstruction
effects

Negligible contribution to gamma precision for B $\rightarrow$ D π as $r_B^{DK} \approx 20 \times r_B^{D\pi}$

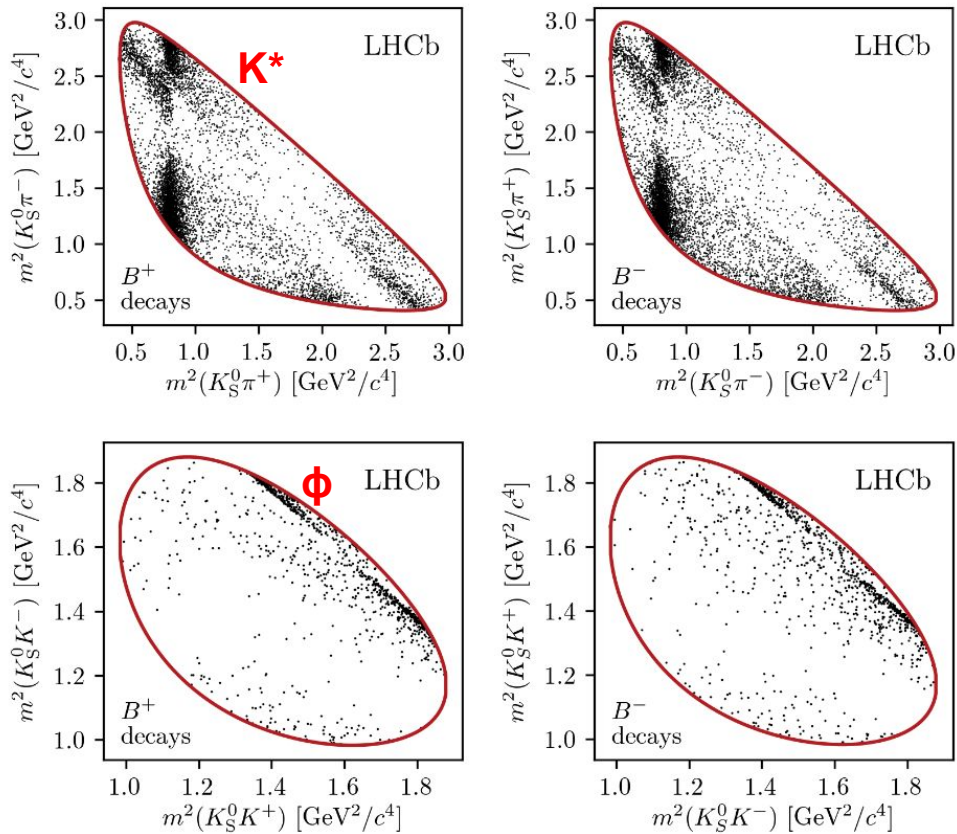


Choice of binning scheme affects only the
statistical sensitivity to gamma - model
independent measurement

Backup - Measurement of gamma in $B^\pm \rightarrow D[Ks0hh]h^\pm$

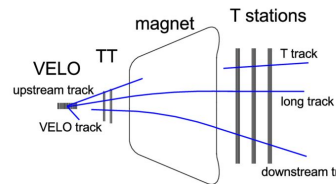
9fb⁻¹

D decay has rich resonance structure - increases sensitivity

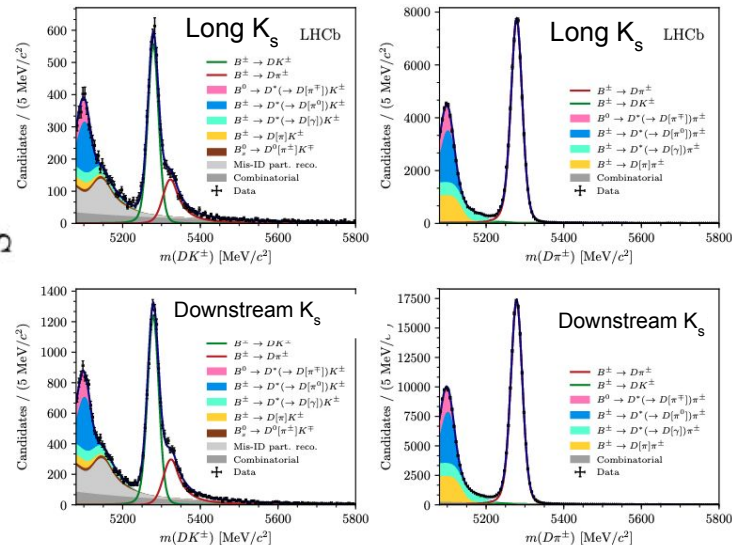


Simultaneous fit to mass distributions in 16 categories

- $B \rightarrow DK$ and $B \rightarrow D\pi$
- B charge
- $D \rightarrow KK$, $D \rightarrow \pi\pi$
- Ks reconstruction

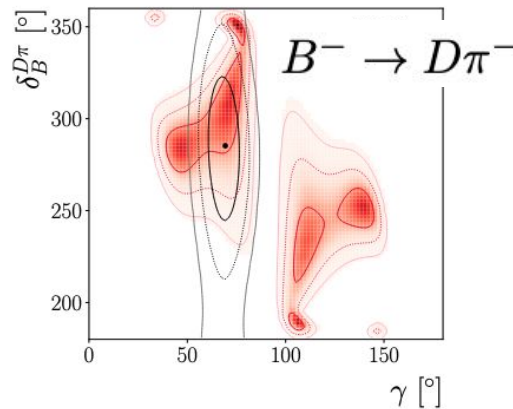
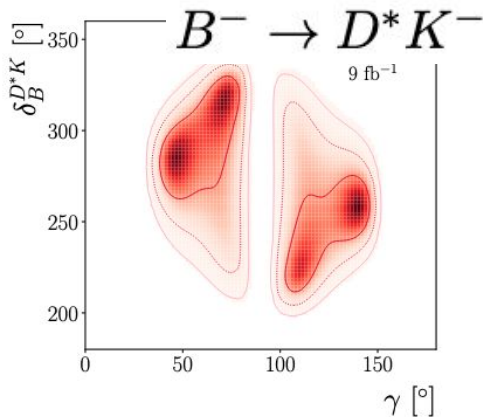
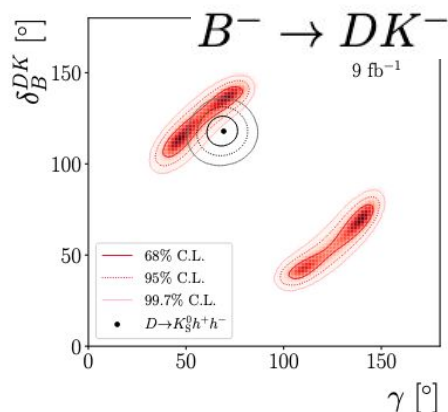


$D \rightarrow K_S^0 \pi^+ \pi^-$



Backup - CP observables in $B^\pm \rightarrow D^{(*)} h^\pm$ decays

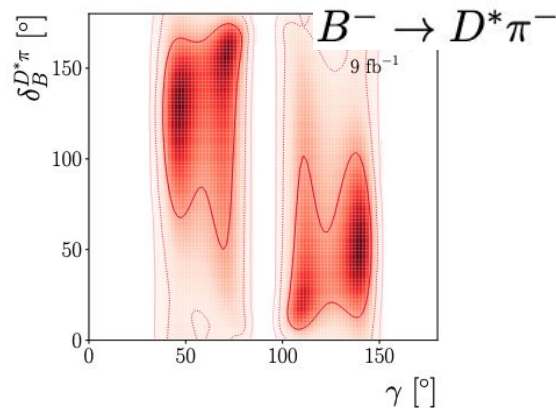
8.7 fb⁻¹



Profile likelihood contours dominated by $B^\pm \rightarrow DK^\pm$ - good agreement with previous LHCb and B factory results

Constructed using charm parameters taken as external constraints from HFLAV average

However $B^\pm \rightarrow D^{(*)} h^\pm$ provide significant sensitivity to these parameters → **motivates combination of gamma and charm results**



Backup - Constraints on γ from $B^\pm \rightarrow Dh^\mp$ with $D \rightarrow h^\pm h^\mp \pi^0$

9fb⁻¹

Eleven experimentally robust CP observables measured through a simultaneous fit to the 16 invariant mass distributions

GLW - asymmetries, double ratios of partial widths

ADS - ratios of partial decay rates, A_{CP} of favoured mode

First observation of $B^\mp \rightarrow [\pi^\mp K^\pm \pi^0]_D K^\mp$ at $>7\sigma!$

Quasi-ADS

Quasi-GLW

