

Exploring dark matter models with global fits

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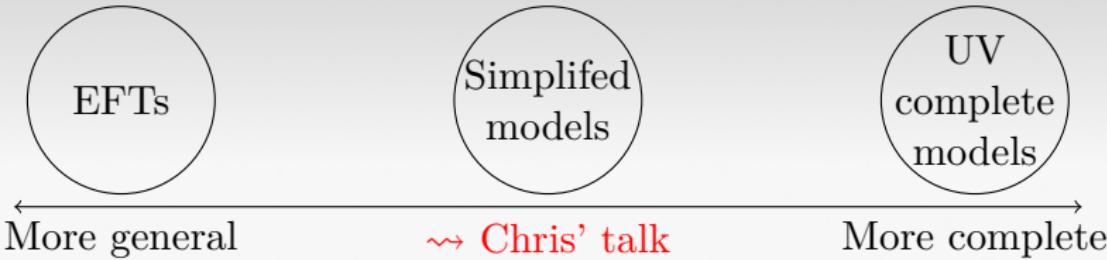
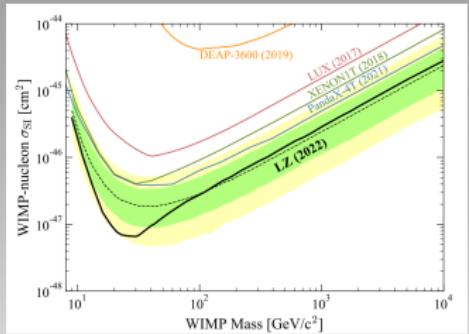
Discrete 2022, 8 Nov 2022

- ① Dark matter and the WIMP miracle
- ② Global fits to explore DM models
- ③ UV complete model: Higgs portal DM
- ④ Effective model: DM EFT
- ⑤ Conclusions and Outlook

Dark matter and the WIMP miracle

DM and WIMPs

- Dark Matter is a thing and there is evidence for it
[Insert obligatory galaxy rotation plot]
- WIMPs are still kicking
 - EW-scale mass, accessible at colliders
 - Just right RD through freeze-out
- Next few years will be critical for the survivability of WIMPs
- Crucial to have rigorous understanding of all possible WIMP models

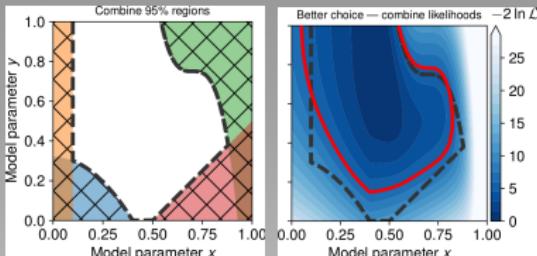


Global fits to explore DM models

Global fits of DM models

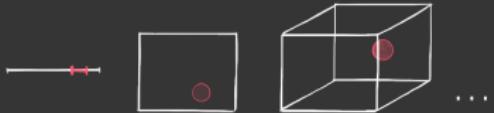
- Multitude of constraints
- Exclusion regions do not properly represent the model predictions
- Composite likelihood

$$\mathcal{L} = \mathcal{L}_{\text{Direct}} \mathcal{L}_{\text{Indirect}} \mathcal{L}_{\text{Collider}} \mathcal{L}_{\text{Astro}} \dots$$



[arXiv:2012.09874 [hep-ph]]

$$\lim_{D \rightarrow \infty} \frac{V_{\text{interesting}}}{V_{\text{total}}} = 0$$



- Multitude of parameters
- Hard to find interesting regions
- Random methods are inefficient
- Need smart sampling strategies (differential, nested, genetic, ...)

- Rigorous statistical interpretations (frequentist / Bayesian)
- Parameter estimation, goodness-of-fit, model comparison, ...

UV complete model: Higgs portal DM

[GAMBIT, Eur.Phys.J.C 77 (2017) 8, 568]

[GAMBIT, Eur.Phys.J.C 79 (2019) 1, 38]

Higgs portal DM

- Scalar DM

$$\mathcal{L}_S = \frac{1}{2}\mu_S^2 S^2 + \frac{1}{2}\lambda_{hS}S^2|H|^2 + \frac{1}{4}S^4 + \frac{1}{2}\partial_\mu S\partial^\mu S$$

- Vector DM

$$\mathcal{L}_V = -\frac{1}{4}W_{\mu\nu}W^{\mu\nu} + \frac{1}{2}\mu_V^2 V_\mu V^\mu - \frac{1}{4!}\lambda_V(V_\mu V^\mu)^2 + \frac{1}{2}\lambda_{hV}V_\mu V^\mu H^\dagger H$$

- Fermionic DM (Dirac)

$$\mathcal{L}_\psi = \bar{\psi}(i\cancel{D} - \mu_\psi)\psi - \frac{\lambda_\psi}{\Lambda_\psi}(\cos\theta\bar{\psi}\psi + \sin\theta\bar{\psi}i\gamma_5\psi)H^\dagger H$$

- Fermionic DM (Majorana)

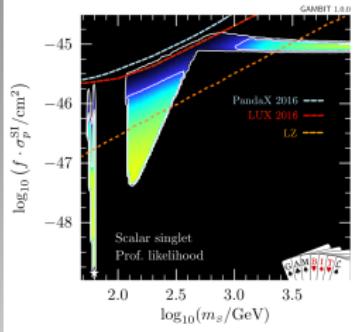
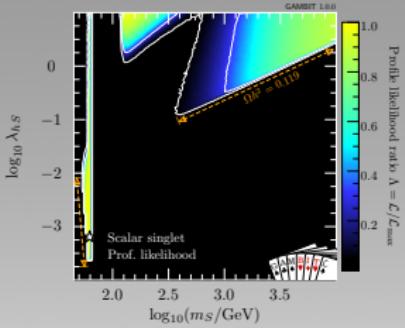
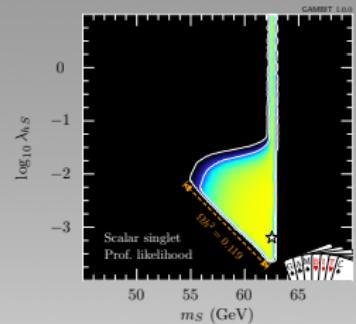
$$\mathcal{L}_\chi = \frac{1}{2}\bar{\chi}(i\cancel{D} - i\mu_\chi)\chi - \frac{1}{2}\frac{\lambda_h\chi}{\Lambda_\chi}(\cos\theta\bar{\chi}\chi + \sin\theta\bar{\chi}i\gamma_5\chi)H^\dagger H$$

Higgs portal DMd

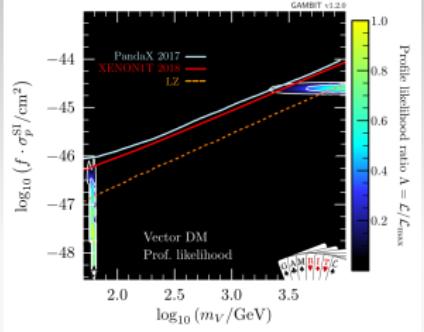
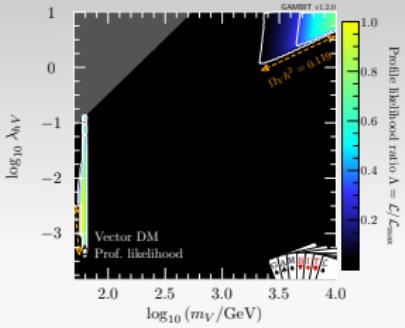
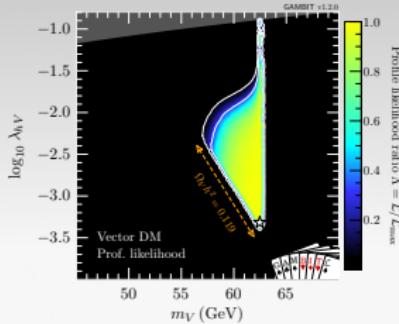
- Direct Detection DDCalc
 → XENON1T, LUX 2016, PandaX 2016-17, CDMSlite, CRESST-II,
 CRESST-III, PICO-60 2017, and DarkSide-50
- Relic abundance DarkSUSY, plc
 → Planck 2015: $\Omega_{\text{DM}} h^2 \leq 0.1188 \pm 0.0010$
- Indirect detection with γ -rays gamLike
 → Pass-8 combined of 15 dSphs from *Fermi*-LAT data
- ID with neutrinos Capt'n General, nulike
 → 79-string IceCube search
- Higgs invisible width
 → $\text{BR}_{\text{inv}}(h \rightarrow \bar{X}X) < 19\% \ (2\sigma)$
- Theoretical constraints
 - Perturbative unitarity $0 \leq \lambda_{hV} \leq \frac{2m_V^2}{v_0^2}$
 - EFT validity $\frac{\lambda_{hX}}{\Lambda_X} \geq \frac{4\pi}{2m_X}$

Higgs portal DM

- Scalar DM

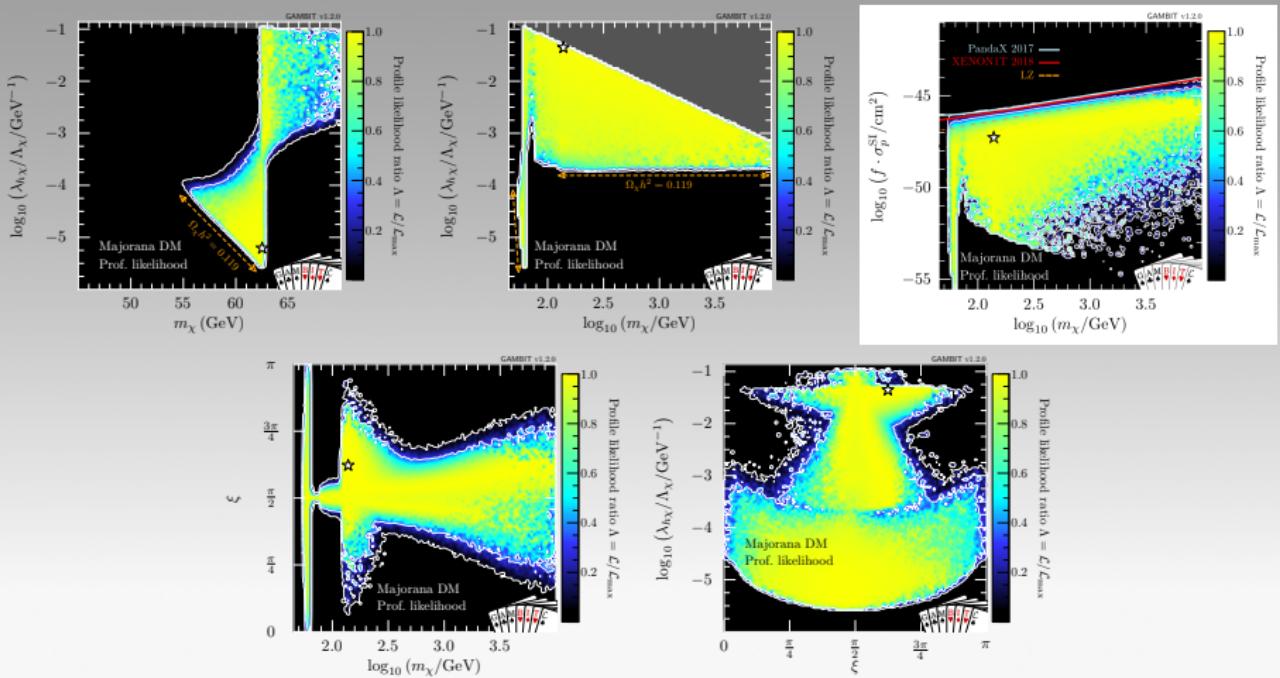


- Vector DM



Higgs portal DM

- Majorana fermion DM (\approx identical to Dirac DM)



Effective model: DM EFT

[GAMBIT, Eur.Phys.J.C 81 (2021) 11, 992]

DM EFT

- Dirac fermionic DM χ : $\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{int}} + \bar{\chi} (i\not{\partial} - m_\chi) \chi$
- Effective interactions (quarks/gluons): $\mathcal{L}_{\text{int}} = \sum_{a,d} \frac{\mathcal{C}_a^{(d)}}{\Lambda^{d-4}} \mathcal{Q}_a^{(d)}$

$$\mathcal{Q}_1^{(5)} = \frac{e}{8\pi^2} (\bar{\chi} \sigma_{\mu\nu} \chi) F^{\mu\nu},$$

$$\mathcal{Q}_2^{(5)} = \frac{e}{8\pi^2} (\bar{\chi} i\sigma_{\mu\nu} \gamma_5 \chi) F^{\mu\nu}$$

$$\mathcal{Q}_{1,q}^{(6)} = (\bar{\chi} \gamma_\mu \chi) (\bar{q} \gamma^\mu q),$$

$$\mathcal{Q}_{2,q}^{(6)} = (\bar{\chi} \gamma_\mu \gamma_5 \chi) (\bar{q} \gamma^\mu q),$$

$$\mathcal{Q}_{3,q}^{(6)} = (\bar{\chi} \gamma_\mu \chi) (\bar{q} \gamma^\mu \gamma_5 q),$$

$$\mathcal{Q}_{4,q}^{(6)} = (\bar{\chi} \gamma_\mu \gamma_5 \chi) (\bar{q} \gamma^\mu \gamma_5 q).$$

$$\mathcal{Q}_1^{(7)} = \frac{\alpha_s}{12\pi} (\bar{\chi} \chi) G^{a\mu\nu} G_{\mu\nu}^a,$$

$$\mathcal{Q}_2^{(7)} = \frac{\alpha_s}{12\pi} (\bar{\chi} i\gamma_5 \chi) G^{a\mu\nu} G_{\mu\nu}^a,$$

$$\mathcal{Q}_3^{(7)} = \frac{\alpha_s}{8\pi} (\bar{\chi} \chi) G^{a\mu\nu} \tilde{G}_{\mu\nu}^a,$$

$$\mathcal{Q}_4^{(7)} = \frac{\alpha_s}{8\pi} (\bar{\chi} i\gamma_5 \chi) G^{a\mu\nu} \tilde{G}_{\mu\nu}^a,$$

$$\mathcal{Q}_{5,q}^{(7)} = m_q (\bar{\chi} \chi) (\bar{q} q),$$

$$\mathcal{Q}_{6,q}^{(7)} = m_q (\bar{\chi} i\gamma_5 \chi) (\bar{q} q),$$

$$\mathcal{Q}_{7,q}^{(7)} = m_q (\bar{\chi} \chi) (\bar{q} i\gamma_5 q),$$

$$\mathcal{Q}_{8,q}^{(7)} = m_q (\bar{\chi} i\gamma_5 \chi) (\bar{q} i\gamma_5 q),$$

$$\mathcal{Q}_{9,q}^{(7)} = m_q (\bar{\chi} \sigma^{\mu\nu} \chi) (\bar{q} \sigma_{\mu\nu} q),$$

$$\mathcal{Q}_{10,q}^{(7)} = m_q (\bar{\chi} i\sigma^{\mu\nu} \gamma_5 \chi) (\bar{q} \sigma_{\mu\nu} q).$$

DM EFT

- Running and mixing

→ For direct detection WCs are needed at $\mu = 2$ GeV (DirectDM)

→ For $\Lambda > m_t(m_t)$:

$$\mathcal{C}_{1,2}^{(5)} = -4 \frac{m_t(m_t)^2}{\Lambda^2} \log \frac{\Lambda^2}{m_t(m_t)^2} \mathcal{C}_{9,10}^{(7)}$$

$$\Delta \mathcal{C}_i^{(7)} = -\mathcal{C}_{i+4,q}^{(7)} \quad (i = 1, 2)$$

$$\Delta \mathcal{C}_i^{(7)} = \mathcal{C}_{i+4,q}^{(7)} \quad (i = 3, 4)$$

- EFT validity, Λ free parameter

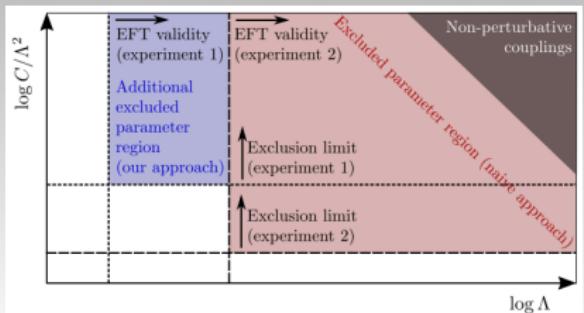
→ DD requires $\Lambda > 2$ GeV

→ Annihilation processes (ID/RD)

require $\Lambda > 2m_\chi$

→ Collider searches $\Lambda > \not{E}_T$

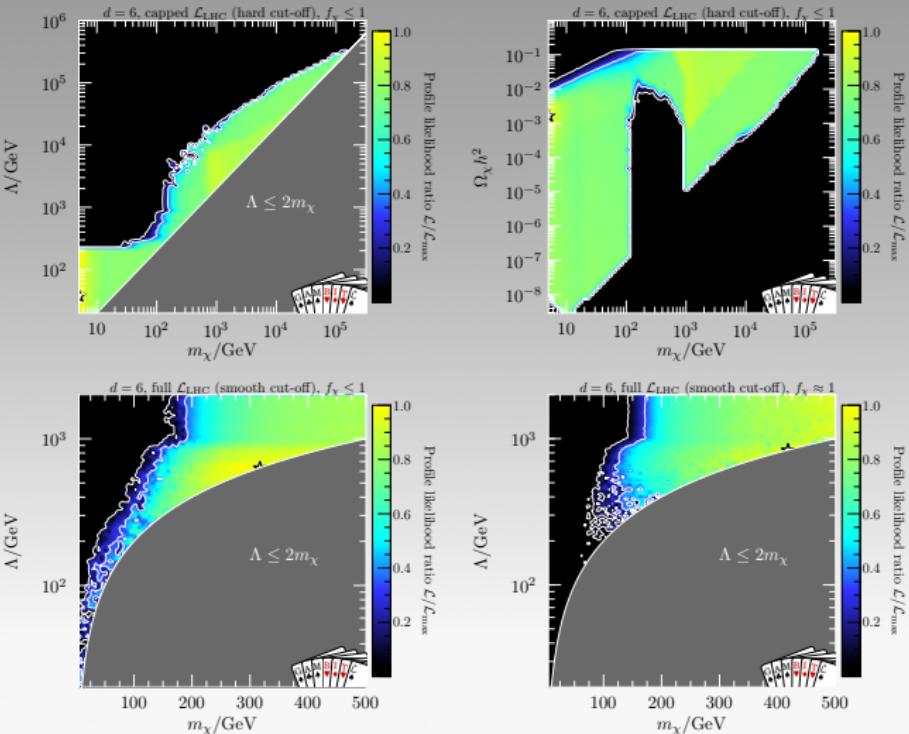
$$\Lambda < \not{E}_T \quad \left\{ \begin{array}{l} \frac{d\sigma}{d\not{E}_T} = 0 \\ \frac{d\sigma}{d\not{E}_T} \rightarrow \frac{d\sigma}{d\not{E}_T} \left(\frac{\not{E}_T}{\Lambda} \right)^{-a} \end{array} \right.$$



DM EFT

- Direct Detection DirectDM, DDCalc
→ XENON1T, LUX 2016, PandaX 2016-17, CDMSlite, CRESST-II,
CRESST-III, PICO-60 2017-19, and DarkSide-50
- Relic abundance CalcHEP, DarkSUSY, plc
→ Planck 2018: $\Omega_{\text{DM}} h^2 \leq 0.120 \pm 0.001$
- Indirect detection with γ -rays CalcHEP, gamLike
→ Pass-8 combined of 15 dSphs from *Fermi*-LAT data
- ID with neutrinos DirectDM, Capt'n General, nulike
→ 79-string IceCube search
- ID constraints from CMB CalcHEP, DarkSUSY, DarkAges
→ 95% CL limit on energy deposition efficiency f_{eff}
- Collider constraints MadGraph_aMC@NLO, Pythia
→ ATLAS 139fb^{-1} mono-jet
→ CMS 36fb^{-1} mono-jet

DM EFT



Conclusions and Outlook

- Conclusions

- WIMP scenarios alive and kicking
- Scalar & Vector DM in trouble
- Fermionic HP DM prefers pseudoscalar interactions
- EFT of DM hints a upper limit on Λ
- Light DM cannot saturate relic abundance

- Outlook

- Next generation DD experiments can constrain very severely HP models
- Very interesting next few years for WIMPs
- WIMPs not the only DM candidate
- See Chris' talk on simplified DM models

Thanks!

Backup

Likelihoods

- Direct Detection

$$\frac{dR}{dE_R} = \frac{\rho}{m_T m_\chi} \int_{v_{\min}}^{\infty} v f(v) \frac{d\sigma}{dE_R} d^3v$$

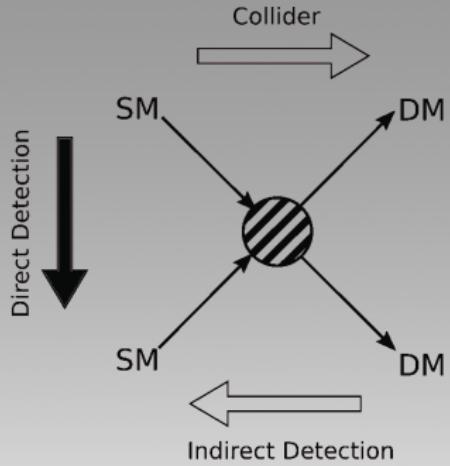
$$v_{\min}(E_R) = \sqrt{\frac{m_T E_R}{2 \mu^2}}$$

→ Non-relativistic operators

$$\mathcal{L}_{\text{NR}} = \sum_{i,N} c_i^N(q^2) \mathcal{O}_i^N ,$$

→ XENON1T, LUX 2016, PandaX 2016-17, CDMSlite, CRESST-II, CRESST-III, PICO-60 2017-19, and DarkSide-50

- Relic abundance $\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle \sigma v_{\text{rel}} \rangle (n_\chi n_{\bar{\chi}} - n_{\chi,\text{eq}} n_{\bar{\chi},\text{eq}})$
 → Planck 2018: $\Omega_{\text{DM}} h^2 \leq 0.120 \pm 0.001$



Likelihoods

- Indirect detection with γ -rays
 - γ -rays from DM annihilation in dSphs

$$\ln \mathcal{L}_{\text{dwarfs}}^{\text{prof.}} = \ln \mathcal{L}_{ki} (\Phi_i \cdot J_k) + \ln \mathcal{L}_J$$

- Pass-8 combined of 15 dSphs from *Fermi*-LAT data

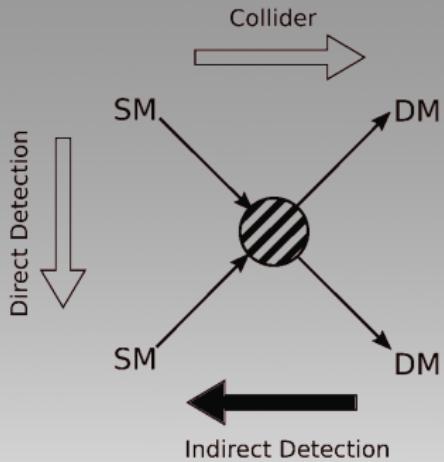
- Indirect detection with ν s

- Solar capture of DM leads to very high energy ν s > solar ν s
- 79-string IceCube search

- Indirect detection constraints from CMB

- Injected energy (γ, e^\pm) changes reion history and optical depth τ
- CMB is sensitive to energy deposition efficiency f_{eff} via combination

$$p_{\text{ann}} = f_\chi f_{\text{eff}} \frac{\langle \sigma v \rangle}{m_\chi}$$



Likelihoods

- Collider constraints
 - Many signatures for DM searches

$$pp \rightarrow \chi\chi j \rightarrow j + \cancel{E}_T$$

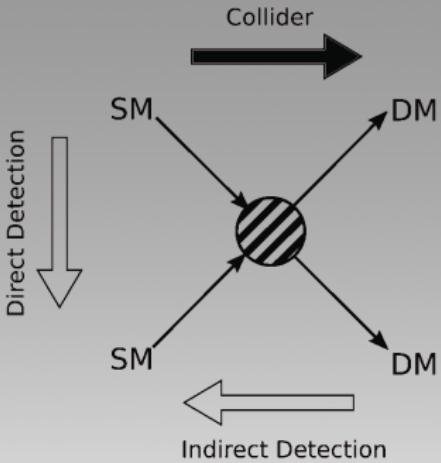
- MadGraph_aMC@NLO \rightsquigarrow Pythia
- Interpolated grids for σ and ϵA
- Events per \cancel{E}_T bin (signal regions)

$$N = L \times \sigma \times (\epsilon A)$$

- ATLAS 139fb^{-1} mono-jet
 - \rightsquigarrow SR with best significance
 - $\rightsquigarrow \mathcal{L}_{\text{ATLAS}}(s_i) \equiv \mathcal{L}_{\text{ATLAS}}(s_i, \hat{\gamma}_i)$

- *Capped* likelihood

$$\mathcal{L}_{\text{cap}}(\mathbf{s}) = \min[\mathcal{L}_{\text{LHC}}(\mathbf{s}), \mathcal{L}_{\text{LHC}}(\mathbf{s} = \mathbf{0})]$$



- CMS 36fb^{-1} mono-jet
 - \rightsquigarrow Profile over systematics
 - $\rightsquigarrow \mathcal{L}_{\text{CMS}}(\mathbf{s}) \equiv \mathcal{L}_{\text{CMS}}(\mathbf{s}, \hat{\gamma})$

Scan framework

- Model parameters

DM mass	m_χ
New physics scale	Λ
Wilson coefficients	$\mathcal{C}_a^{(d)}$

- Nuisance parameters

Local DM density	ρ_0
Most probable speed	v_{peak}
Galactic escape speed	v_{esc}
Running top mass ($\overline{\text{MS}}$ scheme)	$m_t(m_t)$
Pion-nucleon sigma term	$\sigma_{\pi N}$
s -quark contrib. to nucleon spin	Δ_s
s -quark nuclear tensor charge	g_T^s
s -quark charge radius of the proton	r_s^2

- Needs smart sampling to efficiently scan over all parameters and explore interference effects among WCs

Scan framework

GAMBIT: The Global And Modular BSM Inference Tool

gambit.hepforge.org

EPJC **77** (2017) 784

arXiv:1705.07908

- Extensive model database – not just SUSY
 - Extensive observable/data libraries
 - Many statistical and scanning options (Bayesian & frequentist)
 - *Fast* LHC likelihood calculator
 - Massively parallel
 - Fully open-source
- Fast definition of new datasets and theories
 - Plug and play scanning, physics and likelihood packages



Members of:

ATLAS, Belle-II, CLIC,
CMS, CTA, *Fermi*-LAT,
DARWIN, IceCube, LHCb,
SHiP, XENON

Authors of:

DarkSUSY, DDCalc, Diver, FlexibleSUSY, gamlike, GM2Calc,
IsaTools, nulike, PolyChord, Rivet, SoftSUSY, SuperISO, SUSY-
AI, WIMPSim

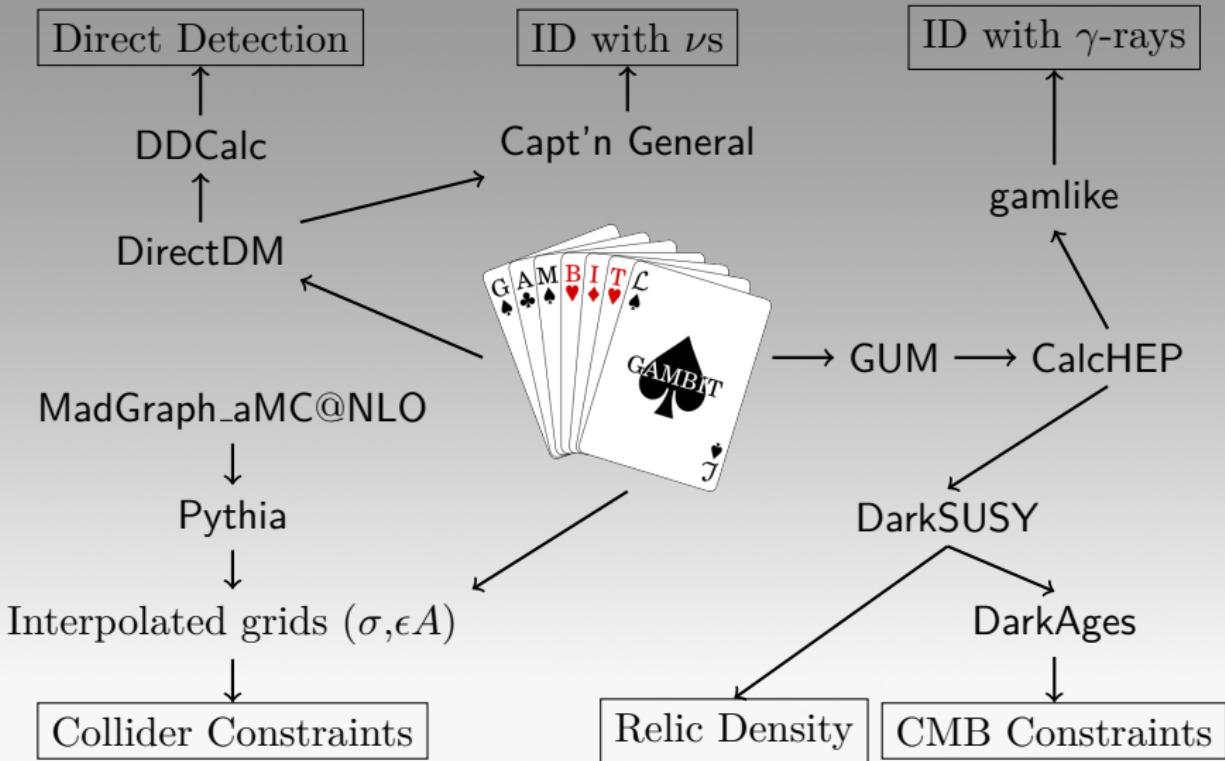


Recent collaborators:

F Agocs, V Ananyev, P Athron, C Balázs, A Beniwal, J Bhom, S Bloor, T Bringmann, A Buckley, J-E Camargo-Molina, C Chang, M Chrzaszcz, J Conrad, J Cornell, M Danninger, J Edsjö, B Farmer, A Fowlie, T Gonzalo, P Grace, W Handley, J Harz, S Hoof, S Hotinli, F Kahlhoefer, N Avis Kozar, A Kvellestad, P Jackson, A Ladhu, N Mahmoudi, G Martinez, MT Prim, F Rajec, A Raklev, J Renk, C Rogan, R Ruiz, I Sáez Casares, N Serra, A Scaffidi, P Scott, P Stöcker, W Su, J Van den Abeele, A Vincent, C Weniger, M White, Y Zhang

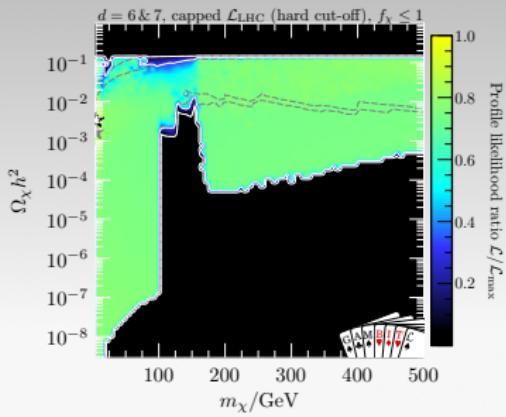
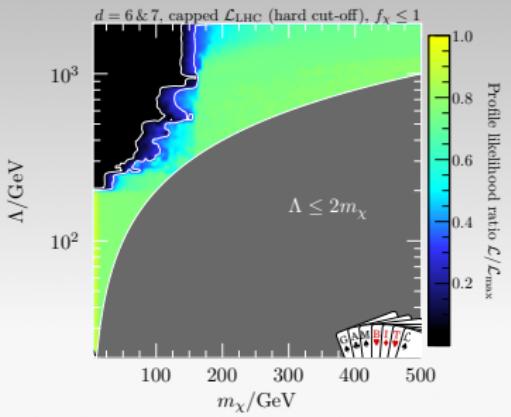
70+ participants in 11 experiments and 14 major theory codes

Scan framework



Results

- Include dim-7 operators, $\Omega_{\text{DM}} h^2$ upper limit, LHC loglike *capped*
 - No change on large Λ - small m_χ region
 - Neither $\mathcal{Q}_{1-4}^{(7)}$ (LHC) nor $\mathcal{Q}_{5-10,q}^{(7)}$ (suppressed) contribute to ann xsec
 - However, RD can be saturated for $m_\chi < 100$ GeV (and small Λ)
 - $\mathcal{Q}_3^{(7)}$ and $\mathcal{Q}_{7,q}^{(7)}$ give unconstrained signals in DD and ID
 - Similar fits to LHC excesses, even when dim-6 ops are zero



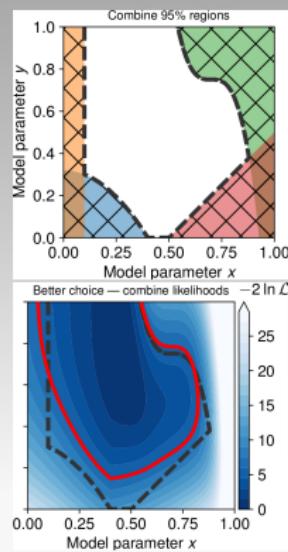
Global fits

- Combine all constraints into a **composite likelihood**

$$\mathcal{L} = \mathcal{L}_{\text{Collider}} \mathcal{L}_{\text{Higgs}} \mathcal{L}_{\text{DM}} \mathcal{L}_{\text{Flavour}} \dots$$

- Perform an extensive **parameter scan**

- Old-school sampling methods (random, grid) are inefficient
- Harder to make statement about statistics
- Need **smart sampling strategies** (differential, nested, genetic, ...)
- **Rigorous** statistical interpretation (frequentist/Bayesian)
 - Goodness-of-fit
 - Parameter estimation
 - Model comparison



[arXiv:2012.09874 [hep-ph]]

Modules (Bits)

- Physics Modules

- **ColliderBit**: collider searches [Eur.Phys.J. C77 (2017) no.11, 795]
- **DarkBit**: relic density, dd,... [Eur.Phys.J. C77 (2017) no.12, 831]
- **FlavBit**: flavour observables [Eur.Phys.J. C77 (2017) no.11, 786]
- **SpecBit**: spectra, RGE running [Eur.Phys.J. C78 (2018) no.1, 22]
- **DecayBit**: decay widths [Eur.Phys.J. C78 (2018) no.1, 22]
- **PrecisionBit**: precision tests [Eur.Phys.J. C78 (2018) no.1, 22]
- **NeutrinoBit**: neutrino likelihoods [Eur.Phys.J.C 80 (2020) no.6, 569]
- **CosmoBit**: cosmological constraints [JCAP 02 (2021) 022]

- **ScannerBit** : stats and sampling

- Diver, GreAT, Multinest, Polychord, ...

[Eur.Phys.J. C77 (2017) no.11, 761]

- **Models**: hierarchical model database

- **Core** : dependency resolution

[Eur.Phys.J. C78 (2018) no.2, 98]

- **Backends** : External tools to calculate observables

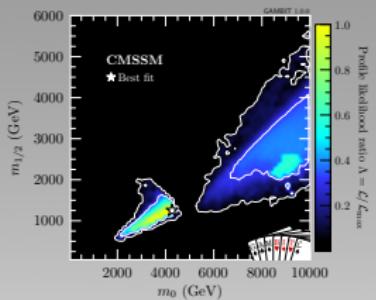
- **GUM**: Autogeneration of code

[S. Bloor, TG, P. Scott et. al., soon]

Examples

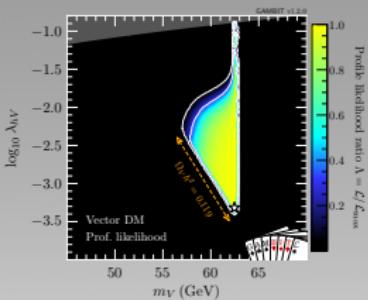
CMSSM

[Eur.Phys.J.C 77 (2017) 12, 824]



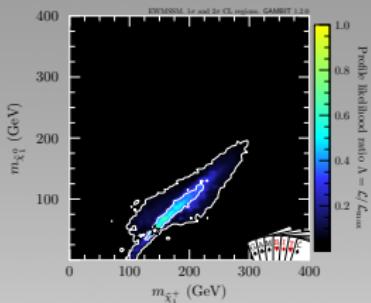
Higgs-portal DM

[Eur.Phys.J.C 79 (2019) 1, 38]



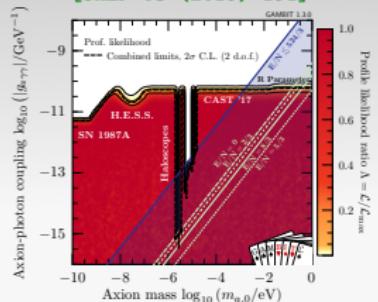
MSSM-EW

[Eur.Phys.J.C 79 (2019) 5, 395]



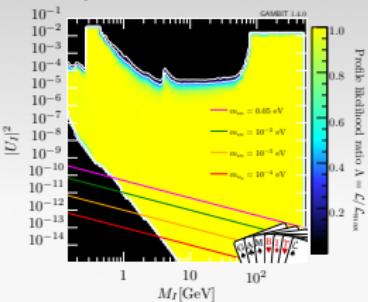
QCD axions

[JHEP 03 (2019) 191]



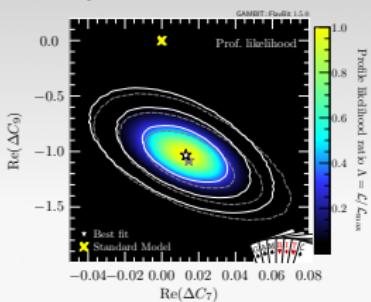
Right-Handed Neutrinos

[Eur.Phys.J.C 80 (2020) 6, 569]



Flavour EFT

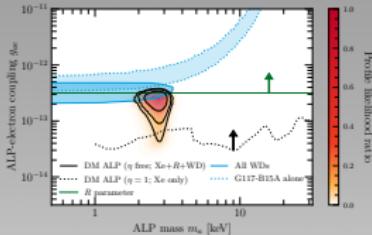
[Eur.Phys.J.C 81 (2021) 12, 1076]



Examples

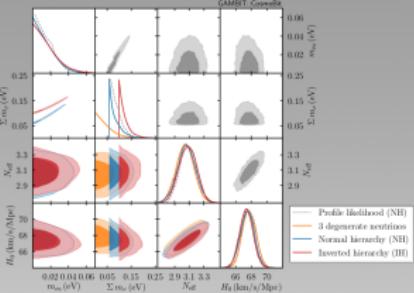
DM ALPs

[JHEP 05 (2021) 159]



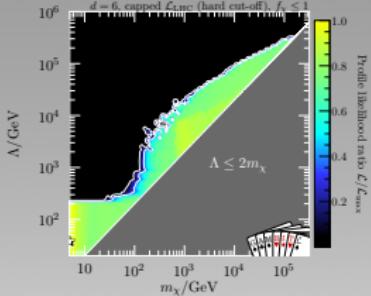
Neutrino Masses

[Phys. Rev. D 103 (2021) 12, 123508]



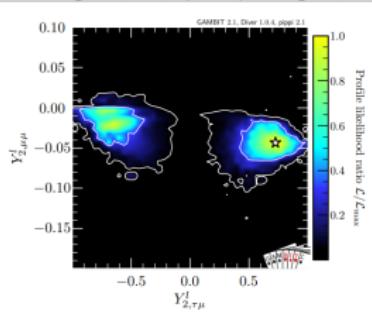
DMEFT

[Eur. Phys. J. C 81 (2021) 11, 992]



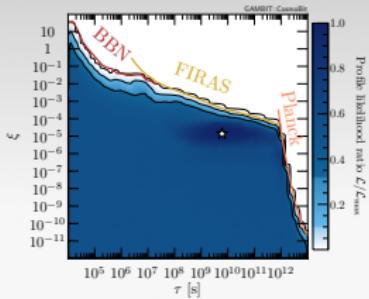
THDM-III

[JHEP 01 (2022) 037]



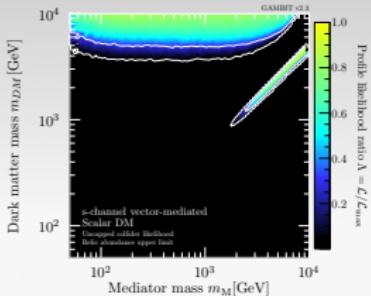
Cosmo ALPs

[arXiv:2205.13549 [astro-ph.CO]]



S-channel DM

[arXiv:2209.13266 [hep-ph]]



Core

- Each module contains a collection of module functions
- Module functions provide a *capability*
- They have dependencies and backend requirements
- Allowed for specific models
- At run time a dependency tree is generated and resolved

```
// SM-like Higgs mass with theoretical uncertainties
#define CAPABILITY prec_mh
START_CAPABILITY

#define FUNCTION FH_HiggsMass
START_FUNCTION(triplet<double>)
DEPENDENCY(unimproved_MSSM_spectrum, Spectrum)
DEPENDENCY(FH_HiggsMasses, fh_HiggsMassObs)
ALLOW_MODELS(MSSM63atQ, MSSM63atMGUT)
#undef FUNCTION

#define FUNCTION SHD_HiggsMass
START_FUNCTION(triplet<double>)
DEPENDENCY(unimproved_MSSM_spectrum, Spectrum)
BACKEND_REQ(SUSYHD_MHiggs, (), MReal, (const MList<MReal>&))
BACKEND_REQ(SUSYHD_DeltaMHiggs, (), MReal, (const MList<MReal>&))
ALLOW_MODELS(MSSM63atQ, MSSM63atMGUT)
#undef FUNCTION

#undef CAPABILITY
```



Models

- Extensive model database

SUSY

CMSSM
NUHM1,2
MSSM63atQ

DM

Scalar Singlet
Fermionic Singlet
Vector Singlet
Axions

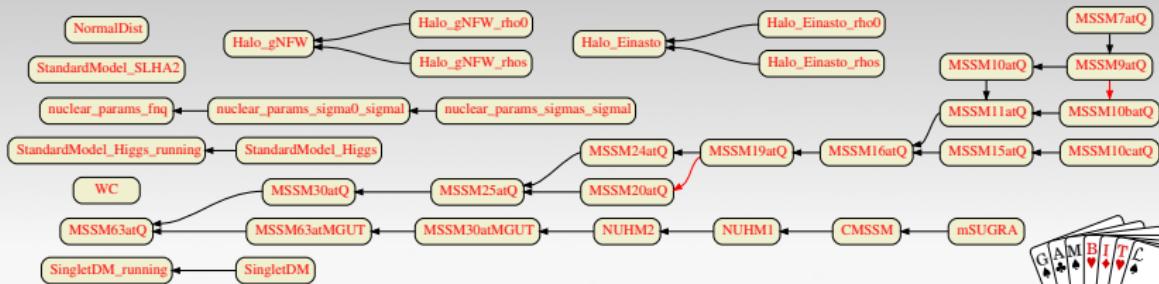
Cosmo

Λ CDM
 ΔN_{eff}
Power-law inflation

Others

SM
RH neutrinos
WC
nuisance models

- Parent-daughter hierarchy
- Module functions are activated for each model



Backends

- C, Fortran \rightsquigarrow POSIX dl
- C++ \rightsquigarrow BOSS + POSIX dl
- Mathematica \rightsquigarrow WSTP
- Python \rightsquigarrow pybind11

CosmoBit

AlterBBN 2.2
 DarkAges 1.2.0
 MontePythonLike 3.3.0
 MultiModeCode 2.0.0
 classy 2.9.4
 plc 3.0

DarkBit

CaptnGeneral 1.0
 DDCalc 2.2.0
 DarkSUSY 6.2.2
 MicrOmegas 3.6.9.2
 gamLike 1.0.1
 nulike 1.0.9

ColliderBit

HiggsBounds 4.3.1
 HiggsSignals 1.4
 Pythia 8.212
 nulike 1.0.9

PrecisionBit

FeynHiggs 2.12.0
 SUSYHD 1.0.2
 gm2calc 1.3.0

SpecBit

FlexibleSUSY 2.0.1
 SPheno 4.0.3

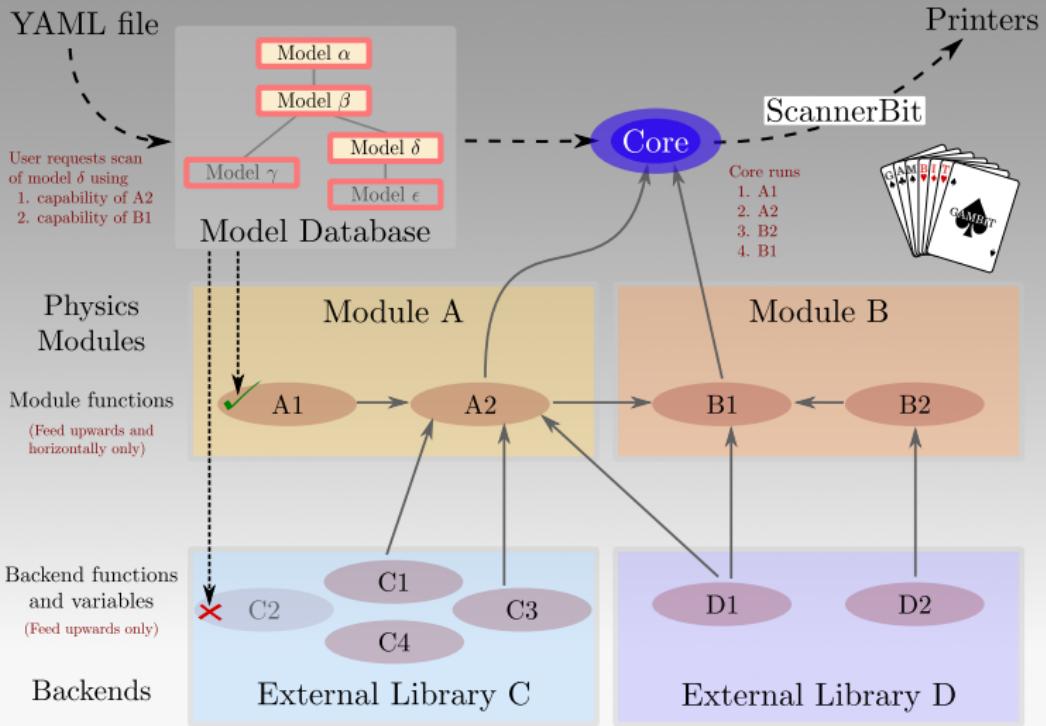
FlavBit

SuperISO 3.6

DecayBit

SUSY_HIT 1.5

An example run



Operators

	SI scattering	SD scattering	Annihilations
$\mathcal{Q}_{1,q}^{(6)} = (\bar{\chi}\gamma_\mu\chi)(\bar{q}\gamma^\mu q)$	unsuppressed	—	<i>s</i> -wave
$\mathcal{Q}_{2,q}^{(6)} = (\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{q}\gamma^\mu q)$	suppressed	—	<i>p</i> -wave
$\mathcal{Q}_{3,q}^{(6)} = (\bar{\chi}\gamma_\mu\chi)(\bar{q}\gamma^\mu\gamma_5 q)$	—	suppressed	<i>s</i> -wave
$\mathcal{Q}_{4,q}^{(6)} = (\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{q}\gamma^\mu\gamma_5 q)$	—	unsuppressed	<i>s</i> -wave $\propto m_q^2/m_\chi^2$
$\mathcal{Q}_1^{(7)} = \frac{\alpha_s}{12\pi}(\bar{\chi}\chi)G^{a\mu\nu}G_{\mu\nu}^a$	unsuppressed	—	<i>p</i> -wave
$\mathcal{Q}_2^{(7)} = \frac{\alpha_s}{12\pi}(\bar{\chi}i\gamma_5\chi)G^{a\mu\nu}G_{\mu\nu}^a$	suppressed	—	<i>s</i> -wave
$\mathcal{Q}_3^{(7)} = \frac{\alpha_s}{8\pi}(\bar{\chi}\chi)G^{a\mu\nu}\tilde{G}_{\mu\nu}^a$	—	suppressed	<i>p</i> -wave
$\mathcal{Q}_4^{(7)} = \frac{\alpha_s}{8\pi}(\bar{\chi}i\gamma_5\chi)G^{a\mu\nu}\tilde{G}_{\mu\nu}^a$	—	suppressed	<i>s</i> -wave
$\mathcal{Q}_{5,q}^{(7)} = m_q(\bar{\chi}\chi)(\bar{q}q)$	unsuppressed	—	<i>p</i> -wave $\propto m_q^2/m_\chi^2$
$\mathcal{Q}_{6,q}^{(7)} = m_q(\bar{\chi}i\gamma_5\chi)(\bar{q}q)$	suppressed	—	<i>s</i> -wave $\propto m_q^2/m_\chi^2$
$\mathcal{Q}_{7,q}^{(7)} = m_q(\bar{\chi}\chi)(\bar{q}i\gamma_5 q)$	—	suppressed	<i>p</i> -wave $\propto m_q^2/m_\chi^2$
$\mathcal{Q}_{8,q}^{(7)} = m_q(\bar{\chi}i\gamma_5\chi)(\bar{q}i\gamma_5 q)$	—	suppressed	<i>s</i> -wave $\propto m_q^2/m_\chi^2$
$\mathcal{Q}_{9,q}^{(7)} = m_q(\bar{\chi}\sigma^{\mu\nu}\chi)(\bar{q}\sigma_{\mu\nu}q)$	loop-induced	unsuppressed	<i>s</i> -wave $\propto m_q^2/m_\chi^2$
$\mathcal{Q}_{10,q}^{(7)} = m_q(\bar{\chi}i\sigma^{\mu\nu}\gamma_5\chi)(\bar{q}\sigma_{\mu\nu}q)$	loop-induced	suppressed	<i>s</i> -wave $\propto m_q^2/m_\chi^2$

Hadronic input parameters

Parameter	Value	Parameter	Value
$\sigma_{\pi N}$	50(15) MeV [1]	μ_p	2.793 - [2]
$Bc_5(m_d - m_u)$	-0.51(8) MeV [3]	μ_n	-1.913 [2]
g_A	1.2756(13) [2]	μ_s	-0.036(21) [4]
m_G	836(17) MeV [1]	g_T^u	0.784(30) [5]
σ_s	52.9(7.0) MeV [6]	g_T^d	-0.204(15) [5]
$\Delta u + \Delta d$	0.440(44) [7]	g_T^s	$-27(16) \cdot 10^{-3}$ [5]
Δs	-0.035(9) [7]	$B_{T,10}^{u/p}$	3.0(1.5) [8]
$B_0 m_u$	0.0058(5) GeV^2 [9]	$B_{T,10}^{d/p}$	0.24(12) [8]
$B_0 m_d$	0.0124(5) GeV^2 [9]	$B_{T,10}^{s/p}$	0.0(2) [8]
$B_0 m_s$	0.249(9) GeV^2 [9]	r_s^2	-0.115(35) GeV^{-2} [4]

[1] [F. Bishara et. al., JHEP 11 (2017) 059] [2] [PDG 2020] [3] [A. Crivellin et. al., Phys. Rev. D 89 (2014) 054021] [4] [D. Djukanovic et. al., Phys. Rev. Lett. 123 (2019) 212001, R. S. Sufian et. al., Phys. Rev. Lett. 118 (2017) 042001] [5] [R. Gupta, et. al., Phys. Rev. D 98 (2018) 091501] [6] [S. Aoki et. al., Eur. Phys. J. C 80 (2020) 113] [7] [J. Liang et. al., Phys. Rev. D 98 (2018) 074505] [8] [B. Pasquini et. al., Phys. Rev. D72 (2005) 094029] [9] [F. Bishara et. al., arXiv:1708.02678.]

Nuisance parameters

Nuisance parameter		Value ($\pm 3\sigma$ range)
Local DM density	ρ_0	0.2–0.8 GeV cm $^{-3}$
Most probable speed	v_{peak}	240 (24) km s $^{-1}$
Galactic escape speed	v_{esc}	528 (75) km s $^{-1}$
Running top mass ($\overline{\text{MS}}$ scheme)	$m_t(m_t)$	162.9 (6.0) GeV
Pion-nucleon sigma term	$\sigma_{\pi N}$	50 (45) MeV
Strange quark contrib. to nucleon spin	Δs	-0.035 (0.027)
Strange quark nuclear tensor charge	g_T^s	-0.027 (0.048)
Strange quark charge radius of the proton	r_s^2	-0.115 (0.105) GeV $^{-2}$

Collider Likelihoods

- ATLAS, Poisson loglike marginalised over nuisance $\xi =$ relative signal/bkg uncertainties

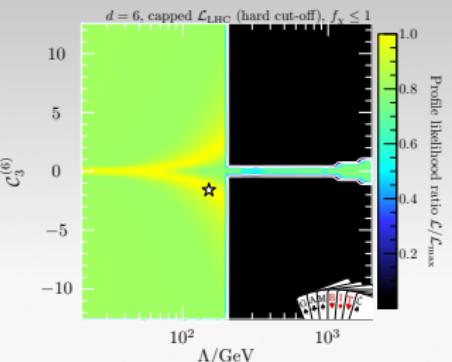
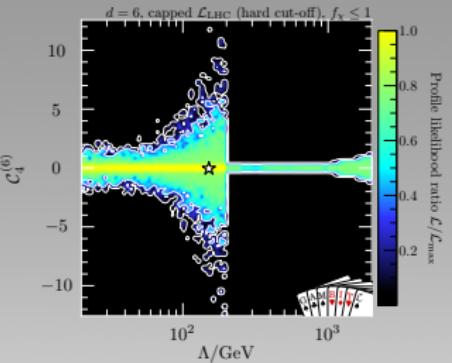
$$\begin{aligned} \mathcal{L}_{\text{marg}}(n|p) &= \int_0^\infty \frac{[\xi p]^n e^{-\xi p}}{n!} \\ &\quad \times \frac{1}{\sqrt{2\pi}\sigma_\xi} \frac{1}{\xi} \exp \left[-\frac{1}{2} \left(\frac{\ln \xi}{\sigma_\xi} \right)^2 \right] d\xi. \end{aligned}$$

- CMS, convolved Poisson-Gaussian, profiled over systematic uncertainties γ on expected background yields with covariance matrix Σ

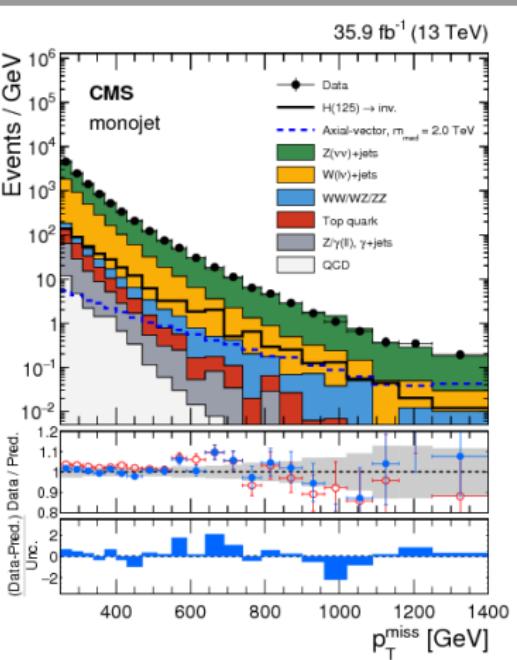
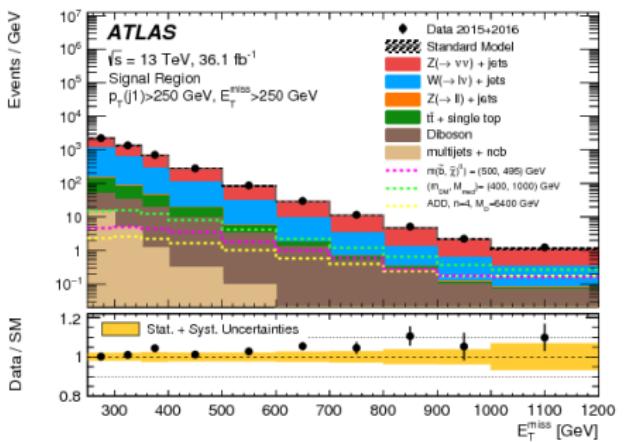
$$\begin{aligned} \mathcal{L}(\mathbf{s}, \gamma) &= \prod_i^{N_{\text{bin}}} \left[\frac{(s_i + b_i + \gamma_i)^{n_i} e^{-(s_i + b_i + \gamma_i)}}{n_i!} \right] \\ &\quad \times \frac{1}{\sqrt{\det 2\pi\Sigma}} e^{-\frac{1}{2} \gamma^T \Sigma^{-1} \gamma}. \end{aligned}$$

Results

- $\mathcal{C}_1^{(6)}$
 - spin-independent scattering
 - strongly constrained \rightsquigarrow very small
- $\mathcal{C}_2^{(6)}$
 - momentum-dependent scattering
 - $\Lambda < 250$ GeV DD constrained
 - $\Lambda > 250$ GeV LHC constrained
- $\mathcal{C}_3^{(6)}$
 - both SD and MD scattering
 - $\Lambda < 250$ GeV weak DD constraints
 - Main contribution to *Fermi – LAT*
 - $\Lambda > 250$ GeV LHC constrained
- $\mathcal{C}_4^{(6)}$
 - spin-dependent scattering
 - identical to $\mathcal{C}_2^{(6)}$



Results



But...

How do I use GAMBIT with my favourite model?

- ~~ Adding a model
- ~~ Sorting out hierarchy
- ~~ Making physics computations work with that model

How do I add a new physical observable or likelihood?

- ~~ Create capabilities
- ~~ Declare dependencies
- ~~ and models
- ~~ and backend requirements

1. Add the model to the **model hierarchy**:

- Choose a model name, and declare any **parent model**
- Declare the model's parameters
- Declare any **translation function** to the parent model

```
#define MODEL HUH1
#define PARENT HUH2
START_MODEL
DEFINEPARM(M0,M12,A0,TanBeta,SignMu)
INTERPRET_AS_PARENT_FUNCTION(HUH1_to_HUH2)
#undef PARENT
#undef MODEL
```

2. Write the translation function as a standard C++ function:

```
void MODEL_NAMESPACE::HUH1_to_HUH2 (const ModelParameters &myP, ModelParameters &targetP)
{
    // Set M0, M12, A0, TanBeta and SignMu in the HUH2 to the same values as in the HUH1
    targetP.setValues(myP,false);
    // Set the values of mH and mHd in the HUH2 to the value of mH in the HUH1
    targetP.setValue("mH", myP["mH"]);
    targetP.setValue("mHd", myP["mH"]);
}
```

3. If needed, declare that existing module functions work with the new model, or add new functions that do.

Adding a new module function is easy:

1. Declare the function to GAMBIT in a module's **rollcall header**

- Choose a capability
- Declare any **backend requirements**
- Declare any **dependencies**
- Declare any specific **allowed models**
- other more advanced declarations also available

```
#define MODULE Flavbit
START_MODULE
// A tasty GAMBIT module.

#define CAPABILITY Rmu
// Observable: RE(X->m_mu nu)/RE(pi->m_mu nu)

#define FUNCTION SI_Rmu
// Name of a function that can compute Rmu
// Function computes the precision weight
// Function computes the precision weight
BACKEND_NER(musin, pmiss, (my_tag), double) // Needs function from a backend
BACKEND_OPTIONC(Superlize, 3.0), (my_tag)) // Backend must be Superlize 3.0
DEPENENCY(Superlize, modelinfo, parameters) // Needs another function to calculate Superlize info
ALLOW_MULTI(MuSMSShutQ, MuSMSShutDUT) // Works with weak/GUT-scale MSSM and descendants
#undef FUNCTION
#undef CAPABILITY
```

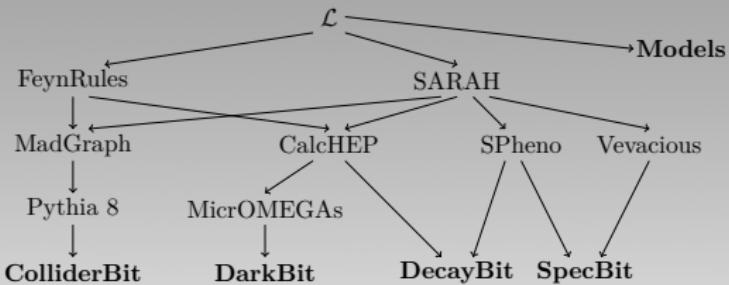
2. Write the function as a standard C++ function
(one argument: the result)

Solution

The GAMBIT Universal Model Machine



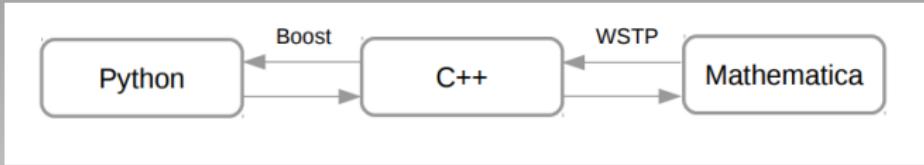
- GUM interfaces LLT SARAH and FeynRules with GAMBIT
- Uses existing HEP toolchains



- GAMBIT-compatible outputs from GUM

Generated output	FeynRules	SARAH	Usage in GAMBIT
CalcHEP	✓	✓	Decays, cross-sections
micrOMEGAs (via CalcHEP)	✓	✓	DM observables
Pythia (via MadGraph)	✓	✓	Collider physics
SPheno	✗	✓	Particle mass spectra, decay widths
Vevacious	✗	✓	Vacuum stability

- Primarily written in Python, with interface to Mathematica via Boost and WSTP



- Automatically generates GAMBIT code
 - Particles → particle database and parameters → Models
 - Module functions for ColliderBit, DarkBit, DecayBit and SpecBit
 - Writes interfaces to requested backends
- GUM will release with GAMBIT 2.0 **VERY SOON**

An example

- Majorana DM χ with scalar mediator Y

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{2}\bar{\chi}(i\not{\partial} - m_\chi)\chi + \frac{1}{2}\partial_\mu Y\partial^\mu Y - \frac{1}{2}m_Y^2 Y^2 - \frac{g_\chi}{2}\bar{\chi}\chi Y - \frac{c_Y}{2}\sum y_f f\bar{f}Y.$$

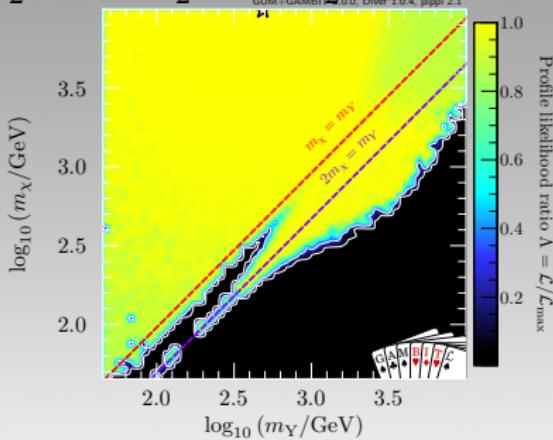
```

math:
# Choose FeynRules
package: feynrules
# Name of the model
model: MDMSM
# Model builds on the Standard Model FeynRules file
base_model: SM
# The Lagrangian is defined by the DM sector (LDM),
# defined in MDMSM.fr, plus the SM Lagrangian (LSM)
# imported from the 'base model', SM.fr
Lagrangian: LDM + LSM
# Make CKM matrix = identity to simplify output
restriction: DiagonalCKM

# PDG code of the annihilating DM candidate in
#<--> FeynRules file
wimp_candidate: 52

# Select outputs for DM physics.
# Collider physics is not as important in this model.
output:
pythia: false
calchep: true
micromegas: true

```



~~> Follow Sanjay's tutorial
3pm Room A