

# Neutrino mass hierarchy from the **Discrete Dark Matter** model

Omar Medina (in collaboration with...)

**IFIC-Universitat de València**

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**CSIC**

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# IN COLLABORATION WITH...



**Cesar Bonilla**  
Universidad Católica del  
Norte  
Antofagasta, Chile



**Johannes Herms**  
Max-Planck-Institut  
für Kernphysik  
Heidelberg, Germany



**Eduardo Peinado**  
IF-UNAM  
México City, México

# Outline



**Motivation**



**Aim of this Model**



**Elements & Properties**



**Phenomenology**



**Conclusions**

# Going BSM

Three drawbacks of the SM that we tackled with this model

**Neutrino Masses**  
(Majorana)

# Going BSM

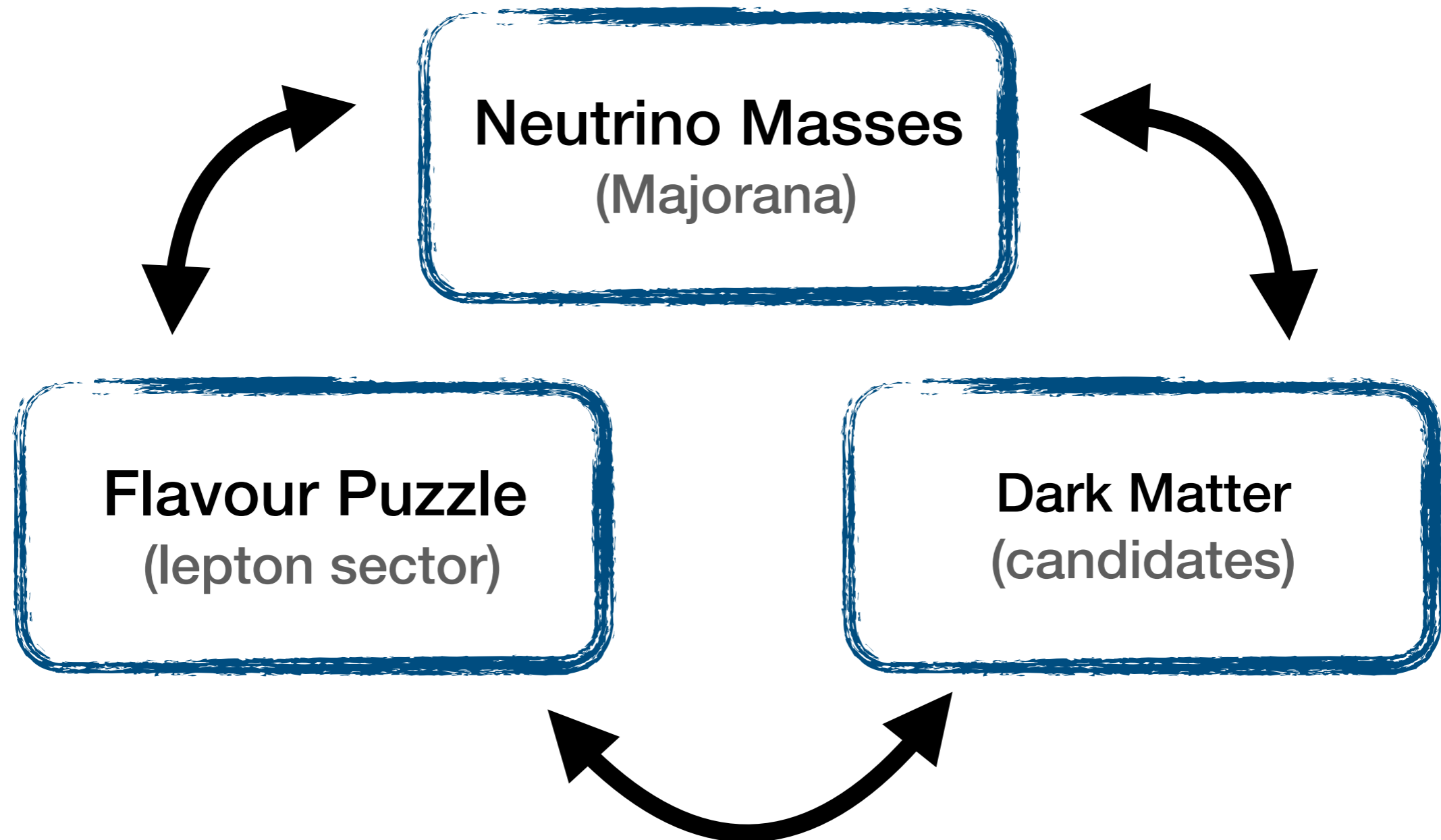
Three drawbacks of the SM that we tackled with this model

**Neutrino Masses**  
(Majorana)

**Flavour Puzzle**  
(lepton sector)

# Going BSM

Three drawbacks of the SM that we tackled with this model





# Neutrino mass mechanisms

## The seesaw mechanism

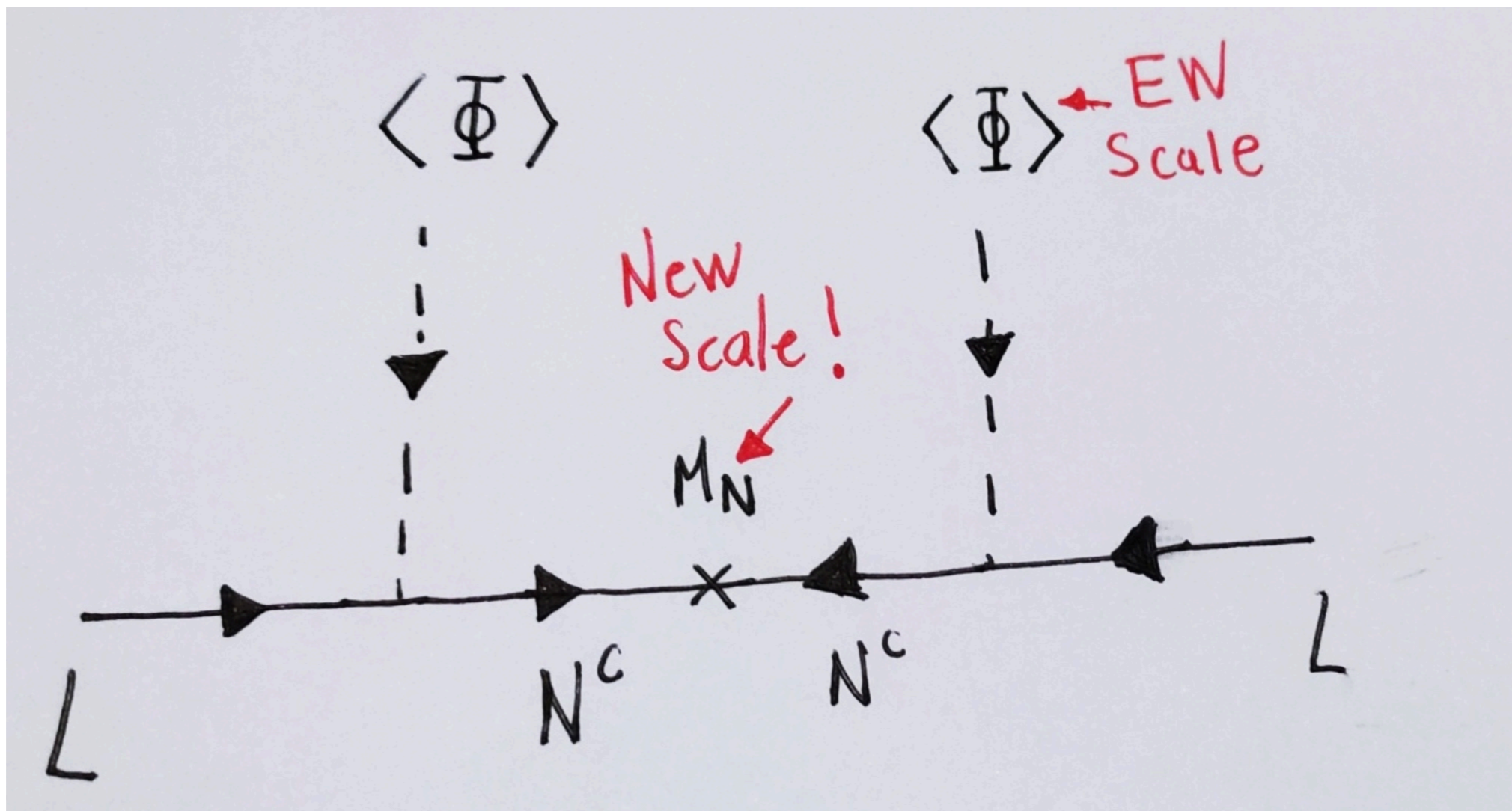
- Explains the lightness of neutrinos
- Introduces Heavy Neutral Leptons  $N^c$
- Introduces a **new physics scale!** (LNV)

[Minkowski,77]

[Yanagida,80]

[Mahopatra and Senjanović ,80]

[Schechter and Valle ,80]



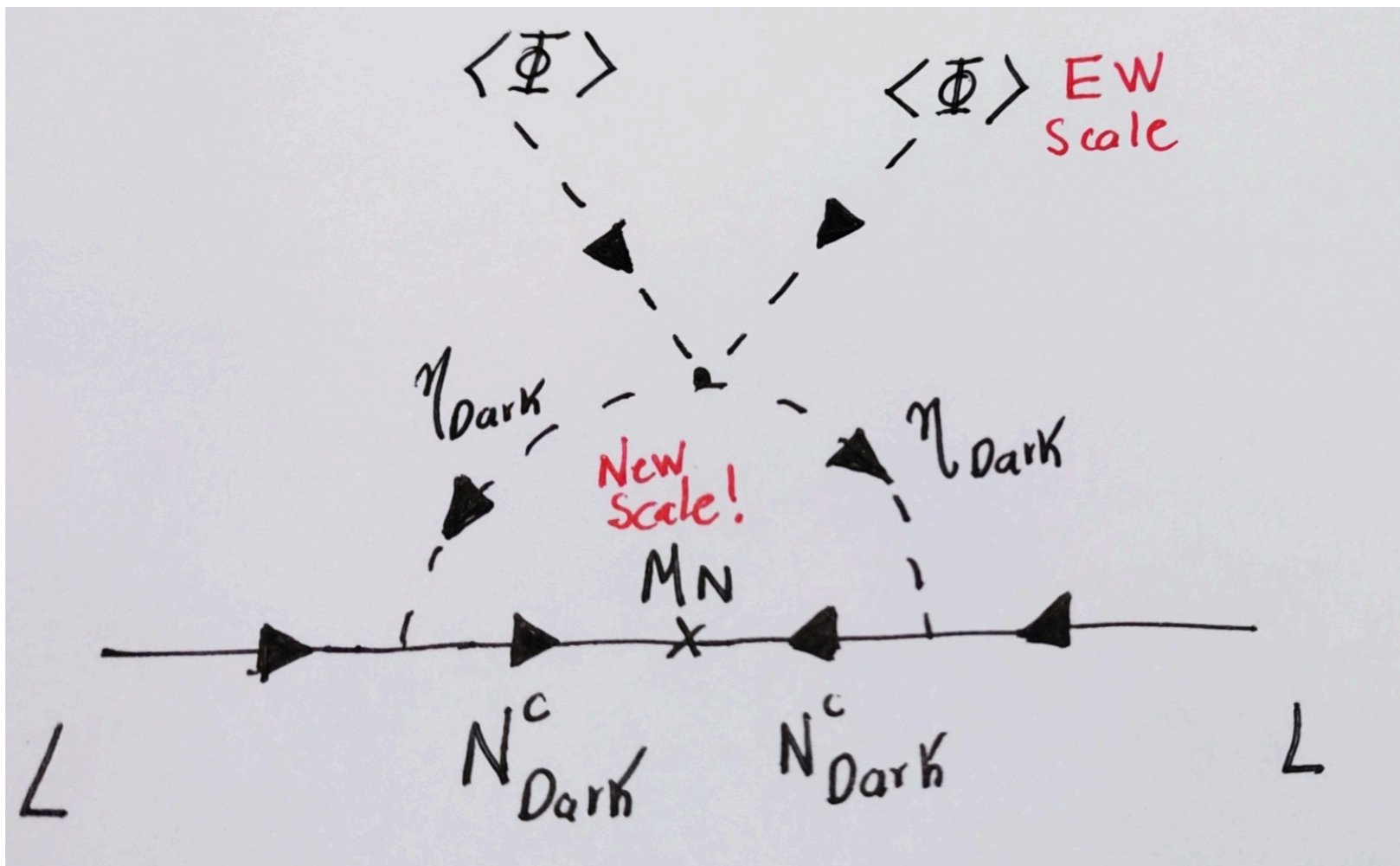
# Neutrino mass mechanisms

## The scotogenic mechanism

[Ma,2006]

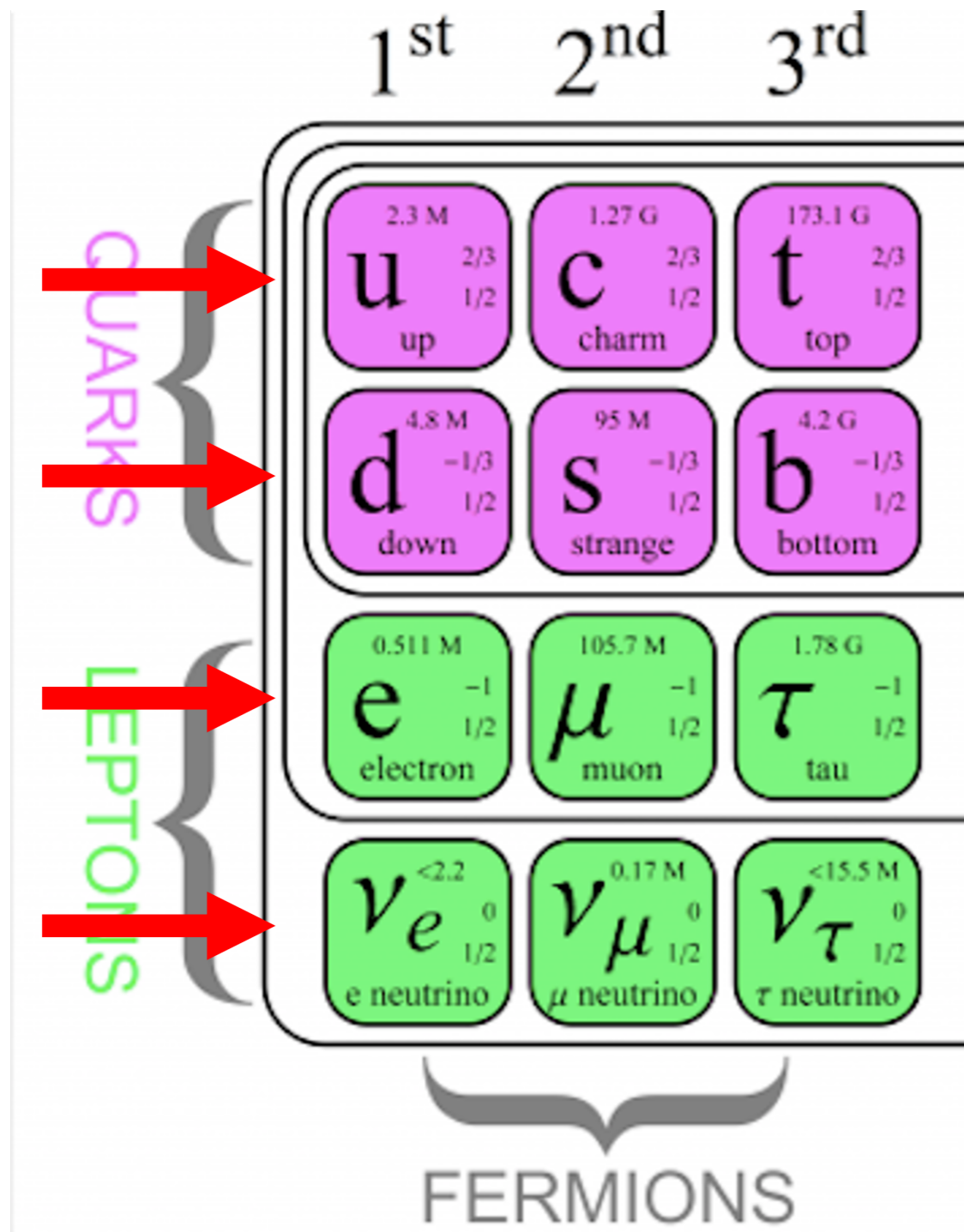
- Additional  $\mathbb{Z}_2$  symmetry and new iso-doublet
- **Dark Matter** generates neutrino mass.

$$N^c \sim -1, \quad \eta \sim -1 \quad \text{under } \mathbb{Z}_2$$





# The Flavour Puzzle



- The  $\mathcal{L}_{SM}$  is **built** to be invariant under

$$SU(3)_c \otimes \underbrace{SU(2)_L \otimes U(1)_Y}_{\text{Electroweak Sector}}$$

- The SM gauge group is **generation blind**, preserves **full flavour symmetry**

# The Flavour Puzzle

- **Yukawa interaction** is **not** based on the gauge principle, and in the SM **breaks** flavour symmetry

$$-\underline{Y_e^{ij}} \overline{L_{iL}^I} \Phi e_{jR}^I, \quad Y_e = \begin{pmatrix} Y_e^{ee} & Y_e^{e\mu} & Y_e^{e\tau} \\ Y_e^{\mu e} & Y_e^{\mu\mu} & Y_e^{\mu\tau} \\ Y_e^{\tau e} & Y_e^{\tau\mu} & Y_e^{\tau\tau} \end{pmatrix}$$

- The CKM matrix and the PMNS matrix **translate flavour symmetry breaking** to the gauge sector

$U_{CKM}$ ,

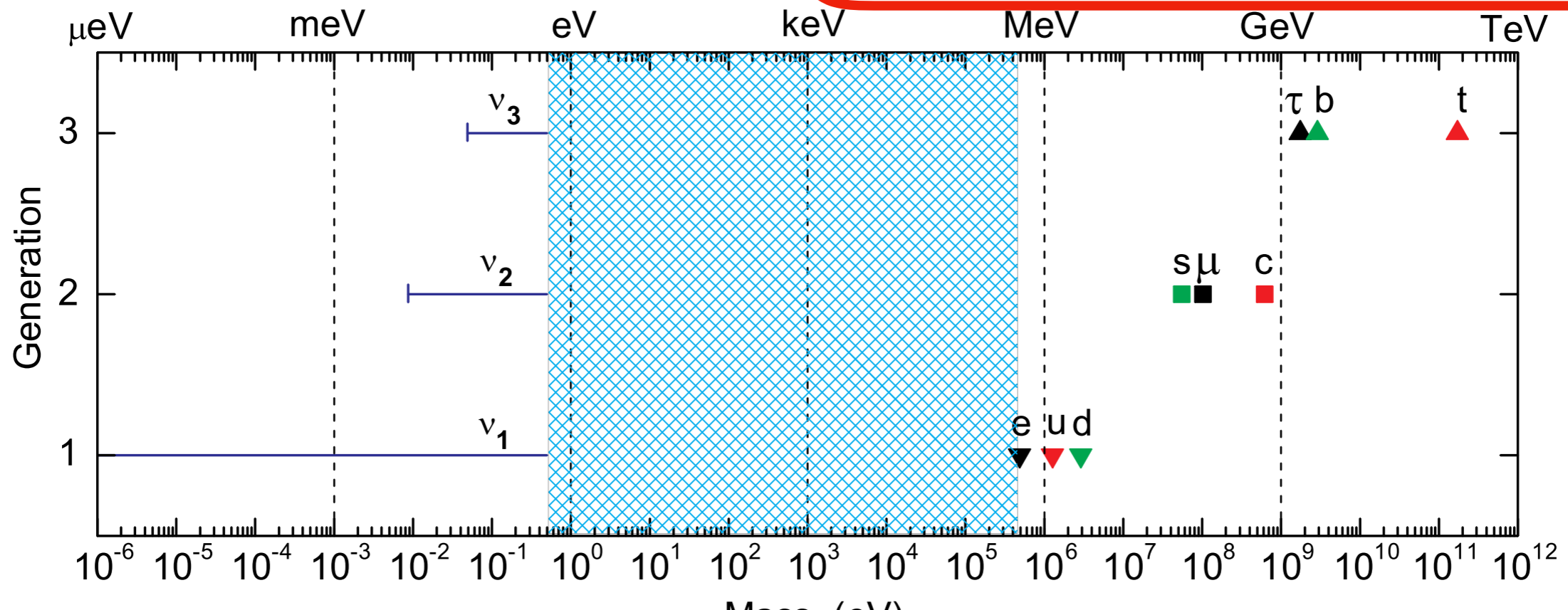
$V_{PMNS}$

# The Flavour Puzzle

- 22 out of the 27 parameters of the SM are in the Yukawa sector. Not constrained by **symmetry**

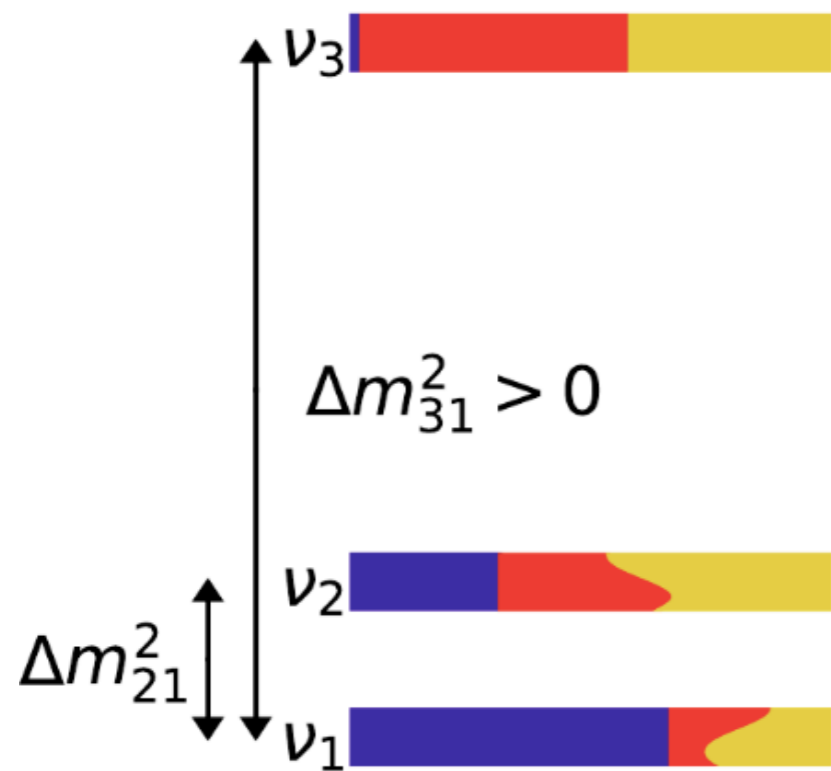
$$\begin{pmatrix} |V_{ud}| & |V_{us}| & |V_{ub}| \\ |V_{cd}| & |V_{cs}| & |V_{cb}| \\ |V_{td}| & |V_{ts}| & |V_{tb}| \end{pmatrix} \sim \begin{pmatrix} \text{large} & \text{small} & \text{tiny} \\ \text{small} & \text{large} & \text{tiny} \\ \text{tiny} & \text{tiny} & \text{large} \end{pmatrix},$$

$$\begin{pmatrix} |U_{e1}| & |U_{e2}| & |U_{e3}| \\ |U_{\mu 1}| & |U_{\mu 2}| & |U_{\mu 3}| \\ |U_{\tau 1}| & |U_{\tau 2}| & |U_{\tau 3}| \end{pmatrix} \sim \begin{pmatrix} \text{large} & \text{medium} & \text{small} \\ \text{small} & \text{medium} & \text{large} \\ \text{small} & \text{medium} & \text{large} \end{pmatrix}.$$



# The Flavour Puzzle

- In the lepton sector: Neutrino Oscillation Parameters



NO

Neutrino Global Fit (Valencia)

$$\begin{pmatrix} |U_{e1}| & |U_{e2}| & |U_{e3}| \\ |U_{\mu 1}| & |U_{\mu 2}| & |U_{\mu 3}| \\ |U_{\tau 1}| & |U_{\tau 2}| & |U_{\tau 3}| \end{pmatrix} \sim \begin{pmatrix} \text{large} & \text{small} & \text{very small} \\ \text{small} & \text{medium} & \text{large} \\ \text{small} & \text{medium} & \text{large} \end{pmatrix}$$

$$\sin^2 \theta_{12} \quad \sin^2 \theta_{13} \quad \sin^2 \theta_{23}$$

$\delta^{CP}$

- Two orders of magnitude hierarchy  $\Delta m_{31}^2 \gg \Delta m_{21}^2$

# Flavour Symmetry

- **Flavour symmetry** at high-energy regime.

$$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y \underbrace{\otimes G}_{\text{Flavour}}.$$

- Constraining, or relating the Yukawa coupling structure

$$\underbrace{G}_{\text{Symmetry}} \xrightarrow{\text{SSB}} \underbrace{V_{\text{CKM}}, U, \text{Mass Hierarchy}}_{\text{Flavour Observables}}.$$

- An appealing option are Discrete and Non-Abelian Groups



# Flavour Symmetry

- The  $A_4$  group

$$A_4 \simeq \left\{ S, T \mid S^2 = T^3 = (ST)^2 = \mathbf{1} \right\},$$

Four Irreps.



$$\mathbf{1}, \mathbf{1}', \mathbf{1}'', \mathbf{3}.$$

$\mathbf{1}:$	$S = 1,$	$T = 1,$
$\mathbf{1}':$	$S = 1,$	$T = \omega,$
$\mathbf{1}'':$	$S = 1,$	$T = \omega^2,$

$$\omega \equiv e^{\frac{2\pi i}{3}}.$$

$$\mathbf{3}: \quad S = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}, \quad T = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix},$$

# Flavour Symmetry

- The  $A_4$  group

$$A_4 \simeq \left\{ S, T \mid S^2 = T^3 = (ST)^2 = \mathbf{1} \right\},$$



## Four Irreps.

$\mathbf{1}, \mathbf{1}', \mathbf{1}'', \mathbf{3}.$

$\mathbf{1}:$	$S = 1,$	$T = 1,$	<b>Generates a <math>\mathbb{Z}_2</math> symmetry</b>
$\mathbf{1}':$	$S = 1,$	$T = \omega,$	
$\mathbf{1}'':$	$S = 1,$	$T = \omega^2,$	

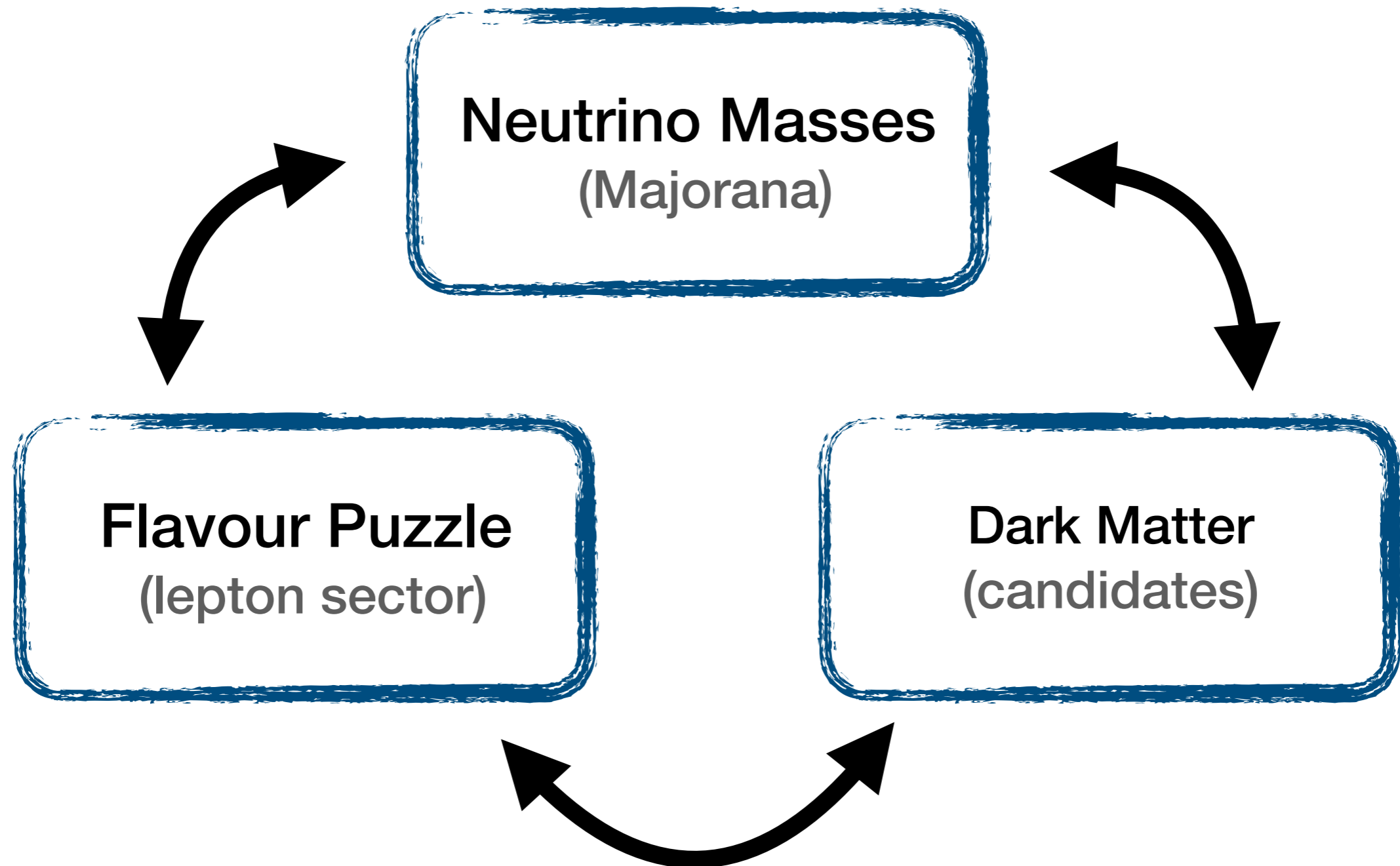
$$\omega \equiv e^{\frac{2\pi i}{3}}.$$

$\mathbf{3}:$

$$S = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}, \quad T = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix},$$

See Ivo's talk

# The Discrete Dark Matter Model

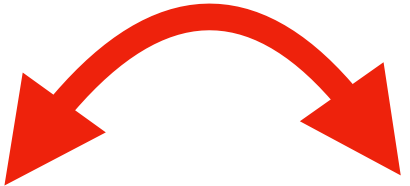


# The Discrete Dark Matter Model

- Fields and Symmetries (Lepton sector only)

[Hirsch,2010]

[Boucenna,2011]



	$L_e$	$L_\mu$	$L_\tau$	$l_e$	$l_\mu$	$l_\tau$	$N_T$	$H$	$\eta$
$SU(2)$	2	2	2	1	1	1	1	2	2
$U(1)_Y$	$-\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$	-1	-1	-1	0	$\frac{1}{2}$	$\frac{1}{2}$
$A_4$	$1'$	1	$1''$	$1'$	1	$1''$	3	1	3

RH-Neutrinos

Scalar iso-doublets

$$N_T = (N_1, N_2, N_3)$$

$$\eta = (\eta_1, \eta_2, \eta_3)$$

Mass degenerate  $m_N$

# The Discrete Dark Matter Model

- Yukawa Lagrangian invariant under  $A_4$

$$\mathcal{L}_{\text{Yukawa}}^H = y_e \bar{L}_e l_e H + y_\mu \bar{L}_\mu l_\mu H + y_\tau \bar{L}_\tau l_\tau H + H.c.$$

$$\mathcal{L}_{\text{Yukawa}}^\eta = y_1^\nu \bar{L}_e [N_T \eta]_{\mathbf{1}} + y_2^\nu \bar{L}_\mu [N_T \eta]_{\mathbf{1}''} + y_3^\nu \bar{L}_\tau [N_T \eta]_{\mathbf{1}'} + M_N [\bar{N}_T^c N_T]_{\mathbf{1}} + H.c.$$



# The Discrete Dark Matter Model

- Yukawa Lagrangian invariant under  $A_4$

$$\mathcal{L}_{\text{Yukawa}}^H = y_e \bar{L}_e l_e H + y_\mu \bar{L}_\mu l_\mu H + y_\tau \bar{L}_\tau l_\tau H + H.c.$$

$$\mathcal{L}_{\text{Yukawa}}^\eta = y_1^\nu \bar{L}_e [N_T \eta]_{\mathbf{1}} + y_2^\nu \bar{L}_\mu [N_T \eta]_{\mathbf{1}''} + y_3^\nu \bar{L}_\tau [N_T \eta]_{\mathbf{1}'} + M_N [\bar{N}_T^c N_T]_{\mathbf{1}} + H.c.$$

- Electroweak and Flavour symmetry breakdown

$$\langle H^0 \rangle = v_H \neq 0, \quad \langle \eta_1^0 \rangle = v_{\eta_1} \neq 0, \quad \langle \eta_{2,3}^0 \rangle = 0,$$

**SM Higgs couple to charged leptons**

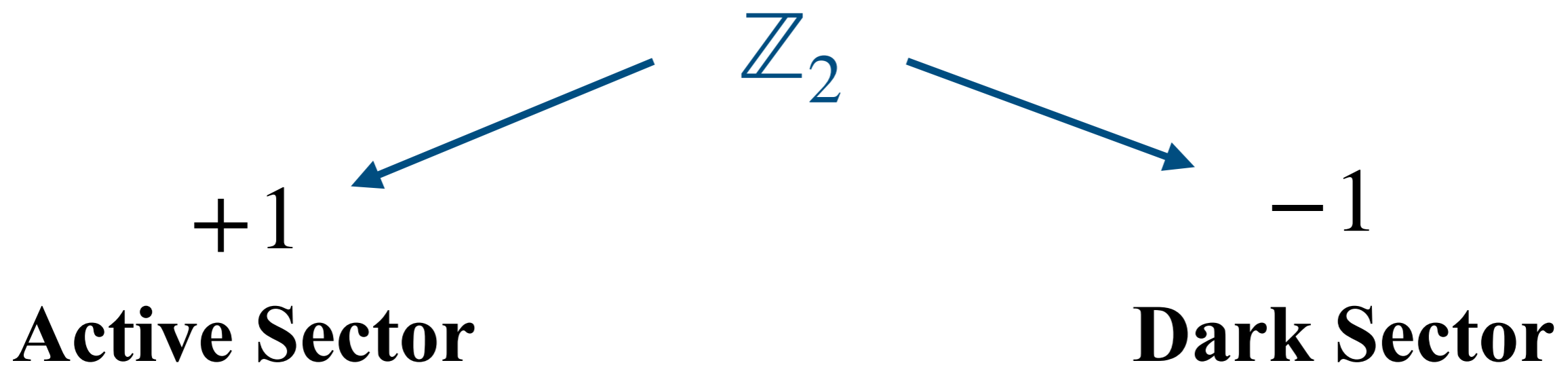
$$v_H Y_l^H = v_H \begin{pmatrix} y_e & 0 & 0 \\ 0 & y_\mu & 0 \\ 0 & 0 & y_\tau \end{pmatrix}$$

$$\langle \eta^0 \rangle = \begin{pmatrix} v_{\eta_1} \\ 0 \\ 0 \end{pmatrix}$$

# The Discrete Dark Matter Model

- Remnant  $\mathbb{Z}_2$  symmetry from  $A_4$

$$\langle \eta^0 \rangle = \begin{pmatrix} v_{\eta_1} \\ 0 \\ 0 \end{pmatrix} \longrightarrow S = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$



# The Discrete Dark Matter Model

- The **Active Sector**

$$N_1, \quad H, \quad \eta_1$$

$$\hat{H} = \begin{pmatrix} H_0'^+ \\ (v_H + H_0' + iA_0')/\sqrt{2} \end{pmatrix}, \quad \eta_1 = \begin{pmatrix} H_1'^+ \\ (v_{\eta_1} + H_1' + iA_1')/\sqrt{2} \end{pmatrix},$$

- **Seesaw mechanism** (rank 1 matrix)

$$m_D = \begin{pmatrix} y_1^\nu v_{\eta_1} & 0 & 0 \\ y_2^\nu v_{\eta_1} & 0 & 0 \\ y_3^\nu v_{\eta_1} & 0 & 0 \end{pmatrix} \quad m_\nu^{\text{Tree}} = -\frac{v_{\eta_1}^2}{M} \begin{pmatrix} y_1^\nu y_1^\nu & y_1^\nu y_2^\nu & y_1^\nu y_3^\nu \\ y_1^\nu y_2^\nu & y_2^\nu y_2^\nu & y_2^\nu y_3^\nu \\ y_1^\nu y_3^\nu & y_2^\nu y_3^\nu & y_3^\nu y_3^\nu \end{pmatrix}$$

# The Discrete Dark Matter Model

- The **Dark Sector**

$$N_2, N_3, \eta_2, \eta_3$$

$$\eta_2 = \begin{pmatrix} H_2'^+ \\ (H_2' + iA_2')/\sqrt{2} \end{pmatrix}, \quad \eta_3 = \begin{pmatrix} H_3'^+ \\ (H_3' + iA_3')/\sqrt{2} \end{pmatrix}$$

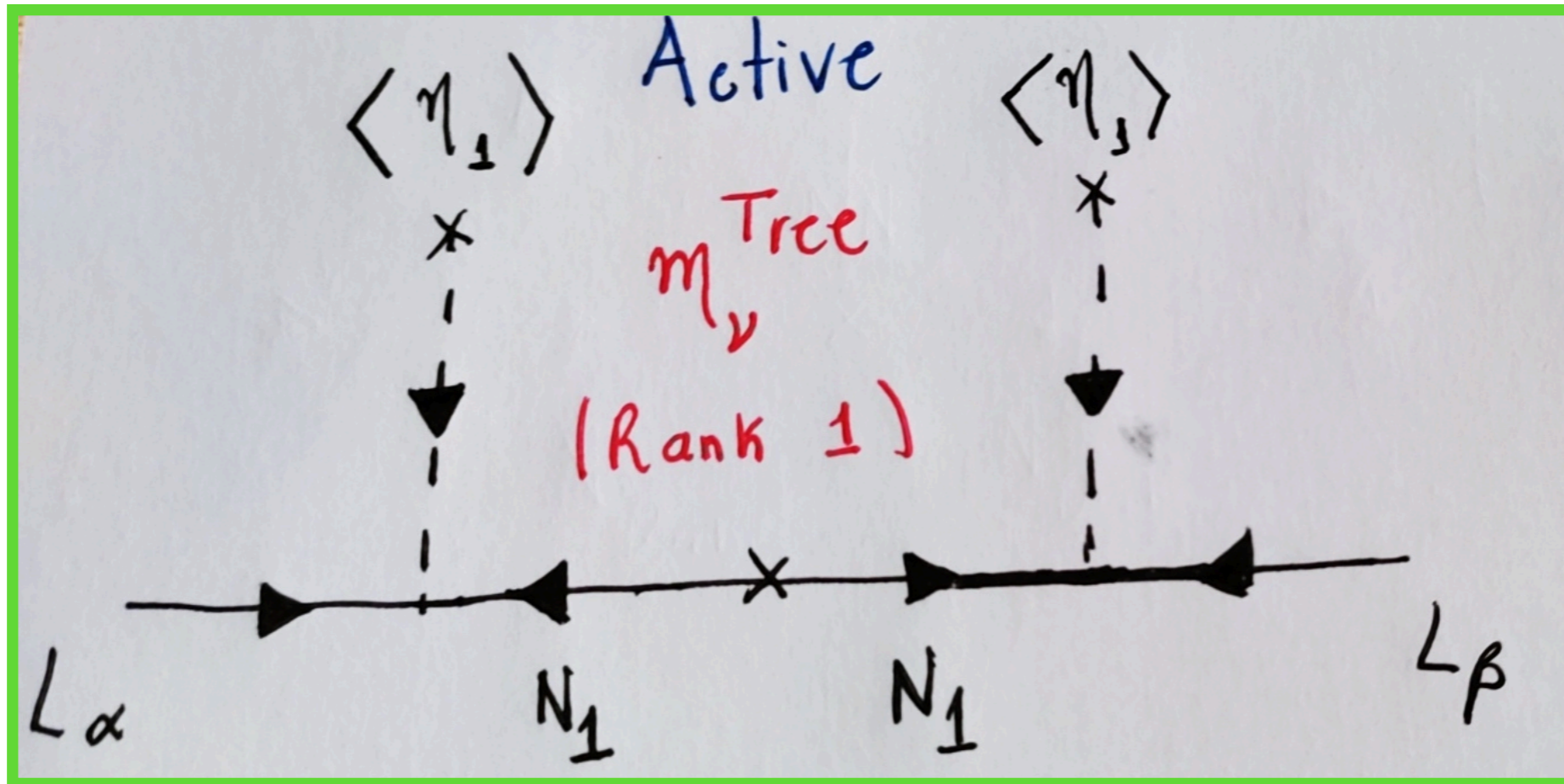
- **Scotogenic mechanism** (rank 2 matrix)

$$Y^{\eta_2} = \begin{pmatrix} 0 & y_1^\nu \omega^2 & 0 \\ 0 & y_2^\nu & 0 \\ 0 & y_3^\nu \omega & 0 \end{pmatrix}, \quad Y^{\eta_3} = \begin{pmatrix} 0 & 0 & y_1^\nu \omega \\ 0 & 0 & y_2^\nu \\ 0 & 0 & y_3^\nu \omega^2 \end{pmatrix}$$

# The Discrete Dark Matter Model

- Active fields (Tree level)

$$N_1, H, \eta_1$$



- Main contribution to:

$$\Delta m_{31}^2$$

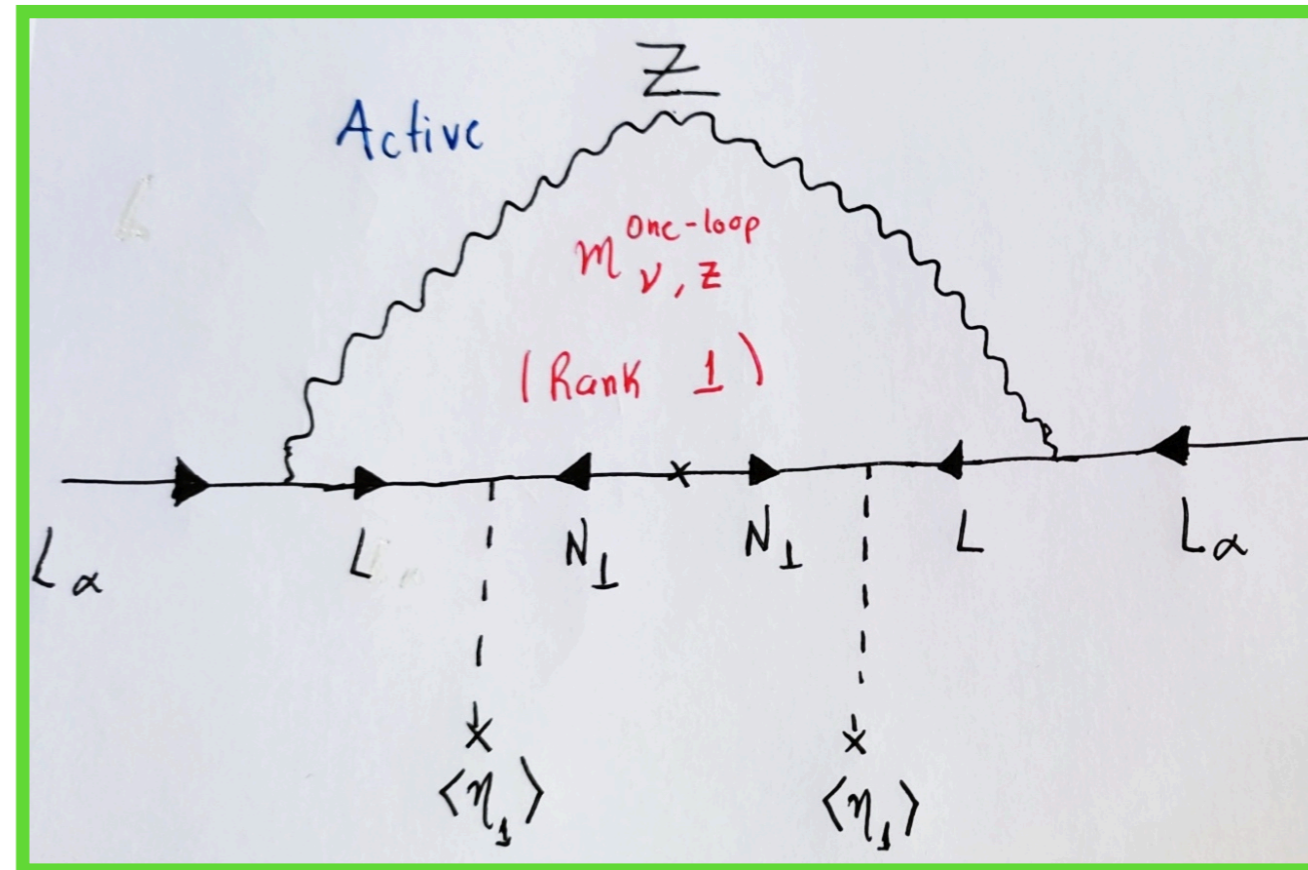
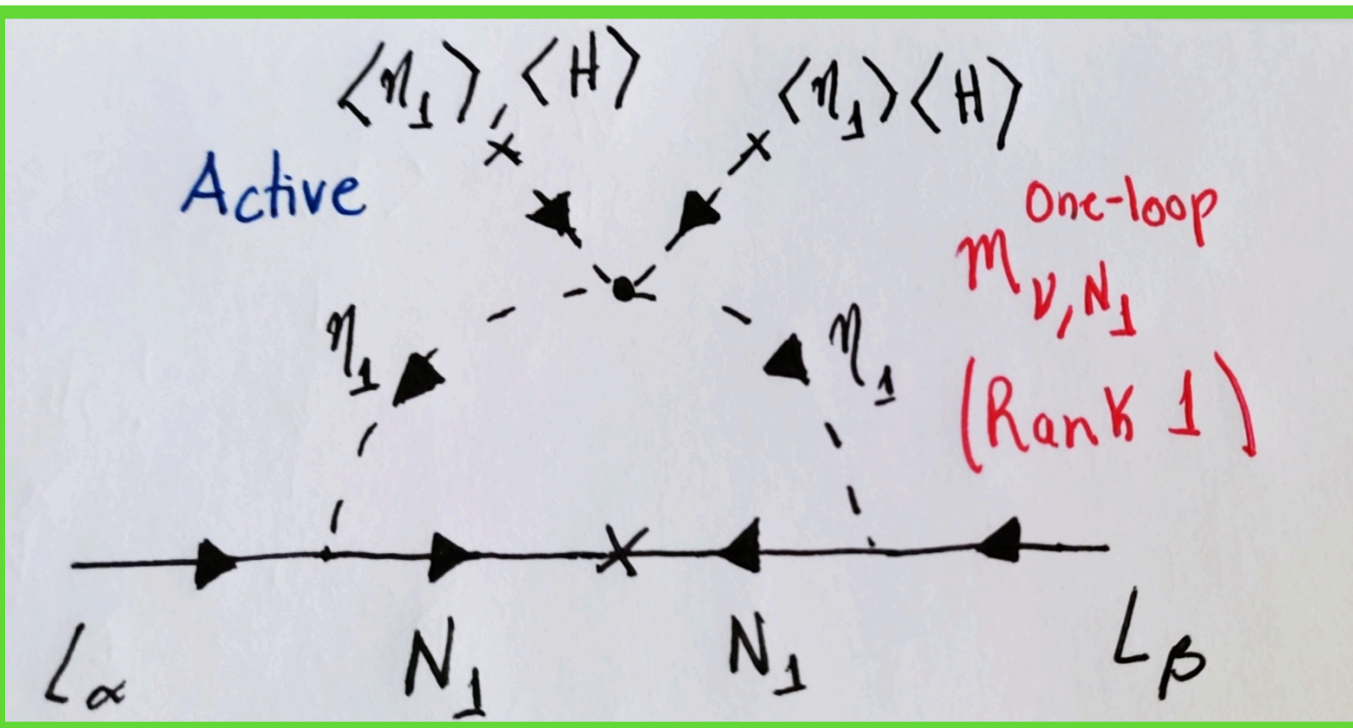
(Normal Ordering)



# The Discrete Dark Matter Model

- Active fields (one-loop)

$$N_1, H, \eta_1$$



- Main contribution to:

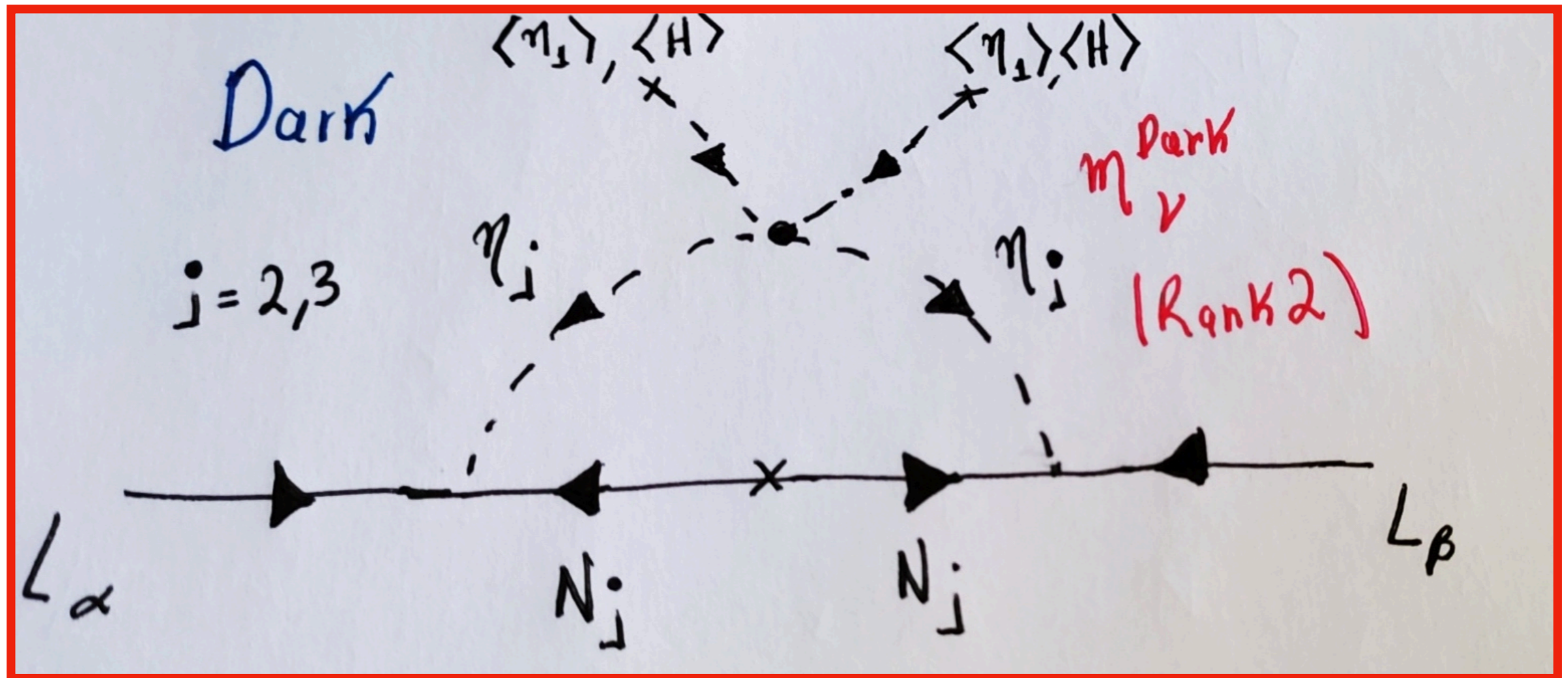
$$\Delta m_{31}^2$$

(Normal Ordering)

# The Discrete Dark Matter Model

- **Dark fields (one-loop)**

$$N_2, N_3, \eta_2, \eta_3$$



- Main contribution to:

$$\Delta m_{21}^2$$

(Normal Ordering)



# The Discrete Dark Matter Model

- **Scalar Sector  $A_4$  invariant** (plays crucial role)

$$\begin{aligned}
 V(H, \eta) = & \mu_H H^\dagger H + \mu_\eta (\eta^\dagger \eta)_{\mathbf{1}} + \lambda_1 (H^\dagger H)^2 + \lambda_2 (\eta^\dagger \eta)_{\mathbf{1}} (\eta^\dagger \eta)_{\mathbf{1}} + \lambda_3 (\eta^\dagger \eta)_{\mathbf{1}'} (\eta^\dagger \eta)_{\mathbf{1}''} \\
 & + \lambda_4 (\eta^\dagger \eta)_{(\mathbf{3}_1)} (\eta^\dagger \eta)_{(\mathbf{3}_2)_1} + \left[ \lambda_5 e^{i\varphi_5} (\eta^\dagger \eta)_{(\mathbf{3}_1)} (\eta^\dagger \eta)_{(\mathbf{3}_1)_1} + \text{H.c.} \right] \\
 & + \lambda_6 (H^\dagger H) (\eta^\dagger \eta)_{\mathbf{1}} + \lambda_7 (H^\dagger \eta)_{(\mathbf{3})} (\eta^\dagger H)_{(\mathbf{3})_1} + \left[ \lambda_8 e^{i\varphi_8} (H^\dagger \eta)_{(\mathbf{3})} (H^\dagger \eta)_{(\mathbf{3})_1} + \text{H.c.} \right] \\
 & + \left[ \lambda_9 e^{i\varphi_9} (\eta^\dagger \eta)_{(\mathbf{3}_1)} (H^\dagger \eta)_{(\mathbf{3})_1} + \text{H.c.} \right] + \left[ \lambda_{10} e^{i\varphi_{10}} (\eta^\dagger \eta)_{(\mathbf{3}_2)} (H^\dagger \eta)_{(\mathbf{3})_1} + \text{H.c.} \right]
 \end{aligned}$$

- **CP-violation** for the dark sector is **necessary** to fit **lepton mixing**

$$M_{\text{neutral}}^2 = \begin{pmatrix} M_{H'_0 H'_1}^2 & 0 & 0 & 0 \\ 0 & M_{A'_0 A'_1}^2 & 0 & 0 \\ 0 & 0 & M_{H'_2 H'_3}^2 & M_{\text{CPV}}^2 \\ 0 & 0 & M_{\text{CPV}}^2 & M_{A'_2 A'_3}^2 \end{pmatrix} \begin{matrix} \sin^2 \theta_{12}^l \\ \sin^2 \theta_{13}^l \\ \sin^2 \theta_{23}^l \\ \Delta m_{21}^2 \quad \Delta m_{31}^2 \end{matrix}$$

# The Discrete Dark Matter Model

## The neutrino mass matrix

$$(m_\nu)_{\alpha\beta} = (m_\nu^{\text{Active}})_{\alpha\beta} + (m_\nu^{\text{Dark}})_{\alpha\beta},$$

### Active contribution

$$m_\nu^{\text{Active}} = m_\nu^{\text{Tree}} + m_{\nu, N_1}^{\text{One-loop}} + m_{\nu, Z}^{\text{One-loop}}.$$

$$m_\nu^{\text{Tree}} = -\frac{v_{\eta_1}^2}{M} \begin{pmatrix} y_1^\nu y_1^\nu & y_1^\nu y_2^\nu & y_1^\nu y_3^\nu \\ y_1^\nu y_2^\nu & y_2^\nu y_2^\nu & y_2^\nu y_3^\nu \\ y_1^\nu y_3^\nu & y_2^\nu y_3^\nu & y_3^\nu y_3^\nu \end{pmatrix}$$

$$(m_{\nu, Z}^{\text{One-loop}})_{\alpha\beta} = \frac{3}{16\pi^2} \frac{m_Z^2}{v^2} \log\left(\frac{m_Z^2}{m_N^2}\right) (m_\nu^{\text{Tree}})_{\alpha\beta},$$

$$(m_{\nu, N_1}^{\text{One-loop}})_{\alpha\beta} = -\frac{1}{32\pi^2} \sum m_N Y_{\alpha N_1}^a Y_{N_1 \beta}^a B_0(0, m_a^2, m_N), \quad a = h, H_0, G, A_0,$$

[Escribano, 2020]

Passarino-Veltman reduction

Mass-eigenstates

# The Discrete Dark Matter Model

## The neutrino mass matrix

$$(m_\nu)_{\alpha\beta} = (m_\nu^{\text{Active}})_{\alpha\beta} + (m_\nu^{\text{Dark}})_{\alpha\beta},$$

## Dark contribution

$$(m_\nu^{\text{Dark}})_{\alpha\beta} = -\frac{1}{32\pi^2} \sum_{n,a} m_N Y_{\alpha n}^a Y_{n\beta}^a B_0(0, m_a^2, m_N), \quad a = \chi_1^D, \dots, \chi_4^D,$$

DM candidate

Scalar mixing plays

Mass-eigenstates

$\chi_4^D$

# Numerical Result

With the **CP-symmetry conservation** in the scalar potential

Parameter	Value
$y_e$	$-2.420 \times 10^{-6}$
$y_\mu$	$5.108 \times 10^{-4}$
$y_\tau$	$8.684 \times 10^{-3}$
$y_1^\nu$	$2.127 \times 10^{-6}$
$y_2^\nu$	$-1.046 \times 10^{-5}$
$y_3^\nu$	$-1.172 \times 10^{-5}$
$v_{\eta_1}/\text{GeV}$	142.243
$v_H/\text{GeV}$	201.079
$M_N/\text{GeV}$	$9.997 \times 10^4$
$\lambda_1$	2.0
$\lambda_2$	2.0
$\lambda_3$	-0.42
$\lambda_4$	0.578
$\lambda_5$	-0.486
$\lambda_6$	2.0
$\lambda_7$	-1.984
$\lambda_8$	-0.369
$\lambda_9$	1.198
$\lambda_{10}$	-1.191
$\varphi_5$	0
$\varphi_9$	0
$\varphi_{10}$	0

Observable	Data		Model best fit
	Central value	$1\sigma$ range	
$\sin^2 \theta_{12}/10^{-1}$	3.18	3.02 $\rightarrow$ 3.34	1.46
$\sin^2 \theta_{13}/10^{-2}$ (NO)	2.200	2.138 $\rightarrow$ 2.269	2.314
$\sin^2 \theta_{23}/10^{-1}$ (NO)	5.74	5.60 $\rightarrow$ 5.88	6.2
$\delta^\ell / \pi$ (NO)	1.08	0.96 $\rightarrow$ 1.21	1.0
$\Delta m_{21}^2/(10^{-5} \text{ eV}^2)$	7.50	7.30 $\rightarrow$ 7.72	7.49
$\Delta m_{31}^2/(10^{-3} \text{ eV}^2)$ (NO)	2.55	2.52 $\rightarrow$ 2.57	2.55
$m_{\text{lightest}}^\nu / \text{meV}$ (NO)			2.03
$m_2^\nu / \text{meV}$			8.89
$m_3^\nu / \text{meV}$			50.54
$\phi_{12}/\pi$			1.5
$\phi_{13}/\pi$			1.5
$\phi_{23}/\pi$			1.0
$m_e / \text{MeV}$	0.486	0.486 $\rightarrow$ 0.486	0.486
$m_\mu / \text{GeV}$	0.102	0.102 $\rightarrow$ 0.102	0.102
$m_\tau / \text{GeV}$	1.746	1.743 $\rightarrow$ 1.747	1.746
$M_H/\text{GeV}$ (Higgs boson)	125.25	125.08 $\rightarrow$ 125.42	125.25
$M_{DM}/\text{GeV}$ (lightest dark scalar)			87.0
$M_N/\text{GeV}$			$9.997 \times 10^4$
$\chi^2$			128.99

$\sin^2 \theta_{12}$   
Excluded

All Other observables in global fit 3-sigma range



# Numerical Result

With the **CP-symmetry breaking** in the scalar potential, all oscillation parameter and scalar sector observables are fitted without tensions.

Parameter	Value	Observable	Data		Model best fit
			Central value	1 $\sigma$ range	
$y_e$	$-2.136 \times 10^{-6}$	$\sin^2 \theta_{12}/10^{-1}$	3.18	3.02 $\rightarrow$ 3.34	3.14
$y_\mu$	$4.509 \times 10^{-4}$	$\sin^2 \theta_{13}/10^{-2}$ (NO)	2.200	2.138 $\rightarrow$ 2.269	2.201
$y_\tau$	$7.665 \times 10^{-3}$	$\sin^2 \theta_{23}/10^{-1}$ (NO)	5.74	5.60 $\rightarrow$ 5.88	5.75
$y_1^\nu$	$5.193 \times 10^{-6}$	$\delta^\ell / \pi$ (NO)	1.08	0.96 $\rightarrow$ 1.21	1.0
$y_2^\nu$	$-3.404 \times 10^{-5}$	$\Delta m_{21}^2 / (10^{-5} \text{eV}^2)$	7.50	7.30 $\rightarrow$ 7.72	7.48
$y_3^\nu$	$-7.540 \times 10^{-5}$	$\Delta m_{31}^2 / (10^{-3} \text{eV}^2)$ (NO)	2.55	2.52 $\rightarrow$ 2.57	2.55
$v_{\eta_1} / \text{GeV}$	93.447	$m_{\text{lightest}}^\nu / \text{meV}$ (NO)			6.21
$v_H / \text{GeV}$	227.799	$m_2^\nu / \text{meV}$			10.65
$M_N / \text{GeV}$	$1.088 \times 10^6$	$m_3^\nu / \text{meV}$			50.87
$\lambda_1$	1.772	$\phi_{12} / \pi$			0.5
$\lambda_2$	2.0	$\phi_{13} / \pi$			0.5
$\lambda_3$	-0.057	$\phi_{23} / \pi$			1.0
$\lambda_4$	1.782	$m_e / \text{MeV}$	0.486	0.486 $\rightarrow$ 0.486	0.486
$\lambda_5$	-1.644	$m_\mu / \text{GeV}$	0.102	0.102 $\rightarrow$ 0.102	0.102
$\lambda_6$	2.0	$m_\tau / \text{GeV}$	1.746	1.743 $\rightarrow$ 1.747	1.746
$\lambda_7$	-1.466	$M_H / \text{GeV}$ (Higgs boson)	125.25	125.08 $\rightarrow$ 125.42	125.25
$\lambda_8$	-0.392	$M_{DM} / \text{GeV}$ (lightest dark scalar)			87.6
$\lambda_9$	1.171	$M_N / \text{GeV}$			$1.09 \times 10^6$
$\lambda_{10}$	-1.154	$M_{H_0} / \text{GeV}$ (Heavy Higgs)			449.57
$\varphi_5$	1.764	$M_{A_0} / \text{GeV}$ (Heavy Pseudoscalar)			435.98
$\varphi_9$	0.059	$M_{H_0}^+ / \text{GeV}$ (Charged Active)			373.5
$\varphi_{10}$	6.228	$M_{H_a}^+ / \text{GeV}$ (Charged Dark)			345.5
		$M_{H_b}^+ / \text{GeV}$ (Charged Dark)			347.7
		$M_{H_a}^0 / \text{GeV}$ (Neutral Dark)			109.47
		$M_{H_b}^0 / \text{GeV}$ (Neutral Dark)			411.36
		$M_{H_c}^0 / \text{GeV}$ (Neutral Dark)			413.1
		$\chi^2$			0.52

$$\Delta m_{31}^2$$

$$\Delta m_{21}^2$$

Both generated

# Conclusions

## The Discrete Dark Matter Model

From a  $A_4$  symmetric origin:

arXiv:2-[soon](#)

- Reproduces lepton masses and mixings.

$$\sin^2 \theta_{12}, \quad \sin^2 \theta_{12}, \quad \sin^2 \theta_{23} \quad \delta_l^{CP} \sim \pi \quad (\text{Normal Ordering})$$

- **Scotogenic-Seesaw** mass mechanism for neutrinos.

$$2 \lesssim m_1^\nu \lesssim 8 [meV] \quad \langle m_{\beta\beta} \rangle \sim 0.2 [meV]$$

- Naturally explains the hierarchy (seesaw dominates over scotogenic):

$$\Delta m_{31}^2 \gg \Delta m_{21}^2$$

- Rich scalar sector: with CP-violation which includes a Scalar Dark Matter candidate stabilized by a remnant

$$A_4 \longrightarrow \mathbb{Z}_2$$

# Back-up Slides

## Dark scalars mass matrices

$$M_{H'_2 H'_3}^2 = \begin{pmatrix} v_{\eta_1}^2 \left(-\frac{3}{2}\lambda_3 + \frac{1}{2}\lambda_4 + \lambda_5 \cos \varphi_5\right) & 6v_{\eta_1} v_s (\lambda_{10} \cos \varphi_{10} + \lambda_9 \cos \varphi_9) \\ 6v_{\eta_1} v_s (\lambda_{10} \cos \varphi_{10} + \lambda_9 \cos \varphi_9) & v_{\eta_1}^2 \left(-\frac{3}{2}\lambda_3 + \frac{1}{2}\lambda_4 + \lambda_5 \cos \varphi_5\right) \end{pmatrix},$$

$$M_{A'_2 A'_3}^2 = \begin{pmatrix} v_{\eta_1}^2 \left(-\frac{3}{2}\lambda_3 + \frac{1}{2}\lambda_4 - \lambda_5 \cos \varphi_5\right) - v_s^2 (8\lambda_8) & 2v_{\eta_1} v_s (\lambda_9 \cos \varphi_9 + \lambda_{10} \cos \varphi_{10}) \\ 2v_{\eta_1} v_s (\lambda_9 \cos \varphi_9 + \lambda_{10} \cos \varphi_{10}) & v_{\eta_1}^2 \left(-\frac{3}{2}\lambda_3 + \frac{1}{2}\lambda_4 - \lambda_5 \cos \varphi_5\right) - v_s^2 (8\lambda_8) \end{pmatrix},$$

$$M_{\text{CPV}}^2 = \begin{pmatrix} -v_{\eta_1}^2 (\lambda_5 \sin \varphi_5) & -2v_{\eta_1} v_s (\lambda_9 \sin \varphi_9 + \lambda_{10} \sin \varphi_{10}) \\ -2v_{\eta_1} v_s (\lambda_9 \sin \varphi_9 + \lambda_{10} \sin \varphi_{10}) & v_{\eta_1}^2 (\lambda_5 \sin \varphi_5) \end{pmatrix}$$

# Back-up Slides

## Yukawa couplings under basis change

$$Y_{n\alpha}^a = (O_{hH_0}^T)^a_k Y_{n\alpha}^k, \quad \text{for } a = h, H_0 \quad \text{and} \quad k = H'_1,$$

$$Y_{n\alpha}^a = (O_{GA_0}^T)^a_k Y_{n\alpha}^k, \quad \text{for } a = G, A_0 \quad \text{and} \quad k = A'_1,$$

$$Y_{n\alpha}^a = (O_{\chi}^T)^a_k Y_{n\alpha}^k, \quad \text{for } a = \chi_1^D, \dots, \chi_4^D, \quad \text{and} \quad k = H'_2, H'_3, A'_2, A'_3.$$