

* talk based on: Electroweak Phase Transition in a Dark Sector with CP Violation

by LB, Margarete Mühlleitner and Jonas Müller [2204.13425]

Electroweak Phase Transition in a Dark Sector with CP Violation

Lisa Biermann¹

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Lisa Biermann (ITP, KIT)

• Electroweak baryogenesis (EWBG) can generate observed baryon asymmetry of the universe (BAU) ($\eta \simeq 6.1 \times 10^{-10}$ [Planck, 2018]) if [A. D. Sakharov, 1967], [D. Morrissey, M. Ramsey-Musolf, 2012]

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 - additional (non-standard) CPV (generation of LH fermion access in front of bubble wall)
 - → creation of BAU: EW sphaleron transitions (triggered by LH fermion access) [F. R. Klinkhammer, N.S. Manton, 1984]
 - sufficiently strong departure from thermal equilibrium, $\xi_c \equiv \frac{v_c}{T_c} \gtrsim 1$ (*conservation* of BAU inside bubble) [M. Quiros, 1994] \Rightarrow **SFOEWPT** (*strong first-order electroweak phase transition*)

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Now .. Can we generate an SFOEWPT within 'CP in the Dark'?

Can the 'hidden' CP violation be translated to the visible sector?

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'CP in the Dark' [D. Azevedo, P. Ferreira, M. Mühleitner, S. Patel, R. Santos, J. Wittbrodt, 2018]

• N2HDM-like extended scalar sector, *one* discrete \mathbb{Z}_2 symmetry

$$\Phi_1 \to +\Phi_1, \quad \Phi_2 \to -\Phi_2, \quad \Phi_S \to -\Phi_S$$

• $SU(2)_L \times U(1)_Y$ and \mathbb{Z}_2 -invariant tree-level potential:

$$V^{(0)} = m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 + \frac{m_S^2}{2} \Phi_S^2 + \left(\mathbf{A} \Phi_1^{\dagger} \Phi_2 \Phi_S + h.c.\right) + \frac{\lambda_1}{2} |\Phi_1|^4 + \frac{\lambda_2}{2} |\Phi_2|^4 + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 |\Phi_1^{\dagger} \Phi_2|^2 + \frac{\lambda_5}{2} \left[\left(\Phi_1^{\dagger} \Phi_2\right)^2 + h.c. \right] + \frac{\lambda_6}{4} \Phi_S^4 + \frac{\lambda_7}{2} |\Phi_1|^2 \Phi_S^2 + \frac{\lambda_8}{2} |\Phi_2|^2 \Phi_S^2$$

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• general vacuum structure @ $T \neq 0$:

r charge-breaking VEV, $\omega_{\rm CB} = 0$

$$\Phi_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} \rho_1 + i\eta_1 \\ \zeta_1 + \omega_1 + i\Psi_1 \end{pmatrix}, \ \Phi_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} \rho_2 + \omega_{\rm CB} + i\eta_2 \\ \zeta_2 + \omega_2 + i(\Psi_2 + \omega_{\rm CP}) \end{pmatrix}, \ \Phi_s = \zeta_s + \omega_s$$

← CP-violating VEV

• general vacuum structure @ T = 0:

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$$\langle \Phi_{1} \rangle |_{T=0} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \nu_{1} \end{pmatrix}, \quad \langle \Phi_{2} \rangle |_{T=0} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \quad \langle \Phi_{S} \rangle |_{T=0} = 0$$

 $\rightarrow \omega_1|_{T=0 \text{ GeV}} = v_1 \equiv v = 246.22 \text{ GeV}$, SM-Yukawa sector and tree-level FCNCs prohibited

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- $\rightarrow \mathbb{Z}_2$ symmetry *unbroken* \Rightarrow conserved quantum number: *dark charge*
 - * Φ_1 (*SM-like particles* with +1): G^{\pm} , G^0 , h
 - * Φ_2 , Φ_S (dark particles with -1): H^{\pm} , h_1 , h_2 , h_3 ($m_{h_1} < m_{h_2} < m_{h_3}$)

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- \Rightarrow **DM**: *stable* particle dark matter candidate h_1
- \Rightarrow explicit CPV: introduced through Im (A) $\neq 0$
 - \rightarrow CPV after SSB, but vacuum is CP-symmetric \Rightarrow CPV is *explicit*
 - \rightarrow solely in the dark sector h_1, h_2, h_3 : states with mixed CP quantum number
 - ⇒ not constrained by EDM constraints

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- \Rightarrow 'CP in the Dark' CPV + DM + SFOEWPT (?)

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- true vacuum state @ finite temperature (FT) including radiative corrections = global minimum of the **effective potential** @ **FT**
- general one-loop effective potential @ FT splits into temperature-dependent and independent part [L. Dolan, R. Jackiw, 1974]



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• V^{CT} absorbs NLO scalar mass and angle shift [P. Basler et al., 2017]

$$0 = \partial_{\phi_i} (V^{CW} + V^{CT}|_{\vec{\omega} = \vec{\omega}_{tree}})$$
$$0 = \partial_{\phi_i} \partial_{\phi_j} (V^{CW} + V^{CT}|_{\vec{\omega} = \vec{\omega}_{tree}})$$

BSMPT [P. Basler, M. Mühlleitner, J. Müller, 2018/20] https://github.com/phbasler/BSMPT

global minimization of the one-loop corrected effective potential @ T ∈ {0, 300} GeV in non-zero FT VEV space \$\vec{\omega} → \text{get}(\$\vec{\omega}_1,\$\vec{\omega}_2,\$\vec{\omega}_{\mathcal{CB}},\$\vec{\omega}_{\mathcal{CP}},\$\vec{\omega}_{\mathcal{S}}\$) @ T

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• temperature-dependent EW VEV v(T):

$$v(T) = \sqrt{\overline{\omega}_1^2 + \overline{\omega}_2^2 + \overline{\omega}_{\mathrm{CB}}^2 + \overline{\omega}_{\mathrm{CP}}^2}$$

• critical temperature T_c : $V^{(1)}(\bar{\omega} = 0, T_c) \equiv V^{(1)}(\bar{\omega}_c \neq 0, T_c)$ \rightarrow bisection method for T_c

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 SFOEWPT:

$$\xi_c \equiv \frac{v_c}{T_c} \gtrsim 1$$

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$$\Rightarrow$$
 SFOEWPT: ξ_c

$$\xi_c \equiv rac{v_c}{T_c} \gtrsim 1$$

- viable parameter points
 - * pass constraints imposed by: ScannerS [R. Coimbra et al., 2013] [M. Mühlleitner et al., 2020] BSMPT [P. Basler, M. Mühlleitner, J. Müller, 2018/20]
 - * BR $(h \to inv.) < 0.11$ [M. Aaboud et al., 2019]
 - * $\mu_{h \to \gamma \gamma} = 1.12 \pm 0.09$ [A. Sirunyan et al., 2021]

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Results: Mass Parameter Distributions for an SFOEWPT



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 \Rightarrow find SFOEWPT (and NLO-stable) points distributed all over allowed^(*) parameter space

(*): by Higgs constraints, DM constraints, theoretical constraints.

- neither requirement of NLO-VEV stability nor SFOEWPT further constrains the parameter space
- restricted $m_{H^{\pm}}$ -range due to $\mu_{\gamma\gamma}$ cut (see Slide 8)

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- tree-level couplings of *h* identical to those of SM Higgs boson
- only presence of dark particles can change $BR(h \rightarrow \gamma \gamma)$



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- increase towards smaller $m_{H^{\pm}}$ (governed by λ_3)
- upper bound in CMS-plot: BFB and unitarity bounds restrict maximal λ_3
- \Rightarrow future increased precision on $\mu_{\gamma\gamma}$ can cut the parameter space on $m_{H^{\pm}}$ substantially

Results: BR $(h \rightarrow inv.)$



Results: BR $(h \rightarrow inv.)$



- SFOEWPT points scattered across allowed ScannerS parameter space
- BR($h \rightarrow \text{inv.}$) strongly correlated with μ_{VV} (V = Z, W) (gauge boson signal strength), agree with results for *fully dark phase* of N2HDM [I. Engeln et al., 2020]
- \rightarrow for $\mu_{VV} \rightarrow 1$, SM-like Higgs BR converges to SM value (invisible decay not allowed)
- ⇒ future precise measurements of BR(h → inv.) and μ_{VV} can constrain parameter space, however *no* further insights into strength of the EWPT





→ @ FT: $|\overline{\omega}_{CP}| \neq 0$ possible for SFOEWPT points \Leftrightarrow @ T = 0 GeV: $\overline{\omega}_{CP}|_{T=0 \text{ GeV}} = 0$ → CPV only possible explicitly (Im(A) $\neq 0$) → no clear correlation - but: $|\overline{\omega}_{CP}| > 0$ only for Im(A) $\neq 0$

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• two different VEV patterns in detail:



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- → @ FT: \mathbb{Z}_2 symmetry is broken → dark charge no longer conserved → dark sector **mixes** with SM-like particles
- ⇒ additional non-standard CPV transferred to the SM-like couplings to fermions @ FT!



$$f_{\chi\chi} \cdot \sigma_{\rm SI,\,DM-nucl.} \equiv \frac{\Omega_{\rm prod}h^2}{\Omega_{\rm obs}h^2} \cdot \sigma_{\rm SI,\,DM-nucl.}$$



Viable SFOEWPT parameter points

- \Rightarrow compatible with *relic density* ($< \Omega h^2$)
- \Rightarrow above neutrino floor
- \Rightarrow testable at future *direct detection* experiments

 $\sigma_{\rm SI, DM-nucl.}$

 $\Omega_{\text{prod}}h^2$

 $f_{\chi\chi} \cdot \sigma_{\text{SI, DM-nucl.}} \equiv$

Conclusion

- dynamical generation of the baryon asymmetry of the universe (BAU) possible if *Sakharov* conditions fulfilled
- electroweak baryogenesis: fulfill Sakharov conditions with
 - BSM models
 - non-standard *CP-violation* (CPV)
 - strong first-order electroweak phase transition (SFOEWPT)
- 'CP in the Dark': special N2HDM + one discrete \mathbb{Z}_2 symmetry
 - dark sector with DM candidate h₁
 - explicit CPV in the dark sector at zero temperature
- \Rightarrow **BSMPT**: global minimization of the one-loop corrected effective potential at finite temperature
- → viable SFOEWPT parameter points for 'CP in the Dark'
 - · within reach of future direct detection experiments
 - ⇒ show spontaneous CPV at finite temperature!

 \rightarrow **Open question:** Can these points successfully generate the BAU?

Thanks for your attention!

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Benchmark Points

All points have: $\lambda_1 \simeq 0.258$, $m_{11}^2 \simeq -7824 \text{ GeV}^2$

	point	$m_{22}^2 [\text{GeV}^2] \qquad m_S^2 [\text{GeV}^2]$		Re (A	Re (A) [GeV] Im (λ_1			
_	no sponCPV sponCPV	96 703.414 32 442.9 65 258.809 36 279.8		949 159.627 847 279.502		-325.391 3.5 -326.645 3.6				
_	point	λ_3	λ_4	λ_5		λ_6	λ_7	λ_8		
	no spon CPV spon CPV	$-0.796 \\ -0.821$	0.787 0.220	-(-(0.055 0.371	10.446 4.715	7.596 7.760	4.683 14.781		
point	<i>m</i> ₁₁ +	<i>m</i> _{<i>h</i>} ,	,	nha	mha		T _c	vc		
no spon spon CP	CPV 269.386 V 200.940	241.71 62.68	18 30 0 21	8.943 8.700	549.265 560.206	14	144.21 189.77		236.53 235.85	
point	ξc	$\overline{\omega}_{CB}$		$\overline{\omega}_1$	$\overline{\omega}_2$	ō	<i>ω</i> _{CP}	$\overline{\omega}_S$		
no spon spon CP	CPV 1.64 V 1.24	-8.977 × -2.212 ×	$ \begin{array}{ccc} 10^{-7} & 23 \\ 10^{-5} & 22 \end{array} $	36.53 26.46	$9.093 \times 10^{-52.72}$	⁻⁷ -3.79	3×10^{-7} 9.52	4.604 × -27.	10 ⁻⁷ 58	

Baryon Asymmetry of the Universe (BAU)

initial: *Big Bang* (symmetric universe) \Leftrightarrow today: **BAU** (asymmetric universe)

$$\eta \equiv rac{n_b - ar{n}_b}{n_\gamma} \simeq rac{n_b}{n_\gamma} \simeq 6.1 imes 10^{-10}$$
 [Planck, 2018]

How can we generate a non-zero baryon asymmetry of the universe?

[Sakharov, 1967]: dynamical generation of a BAU with an initially symmetric state possible if

condition								
existence of <i>B</i> violating processes	⇒	sphaleron-mediated @ $T > T_{EW} = 100 \text{ GeV}$ [N. Manton, 1983], [F. Klinkhammer, N. Manton, 1984]						
$\mathcal C$ and $\mathcal {CP}$ violation (CPV)	\Rightarrow	<i>Cabibbo-Kobayashi-Maskawa</i> mechanism (?) [N. Cabibbo, 1963], [M. Kobayashi, T. Maskawa, 1973]						
departure from thermal equilibrium	\Rightarrow	electroweak phase transition (EWPT) [D. Kirznits, 1972], [L. Dolan, R. Jackiw, 1974]						
Lisa Biermann (ITP KIT) EV	VPT in 'CP i	in the Dark' 08.11.2022 2/3						

Electroweak Baryogenesis (EWBG) [D. Morrissey, M. Ramsey-Musolf, 2012] V_{eff}

- EWBG takes place around $T \sim T_{\rm EW}$
- EWPT happens and bubbles with non-zero vacuum expectation value (VEV) are created and expand
- necessary departure from thermal equilibrium achieved through strong first-order EWPT (SFOEWPT)



$$V(v = 0, T_c) = V(v \neq 0, T_c)$$



- How do we see this in the potential? \rightarrow global minimum jumps from symmetric to broken minimum @ T_c
- \rightarrow 'strong': conservation of BAU through sufficient suppression of the sphaleron rate inside the bubbles

irvon-wash-out condition* I. Quiros, 1994]

 $T = T_2 > T_C$

 $T = T_C$

 $T = T_1 < T_C$

- EWPT in SM only smooth cross-over [K. Kajantie et al., 1996]
- need BSM models that enable an SFOEWPT* + non-standard CPV

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