

Thermal Friction as a Solution to the Hubble Tension

**Kim V. Berghaus,
California Institute of Technology**

based on

Phys. Rev. D 101, 083537 (2020) (1911.06281)

Phys. Rev. D 107, 103515 (2023) (2204.09133)

With Tanvi Karwal, KICP U Chicago

The Hubble Tension

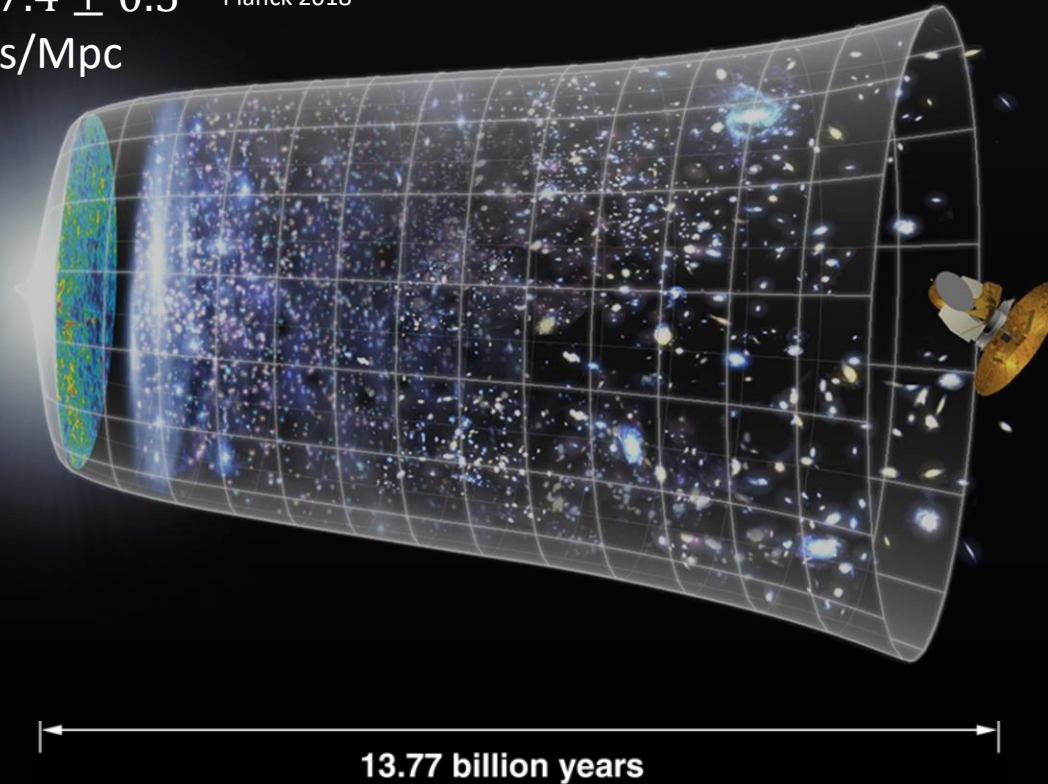
Λ CDM +
CMB measurement

$$H_0 = 67.4 \pm 0.5 \text{ km/s/Mpc} \quad \text{Planck 2018}$$

$\sim 5 \sigma$

Direct
measurement

$$H_0 = 73.3 \pm 1.0 \text{ km/s/Mpc} \quad \text{SH0ES 2022}$$



The Large-Scale Structure Tension

Λ CDM +
CMB measurement

$$S_8 = 0.82 \pm 0.01 \quad \text{Planck 2018}$$

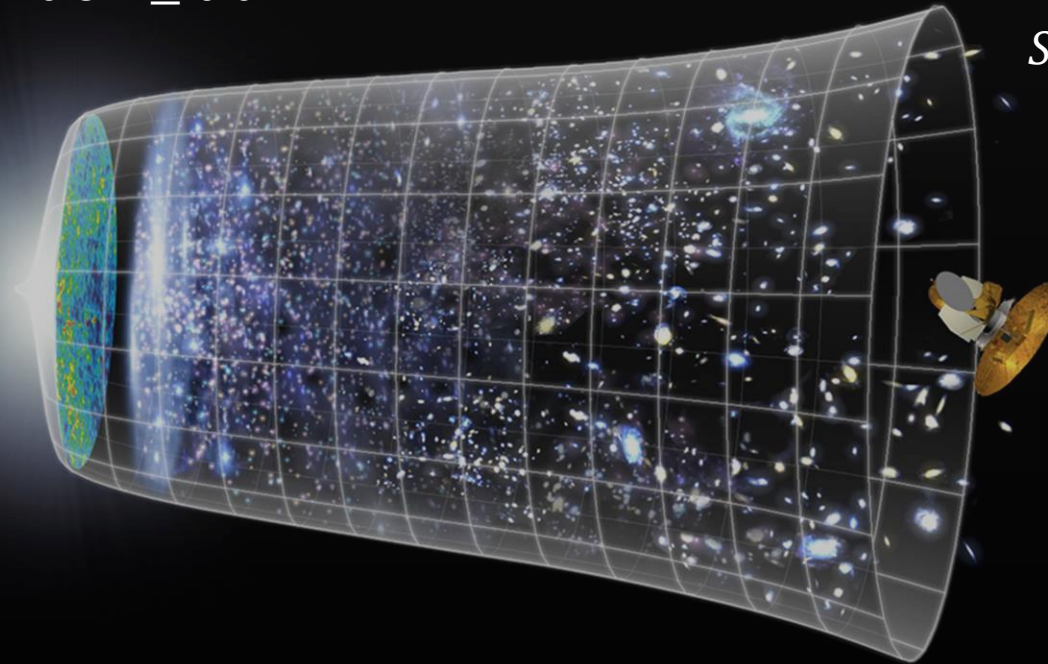
$\sim 2 - 3 \sigma$

Galaxy clustering
+ weak lensing
(DES Y1)

$$S_8 = 0.77 \pm 0.02 \quad \text{DES 2017}$$

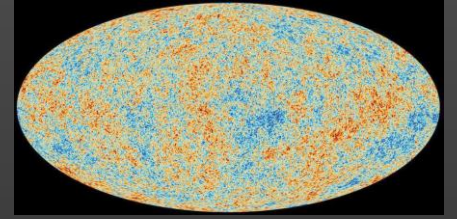
$$S_8 = \sigma_8 \sqrt{\frac{\Omega_m}{0.3}}$$

Amplitude of matter
fluctuations



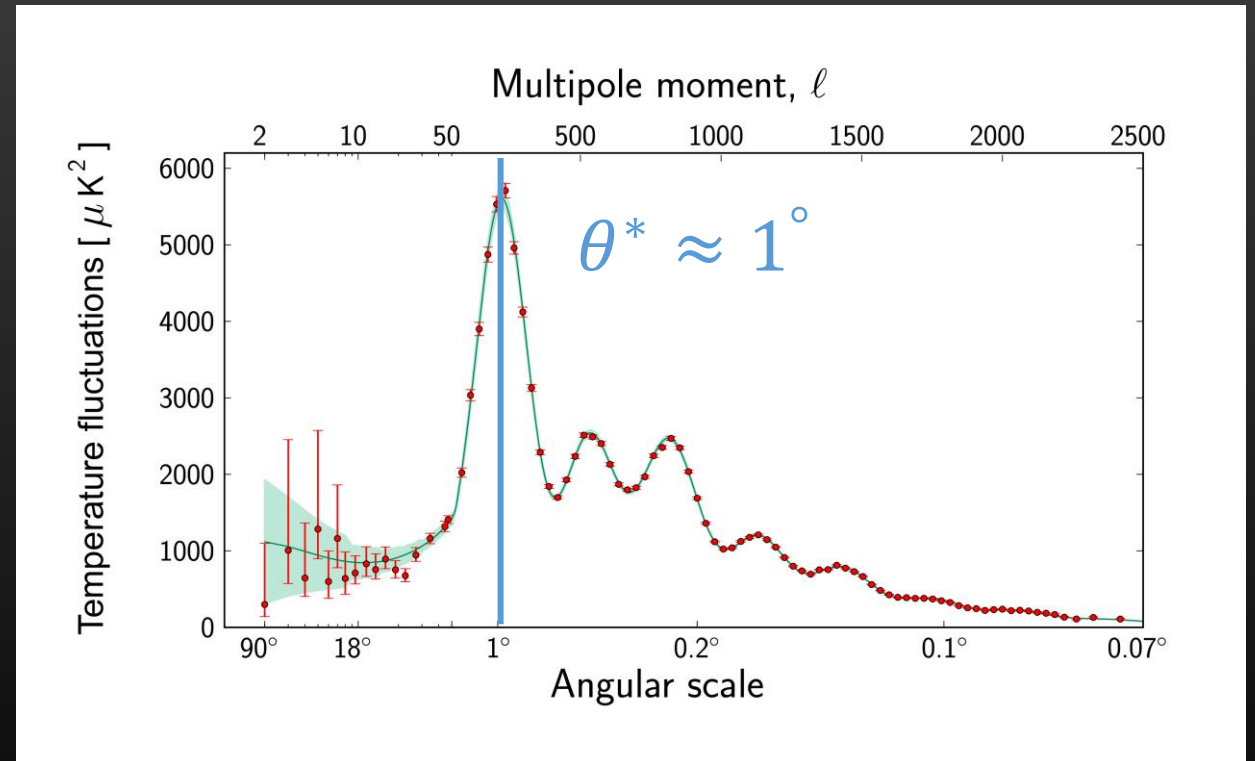
13.77 billion years

The Hubble Measurement with the CMB



$$\theta^* \propto r_s H_0$$

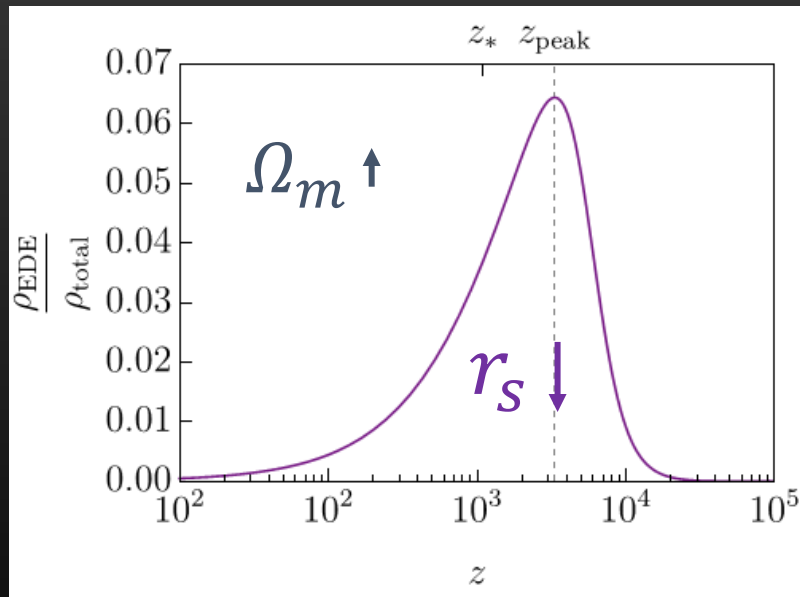
- r_s depends only on physics before formation of CMB
- Lowering r_s increases H_0



Planck 2018

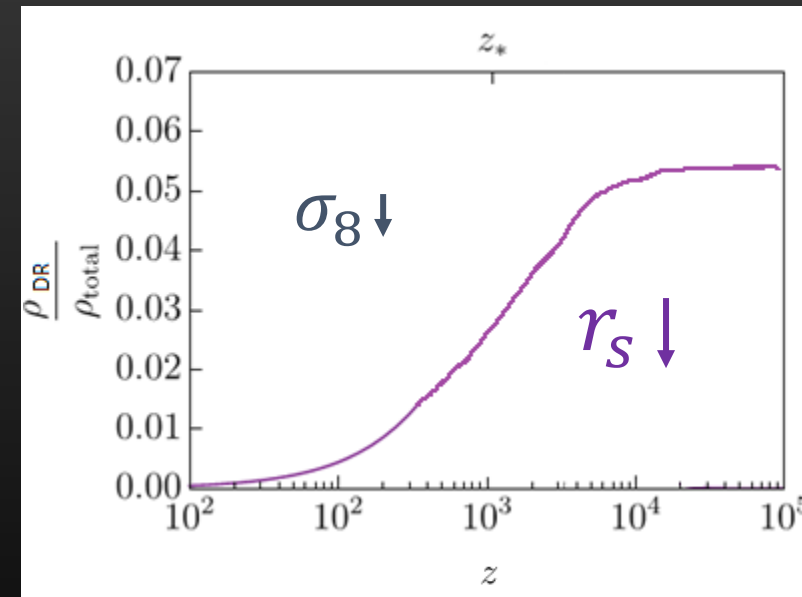
New physics that can lower the sound horizon

Early Dark Energy



$$S_8 = \sigma_8 \sqrt{\frac{\Omega_m}{0.3}}$$

Extra Radiation



Maintains great fit to CMB

Resolves H_0 but exacerbates LSS tension (Hill et. al., 2020, Ivanov et al. 2020)

Is fine-tuned

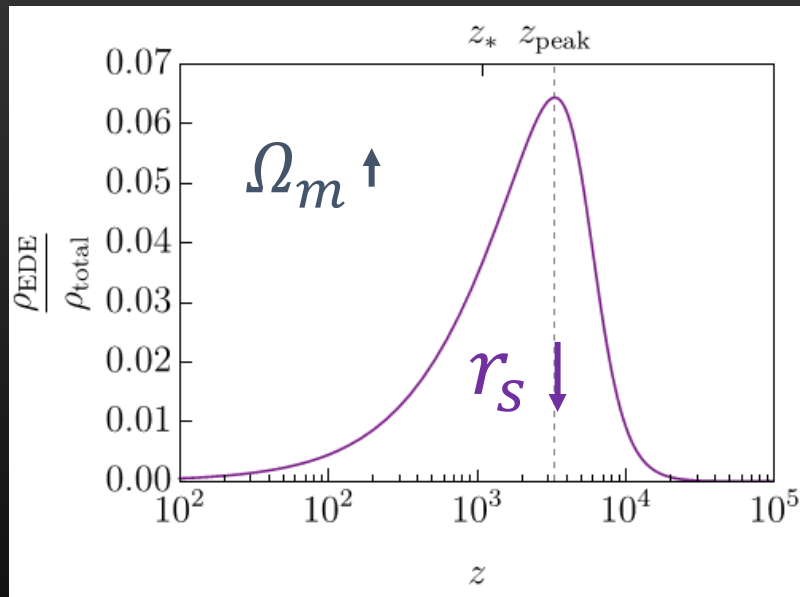
Worsens fit to CMB

Can ease LSS and Hubble tension

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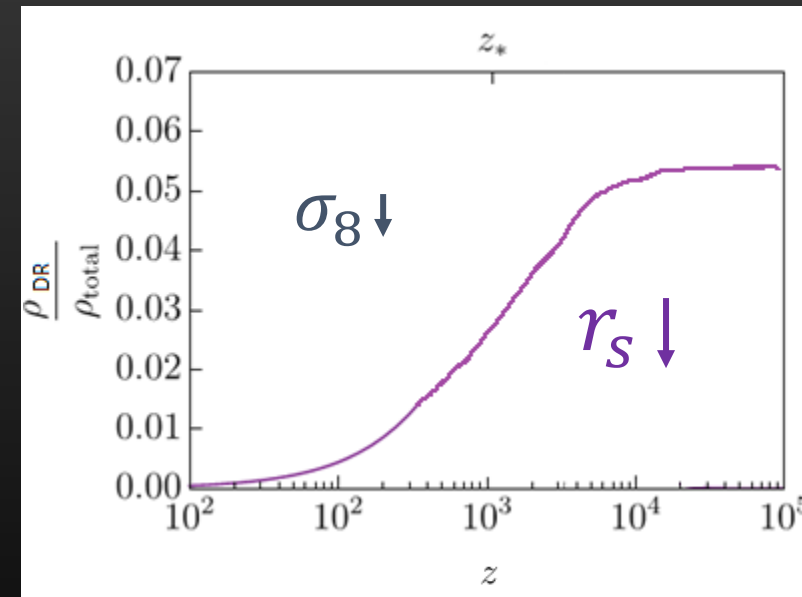
Early Dark Energy with **thermal friction**

Extra Radiation



combines

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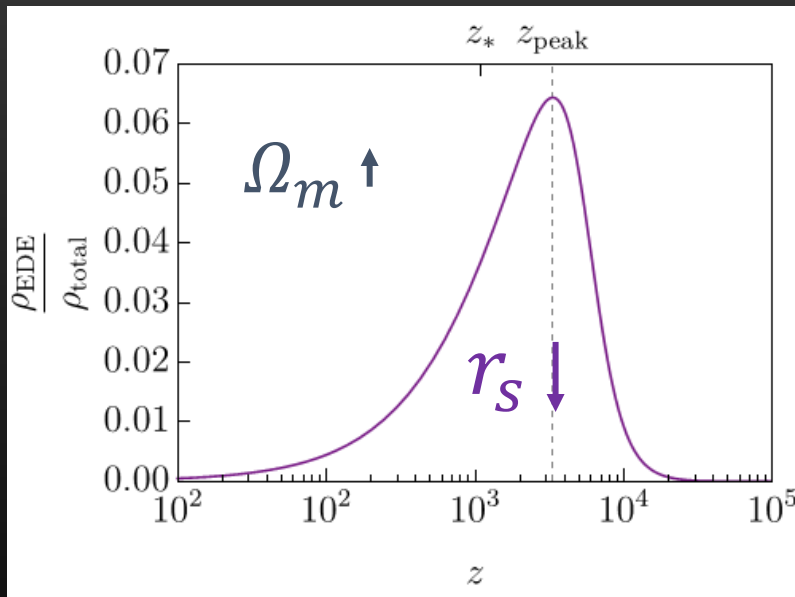
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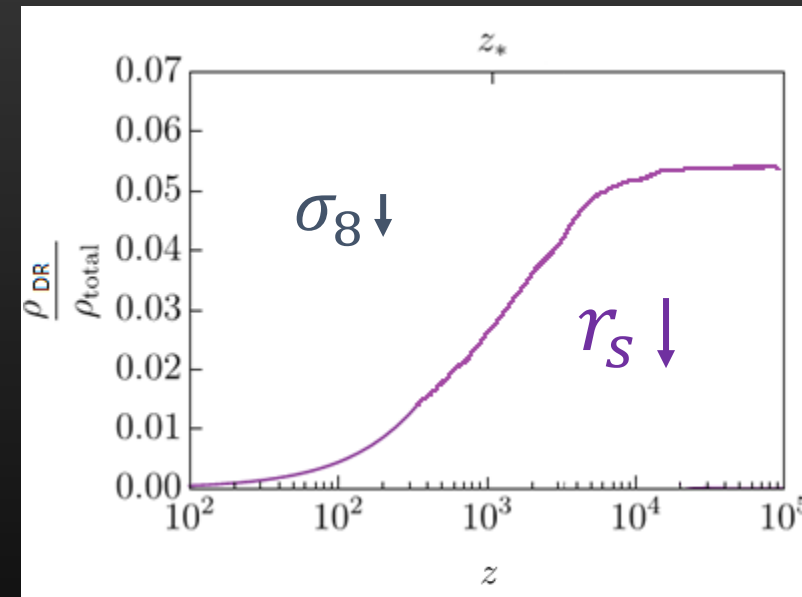
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Maintains great fit to CMB ?

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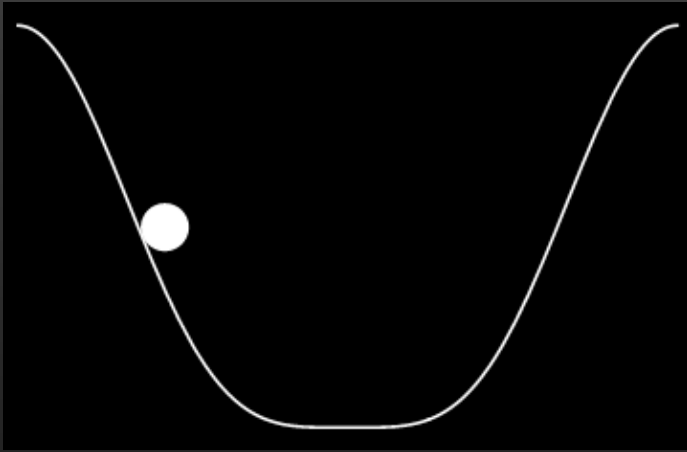
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Worsens fit to CMB

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Early Dark Energy



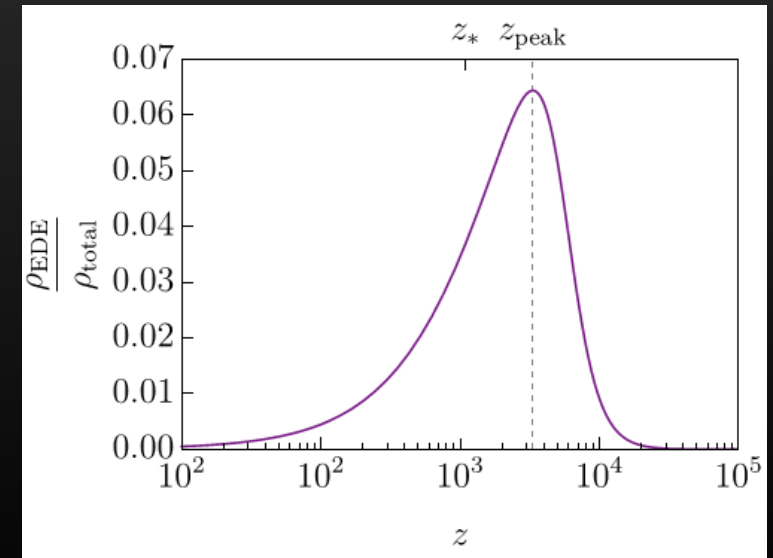
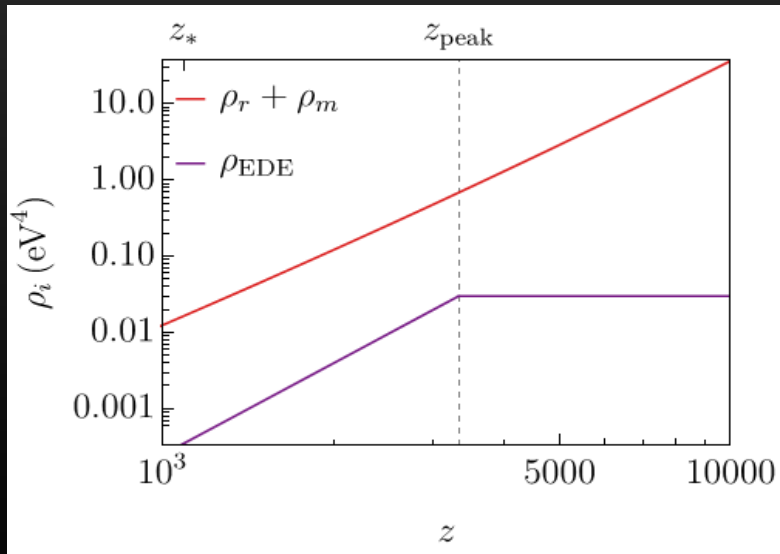
Frozen at early times by Hubble friction

$$\ddot{\phi} + 3H\dot{\phi} + V' = 0$$

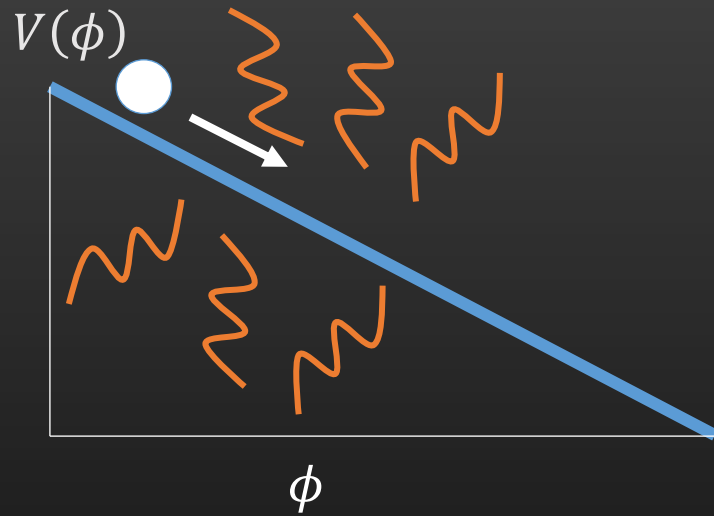
Dilutes away as radiation or faster when axion starts oscillating at critical redshift z_c

$$V \propto \left(1 - \cos \frac{\phi}{f}\right)^n$$

$$n \geq 2$$



Early Dark Energy with **thermal friction**



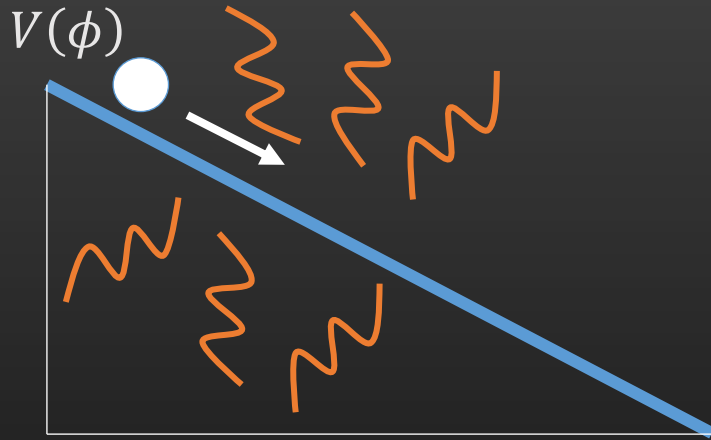
Frozen at early times by **thermal friction**

$$\ddot{\phi} + (3H + \gamma)\dot{\phi} + V' = 0$$

$$\dot{\rho}_{DR} + 4H\rho_{DR} = \gamma\dot{\phi}^2$$

Overdamped at all time. Axion converts its energy into dark radiation at critical redshift z_c

Early Dark Energy with **thermal friction**

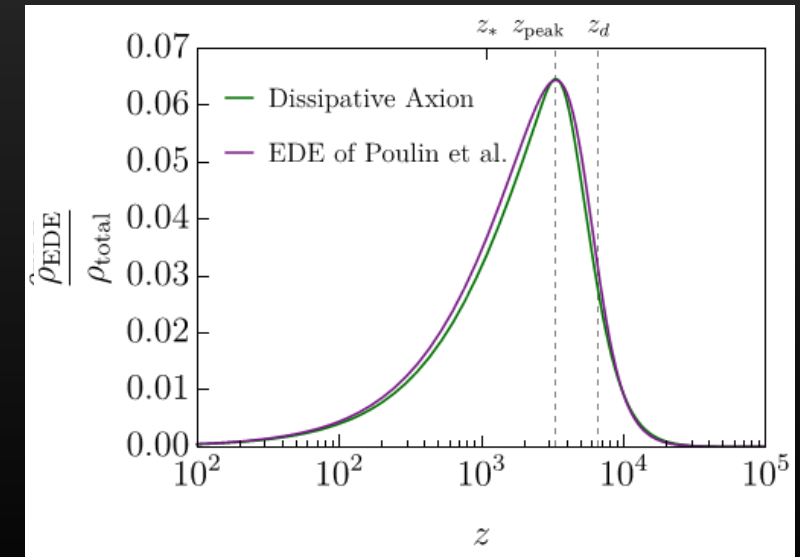
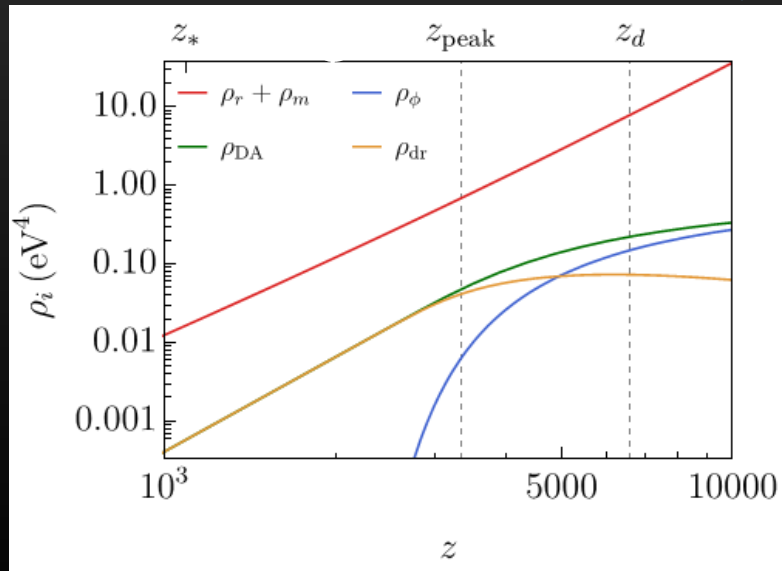


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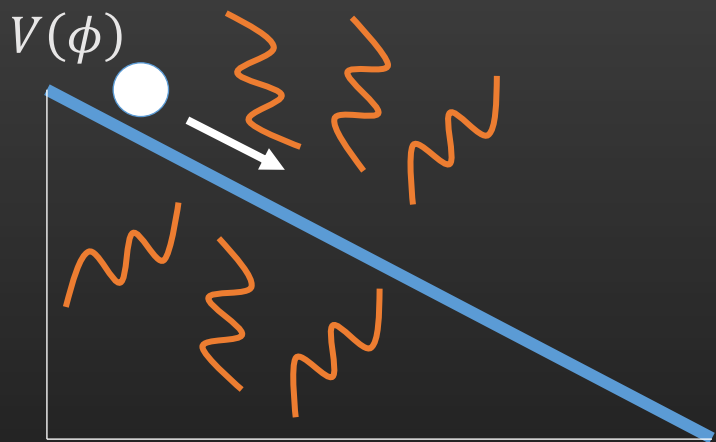
$$\ddot{\phi} + (3H + \Upsilon)\dot{\phi} + V' = 0$$

$$\dot{\rho}_{DR} + 4H\rho_{DR} = \Upsilon\dot{\phi}^2$$

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Early Dark Energy with **thermal friction**



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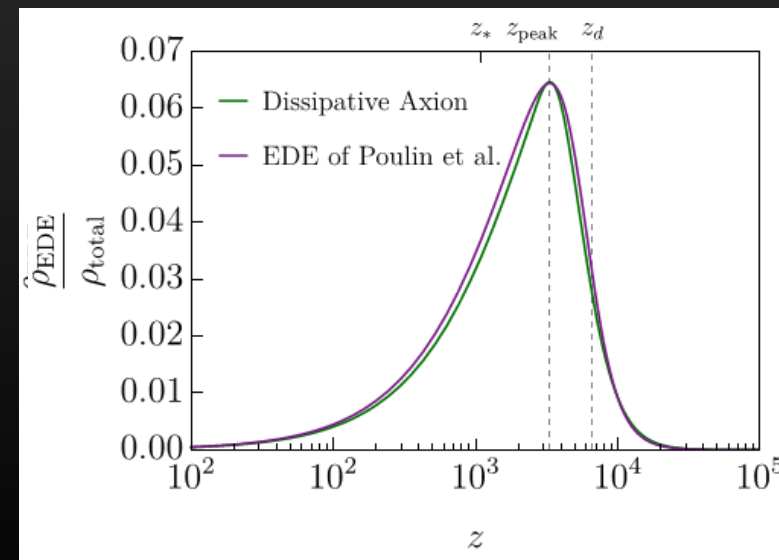
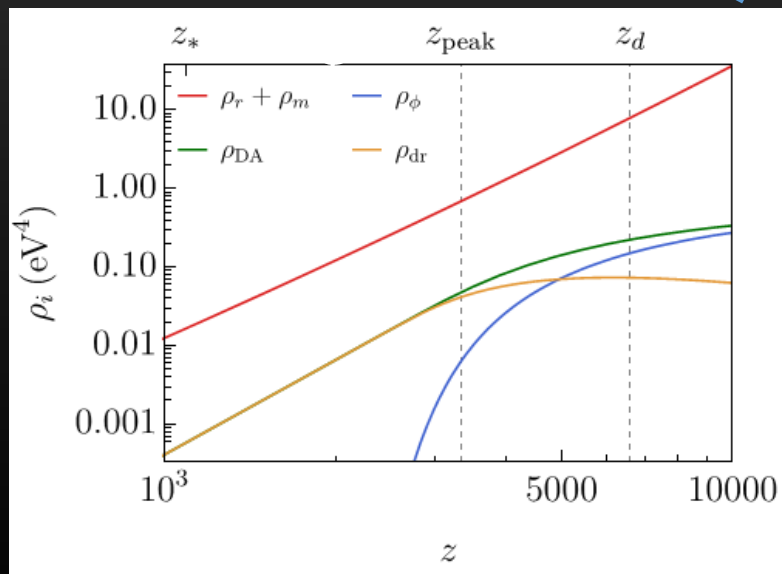
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Overdamped at all time. Axion converts its energy into dark radiation at critical redshift z_c

Thermal friction generically arises for axions couplings to gauge fields

$$L_{\text{int}} = -\phi \frac{\alpha}{16\pi f} \tilde{G}G$$



Analysis

- Put **thermal friction** into CLASS: $V(\phi) = \frac{1}{2}m^2\phi^2$
- Effective parameters: f_{ede}

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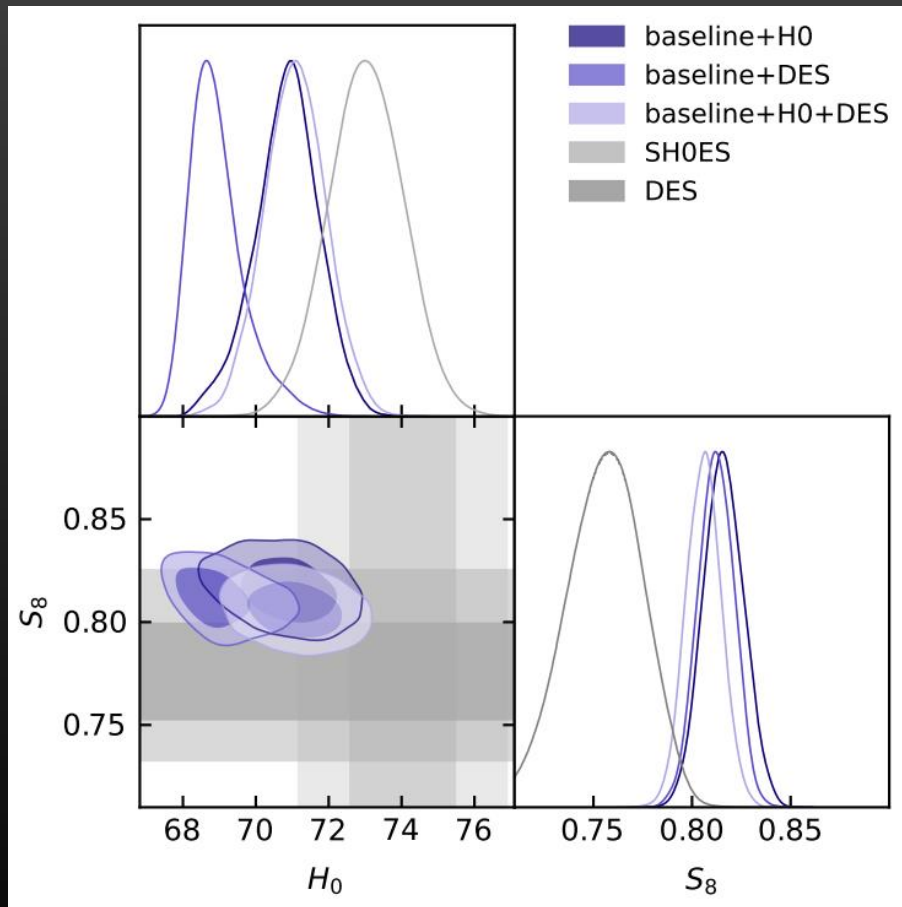
- Put **thermal friction** into CLASS: $V(\phi) = \frac{1}{2}m^2\phi^2$
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Analysis

- Put **thermal friction** into CLASS: $V(\phi) = \frac{1}{2}m^2\phi^2$
- Effective parameters: $f_{ede}, z_c(m, \mathbf{Y}), m$
- Data sets:
 - Planck 2018 CMB (TTTEEE) + lensing
 - BAO (BOSS DR12,, SDSS Main Galaxy Sample, 6dFGS)
 - Pantheon Supernovae sample
 - SH0ES measurement $H_0 = 73.04 \pm 1.04$ km/s/Mpc
 - Dark Energy Survey Year 1 galaxy lensing and clustering

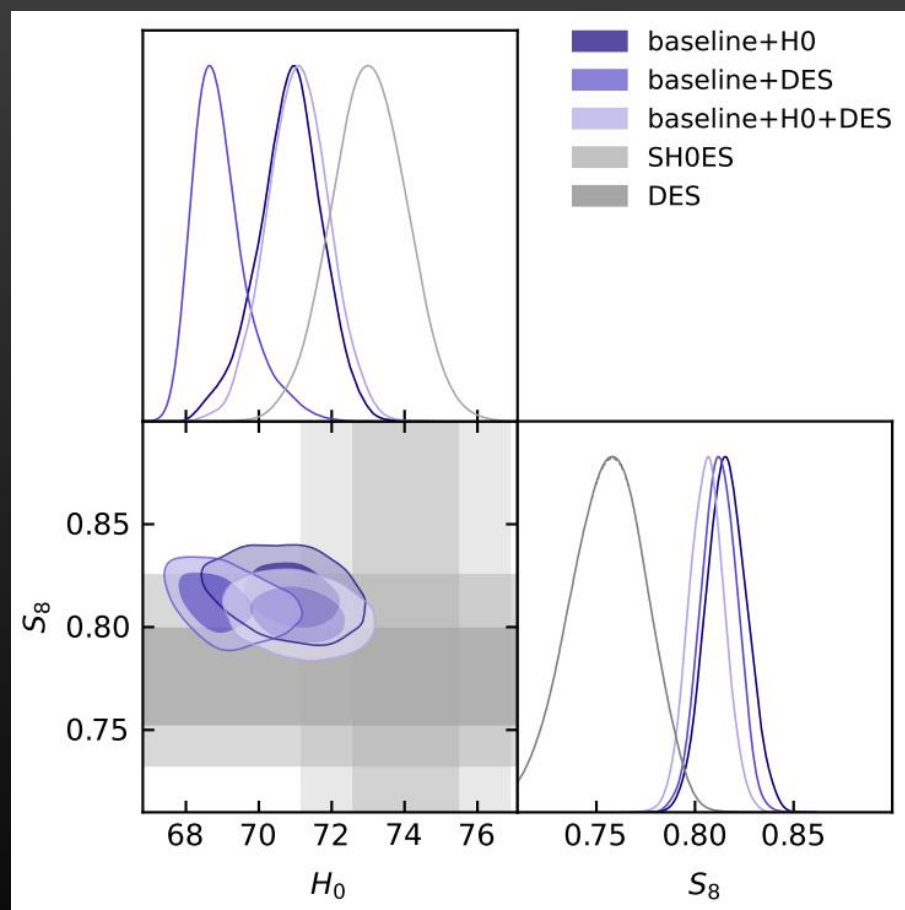
} baseline

Results



Berghaus, Karwal 2022

Results



Berghaus, Karwal 2022

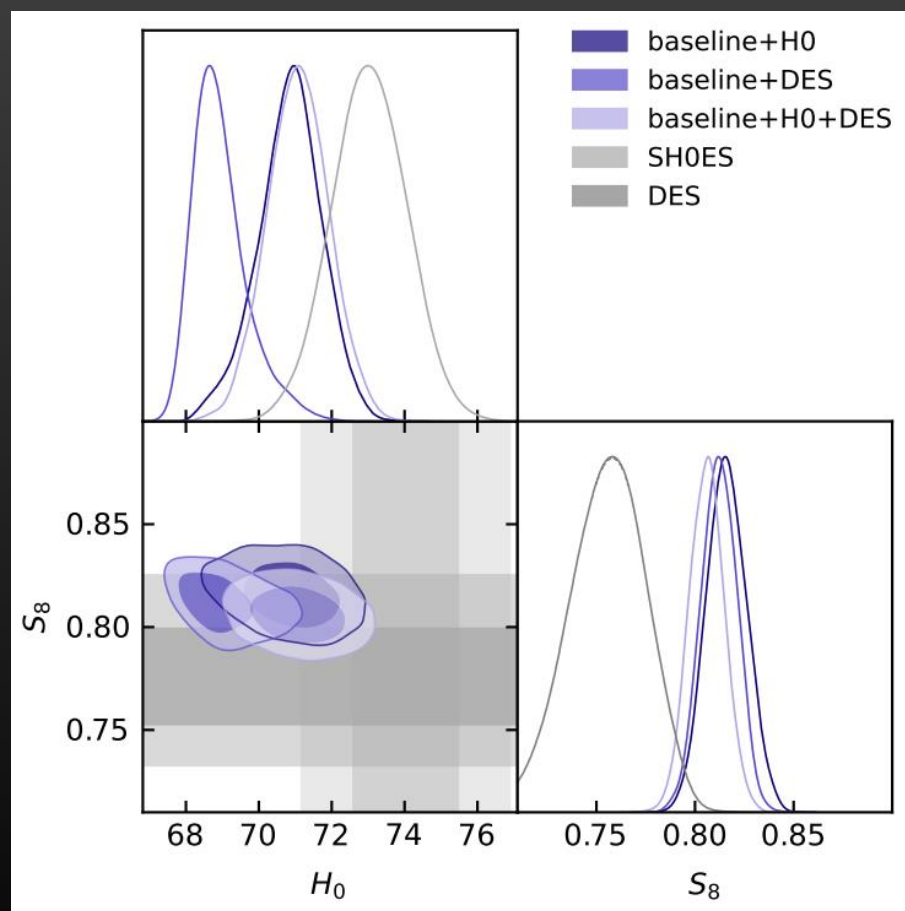
Model	H_0 [km/s/Mpc]	S_8
Λ CDM	$68.76(68.63) \pm 0.36$	$0.8013(0.8055) \pm 0.0087$
DA EDE	$71.08(71.06) \pm 0.85$	$0.8058(0.8075) \pm 0.0089$
N_{eff}	$70.50(70.86) \pm 0.78$	$0.8102(0.8106) \pm 0.0096$

Table VI. 1D marginalized posteriors of measurements quantifying the two cosmological tensions, fitting to *baseline*+ H_0 +DES.

Model	χ^2_{CMB}	$\chi^2_{H_0}$	χ^2_{DES}
Λ CDM	2778.4	18.0	508.0
DA EDE	2778.7	3.6	508.3
N_{eff}	2783.3	4.4	508.8

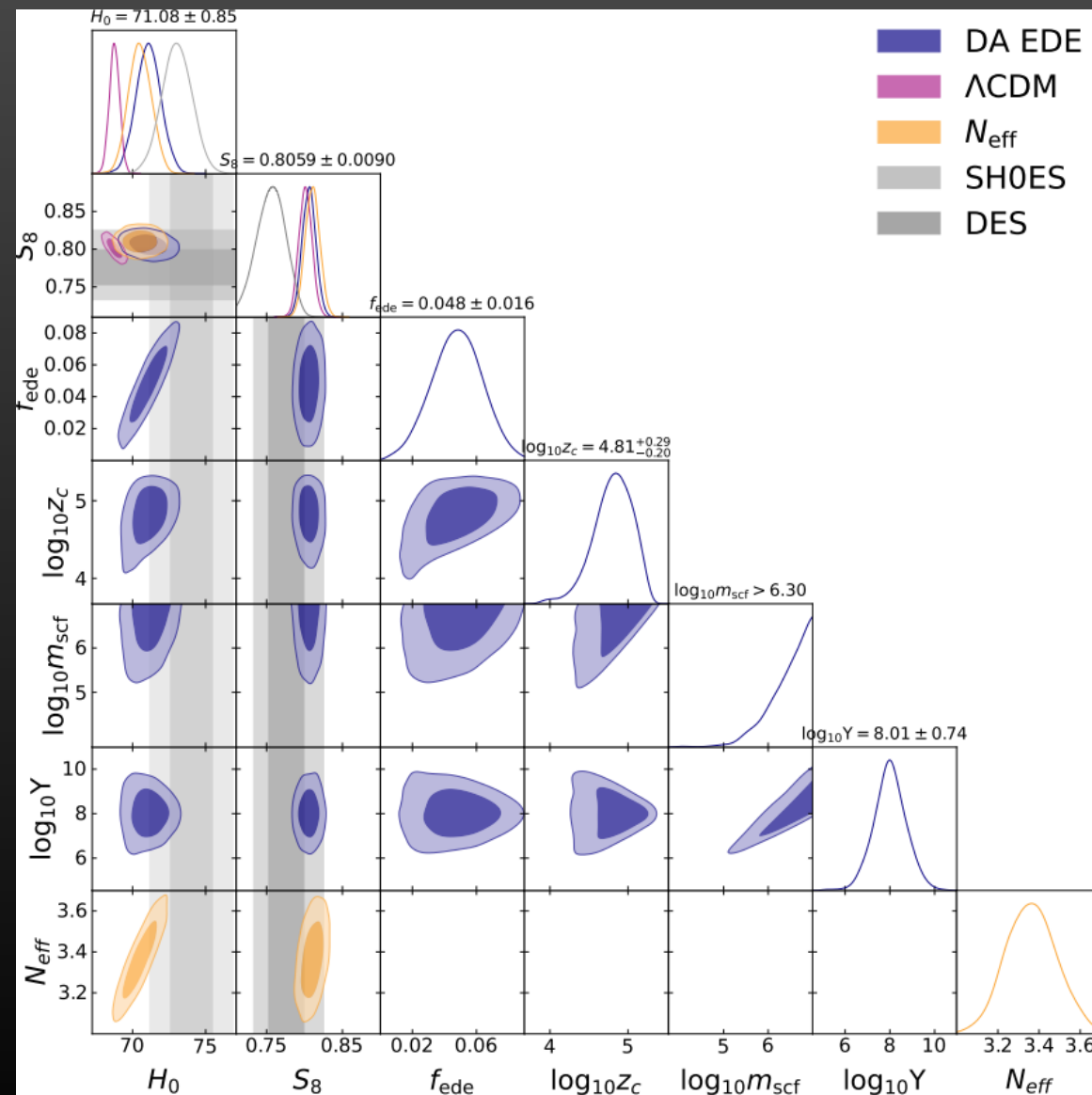
Table VII. The goodness of fit to CMB and DES data, while cumulatively fitting to *baseline*+ H_0 +DES.

Results



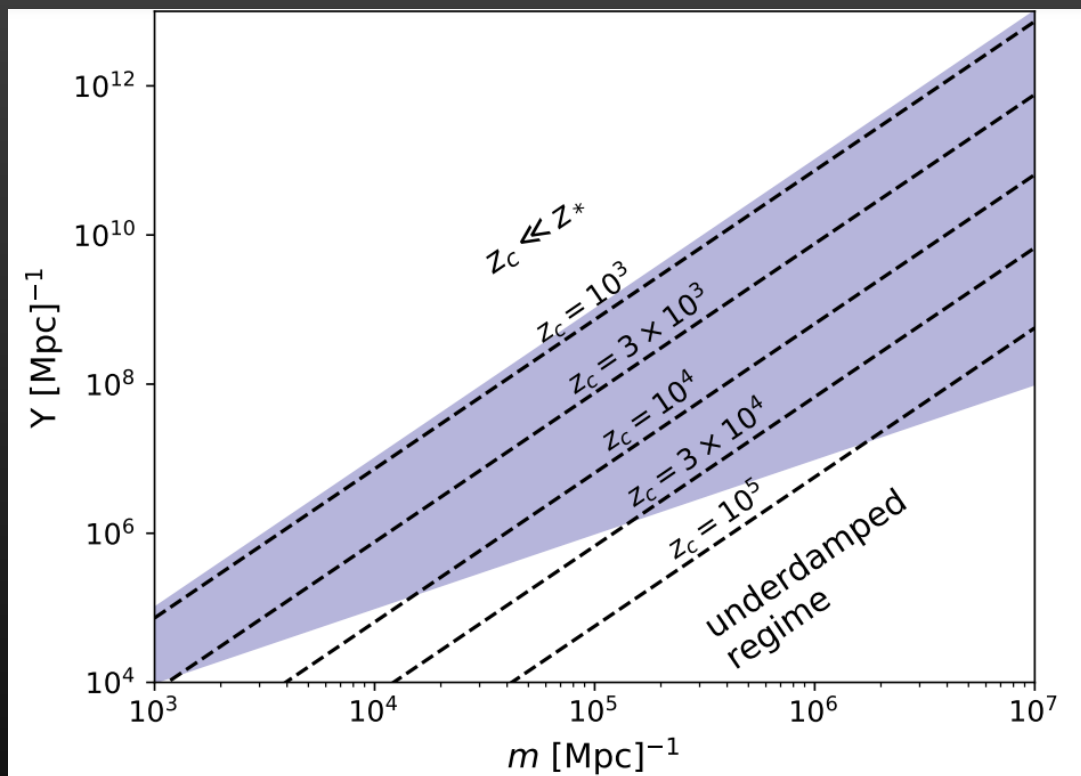
Berghaus, Karwal 2022

Baseline + SH0ES + DES



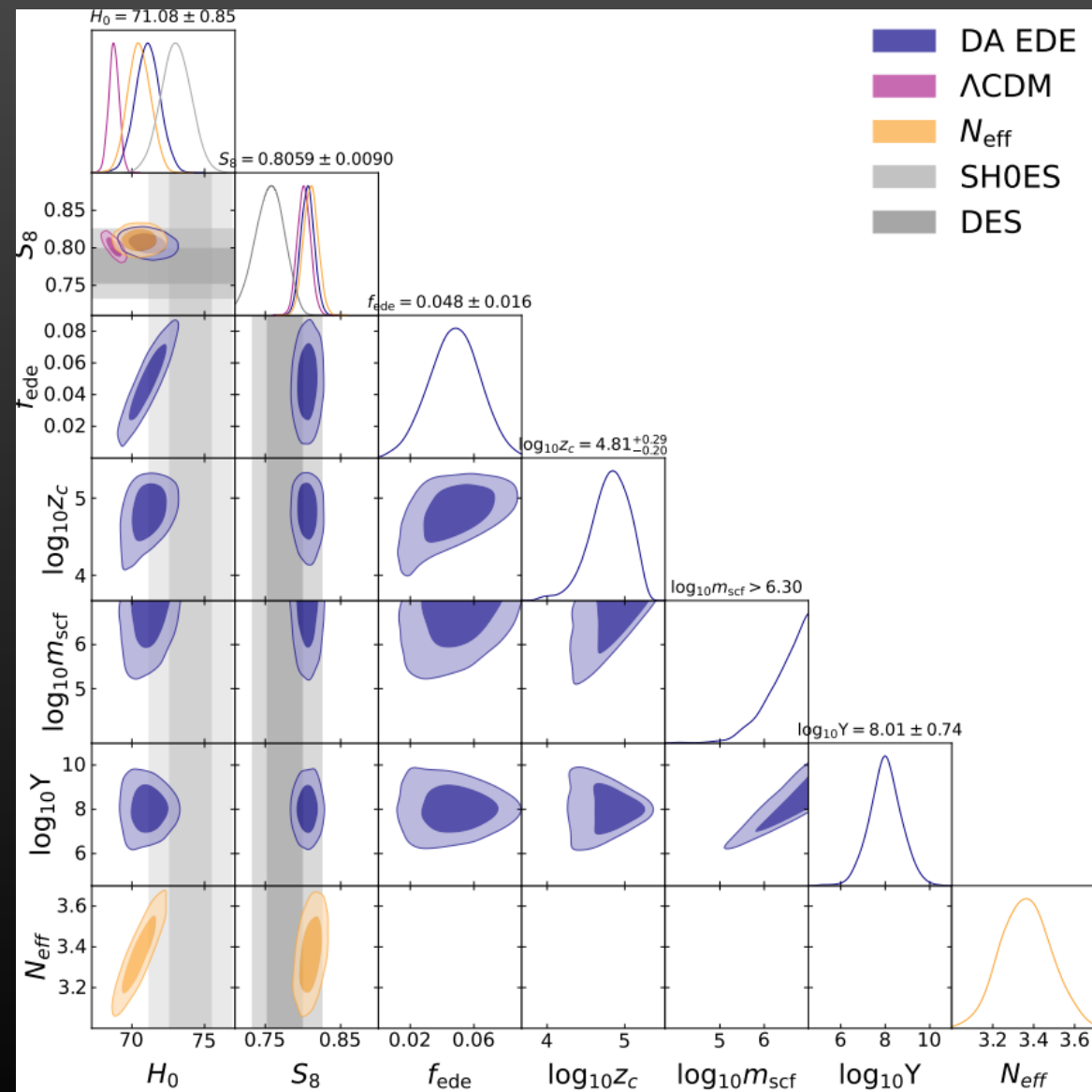
Berghaus, Karwal 2022

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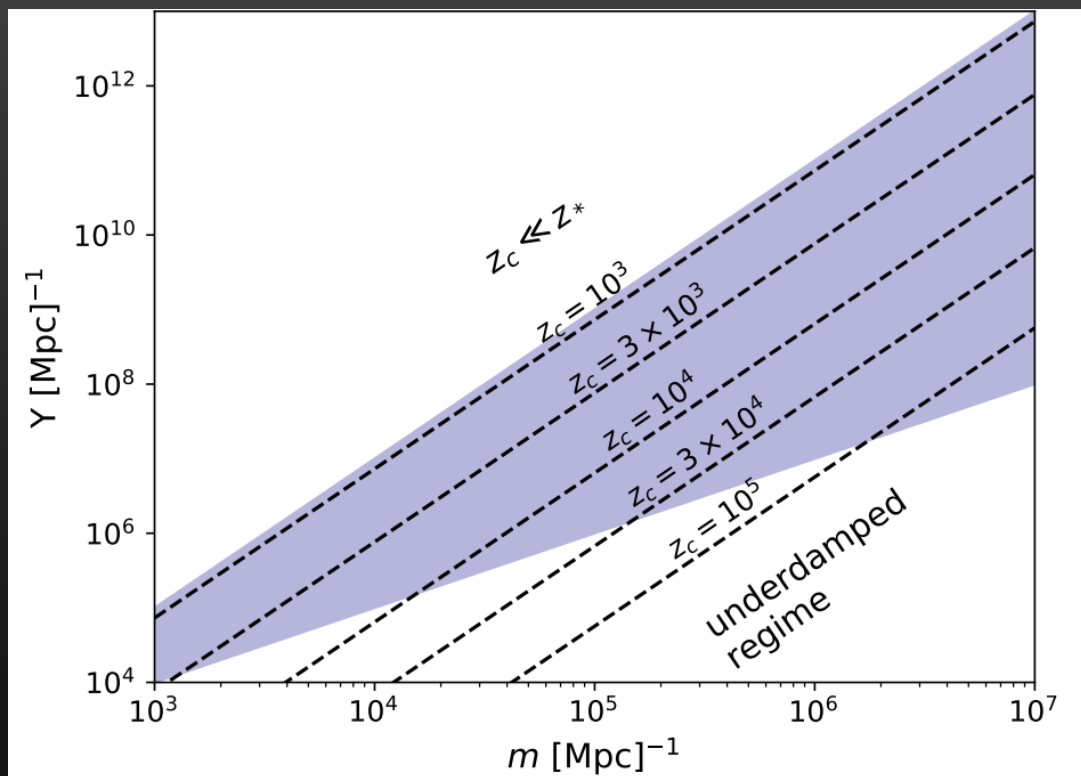
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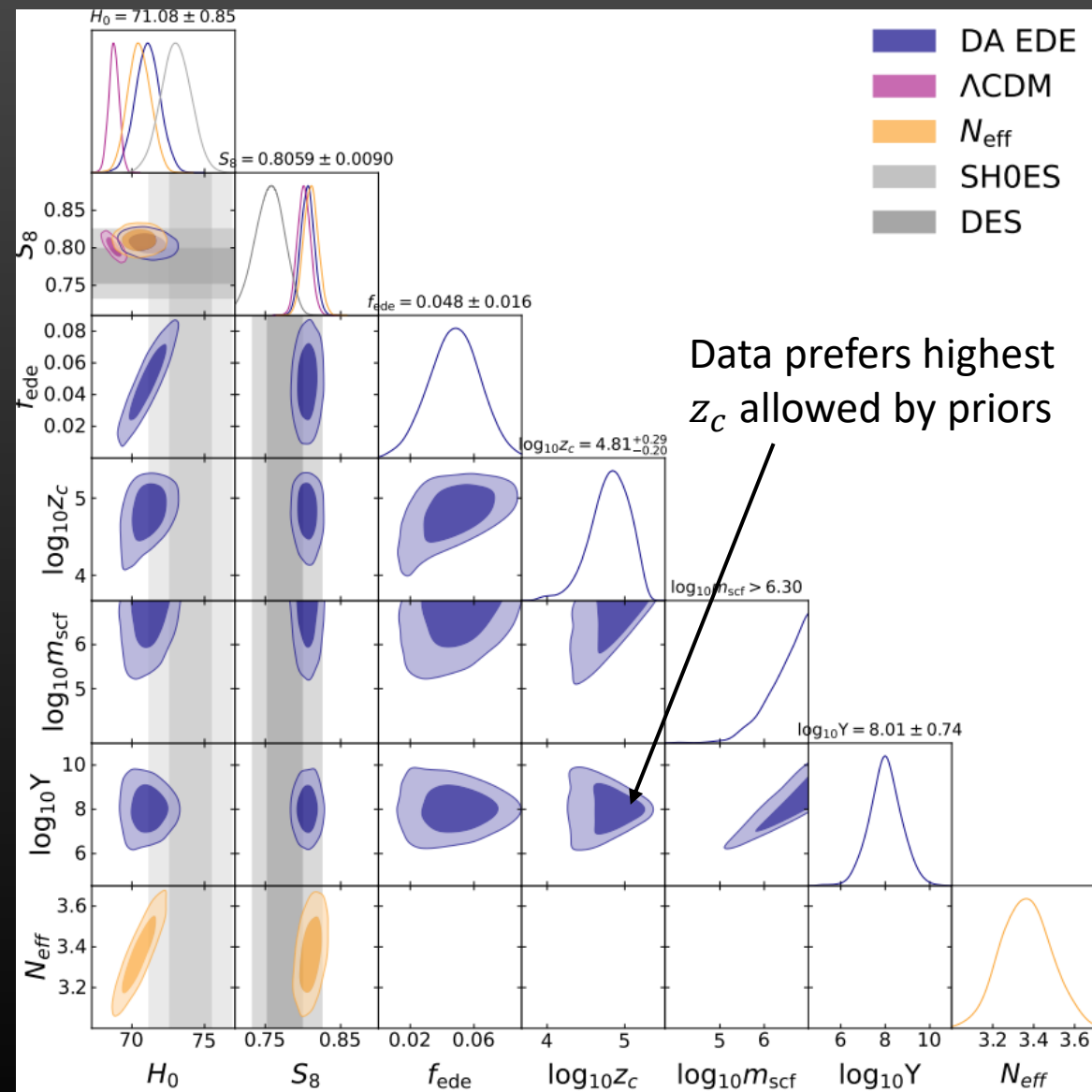
Berghaus, Karwal 2022

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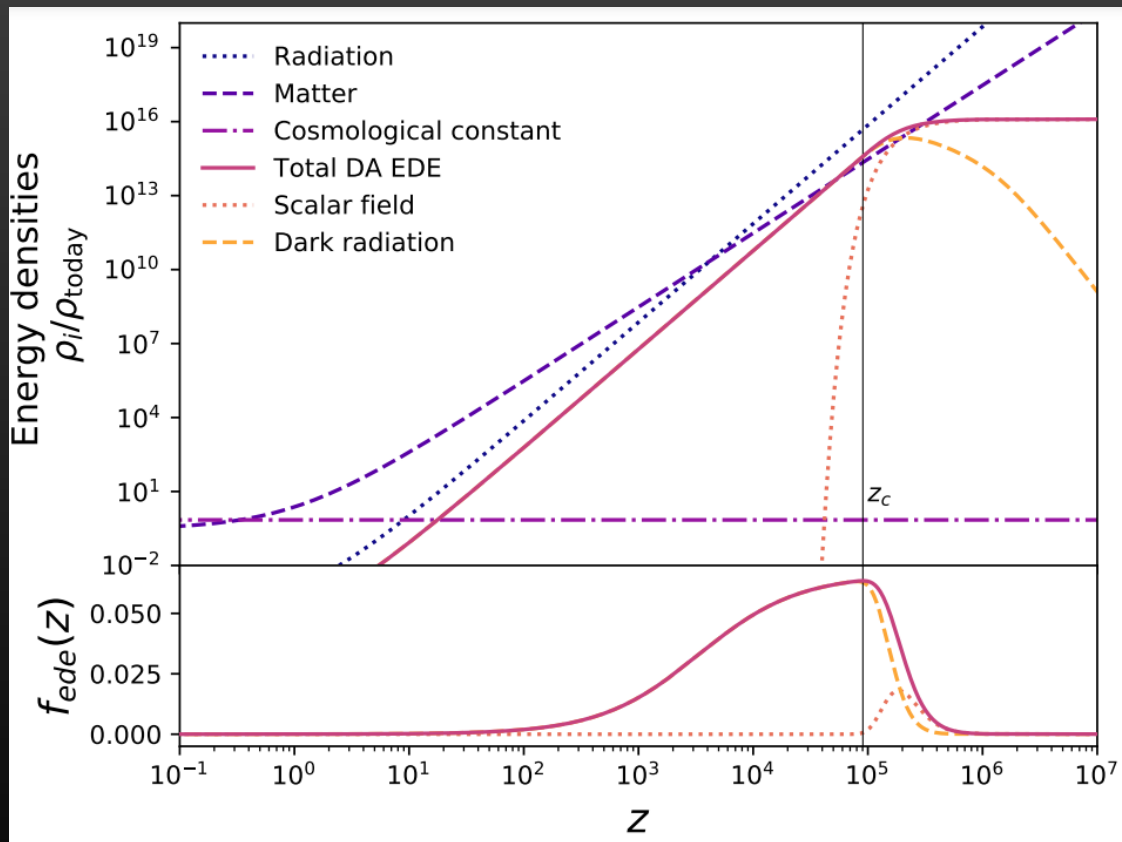
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Baseline + SH0ES + DES

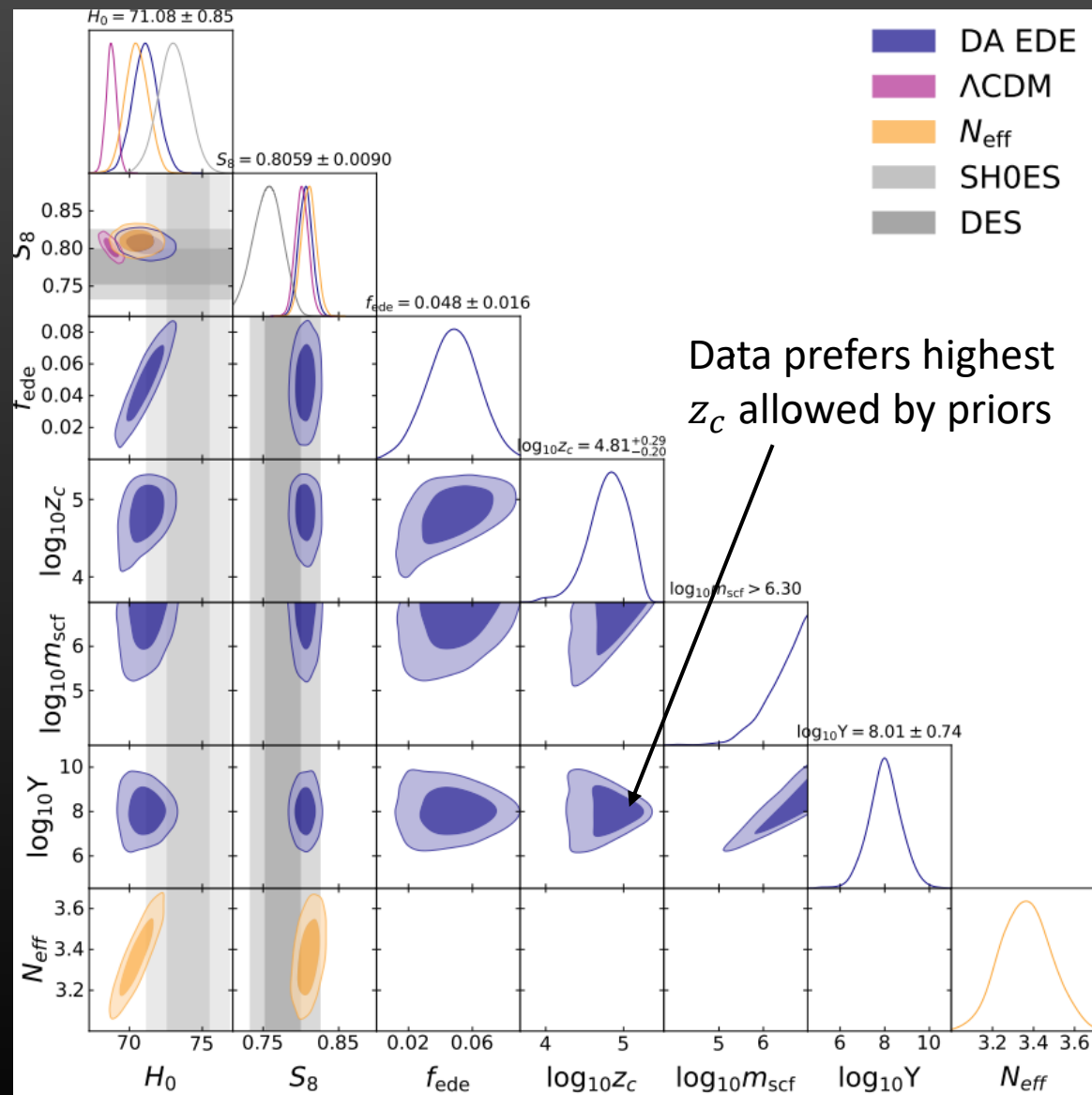


Berghaus, Karwal 2022

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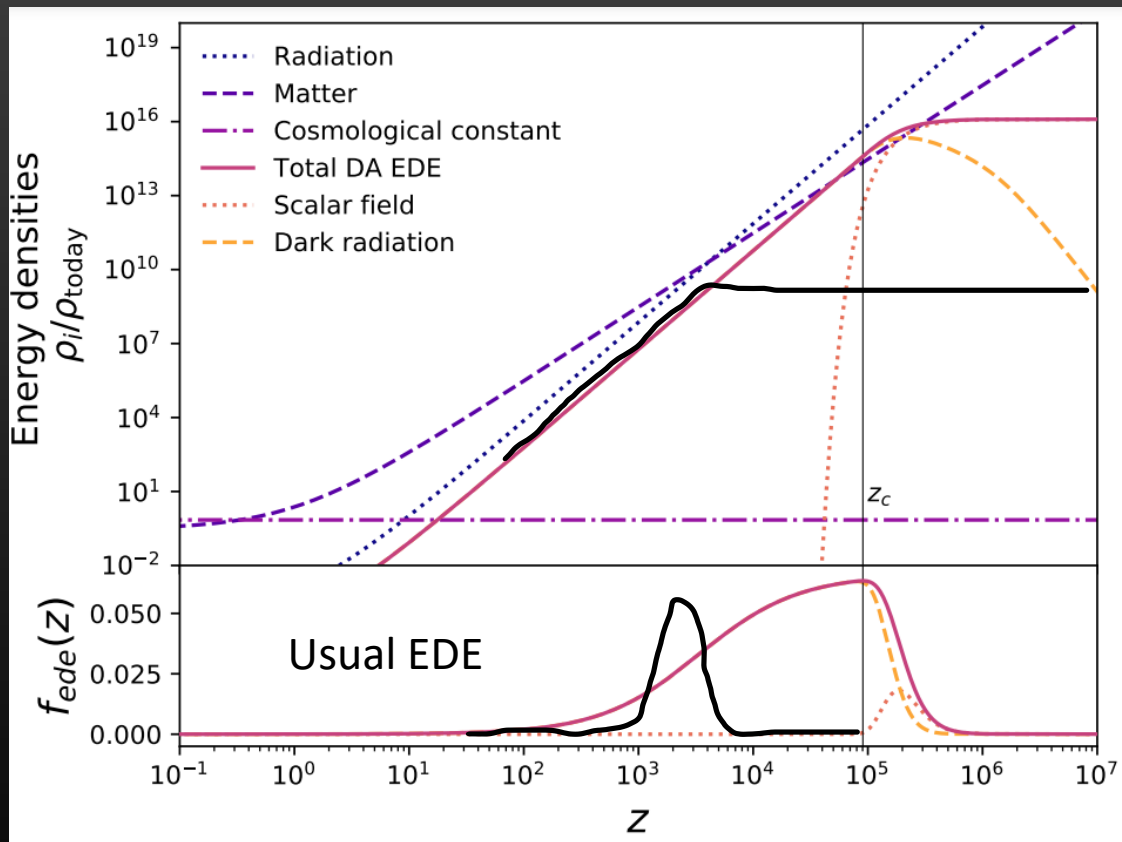


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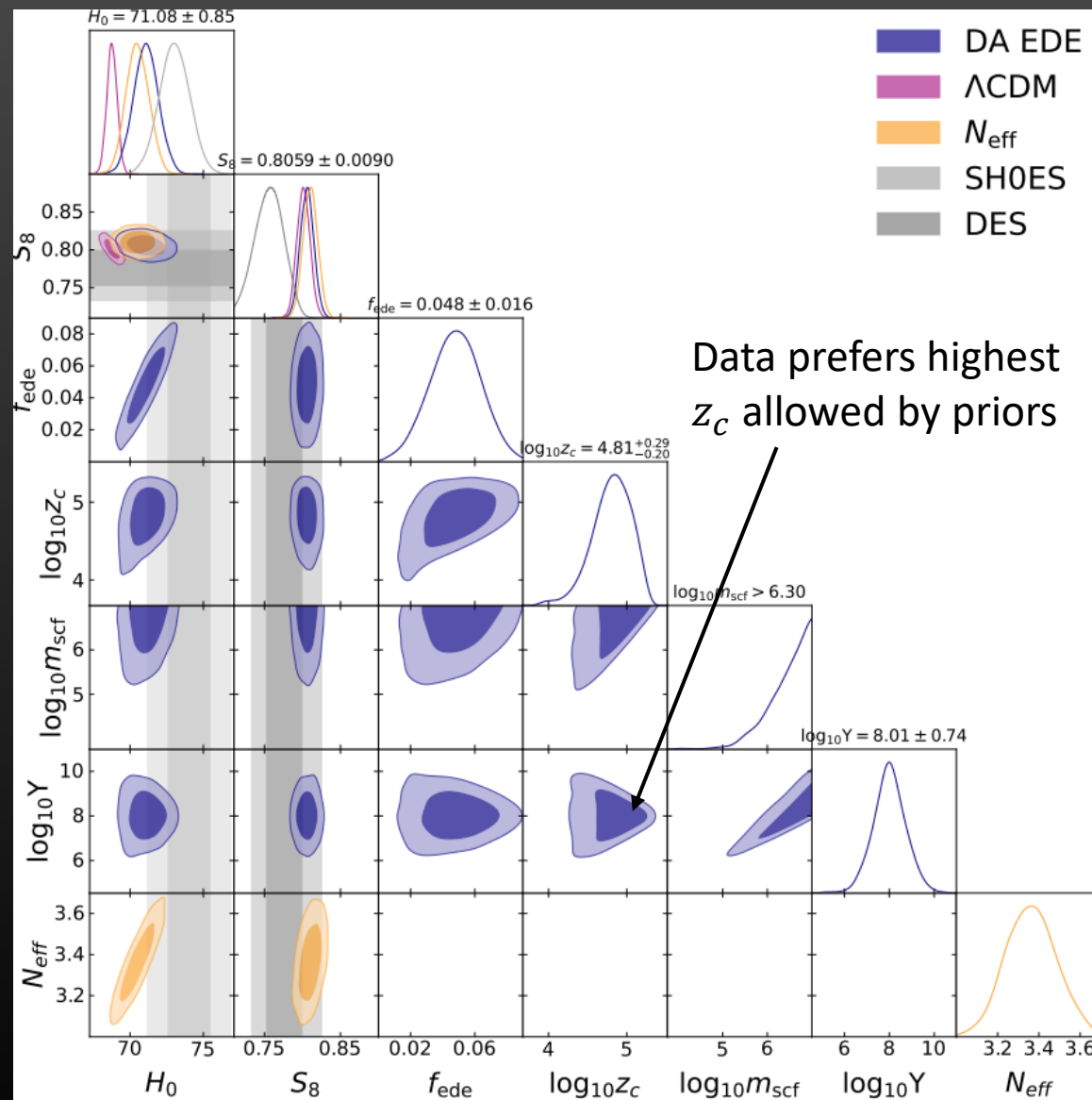


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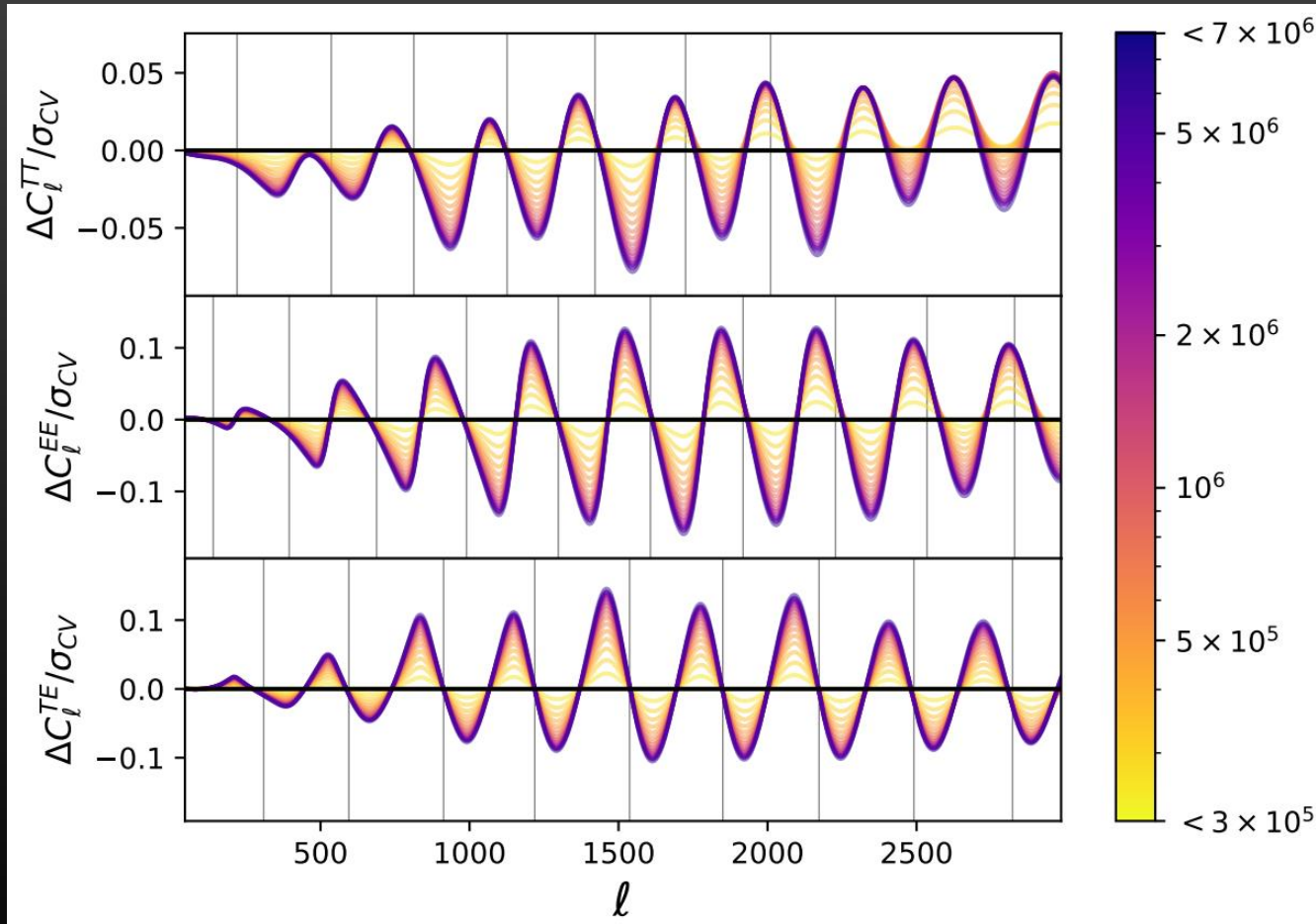


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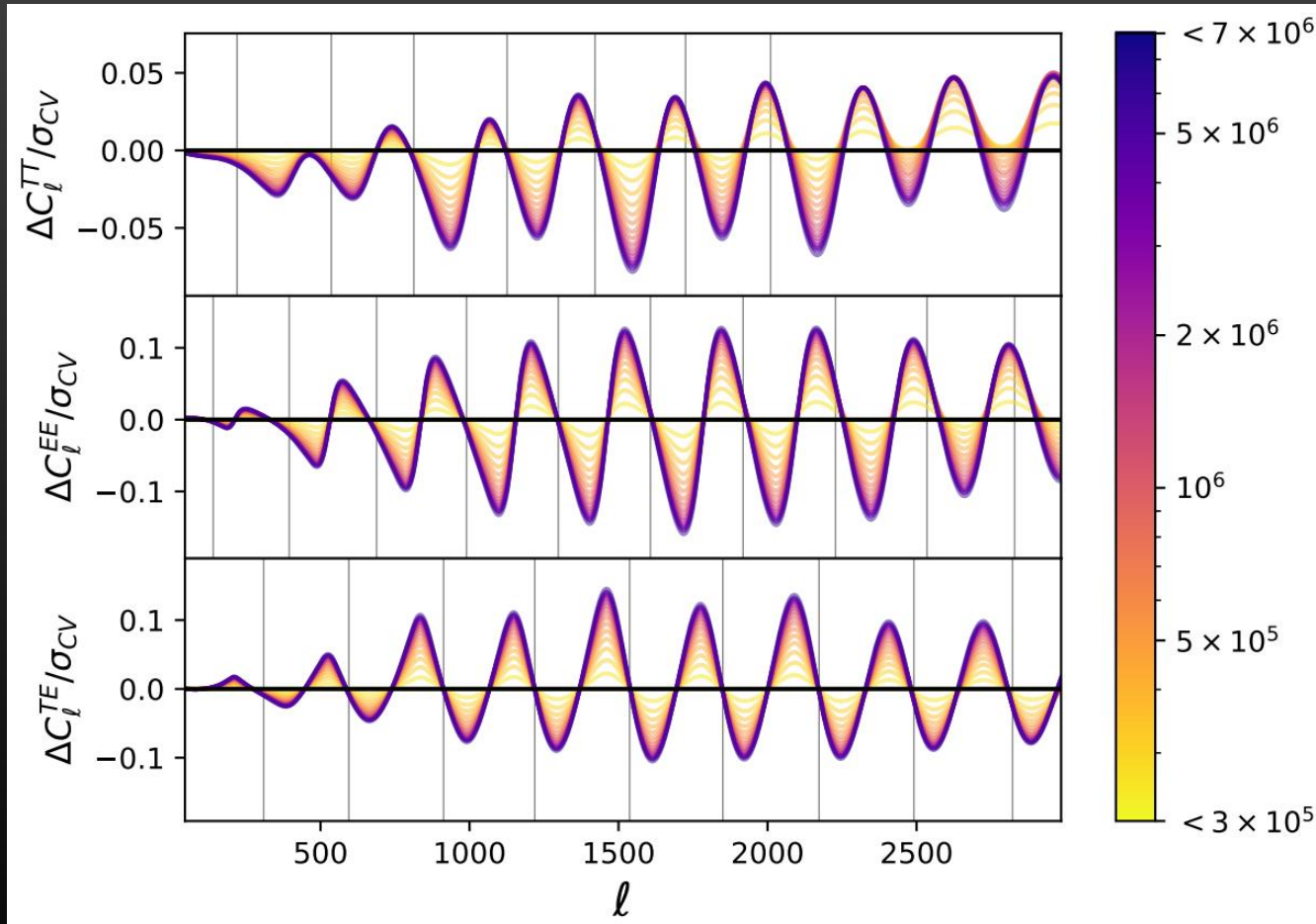
Going to higher redshifts



Berghaus, Karwal 2022

Data has sensitivity when $\frac{\Delta C_{el}^{XX}}{\sigma_{CV}} > 1$

Going to higher redshifts

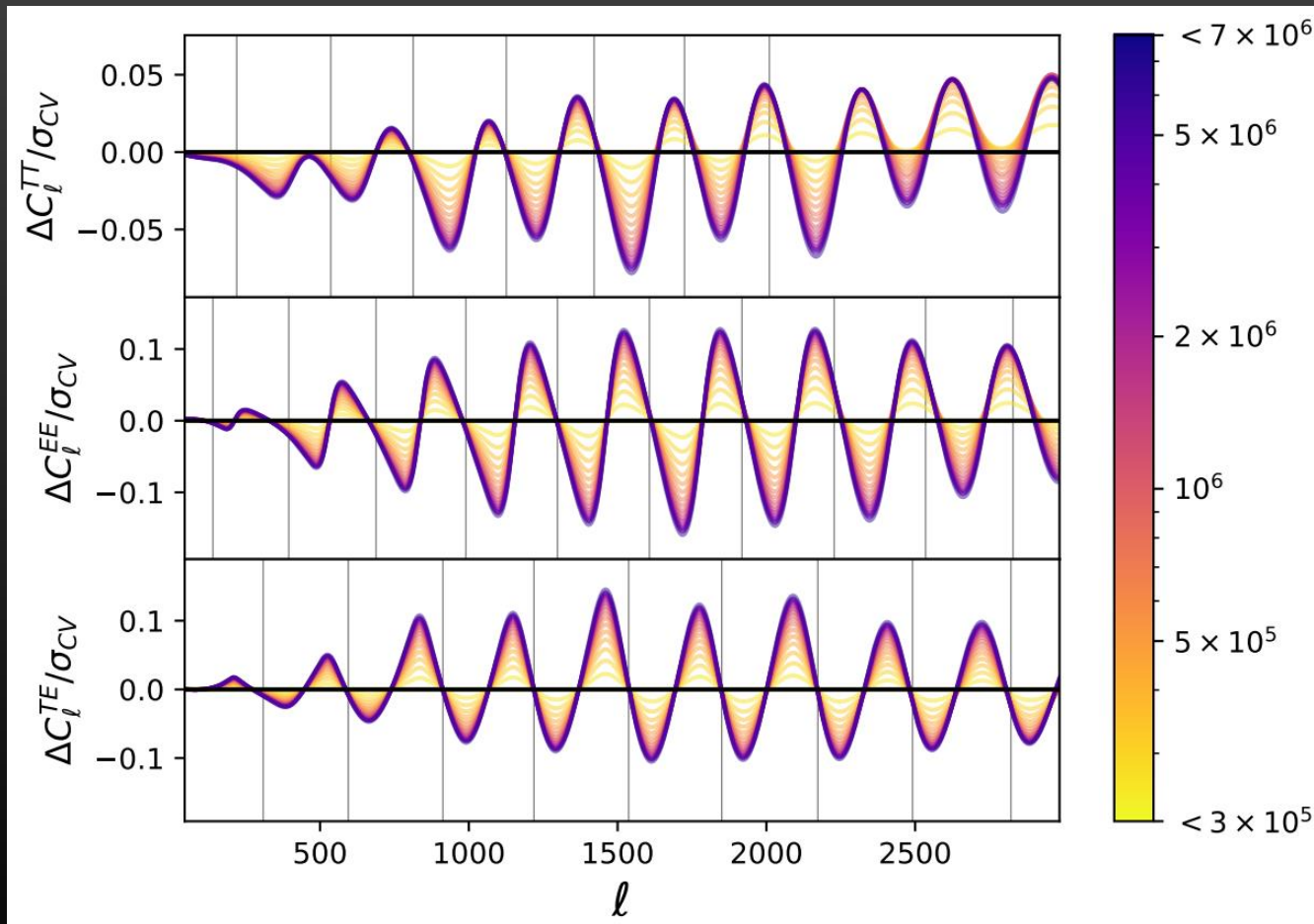


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Differences in predictions of the theory at larger z_c are not resolvable

Going to higher redshifts



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Thermal friction asymptotes to an extra-radiation solution

Why does data not prefer EDE thermal friction?

Answer must be in perturbations since background looks favorable

$$\delta\phi'' + 2aH\delta\phi' + (k^2 + a^2V''(\phi))\delta\phi = -\frac{h'\phi'}{2}$$

Regular EDE perturbations

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Axion perturbations with thermal friction

$$\delta'_{DR} = -\frac{2h'}{3} - \frac{4}{3}\theta_{DR} + \frac{2\Upsilon}{a\rho_{DR}}\delta\phi'\phi' - \frac{\Upsilon\phi'^2}{a\rho_{DR}}\delta_{DR}$$

Dark radiation perturbations sourced by thermal friction

$$\delta_{DR} \equiv \frac{\delta\rho_{DR}}{\rho_{DR}} \quad \theta_{DR} \equiv ik^j v_j$$

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Dark radiation perturbations sourced by thermal friction

$$\delta_{DR} \equiv \frac{\delta\rho_{DR}}{\rho_{DR}} \quad \theta_{DR} \equiv ik^j v_j$$

If injection occurs early the axion decays away completely. Only radiation left without friction terms.

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If injection occurs around $z_c \sim z_{eq}$ then $\delta\phi \ll \delta\rho_{DR}$ shortly after z_c .

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Dark radiation perturbations sourced by thermal friction

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Sourcing of DR
smoothes anisotropies

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Conclusions

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- Data has no sensitivity to step-like transitions at redshifts $z > 10^5$
- Data disfavors new physics with smoothed anisotropies

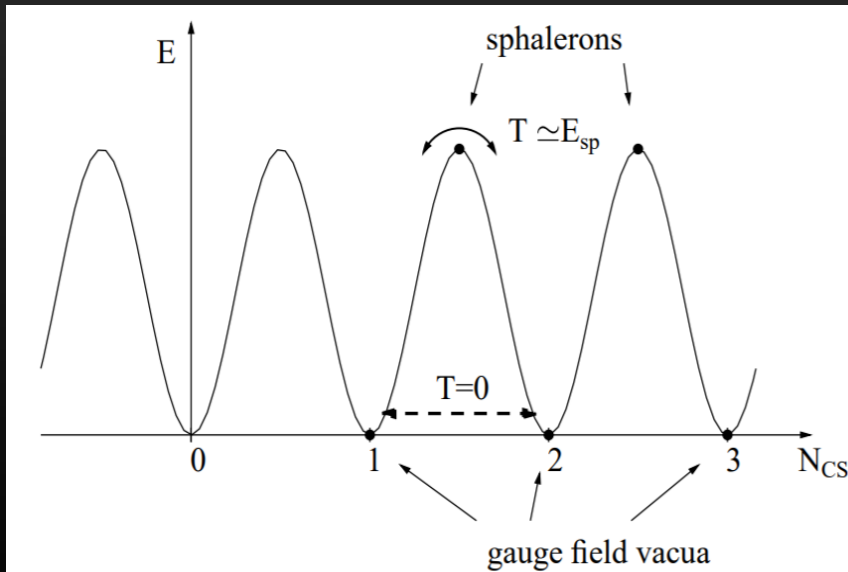
Thank You

Dissipative Axion

$$\frac{\partial L}{\partial \phi} - \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} = 0$$

- Couple scalar field to light degrees of freedom $L_{\text{int}} = -\phi \frac{\alpha}{16\pi f} \tilde{G} G$

$$\ddot{\phi} + 3H\dot{\phi} + V' = - \left\langle \frac{\alpha}{16\pi f} \tilde{G} G \right\rangle_{\text{non-eq}}(\phi)$$



$$\left\langle \frac{\alpha}{16\pi f} \tilde{G} G \right\rangle_{\text{non-eq}}(\phi) \approx \cancel{m_{\text{eff}}^2} \phi + \gamma \dot{\phi} + O(\ddot{\phi})$$

Not allowed by symmetry