



How to Measure the Spin of Dark Matter in $e^+ e^- \rightarrow \gamma + X$

Sofie Nordahl Erner

IPPP @ Durham University

sofie.n.erner@durham.ac.uk

Bauer & Erner [arXiv:2308.09746]

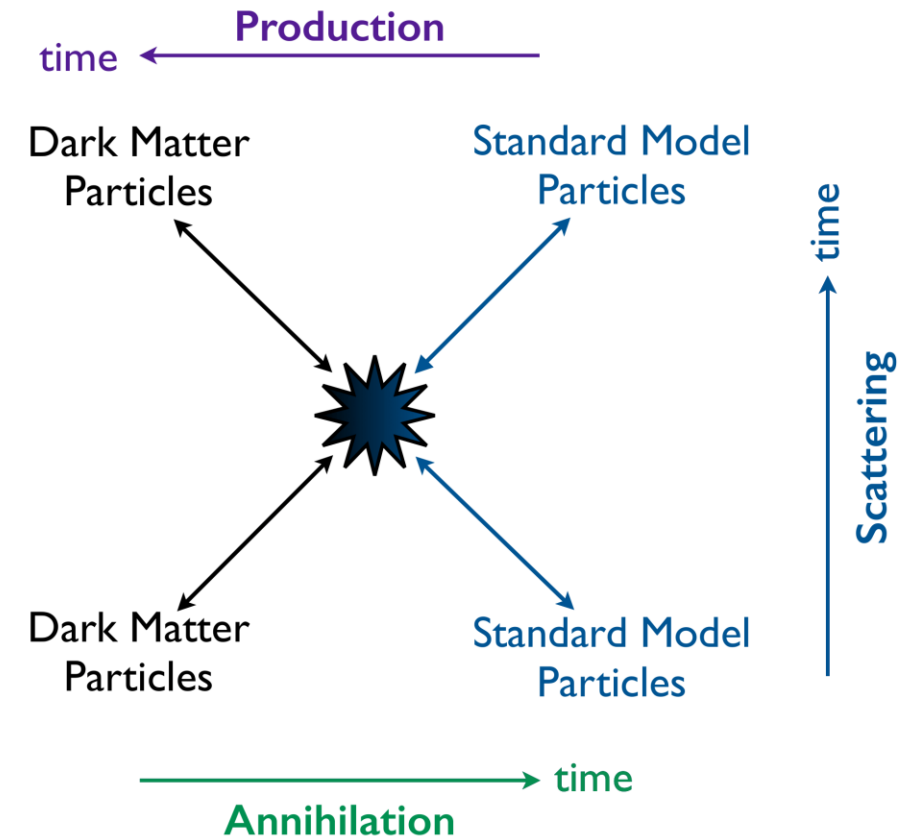


Introduction

With current collider experiments

- Can determine mass and coupling strength
- Cannot determine quantum numbers such as spin

⇒ Dark matter models are indistinguishable



<https://particleastro.brown.edu/dark-matter/>

Outline

Goal

Measure the Spin of Dark Matter

Process

Standard Model Background

Dark Matter Models

Currently in Detectors

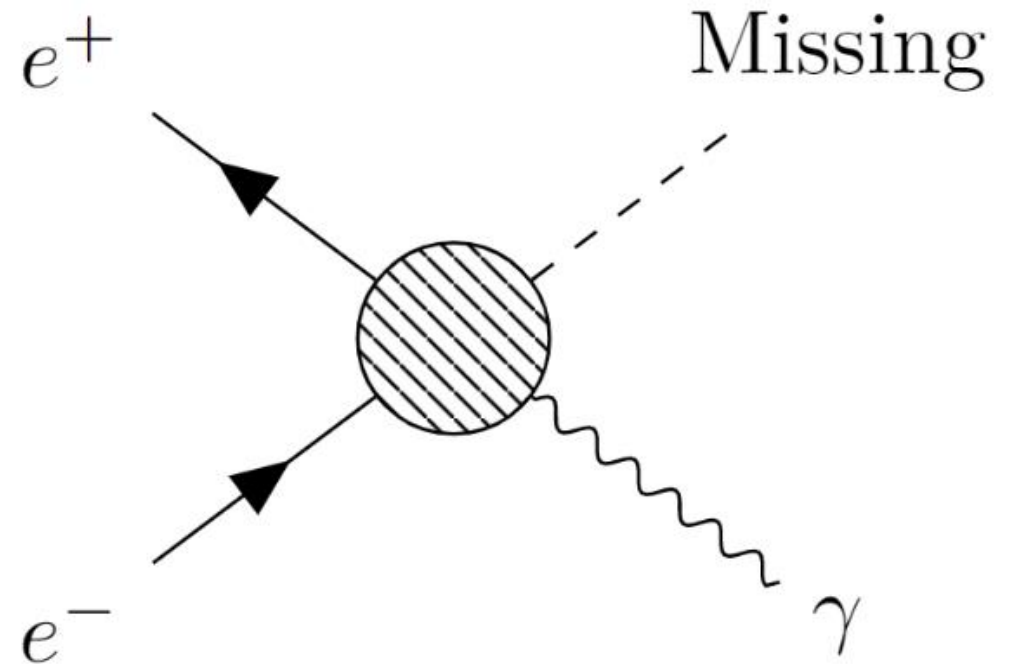
Polarised Beams

Conclusion

Process

- Mono-photon signal with missing energy

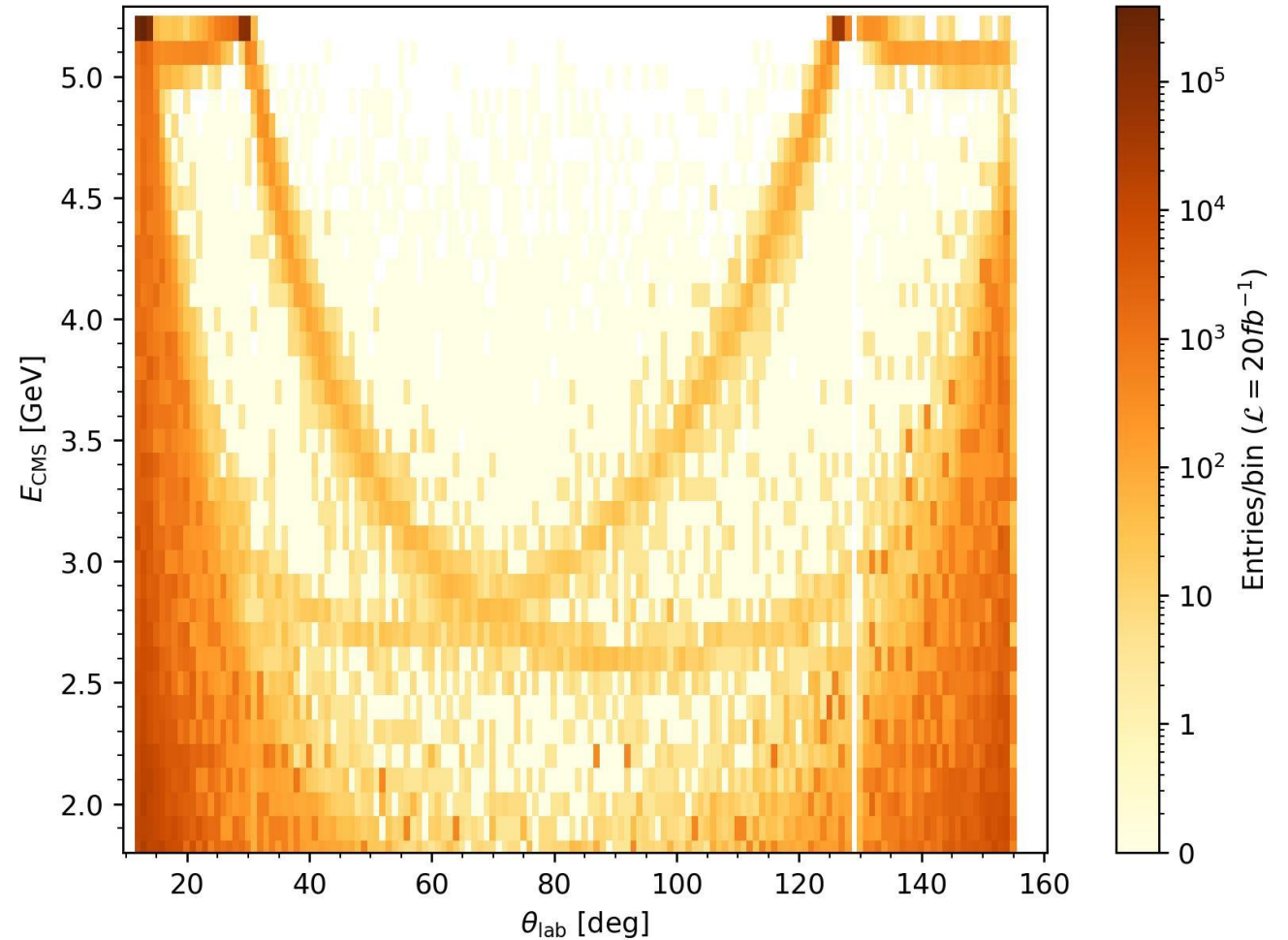
- Collider: Belle II
 - Beam Energies: 4 & 7 GeV
 - Asymmetric Angular Coverage



Standard Model Background

Processes:

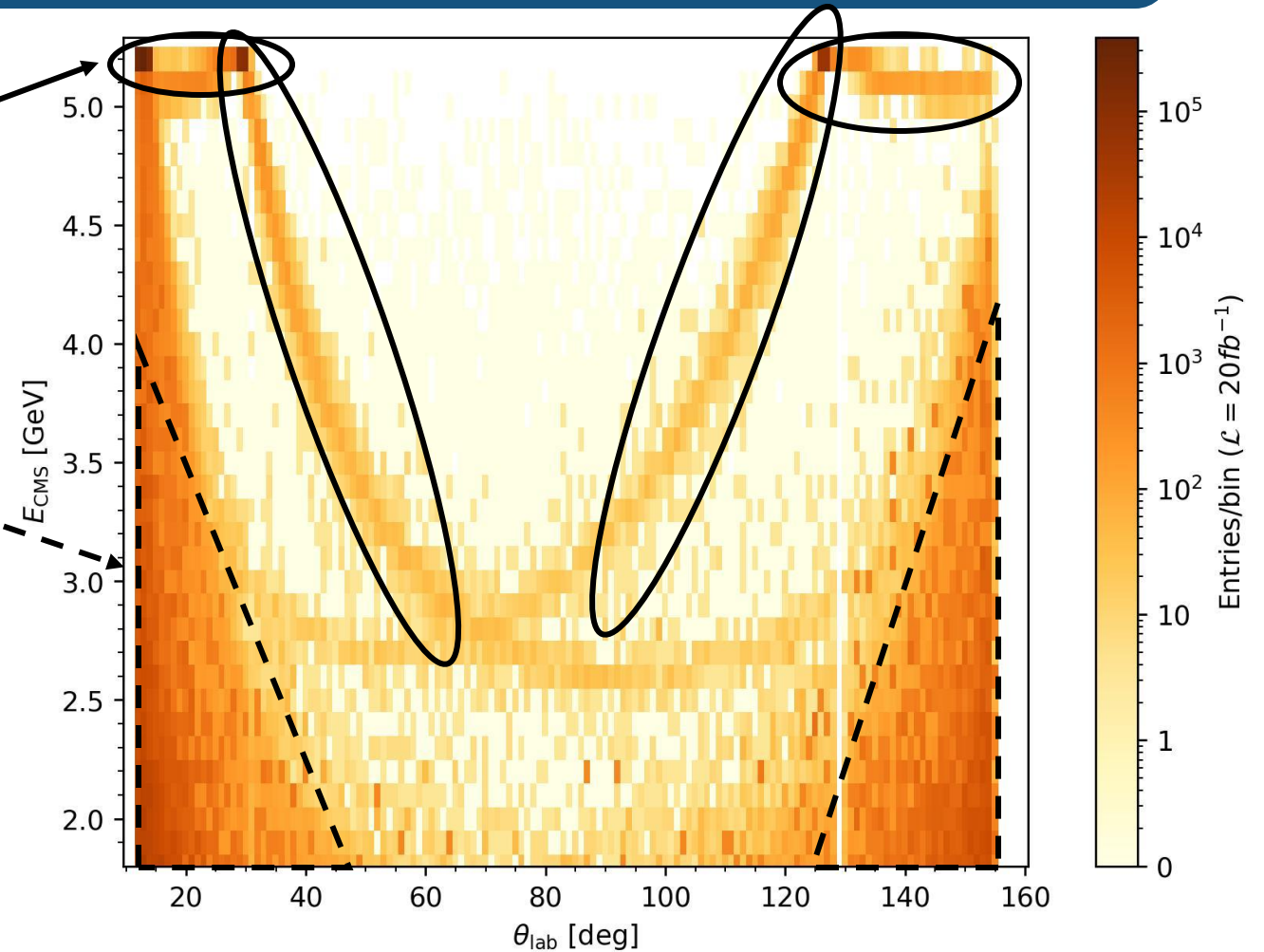
- $e^+ e^- \rightarrow$
 - $\gamma\gamma(\gamma)$
 - $e^+ e^- \gamma(\gamma)$
 - $\bar{\nu}\nu\gamma(\gamma)$
- All particles but one photon undetected
 - Asymmetric angular coverage
 - Gaps between barrel and endcaps



Standard Model Background

Processes:

- $e^+ e^- \rightarrow$
 - $\gamma\gamma(\gamma)$
 - $e^+ e^- \gamma(\gamma)$
 - $\bar{\nu}\nu\gamma(\gamma)$
- All particles but one photon undetected
 - Asymmetric angular coverage
 - Gaps between barrel and endcaps
- Neutrino background irreducible & everywhere

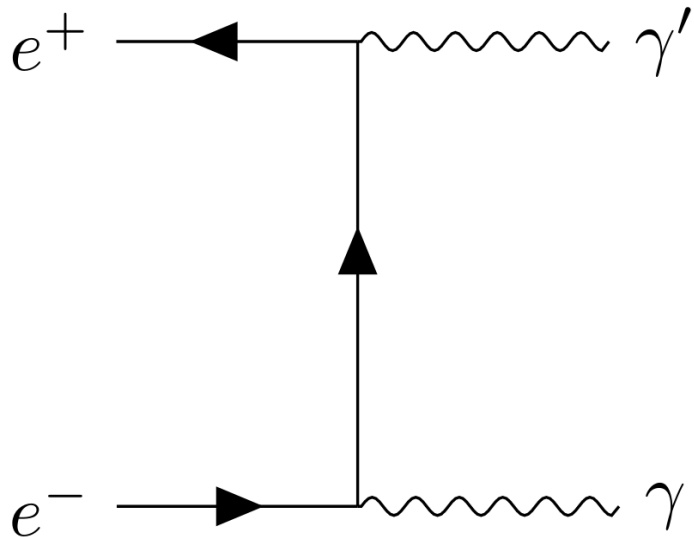


Dark Matter Models

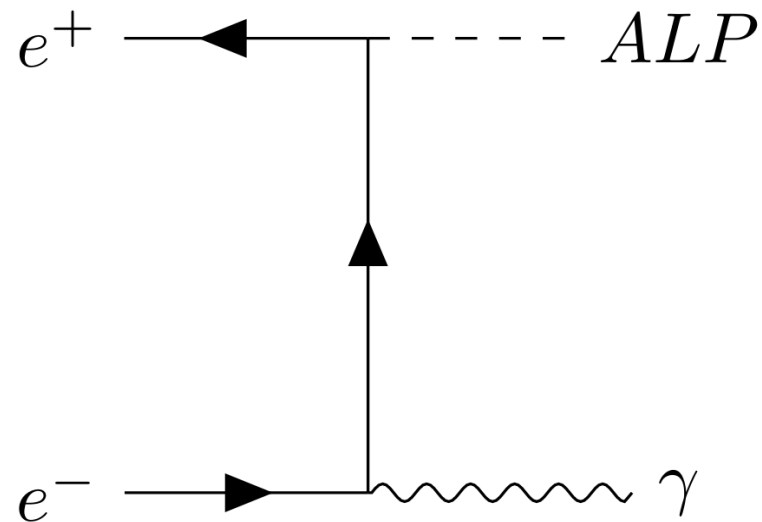
Type & Spin	Fermion Coupling	Photon Coupling
Axion-Like-Particles (ALPs)		
Pseudo-Scalar spin-0	$-\frac{c_{ff}}{f_a} a \sum_{\psi} m_{\psi} \bar{\psi} \gamma^5 \psi$	$-\frac{1}{4} g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$ $g_{a\gamma\gamma} = 4/f_a (C_{BB} c_{\theta}^2 + C_{WW} s_{\theta}^2)$
Dark Photon		
Vector spin-1	$-g_X X_{\mu} \sum_{e,\mu,\tau} (\bar{L} \gamma^{\mu} L + \bar{I} \gamma^{\mu} I)$	NA

Dark Matter Models

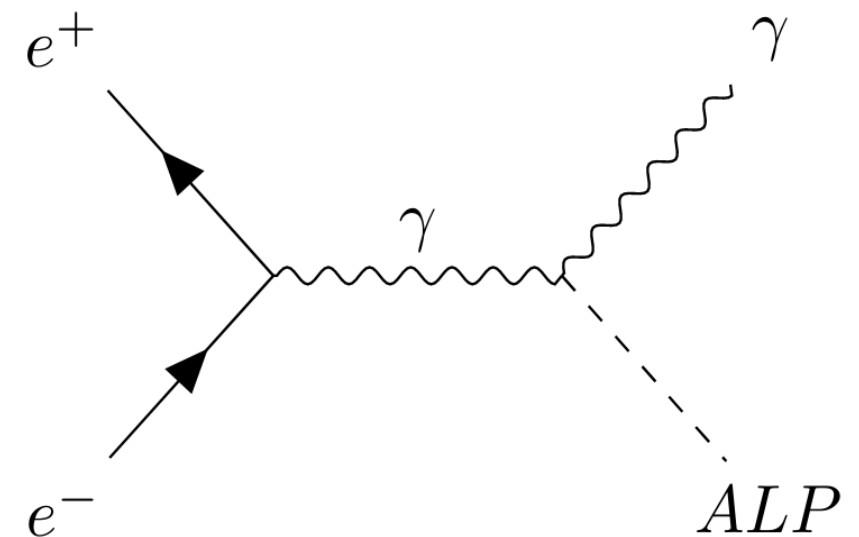
Dark Photon



ALP t/u-channel

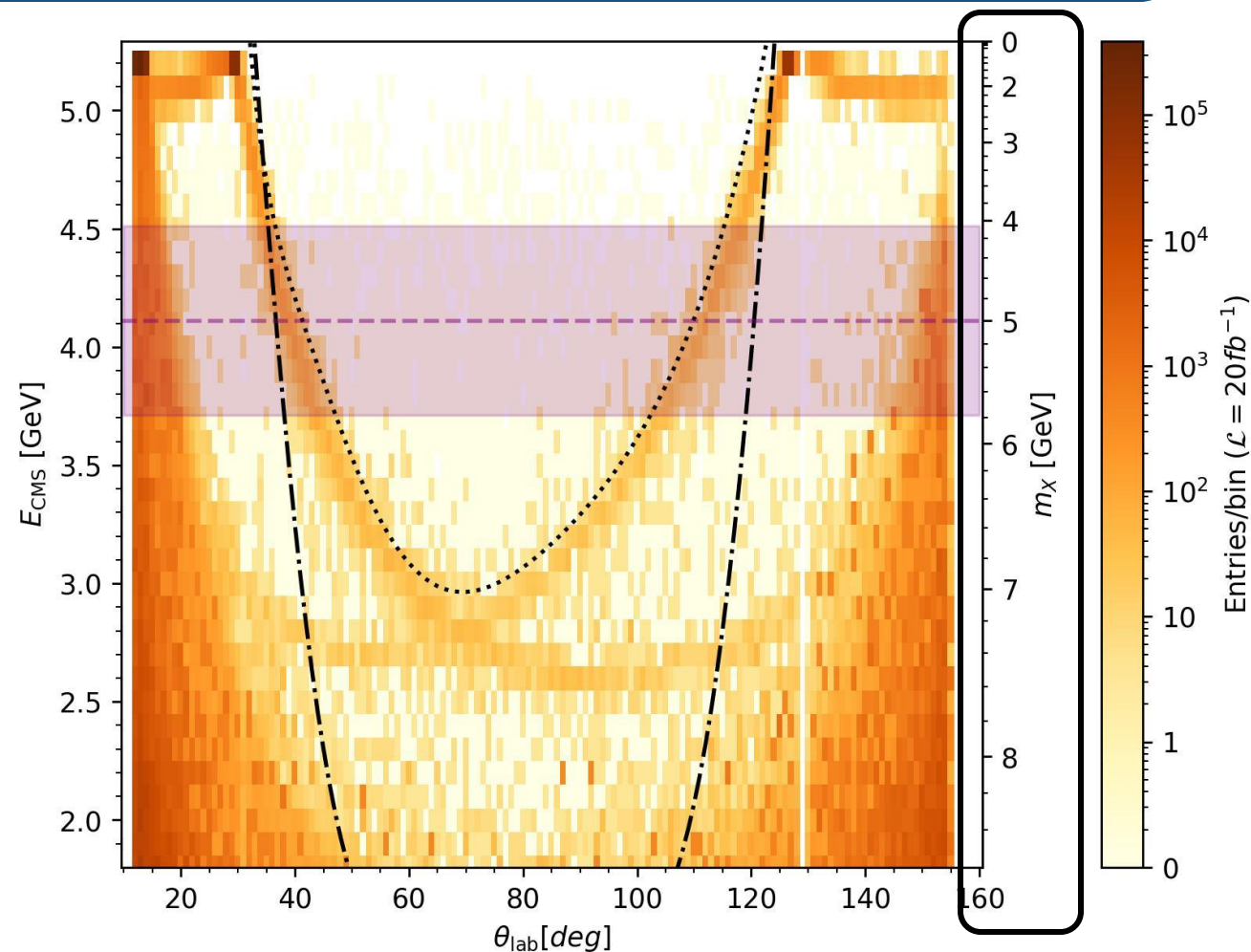


ALP s-channel

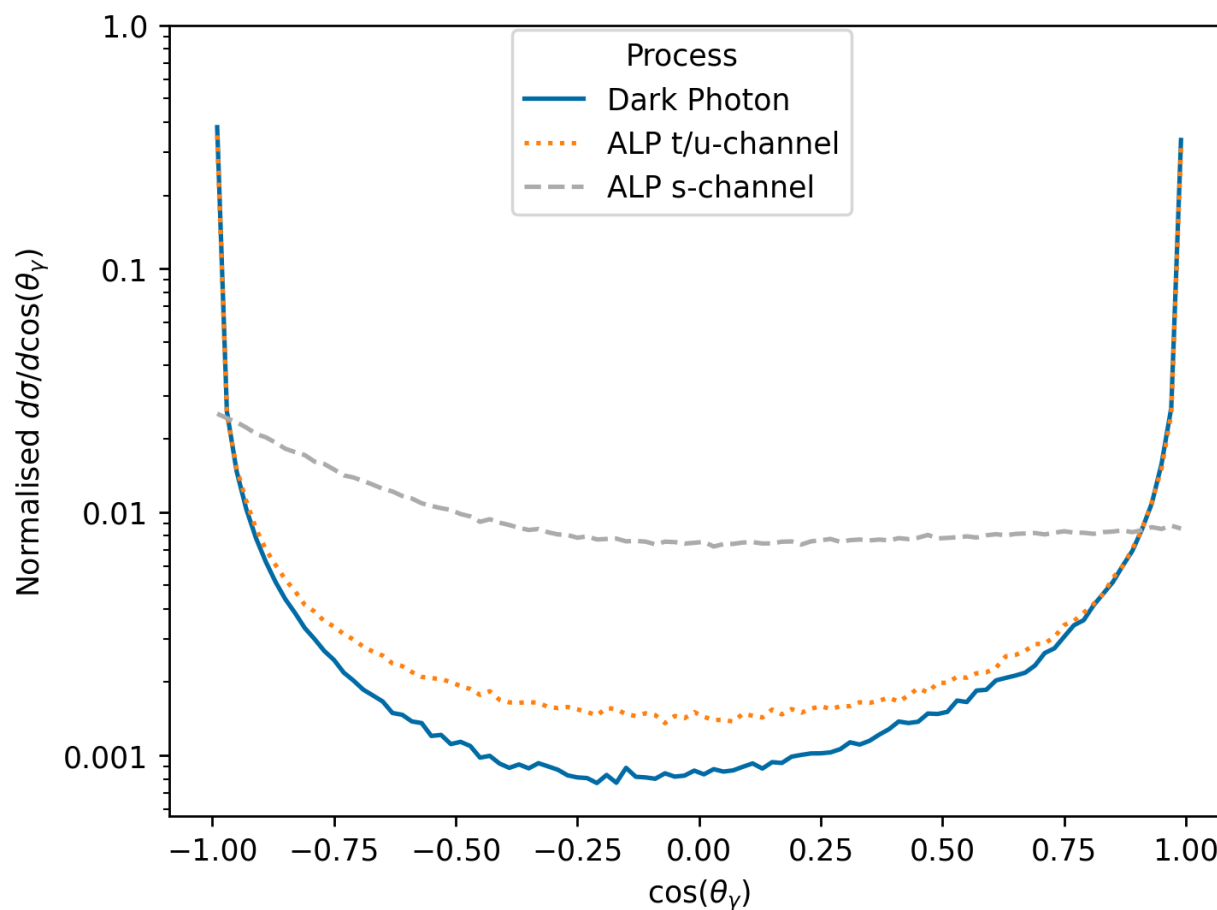


Currently in Detectors

- $2 \rightarrow 2$ scattering
⇒ Final photon E_{CMS} function of dark matter mass
 - Not including initial-state-radiation (ISR)
- Same for all DM contributions



Currently in Detectors



Angular Distribution

- Difference between s- and t/u-channel contributions
 - Distinguish between ALP contributions
 - Cannot distinguish between ALP and Dark Photon t/u-channel contributions

Polarised Beams



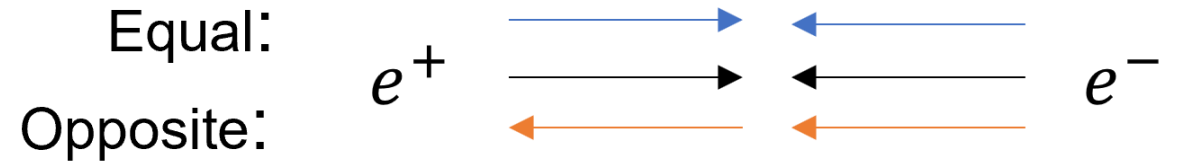
Planned: 70% electron beam polarisation at Belle II



Interest: 100% polarise both beams

Polarised Beams

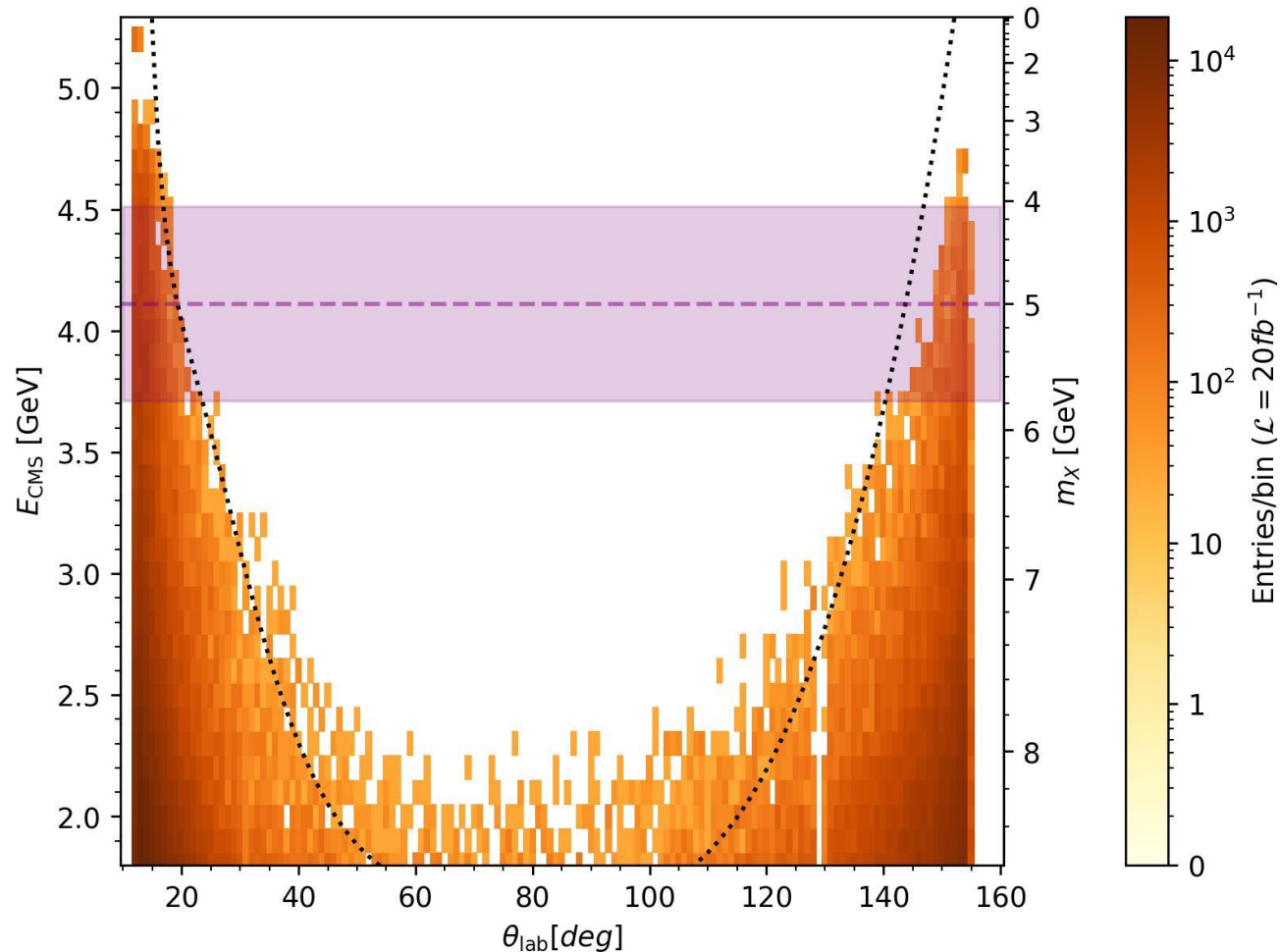
Allows for "selection" of processes as they require specific spin structures



Polarisation	Opposite	Equal
DM Signal present	Dark photon & ALP s-channel	ALP t/u-channel

Distinguish t/u-channel Dark Photon and ALPs contributions

Polarised Beams

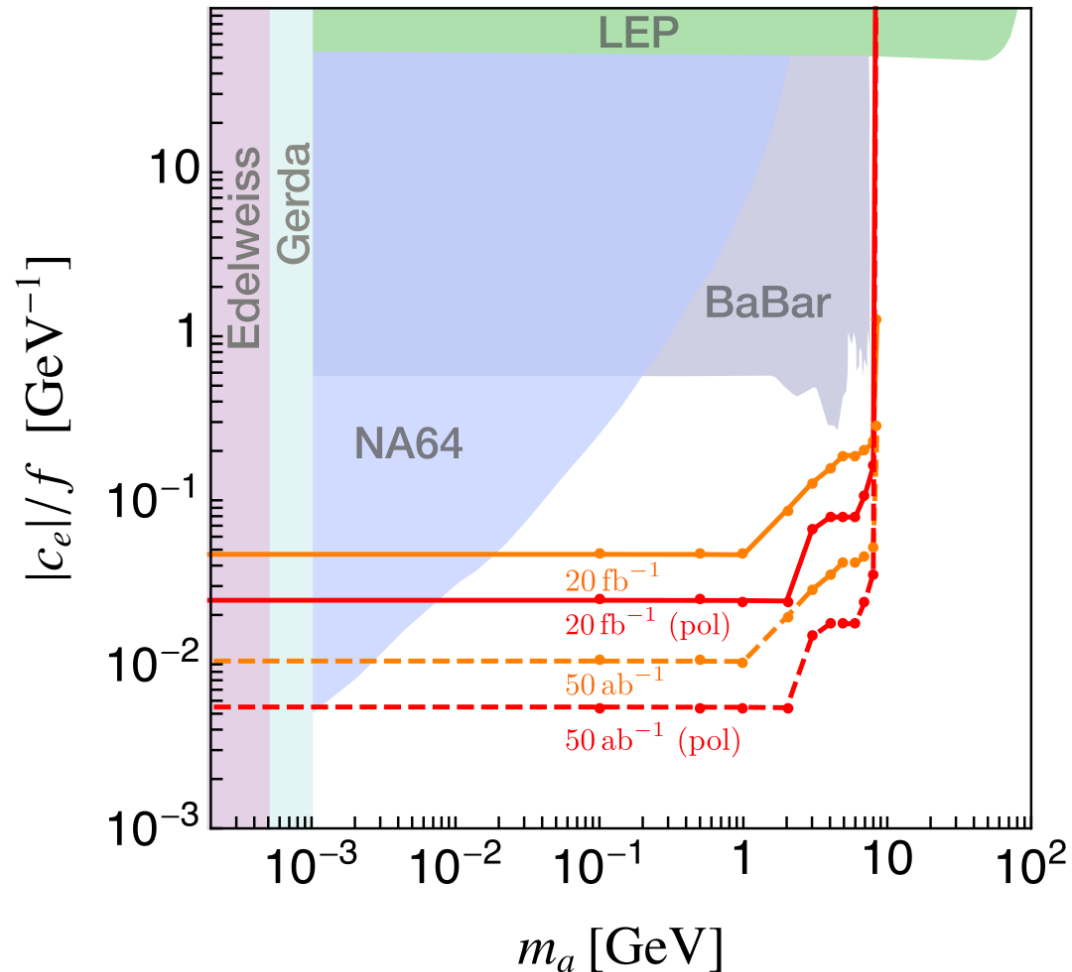


Reduced SM background for
equally polarised beams

- Photon-only final states drop out

⇒ Provide significantly better
exclusion limits on ALP-fermion
coupling

Polarised Beams



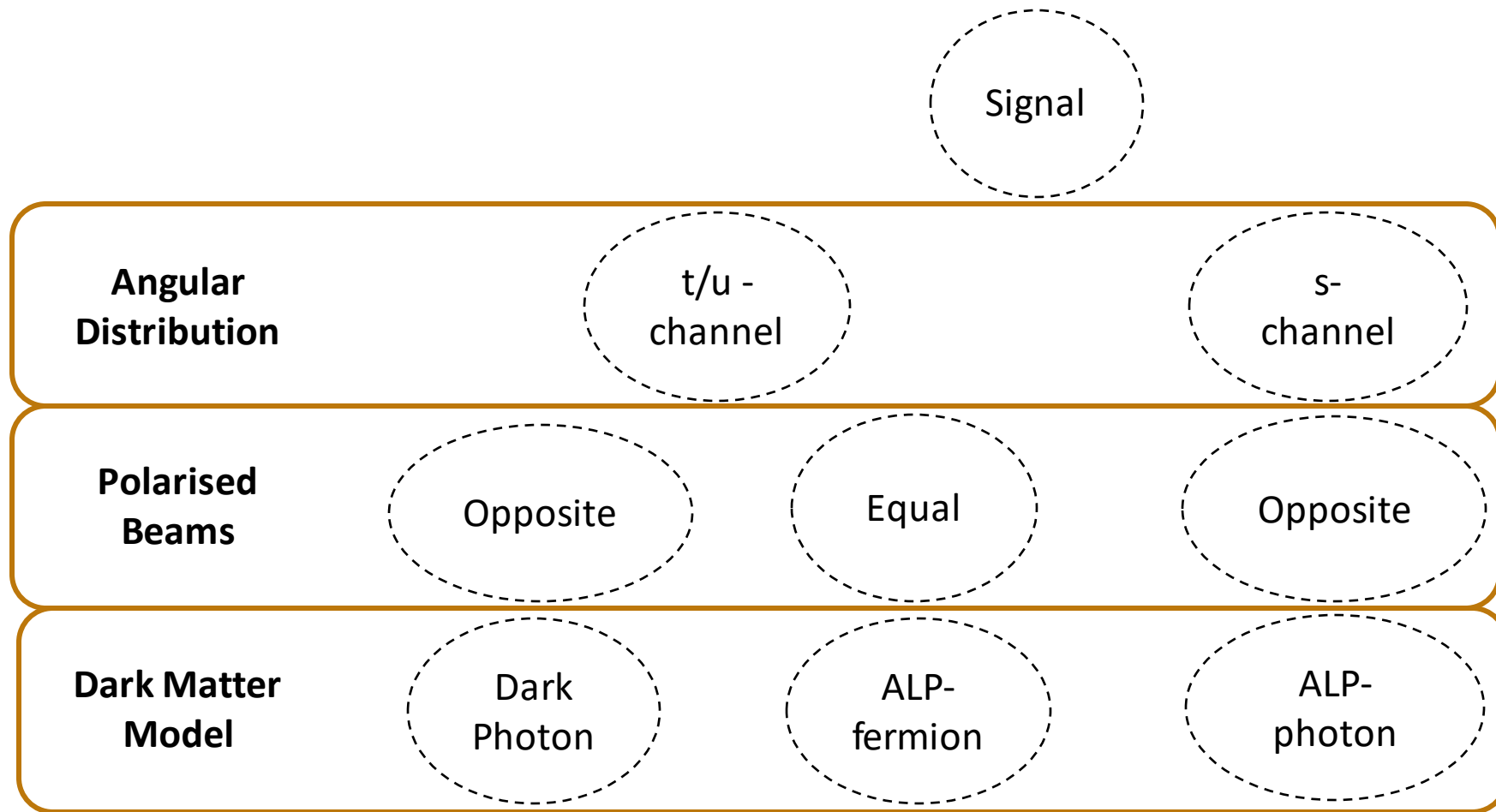
Reduced SM background for
equally polarised beams

⇒ Provide significantly better
exclusion limits on ALP-fermion
coupling

Orange: Unpolarised beams

Red: Beams with equal polarisation

Conclusion



Conclusion

- Separate the dark matter contributions
 - ⇒ Distinguish between spin-0 ALP and spin-1 Dark Photon
 - ⇒ **Measure the Spin of Dark Matter**
- Reduced SM background for equally polarised beams
 - ⇒ Provide significantly better exclusion limits on ALP-fermion coupling

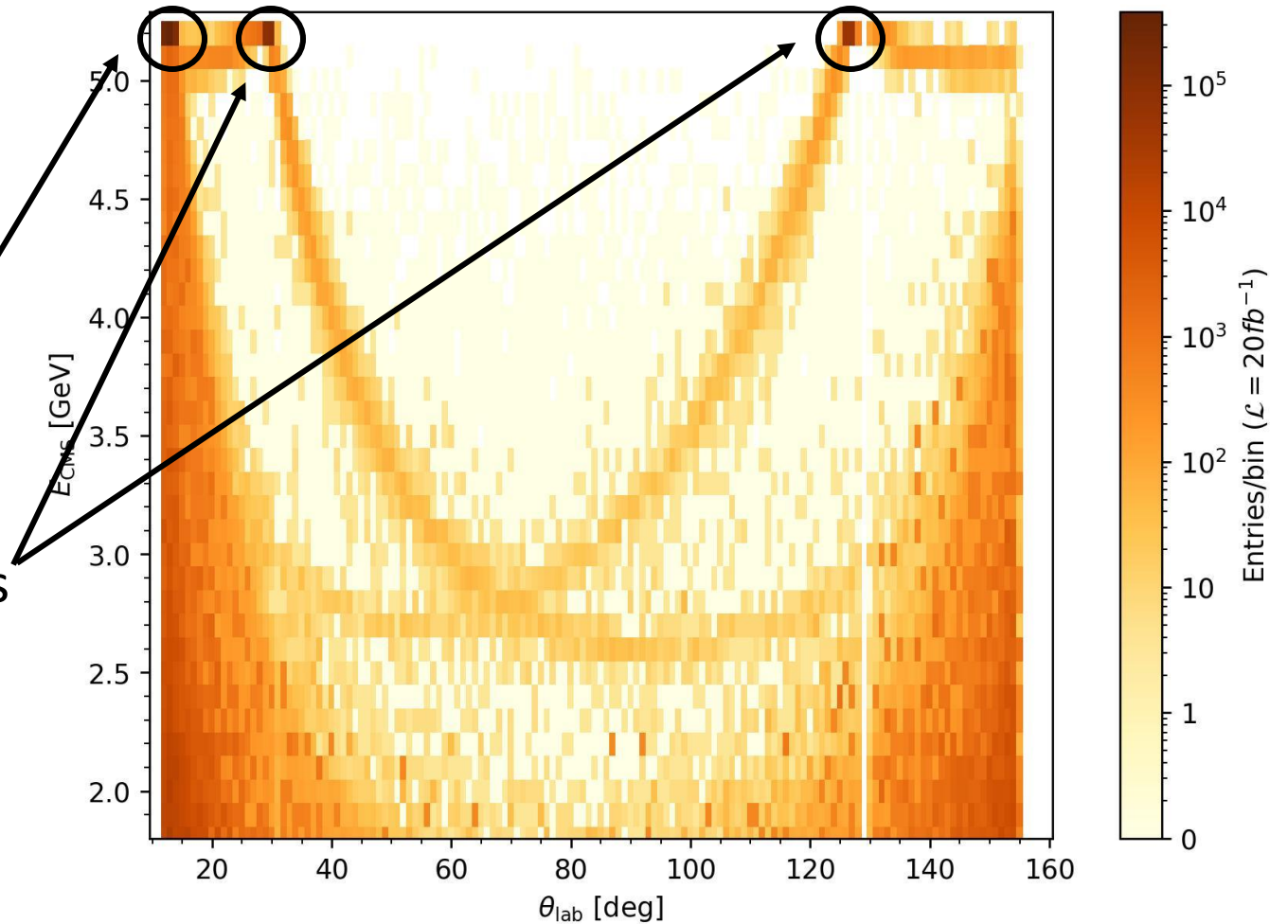
Thanks for listening!

Backup Slides

Standard Model Background

Processes:

- $e^+ e^- \rightarrow$
 - $\gamma\gamma(\gamma)$
- Asymmetric angular coverage
- Gaps between barrel and endcaps

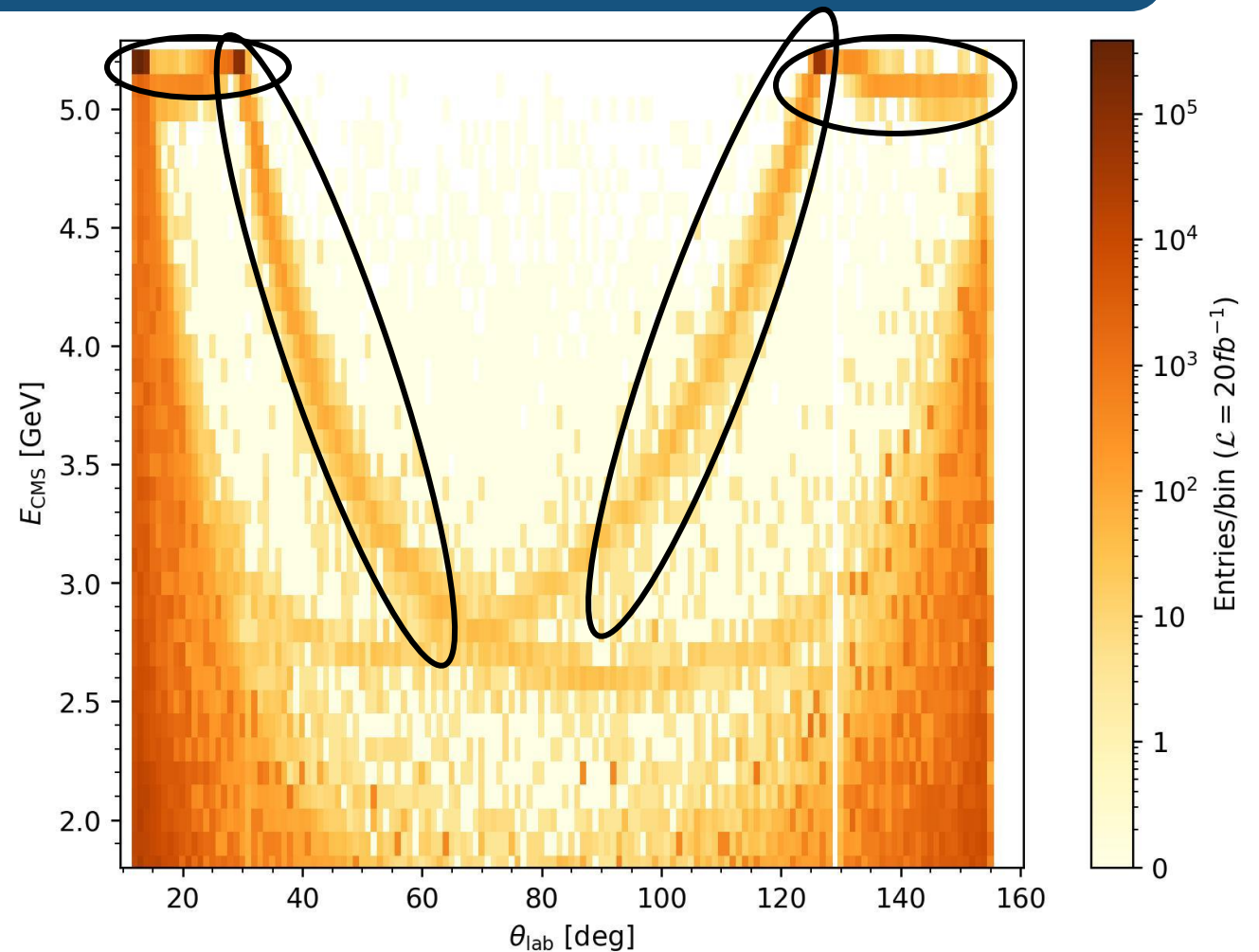


Standard Model Background

Processes:

- $e^+ e^- \rightarrow$
 - $\gamma\gamma(\gamma)$

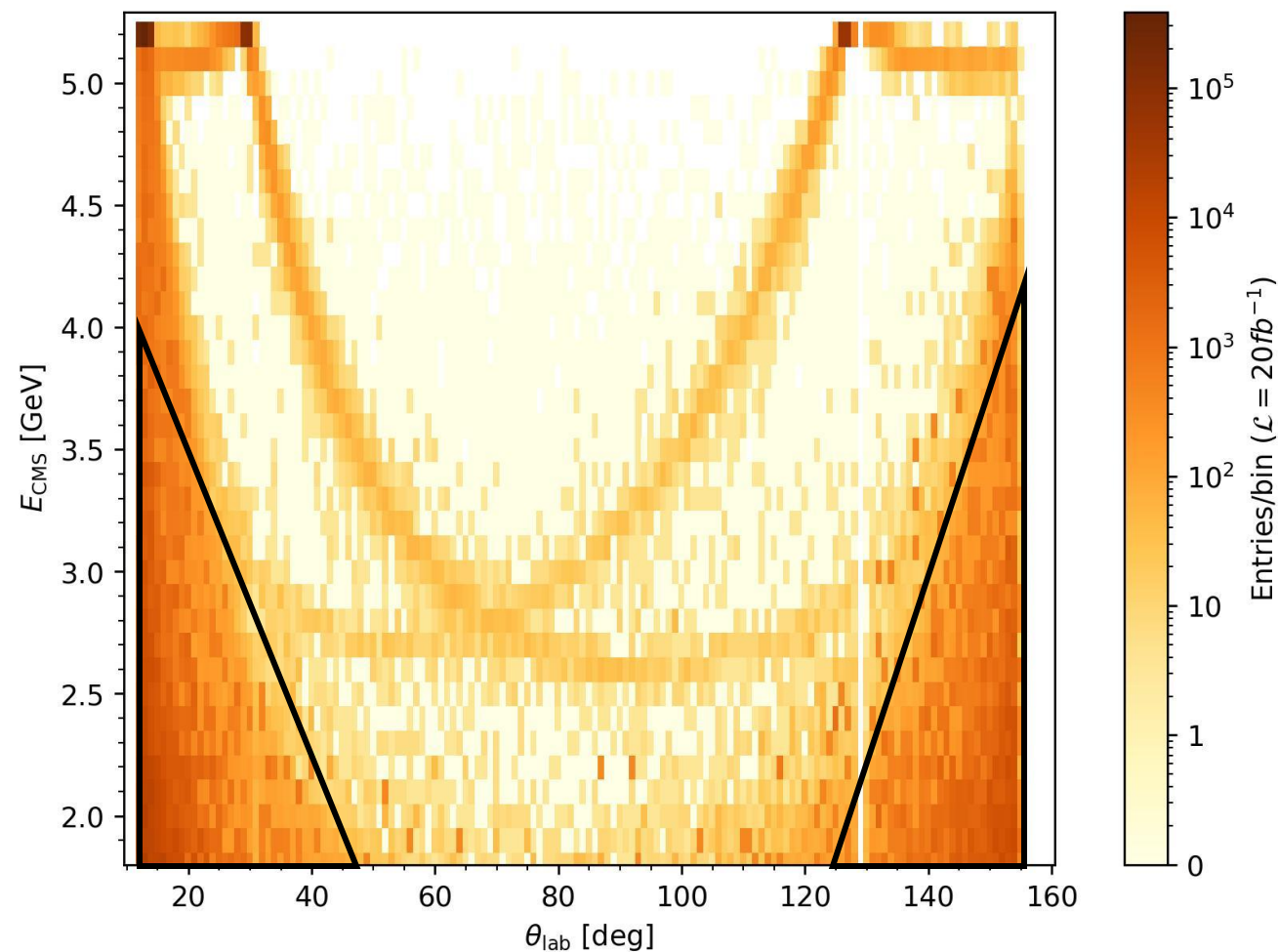
Gaps between barrel and endcaps



Standard Model Background

Processes:

- $e^+ e^- \rightarrow$
 - $\gamma\gamma(\gamma)$
 - $e^+ e^- \gamma(\gamma)$

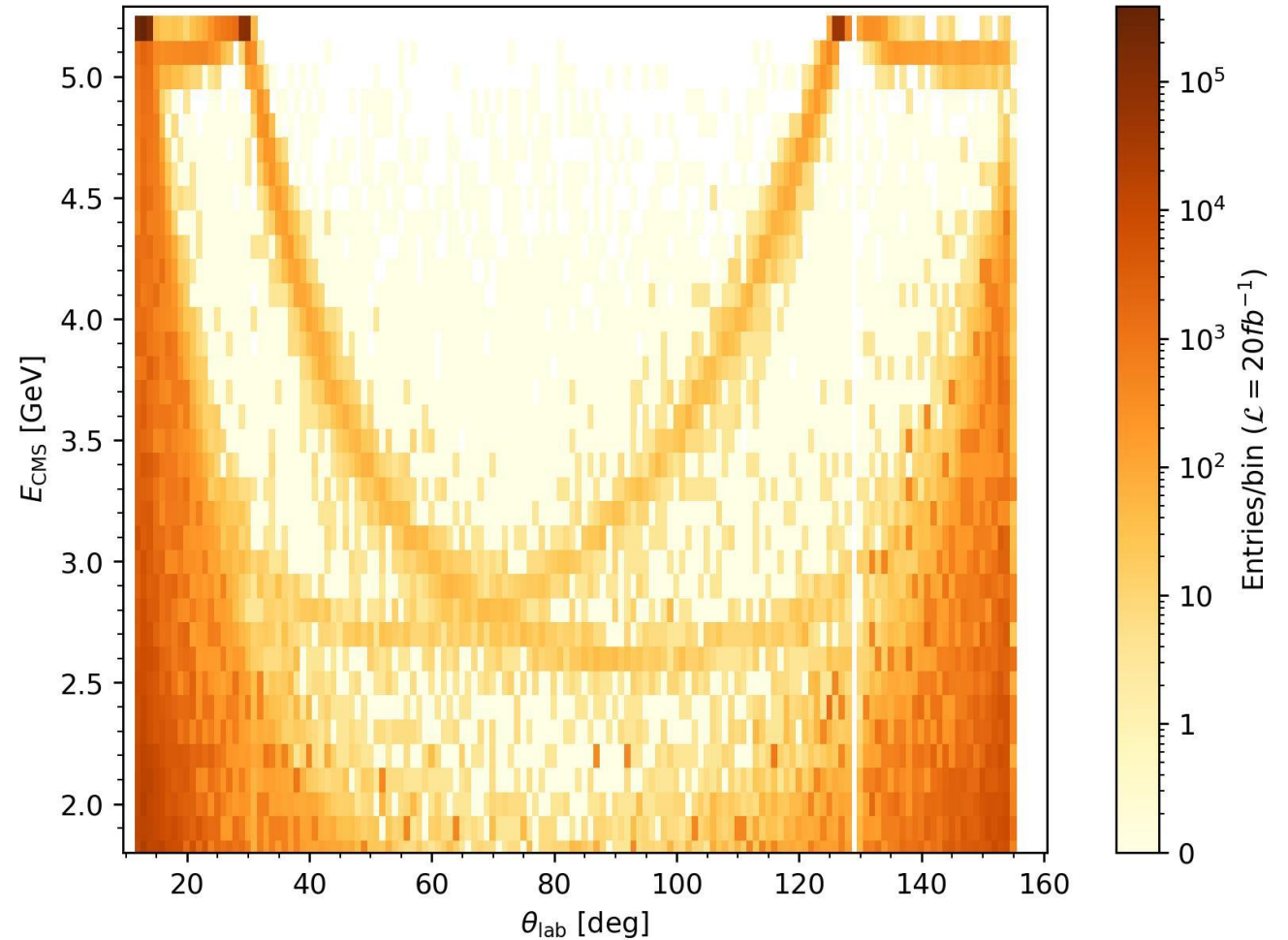


Standard Model Background

Processes:

- $e^+ e^- \rightarrow$
 - $\gamma\gamma(\gamma)$
 - $e^+ e^- \gamma(\gamma)$
 - $\bar{\nu}\nu\gamma(\gamma)$

Irreducible & everywhere



Dark Photon

Introduce a new $U(1)_X$ gauge symmetry with vector boson X_μ which kinetically mixes with the hypercharge boson B_μ .

$$\mathcal{L} = -\frac{1}{4}\hat{F}_{\mu\nu}\hat{F}^{\mu\nu} - \frac{\epsilon'}{2}\hat{F}_{\mu\nu}\hat{X}^{\mu\nu} - \frac{1}{4}\hat{X}_{\mu\nu}\hat{X}^{\mu\nu} - g'j_\mu^Y\hat{B}_\mu \\ - g_Xj_\mu^X\hat{X}_\mu - \frac{1}{2}\hat{M}_X^2\hat{X}_\mu\hat{X}^\mu$$

Other options are gauging a global symmetry of SM:

- $U(1)_{L_i-L_j}$ lepton number difference
- $U(1)_{B-L}$ baryon - lepton number difference (Used)

[Bauer, Foldenauer, and Jaeckel arXiv:1803.05466]

Axion-Like-Particles (ALPs)

Pseudo-scalar particle with linear effective Lagrangian

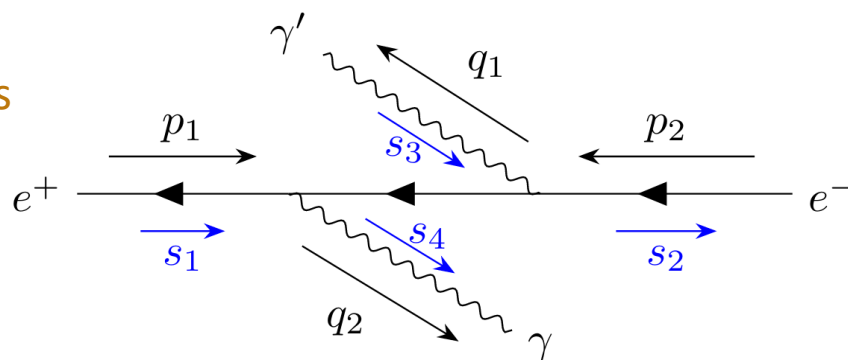
$$\begin{aligned}\mathcal{L}_{\text{eff}} = & \frac{1}{2} (\partial_\mu a) (\partial^\mu a) - \frac{m_a^2}{2} a^2 + \sum_f \frac{c_{ff}}{2} \frac{\partial^\mu a}{f_a} \bar{f} \gamma_\mu \gamma_5 f \\ & + g_s^2 C_{GG} \frac{a}{f_a} G_{\mu\nu}^A \tilde{G}^{\mu\nu,A} + g^2 C_{WW} \frac{a}{f_a} W_{\mu\nu}^A \tilde{W}^{\mu\nu,A} + g'^2 C_{BB} \frac{a}{f_a} B_{\mu\nu} \tilde{B}^{\mu\nu}\end{aligned}$$

[I. Brivio et al arXiv:1701.05379 & Bauer et al arXiv:1808.10323]

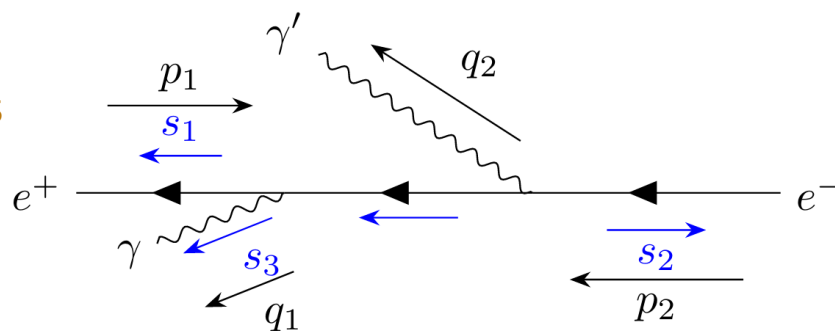
Spin Structure Diagrams

Dark Photon

Small DM mass

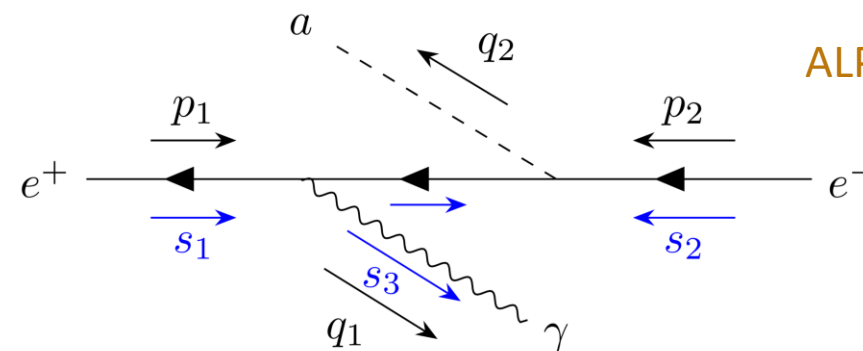


Large DM mass

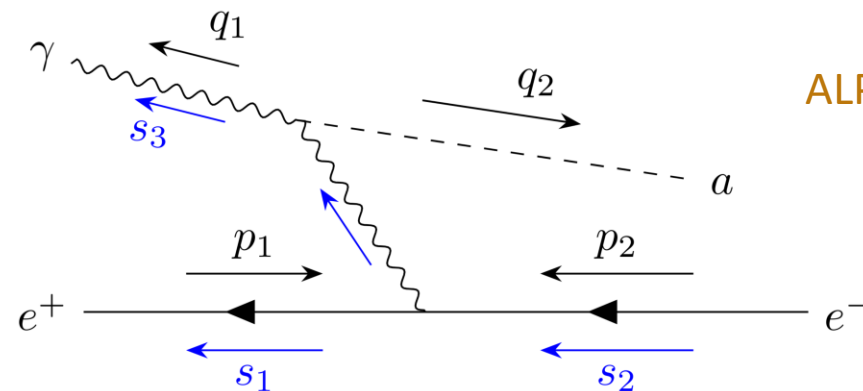


Axion-Like-Particles

ALP-Fermion



ALP-Photon



Polarisation of the Photon



Not Planned!



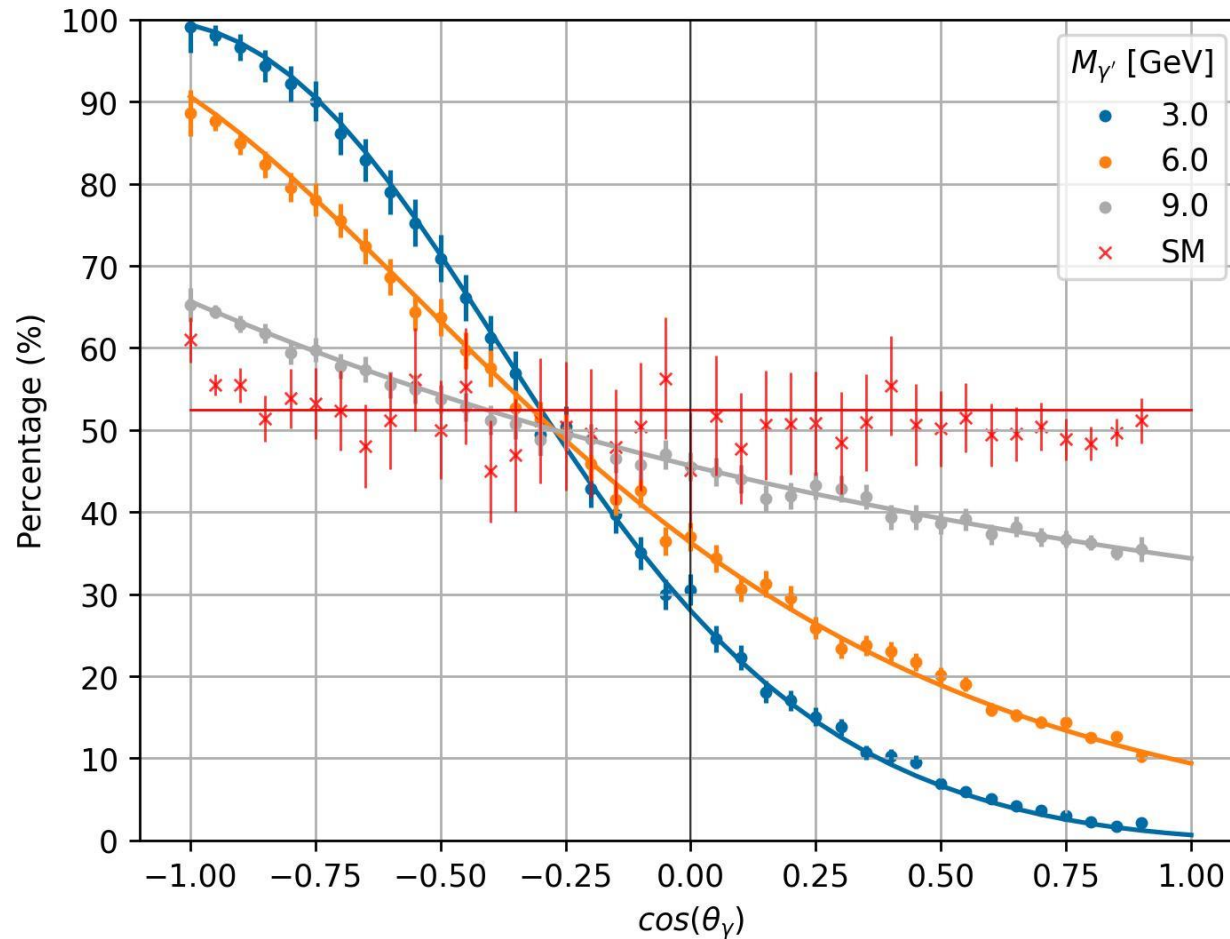
Speculative

Polarisation of the Photon

Percentage of the photon helicity being equal to the incoming electron for 100% polarised beams

$$\text{Percentage} = \frac{|\mathcal{M}|^2(2\lambda_{e^-} = \lambda_\gamma = 1)}{|\mathcal{M}|^2(\lambda_{e^\pm} = \lambda_\gamma = 0)}$$

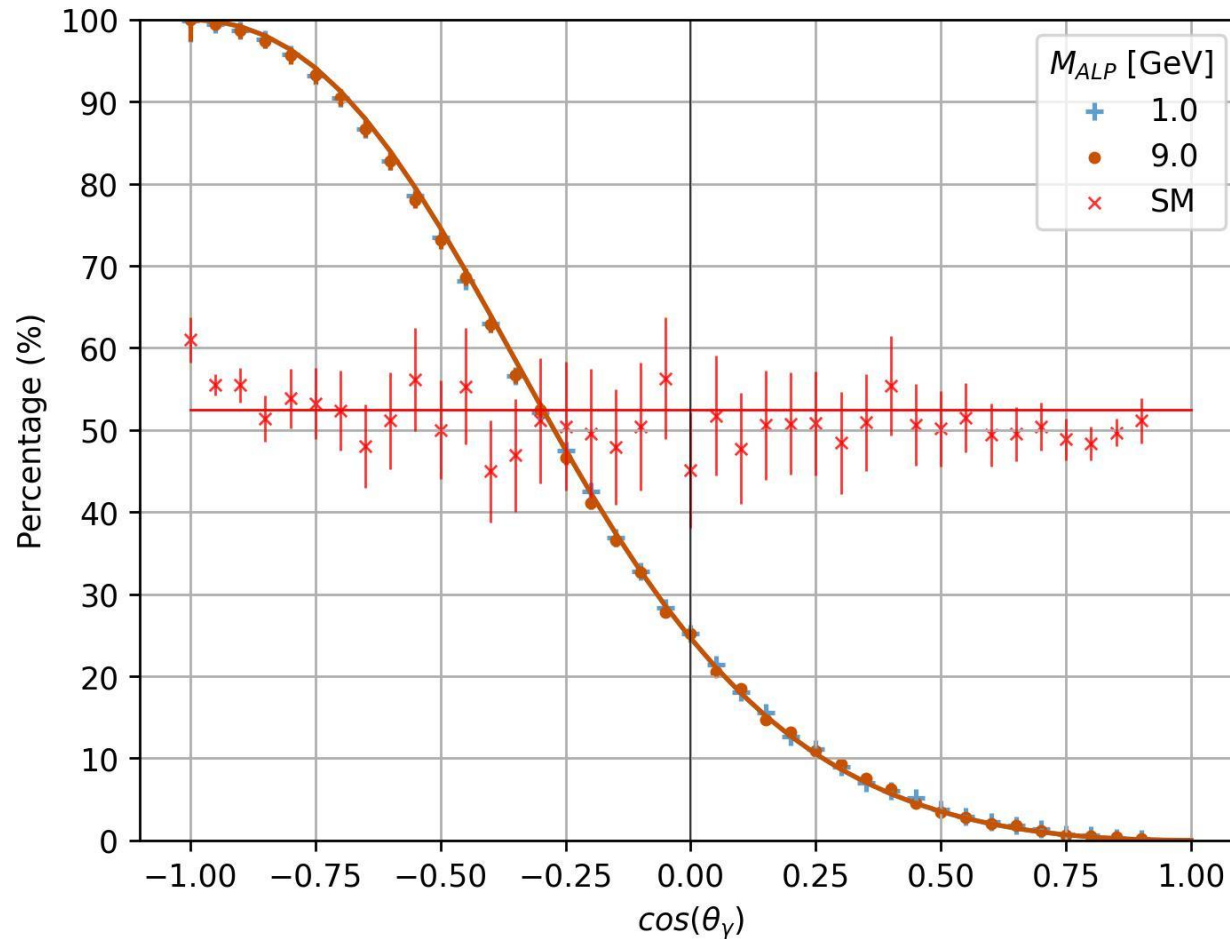
Polarisation of the Photon



Dark Photon

- Angular correlation
 - Follows incoming beams
- As mass increases, direction of final photon becomes more random
⇒ Percentage tends towards 50 %
- Skew from unequal beam energies

Polarisation of the Photon



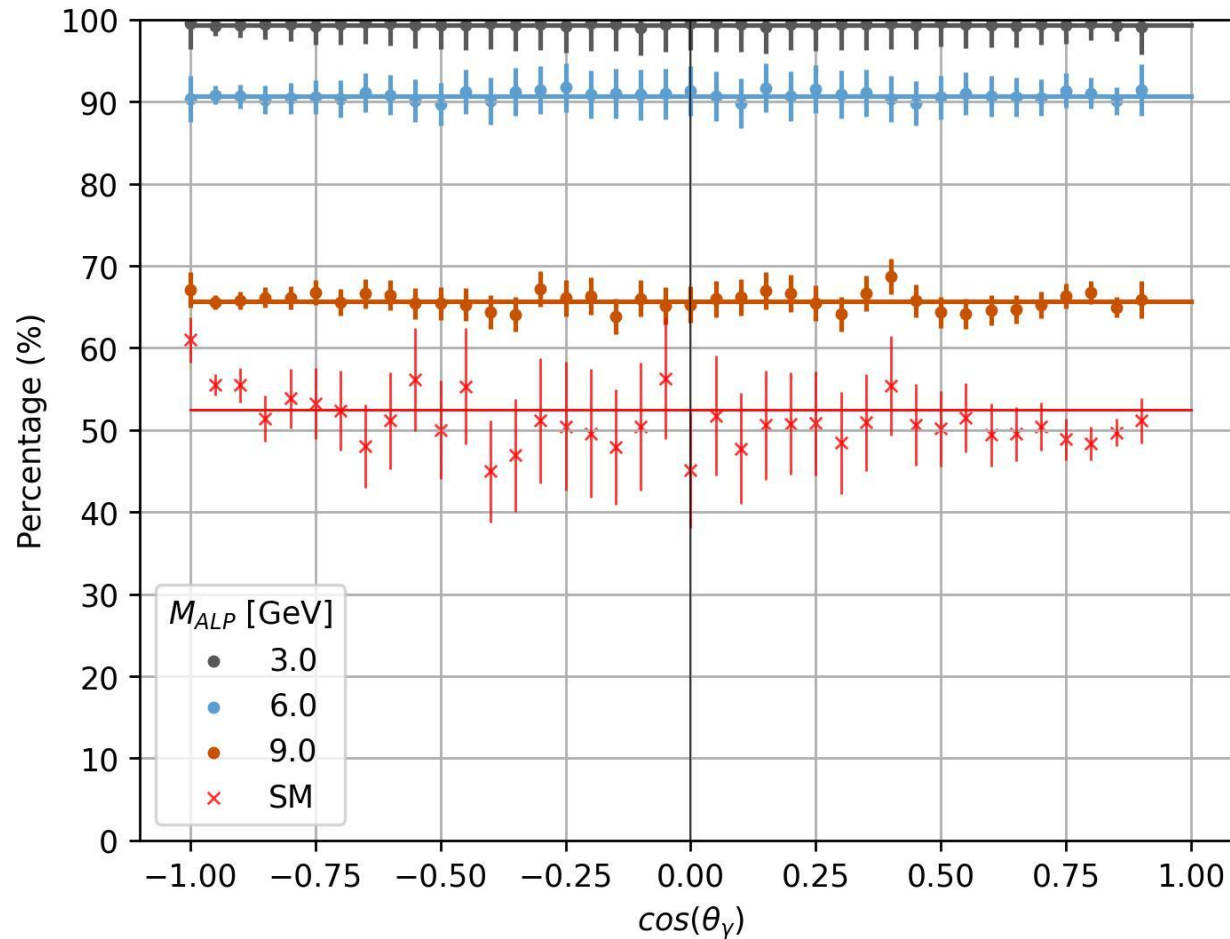
ALP-Photon Coupling

- Same shape as Dark Photon
- ALP emission independent of incoming beams
⇒ ALP mass factors out and has no influence

Significantly different to SM Background

⇒ Better exclusion limits

Polarisation of the Photon



ALP-Fermion Coupling

- ALP-fermion interaction causes helicity flip
⇒ Helicity match both incoming beams
⇒ No angular dependence
- Same shape as SM Background

Different from dark photon and ALP-photon coupling

Polarisation of the Photon

Not Planned! This was speculative

- Do not know how it would be implemented
⇒ Not going to try and provide exclusion limits