

Looking into the early Universe with a solid-state device

Quantum measurement device
with quantum impurities
for particle physics
and physics of the early universe

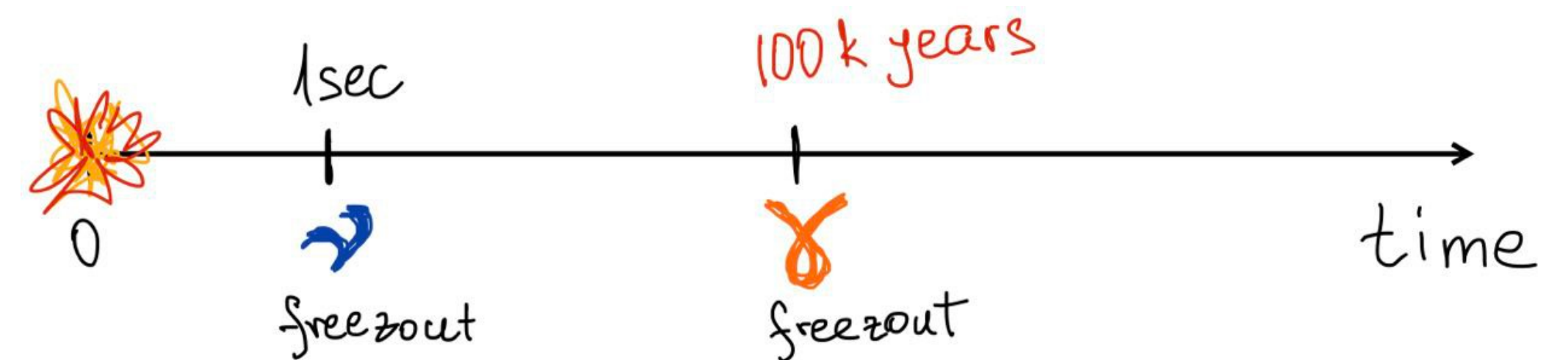
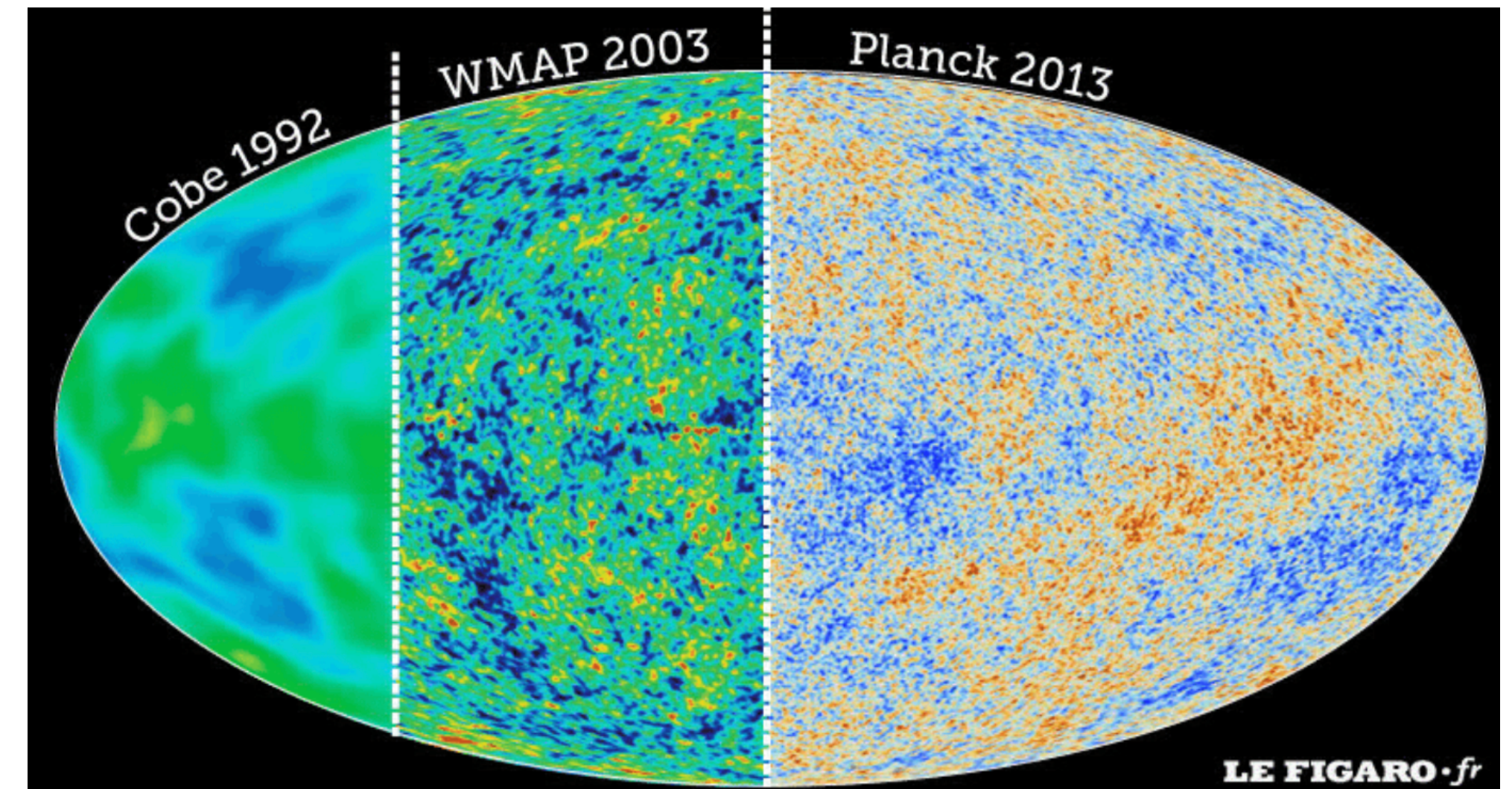


Oleksiy Boyarsky, Yevheniia Cheipesh, Oleksiy Mikulenko, Zhiyang Tan, Vadim Cheianov



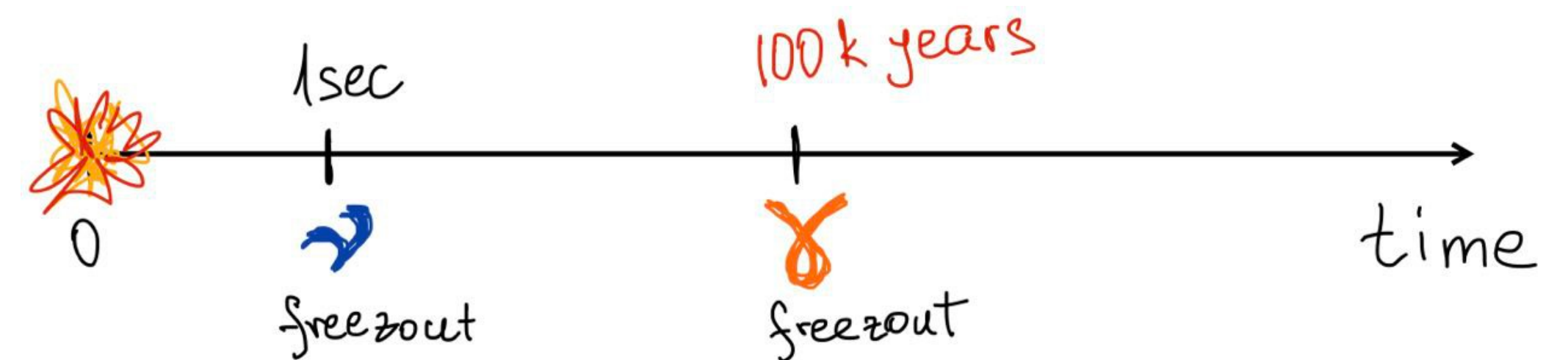
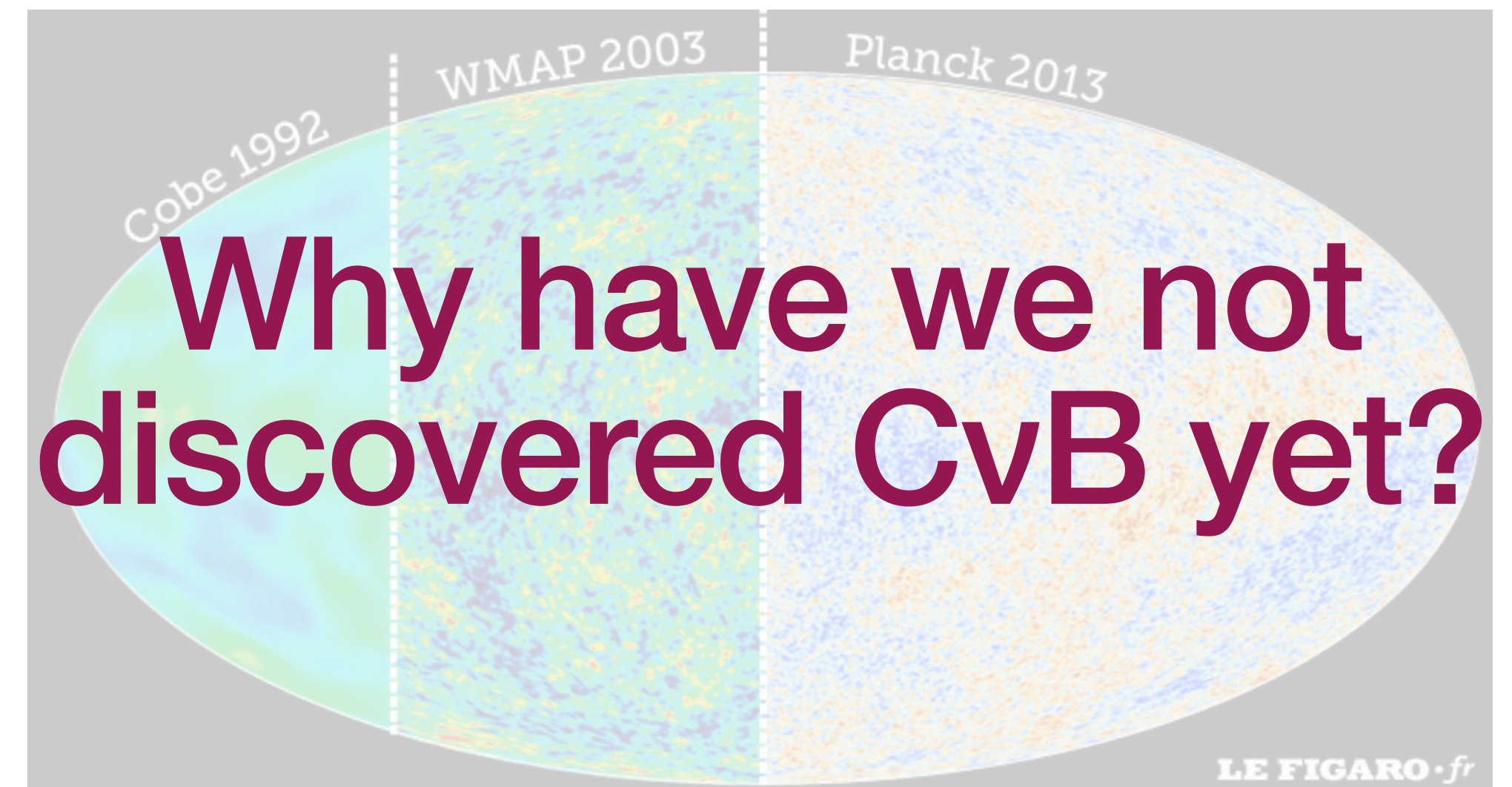
How to probe the early Universe? $C\nu B$ and CMB

- **Gamov:** Early Universe is radiation dominated.
 $\rho_{\text{rad}}/\rho_{\text{matter}} \sim 10^{10}$
- For cosmology, any relativistic particle is “radiation” !
- In fact, early Universe has equal populations of ν/γ .
- As Universe expands, ν/γ decouple, relic backgrounds (CMB / $C\nu B$), keep a “**frozen**” pictures of the Universe.
- Right now, in your room, there are **411 relic photons** and **339 relic neutrinos** in every cm^3 !
- The “ ν - **freezout**” is much earlier than photons
- $C\nu B$: one of the few **yet untested predictions of the SM**
- **Detecting $C\nu B$ is a strategic goal for fundamental physics.** [Weinberg , 1962]



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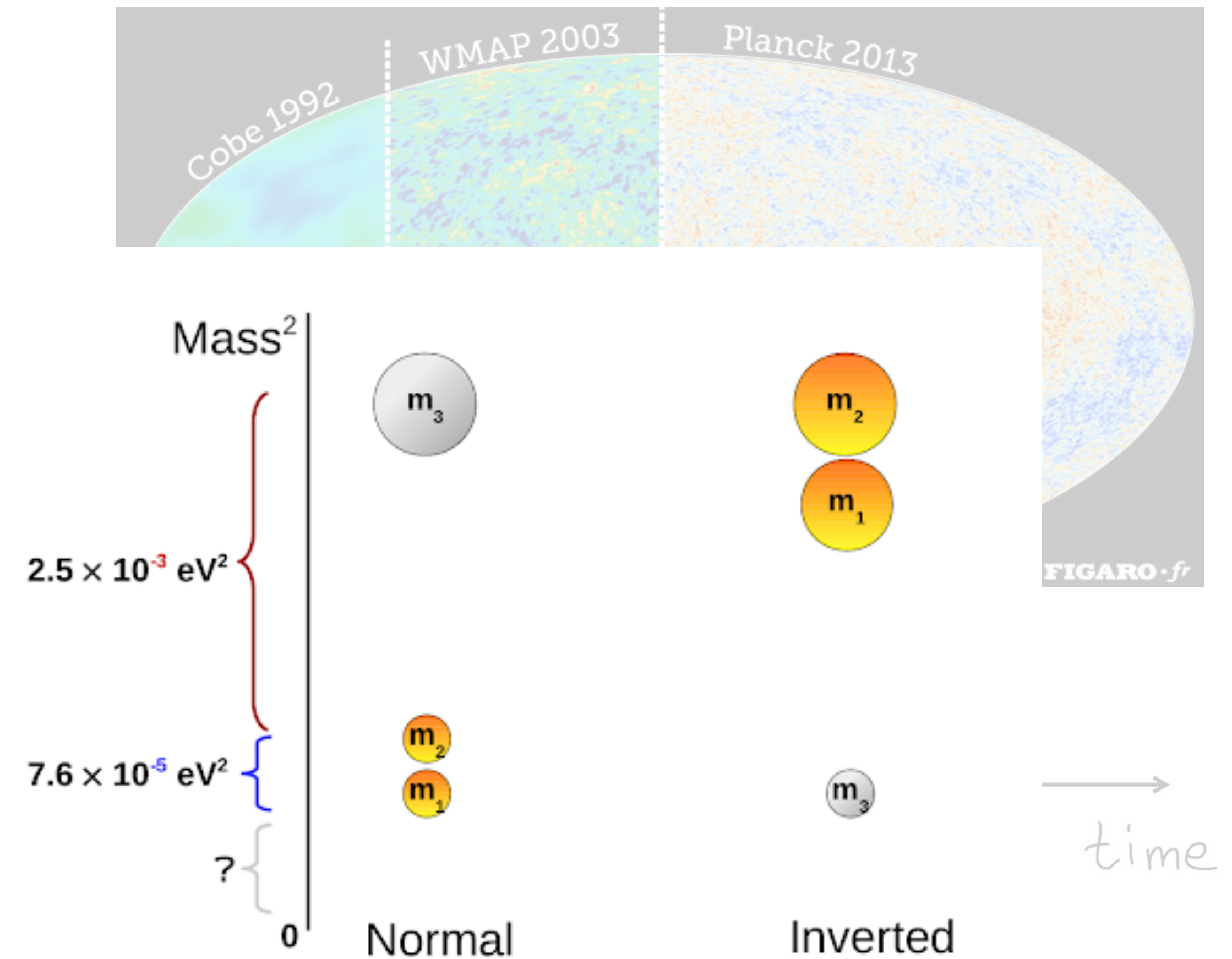
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Observation of the cosmological neutrinos would then provide a window into the 1st second of creation

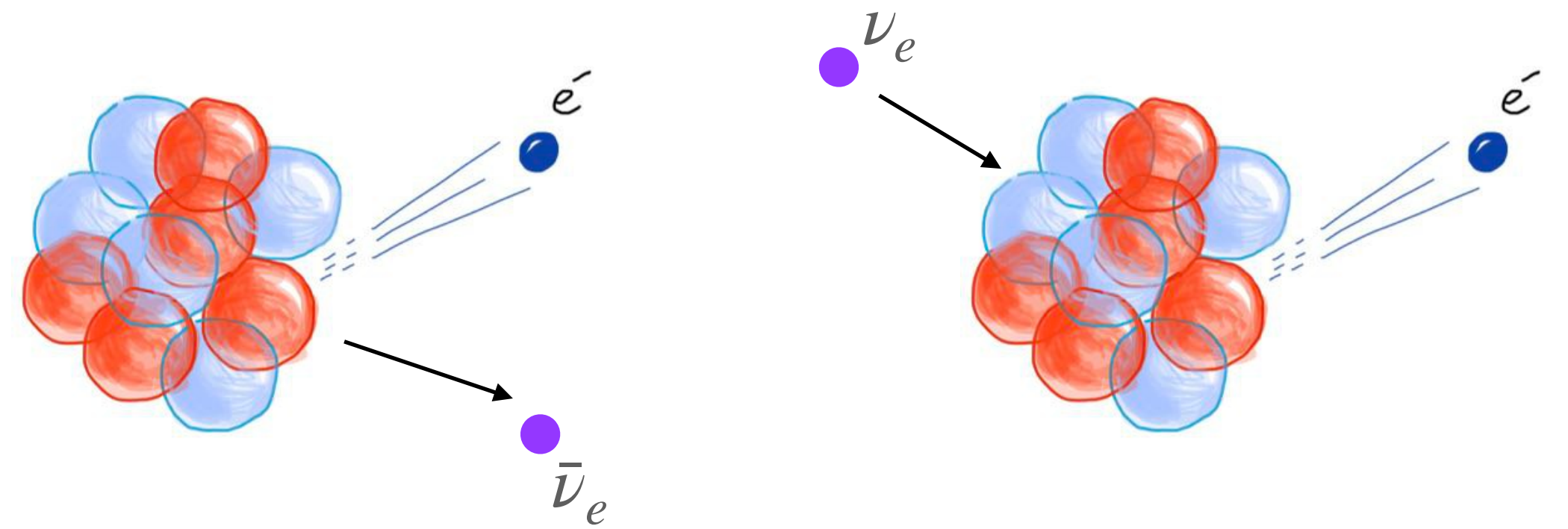
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- **Detecting $C\nu B$ is a strategic goal for fundamental physics.** [Weinberg , 1962]
- Neutrino oscillations can only measure Δm and hierarchy



Detecting relic neutrinos via β decay

- Neutrino capture is **threshold-less** – soft relic neutrino detection
[Weinberg, 1962]



β decay

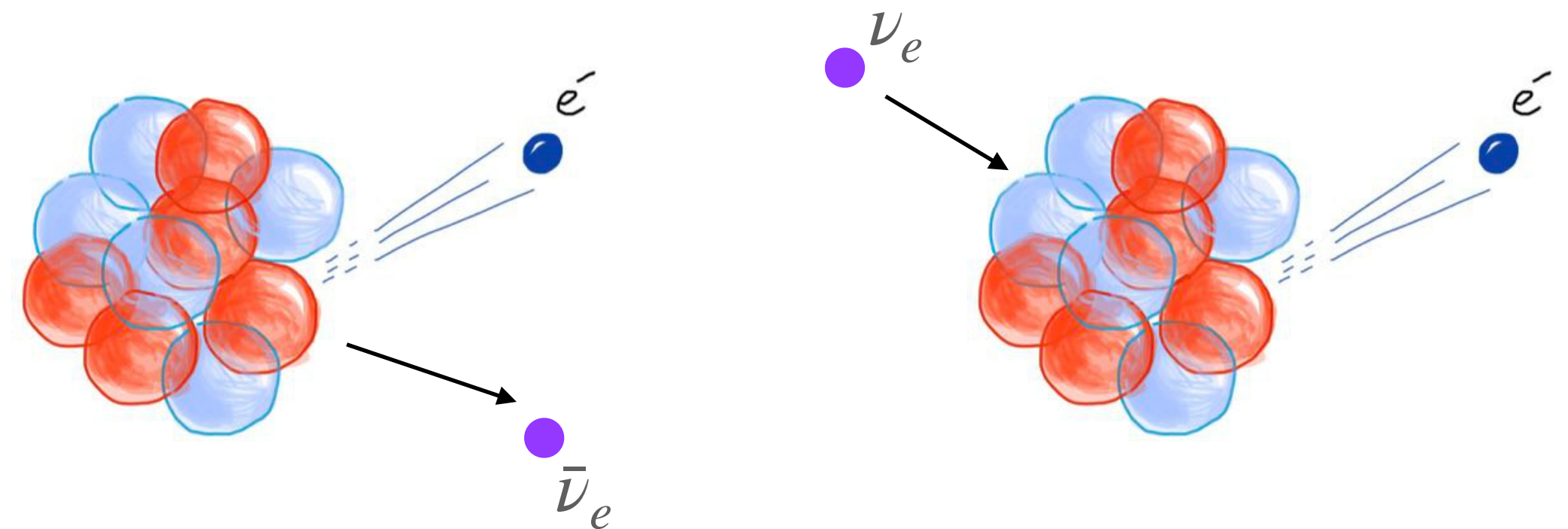
Neutrino capture

Detecting relic neutrinos via β decay

- Neutrino capture is **threshold-less** – soft relic neutrino detection [Weinberg, 1962]
- The **2 parts of the spectrum** are separated by $2m_\nu$

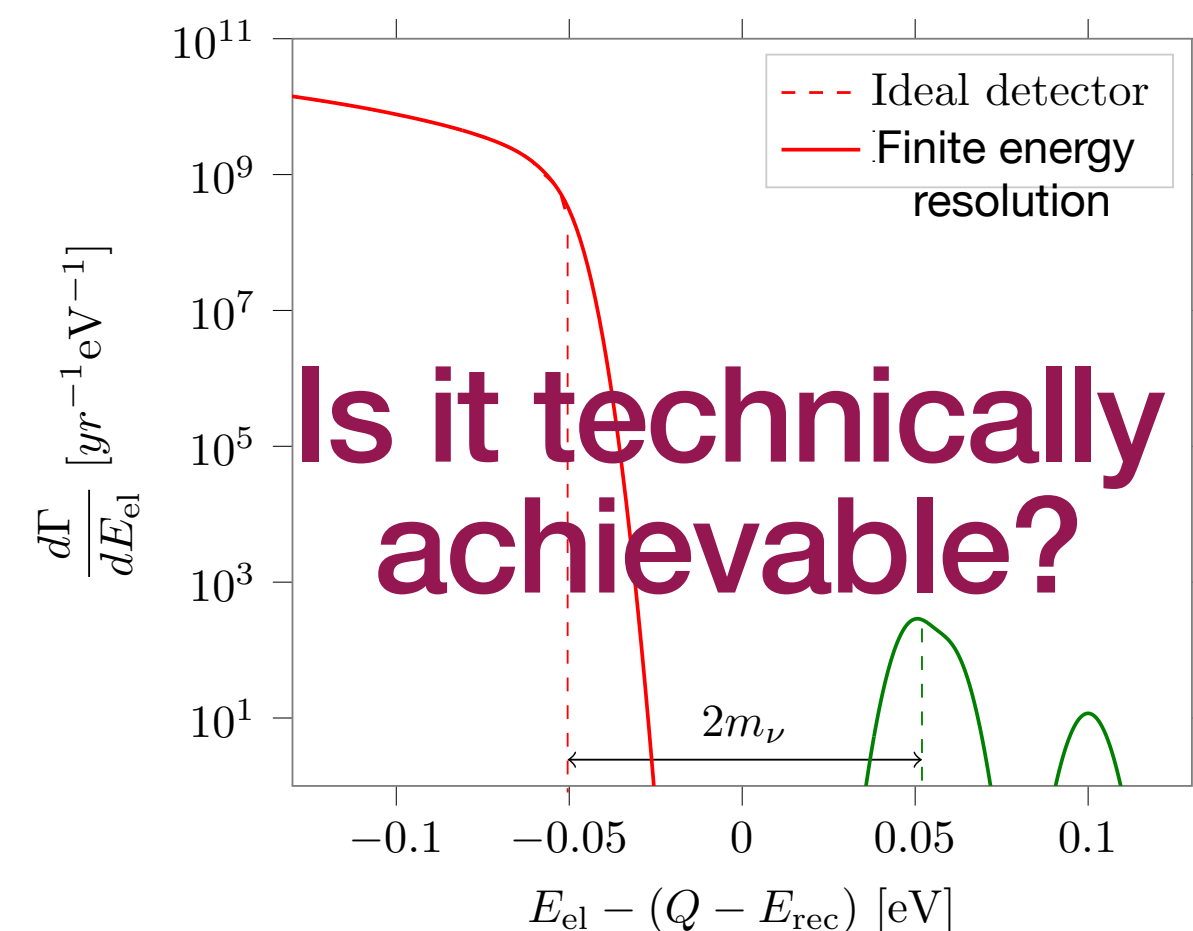
Challenges

- High **energy precision** (order of $m_\nu \sim 10$ meV)
- Sufficient **activity rate** (several events per year)

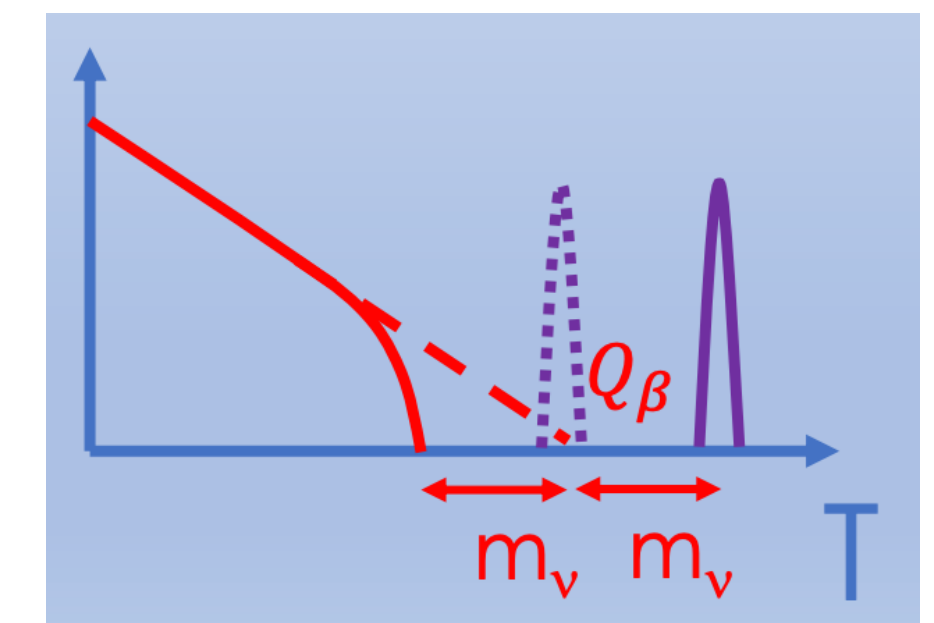
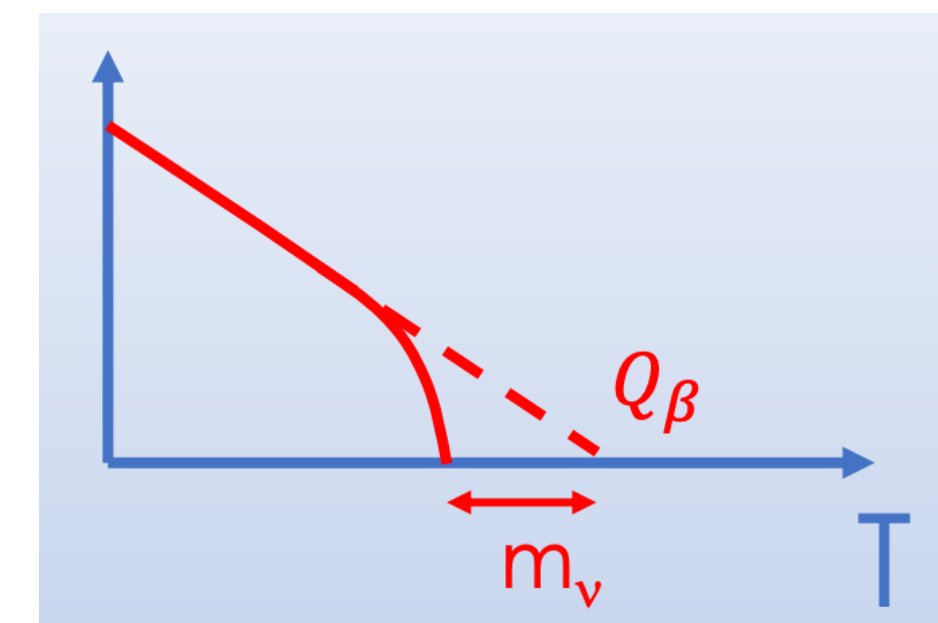


β decay

Neutrino capture



$$(A, Z) \rightarrow (A, Z + 1) + e^- + \bar{\nu}_e \quad \nu_e + (A, Z) \rightarrow (A, Z + 1) + e^-$$



Requirements to the experiment

High enough activity

- Low emitter **Q-value**
[Cocco et.al., 2007]
- **Lifetime of emitter:**
small enough to have a high decay rate, but large enough not to decay instantly

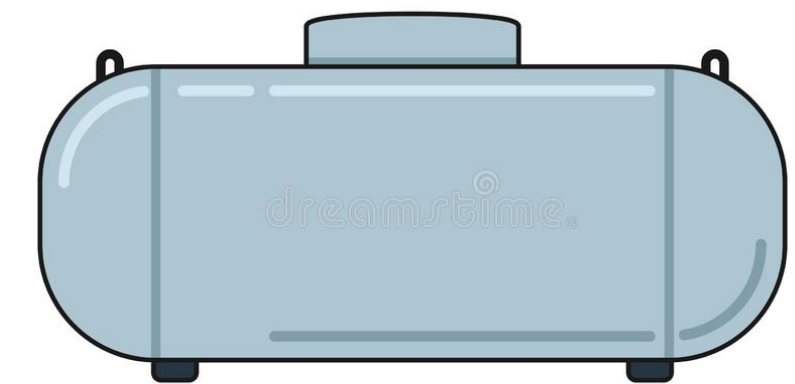
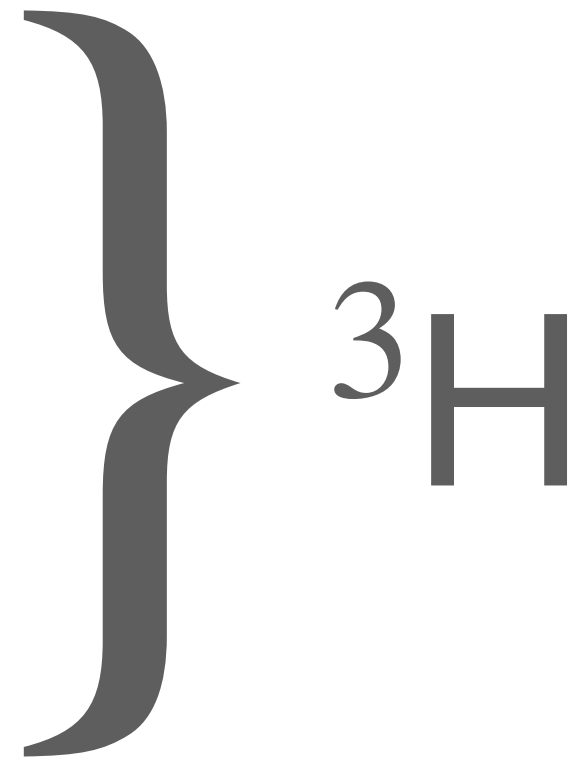
} ^3H

$$\sigma_\nu \sim \frac{1}{\tau Q^3}$$

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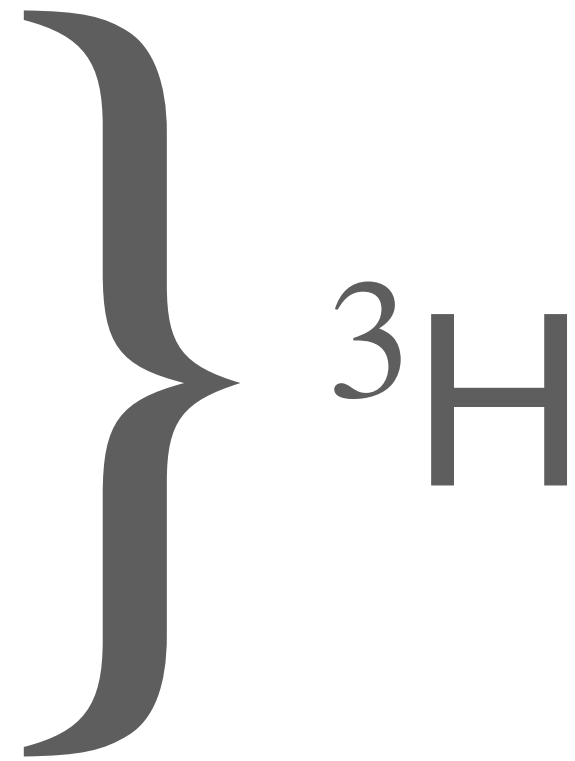


In gaseous form?

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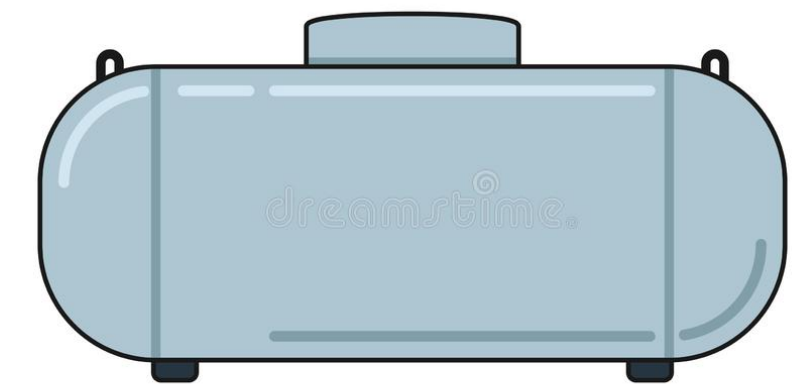
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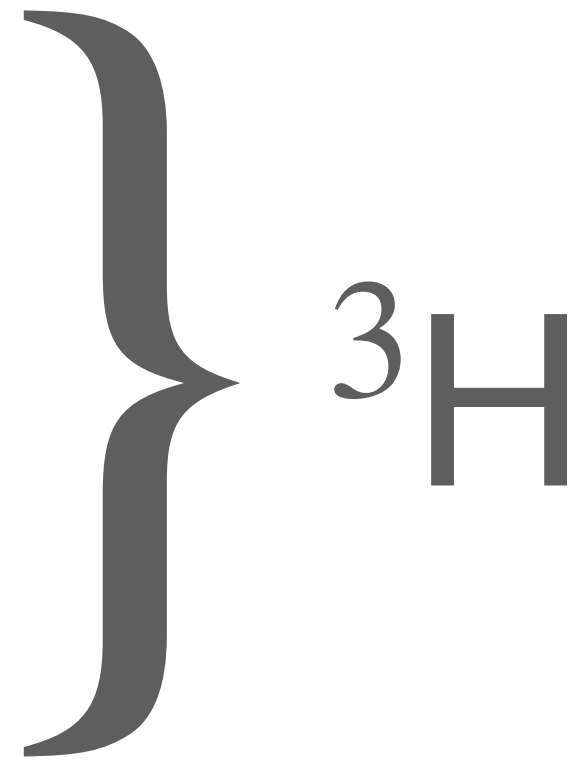


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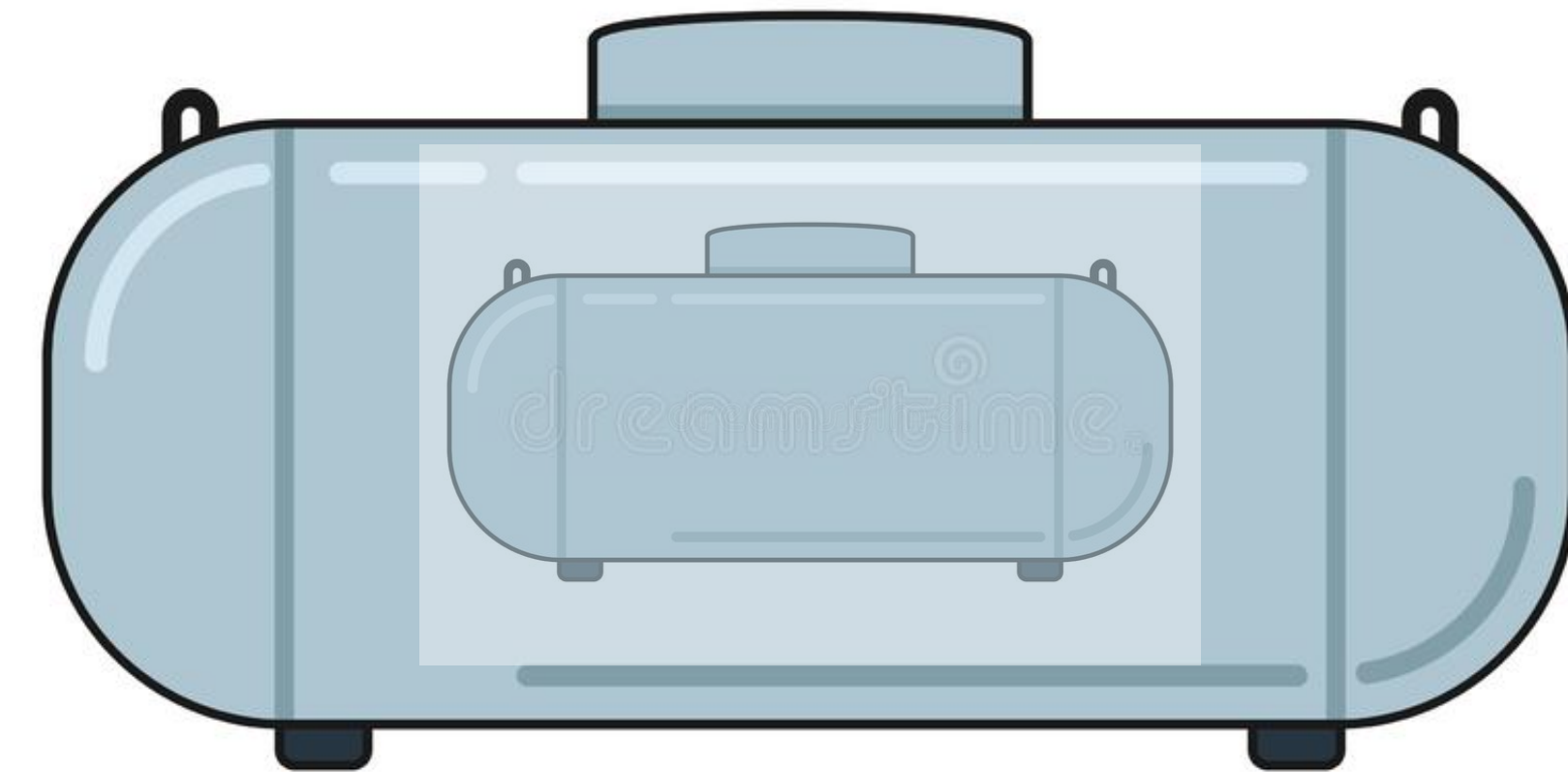
- Low emitter **Q-value**
- Low **emitter densities**
electron free path bigger than the system size

Cross section

$$\lambda = \left(R_{\text{atom}}^2 \frac{N}{L^3} \right)^{-1} > L$$

$$L \sim 1\text{km}$$

Very naive! In reality much bigger

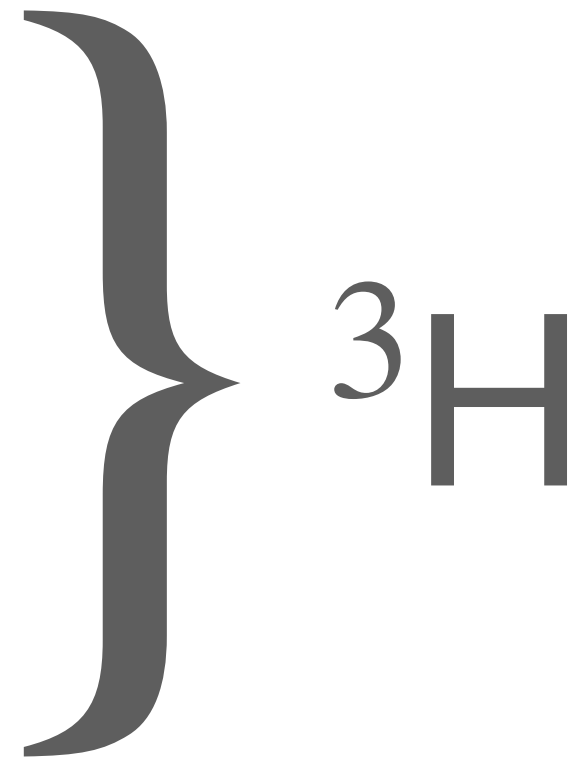


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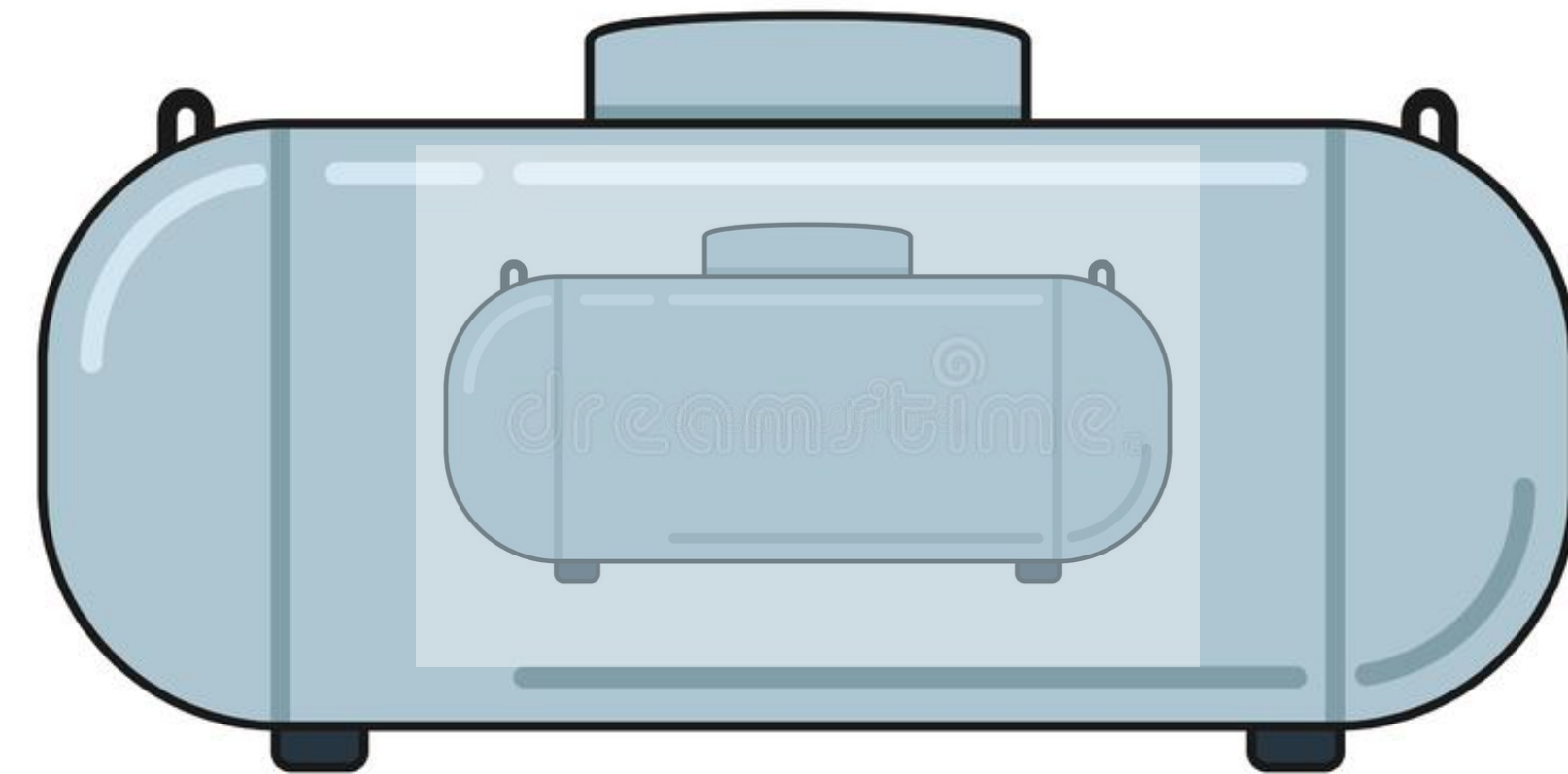
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High energy precision

- Low emitter **Q-value**
- Low **emitter densities**
electron free path bigger than the system size
- Low **volume**

$$\Delta E \sim \frac{V_{\text{source}}}{V_{\text{detector}}}$$

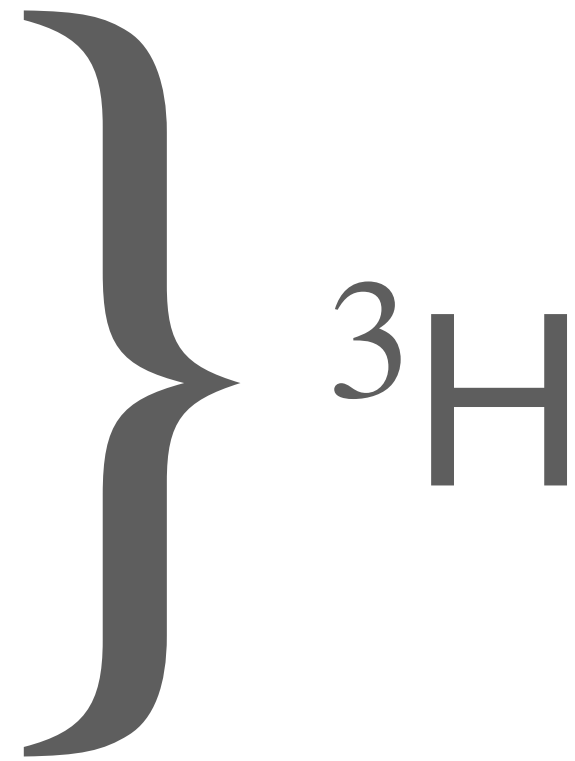


In gaseous form?

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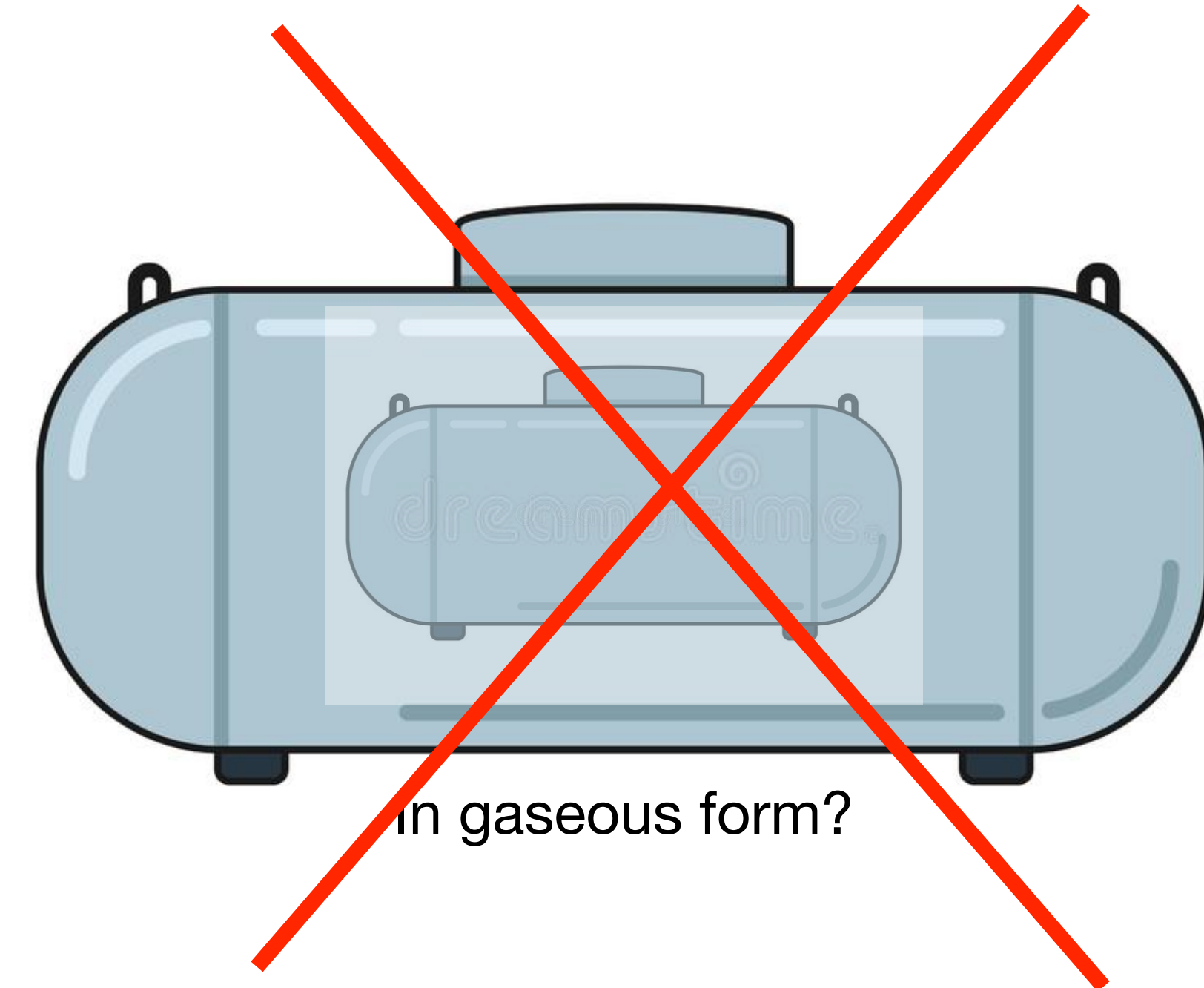
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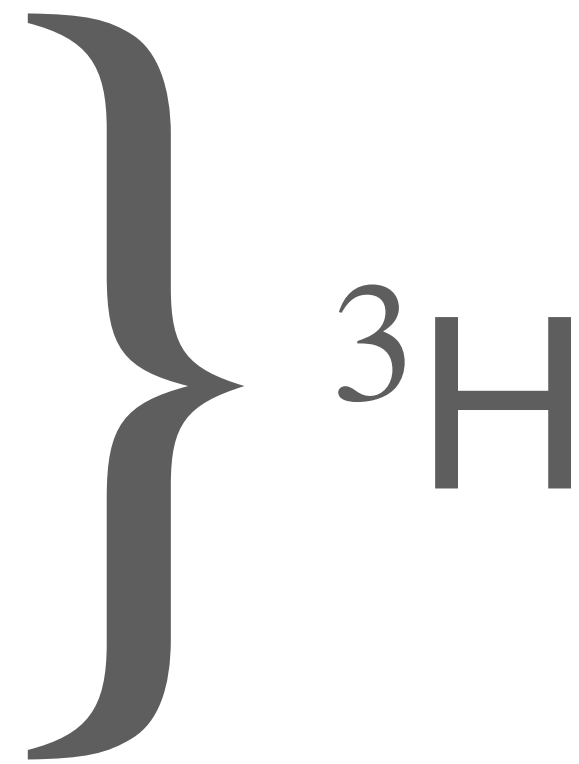
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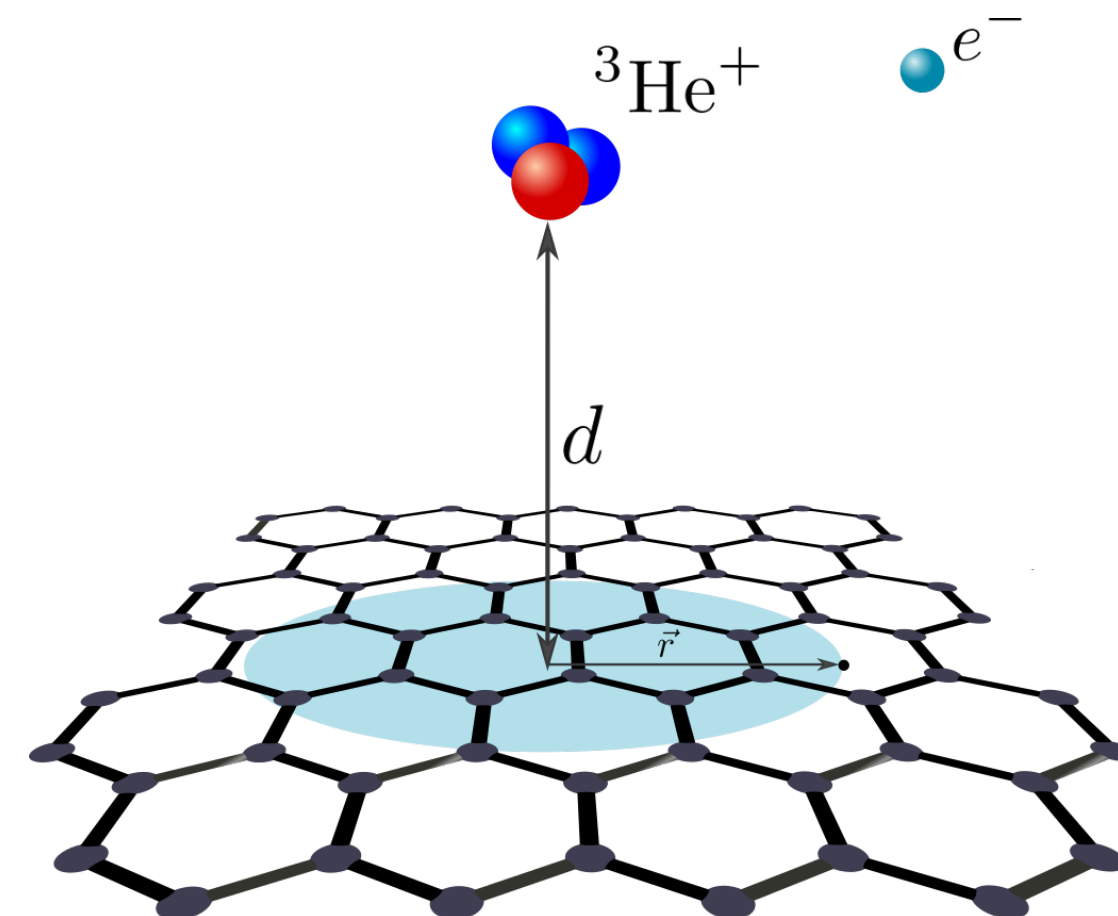
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- Radioactive material in gaseous form does not suit
- Need in the **solid-state based experiment**

POLLICA
SUMMER
WORKSHOP

CASTELLO DEI PRINCIPI CAPANO, POLLICA, ITALY

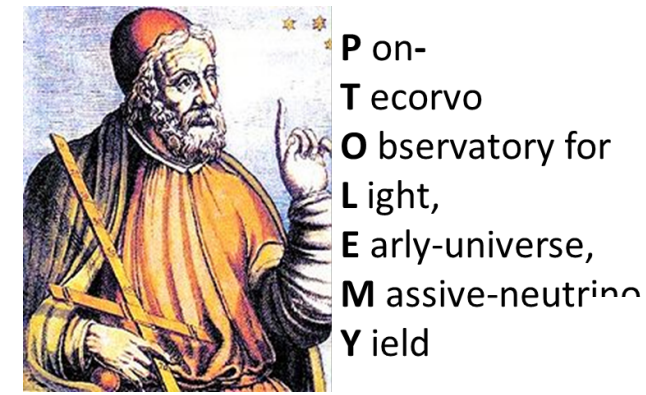
Overview Poster Application Timetable Participant List Local logistics Shuttle service

High energy physics meets low energy phenomena

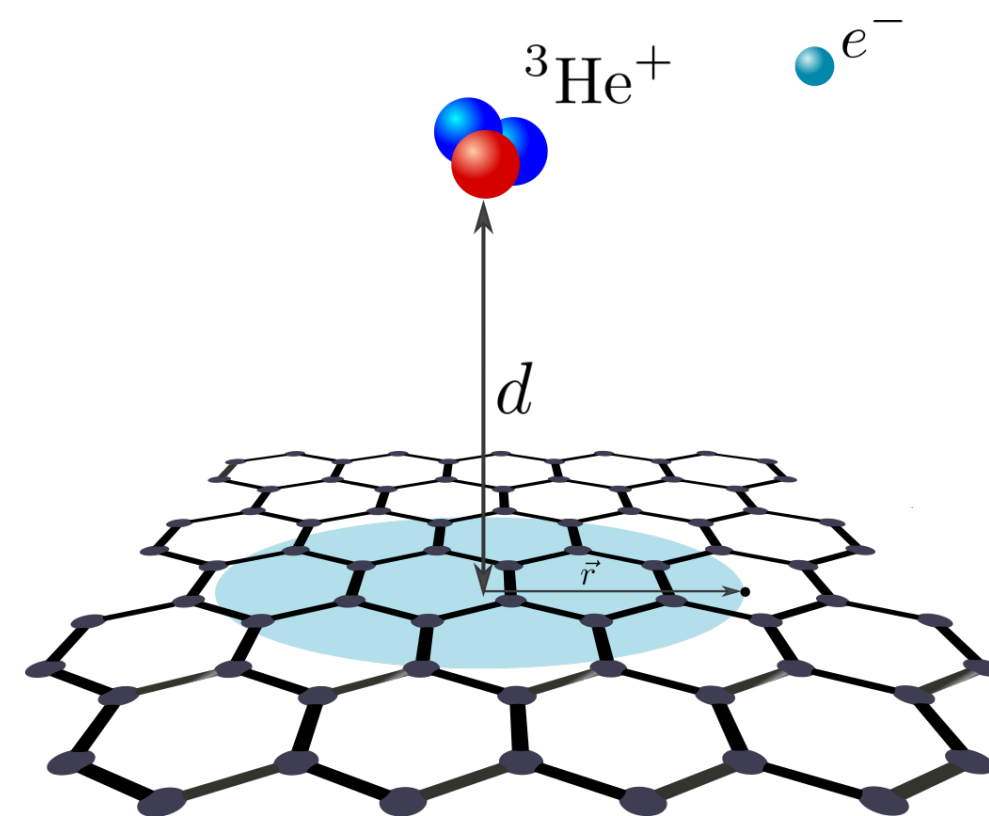
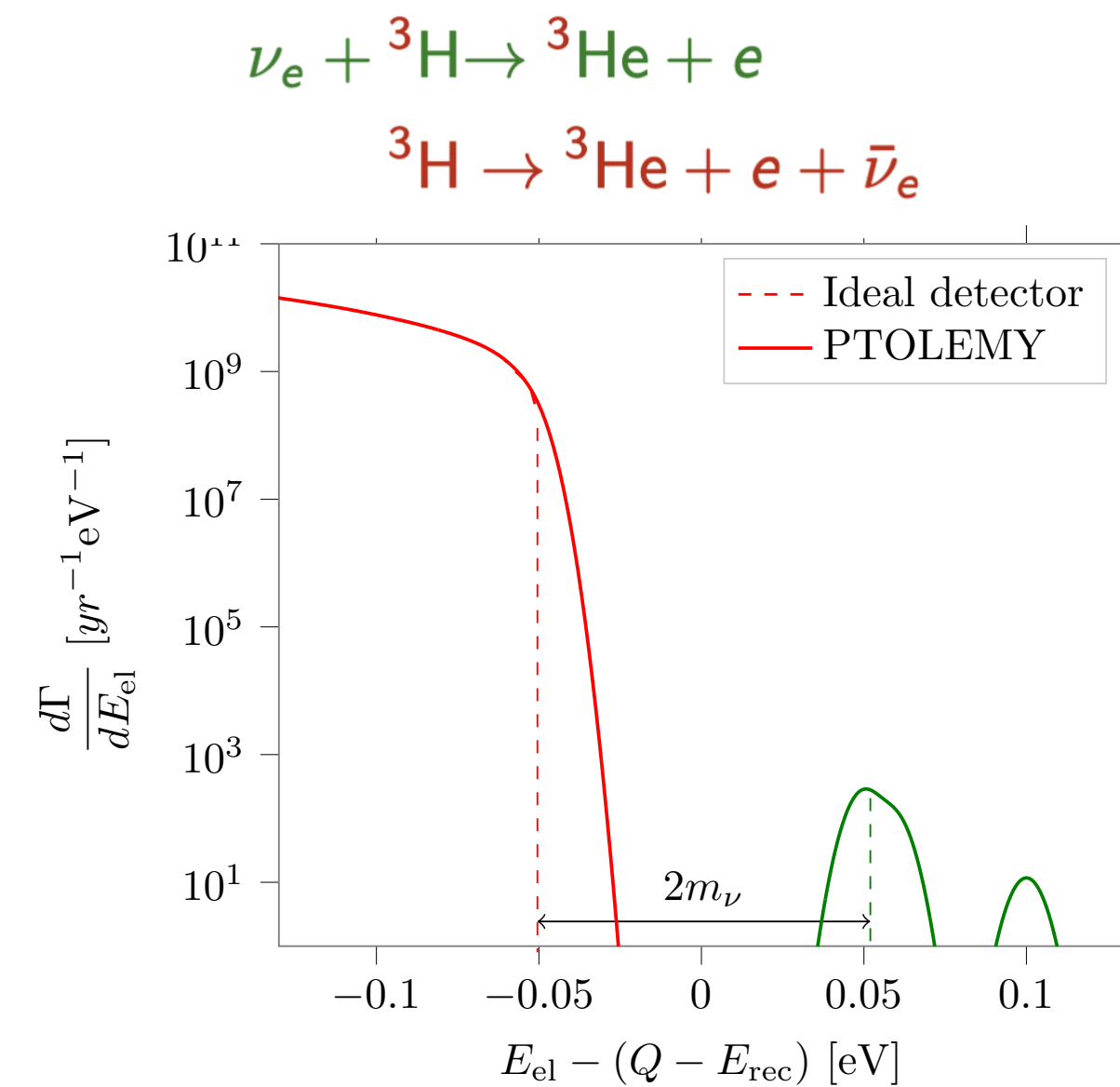
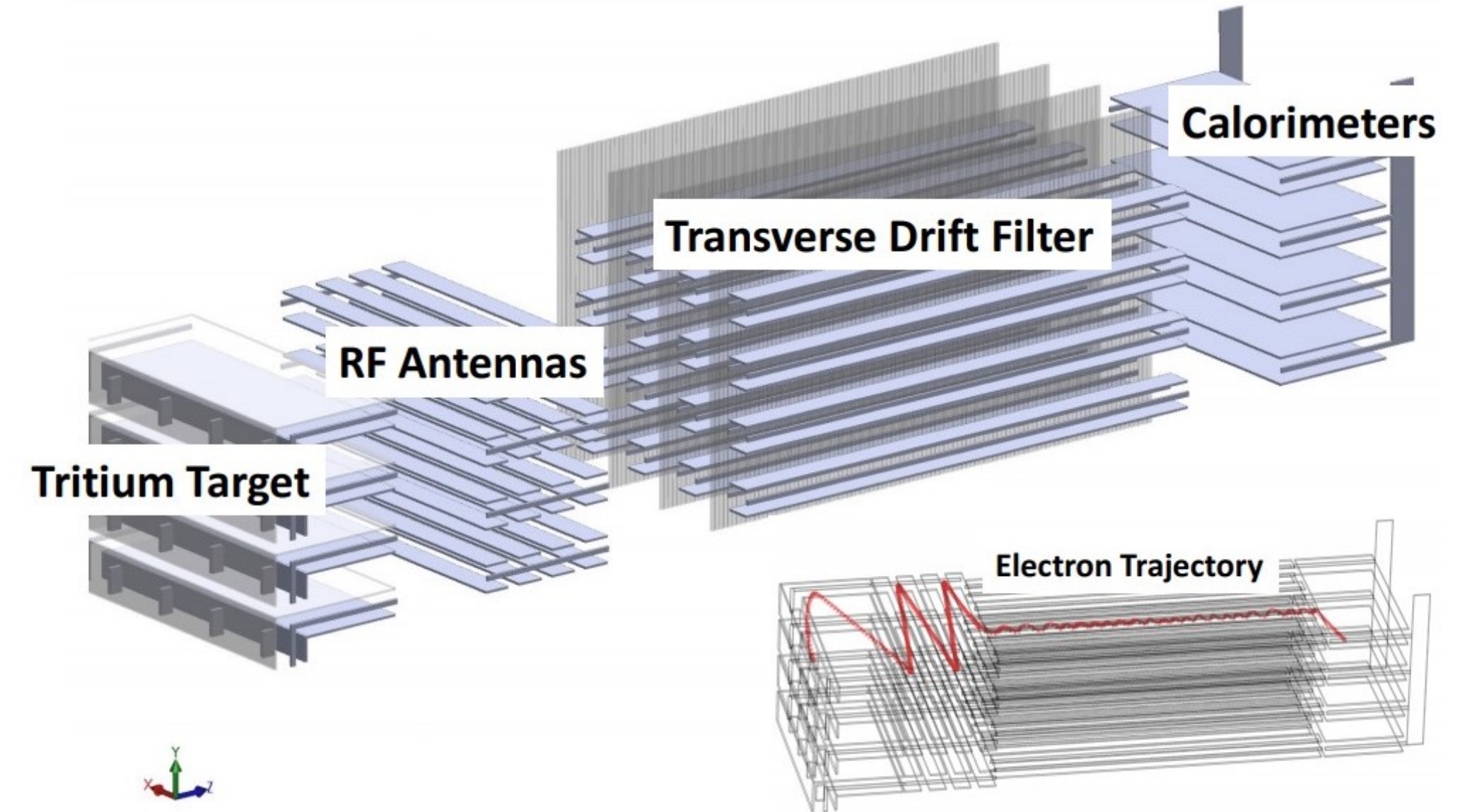
Throughout the history of quantum field theory there has been a rich cross-pollination between **high energy** and **condensed matter physics**. From the theory of renormalization to the consequences of spontaneous symmetry breaking, this interaction has brought tremendous progress in our understanding of natural phenomena. The recent developments in our systematic study of conformal field theories and new phases of matter, the conformal bootstrap, gauge/gravity and other types of dualities, have re-ignite a fervent exchange between high energy and condensed matter theorists which has brought already remarkable progress.

PTOLEMY project

State of the art

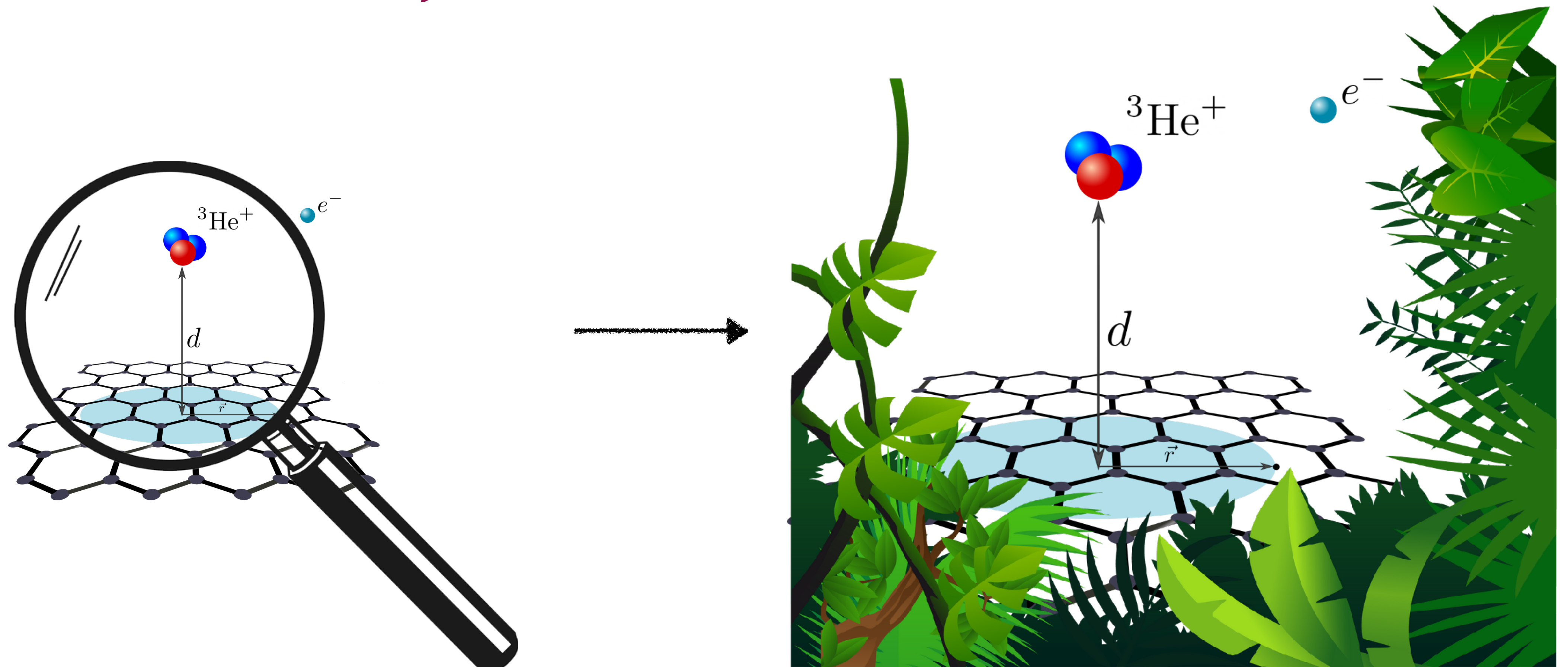


- **Tritium** as a β -decay emitter.
- Tritium is deposited on **graphene** sheets
- ≈ 4 CvB **events** per year.
- Outstanding **energy resolution** of the apparatus ≈ 10 meV.



Jungle of many-body and chemical effects

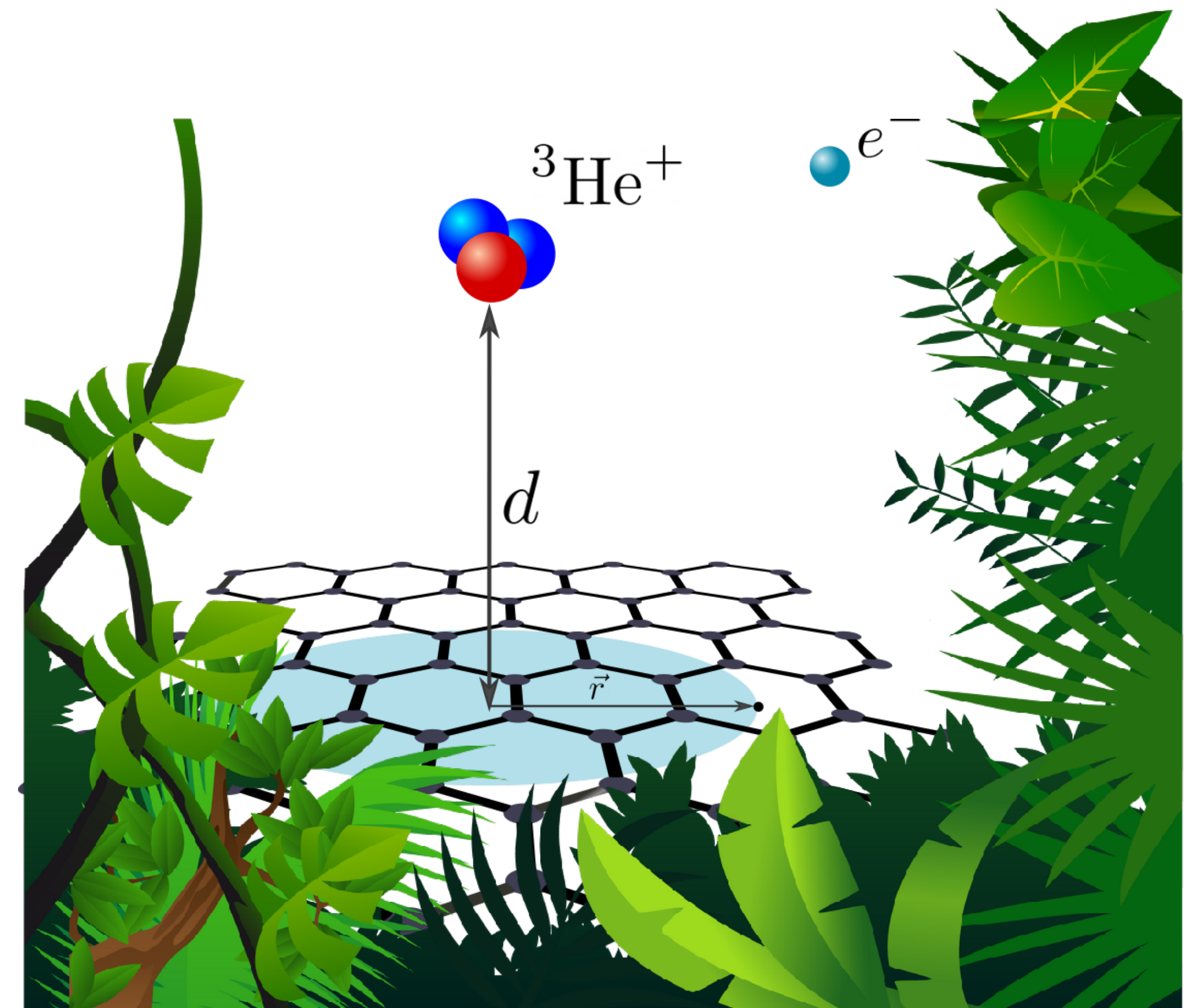
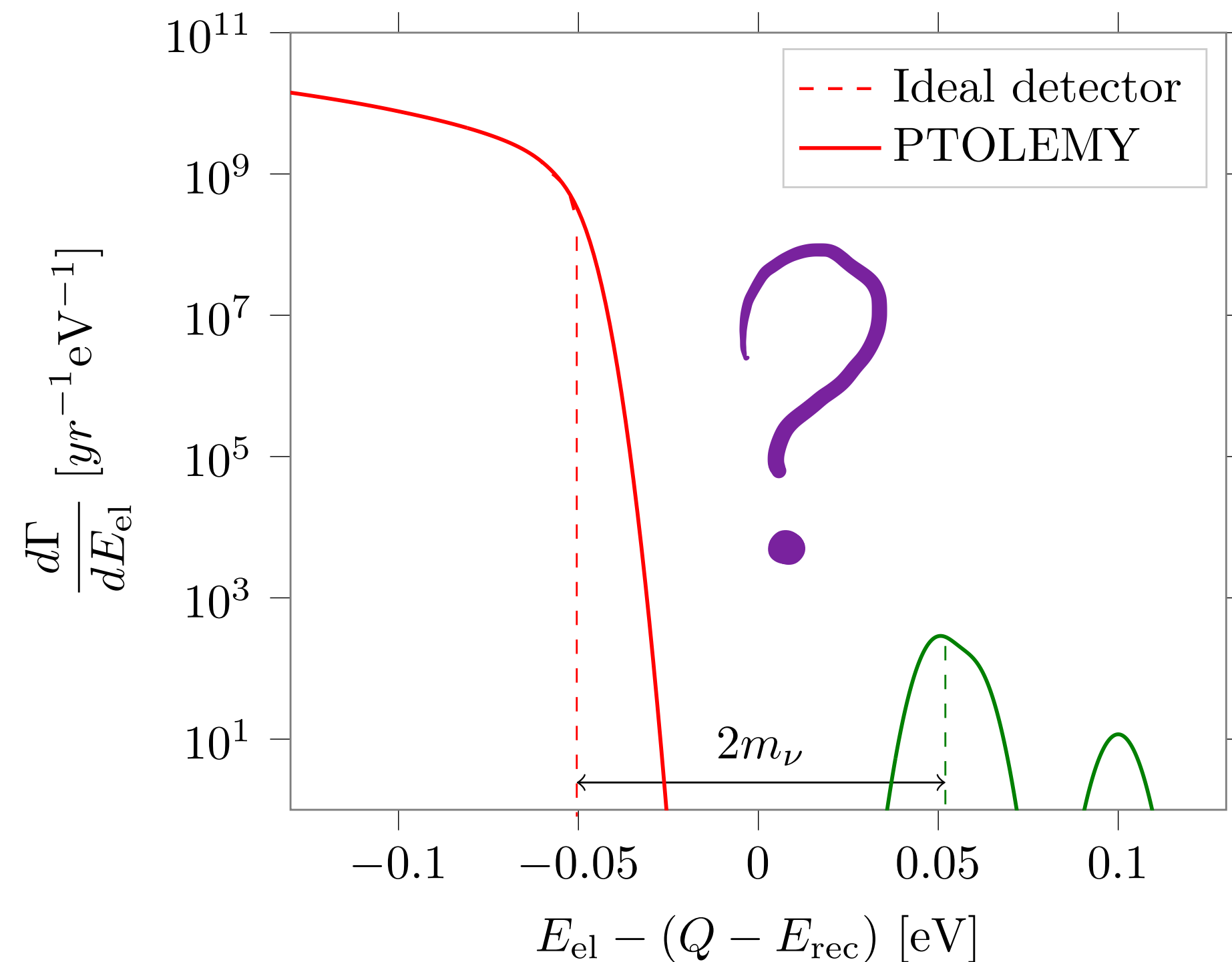
We need energy resolution $> m_\nu \sim 10$ meV



Jungle of many-body and chemical effects

The **width of the peak** that serves as a signature of $C\nu B$ is defined by

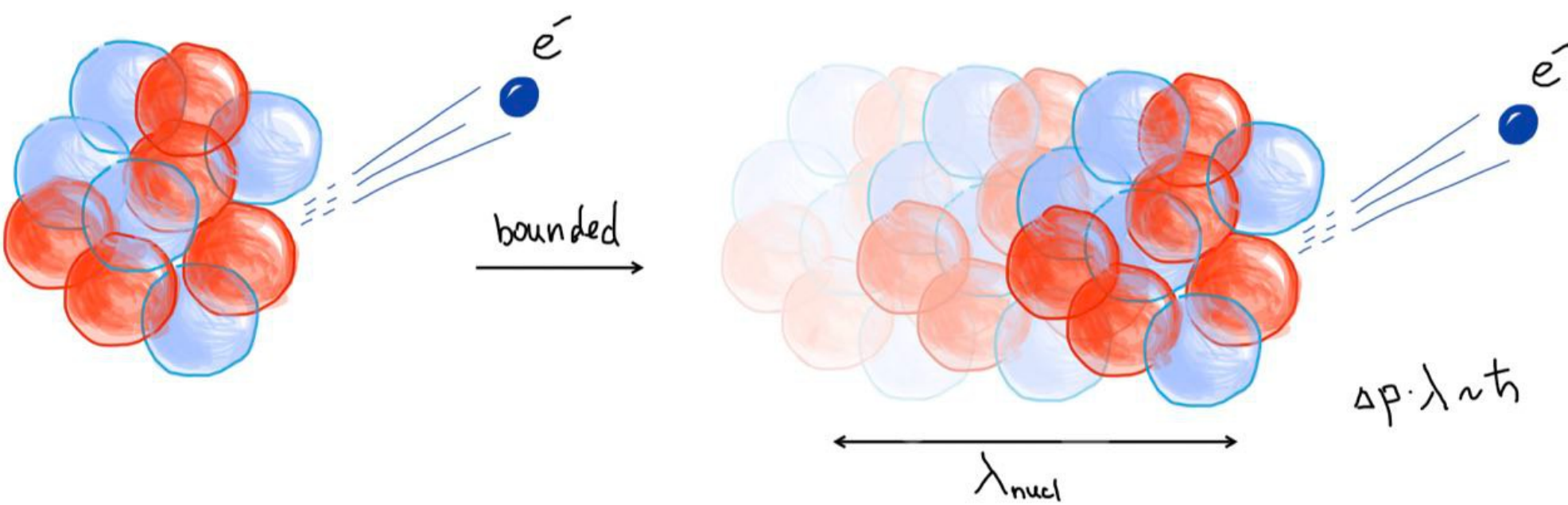
- The energy **resolution of the apparatus**
- **Intrinsic** physical effects



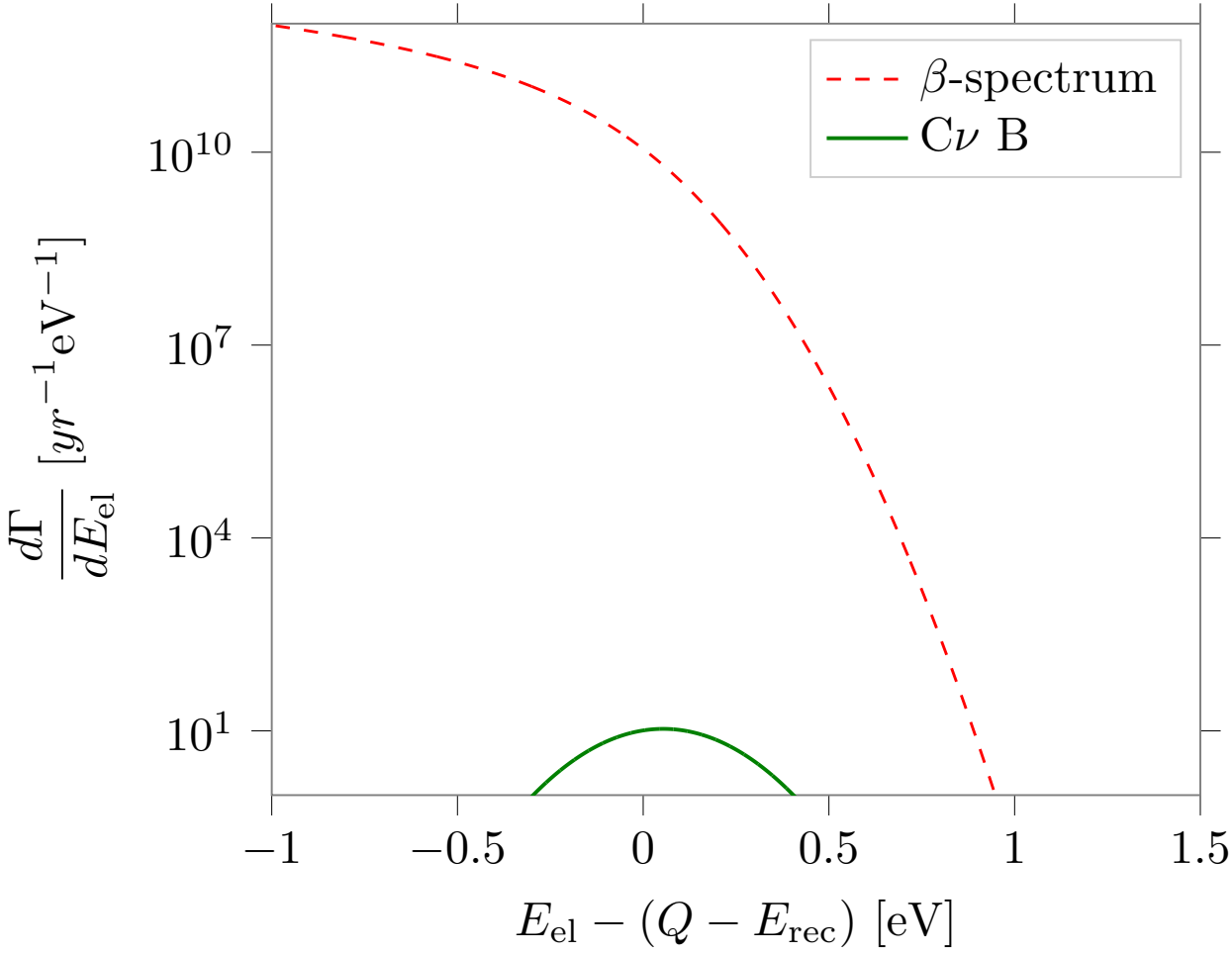
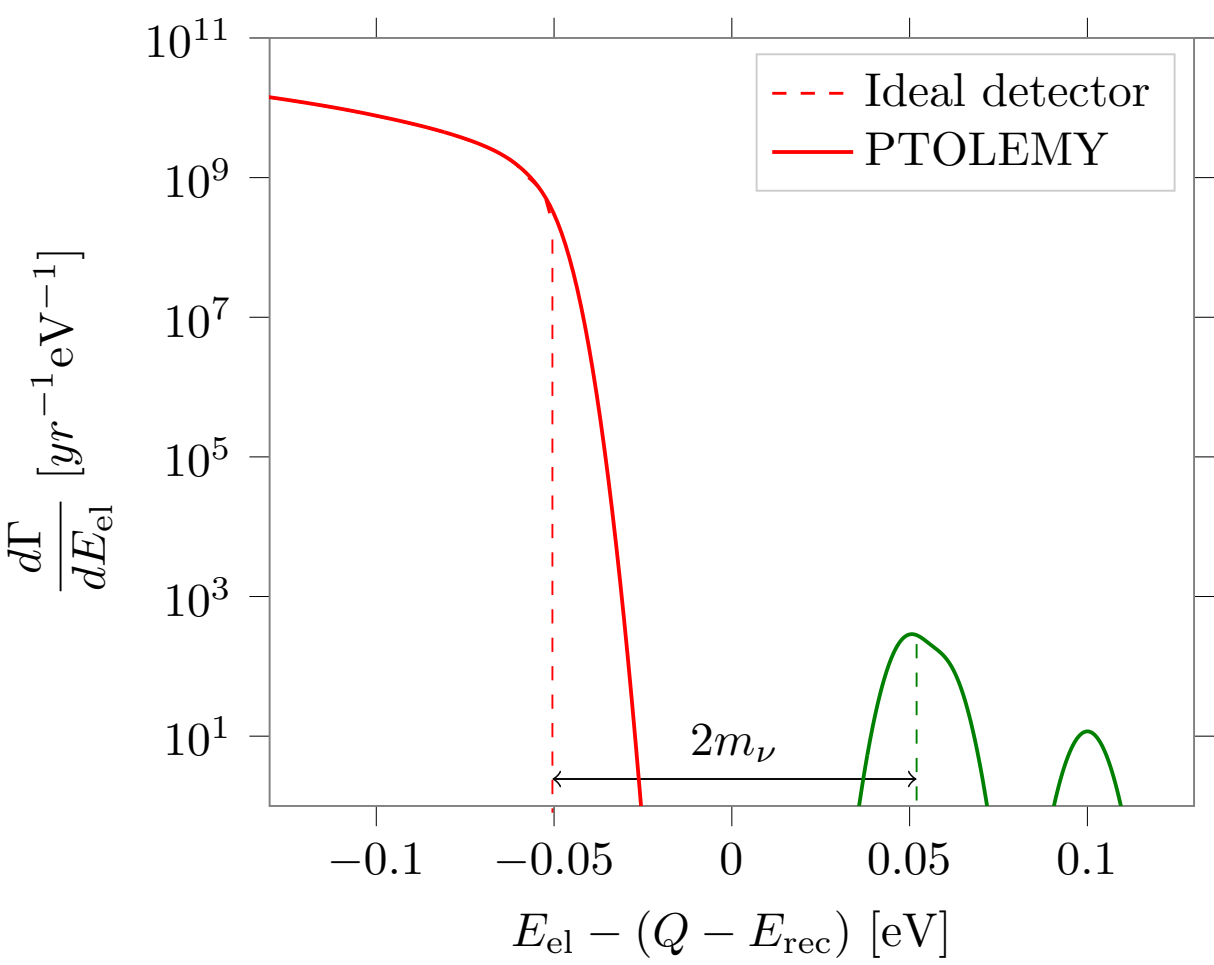
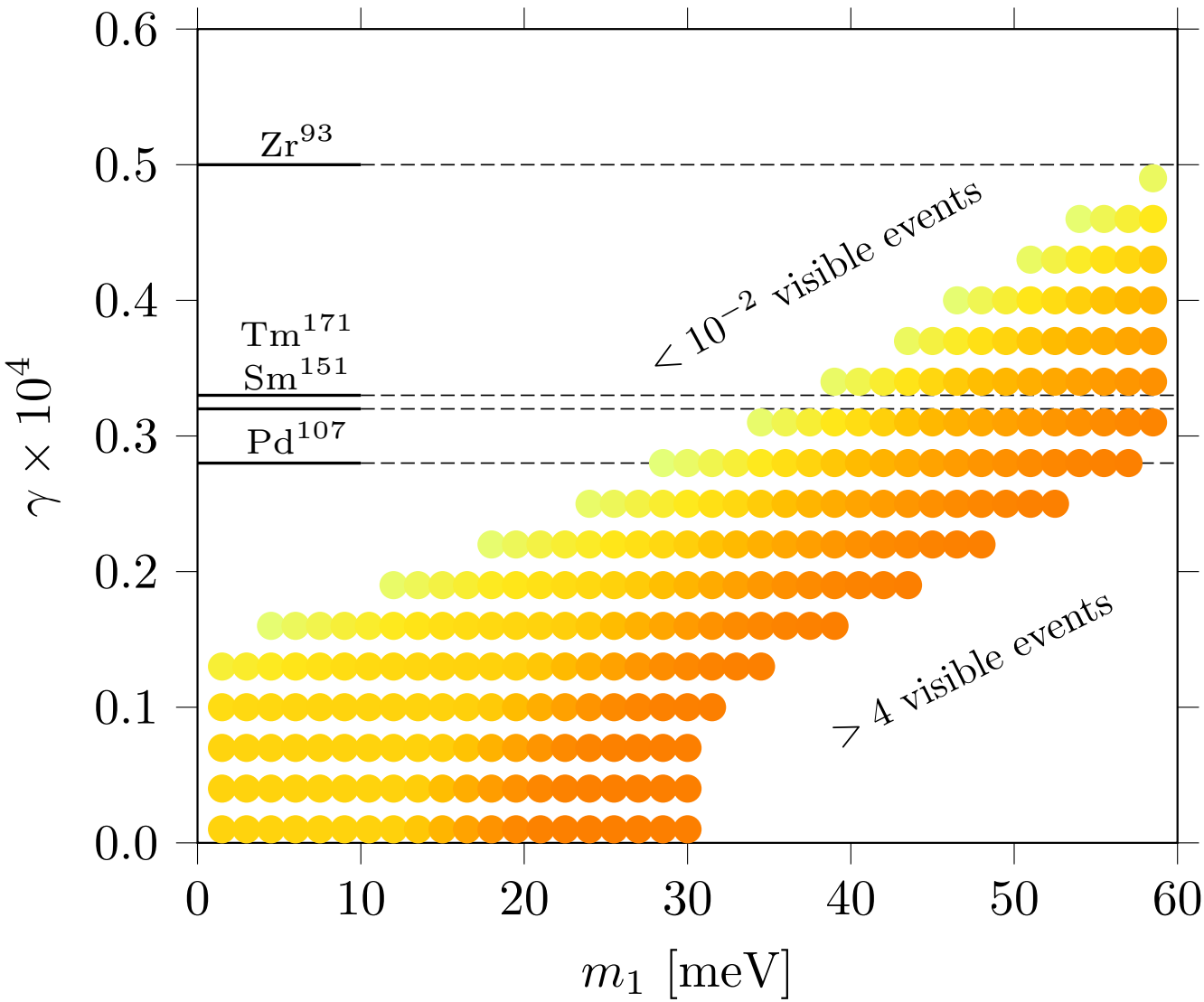
Heisenberg uncertainty

The uncertainty in energy of the emitted electron ΔE

- Is of the **order of 0.5 eV**
- Is 2 orders of magnitude greater than the resolution needed
- Weakly depends on the potential stiffness.
- Strongly depends on the radioactive nucleus.



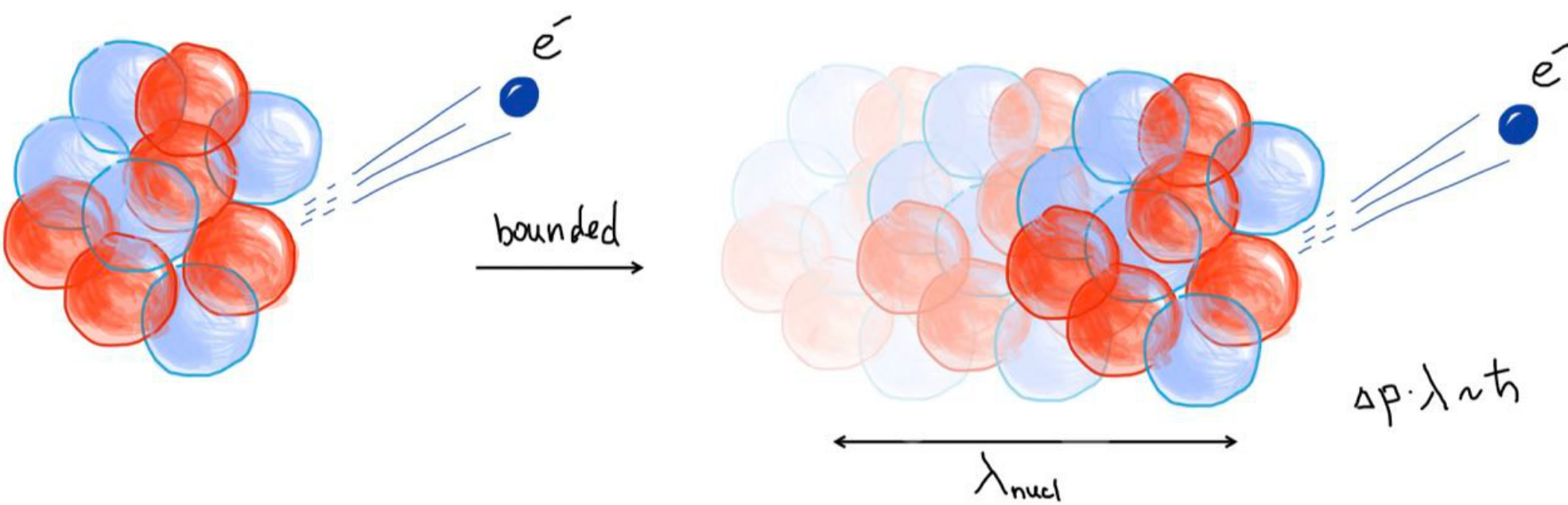
$$\frac{\Delta E}{\sqrt{\hbar m_e}} \approx \underbrace{\pi^{1/4}}_{\text{potential}} \underbrace{\sqrt{\frac{Q}{m_{\text{nucl}}^{3/2}}}}_{\text{nucleus}}$$



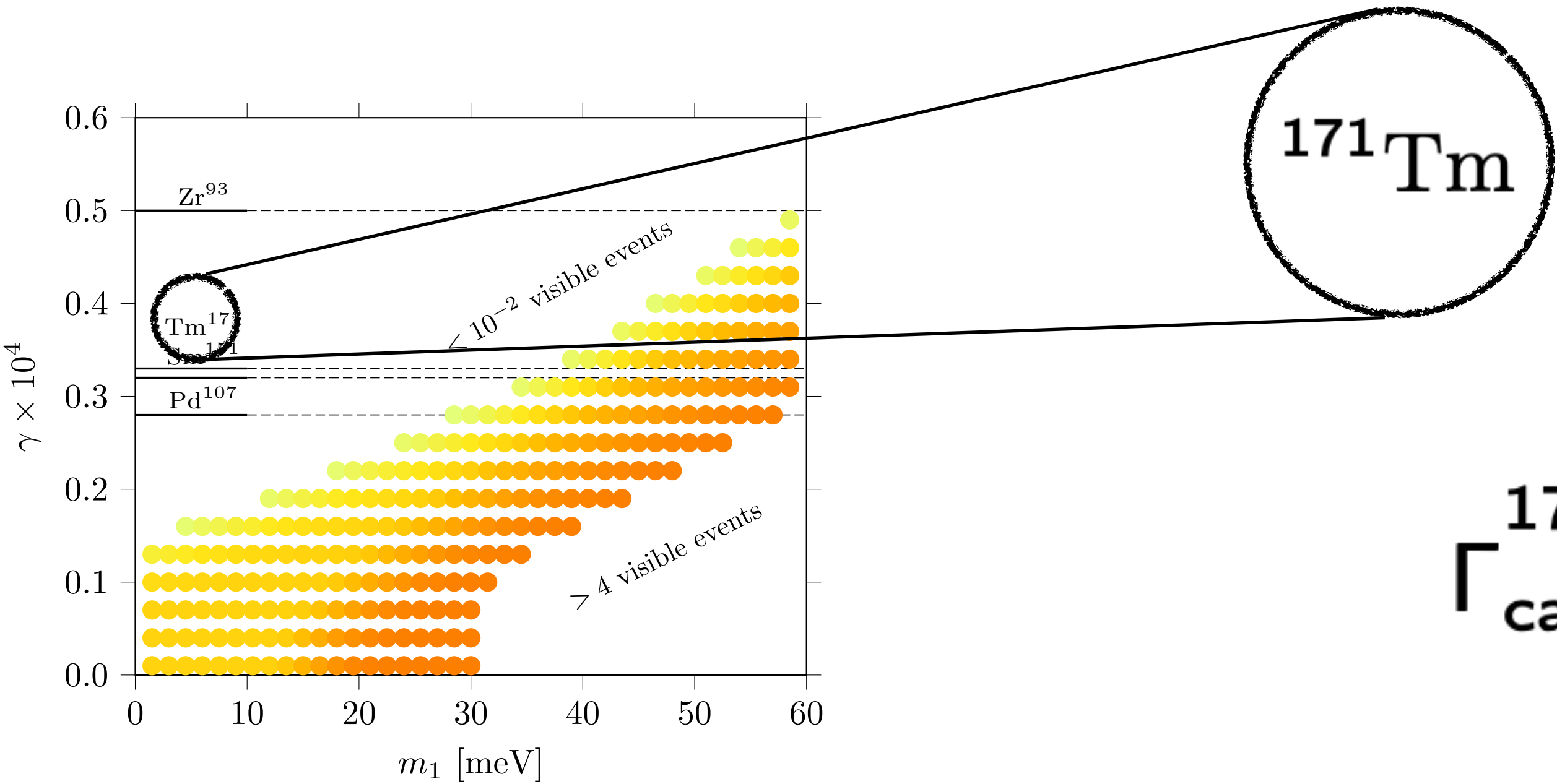
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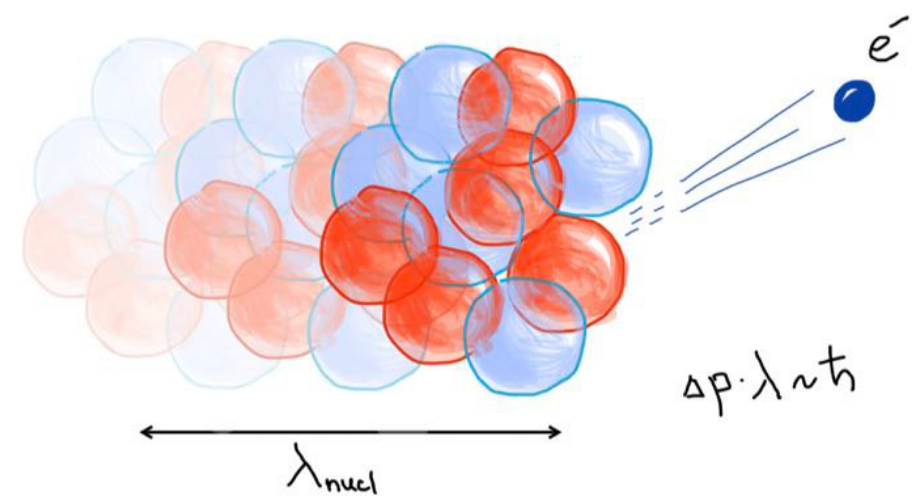


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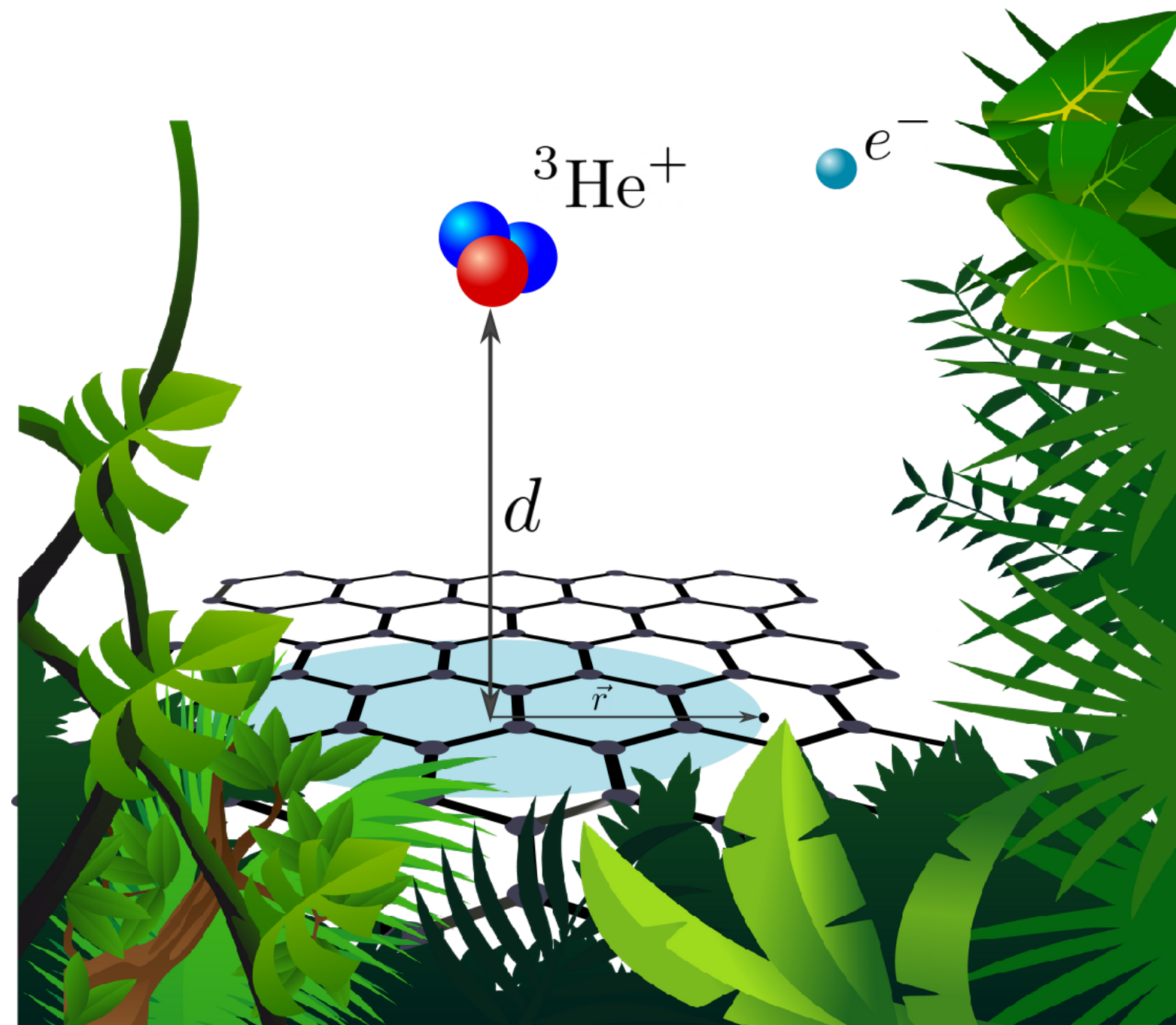


$$\Gamma_{\text{capture}}^{171\text{Tm}} = 3 \cdot 10^{-2} \Gamma_{\text{capture}}^3\text{H},$$

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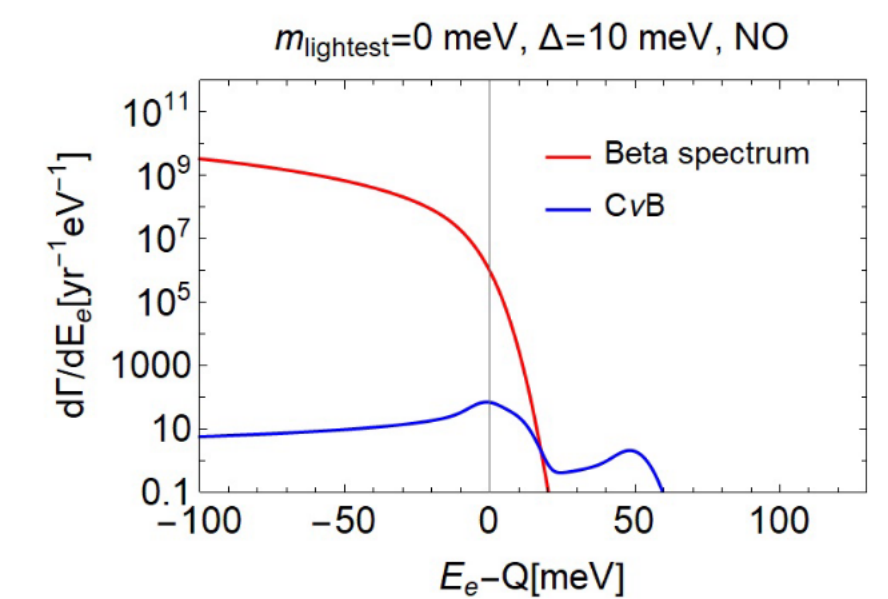
Zero-point motion



Electric effects

Etc...

Finite lifetime of the daughter ion



X-ray edge singularity

Generalised smeared spectrum

The bare spectrum of the β decay

$$\frac{d\Gamma^{(0)}}{dE_e} = \frac{4E_e p(E_e)}{(2\pi)^4 \hbar} \int E_\nu k(E_\nu) dE_\nu \left| \int dx j_{\text{lept}}^\mu(x, E_e, E_\nu) J_\mu^{\text{nucl}}(x) \right|^2 \delta(E_e + E_\nu - E_{\alpha_0}).$$

in the presence of the substrate becomes smeared

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where the correlation function $\mathcal{F}(\omega)$ represents the **finite lifetime** of the system

$$\frac{1}{2\pi} \int d\tau \langle \alpha_0, z | \hat{\chi} e^{i\tau \hat{H}_{z+1}^\alpha} \hat{\chi}^\dagger | \alpha_0, z \rangle e^{i\tau(E_e + E_\nu - E_{\alpha_0})} = \mathcal{F}(E_e + E_\nu - E_{\alpha_0})$$

and maybe can be probed on the experiment