

The Path to Predict the Axion Mass with large String Tension

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Klaer and Moore, arXiv:1708.07521 and arXiv:1707.05566

The Axion Quest!

Dark matter is still a mystery, it is

- ▶ **matter** : makes up 25% of the density of the Universe
- ▶ **dark** : interaction is feeble (except gravitationally)
- ▶ **cold** : almost pressureless

The Axion could be a likely candidate, would you like to investigate the Axion?

Yes

No



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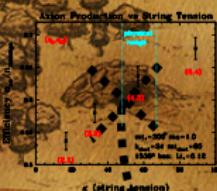
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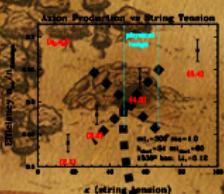




$$L_a = \partial^\mu \varphi^* \partial_\mu \varphi + \frac{\Lambda}{8} (2\varphi^* \varphi - f_a^2)^2 + \chi(T)$$

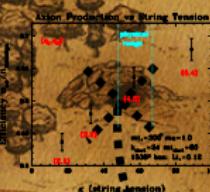
$$\begin{aligned} L(\Phi_1, \Phi_2, A_\mu) = & \frac{1}{4} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2 \\ & + \frac{\lambda}{8} [(2\Phi_1^* \Phi_1 - f^2) + (2\Phi_2^* \Phi_2 - f^2)] \\ & + [(\partial_\mu - i\eta_\mu e A_\mu) \Phi_1]^2 + [(\partial_\mu - i\eta_\mu e A_\mu) \Phi_2]^2 \end{aligned}$$





$$L_a = \partial^\mu \phi^* \partial_\mu \phi + \frac{\lambda}{8} (2\phi^* \phi - f)$$

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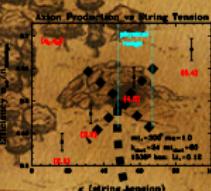


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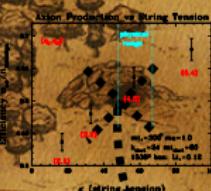
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z (string tension)

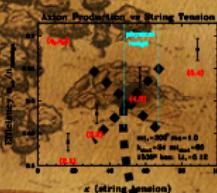


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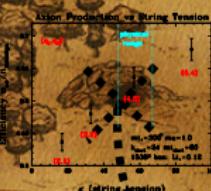


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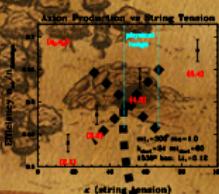


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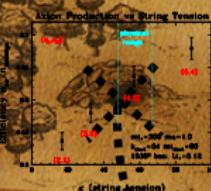
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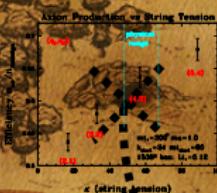
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$$L_a = \partial^\mu \phi^* \partial_\mu (\frac{\lambda}{\beta} (\phi^* \phi - f_c^2)^2 + \chi(T) \text{Re}\phi)$$

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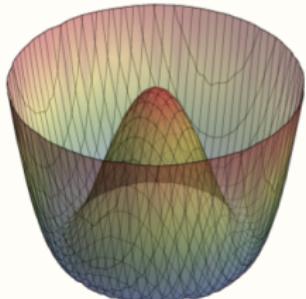


The Axion Lagrangian

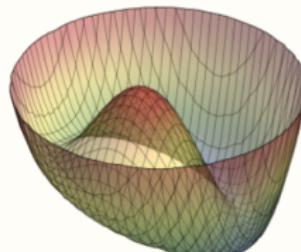
You found a Lagrangian

$$\mathcal{L}_a = \partial^\mu \phi^* \partial_\mu \phi + \frac{\lambda}{8} (2\phi^* \phi - f_a^2)^2 + \chi(T) \text{Re}\phi$$

$$T_{PQ} > T > T_c$$



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Would you like to take this Lagrangian to describe your Axion dynamics?

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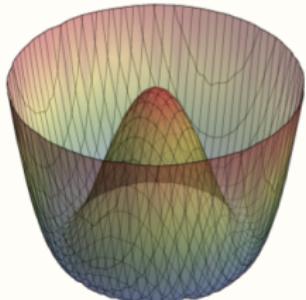
No

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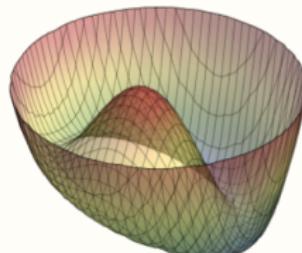
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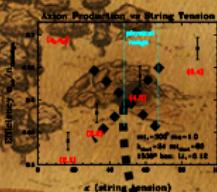
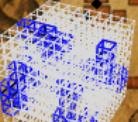
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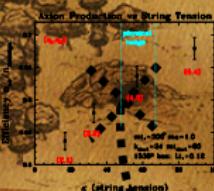
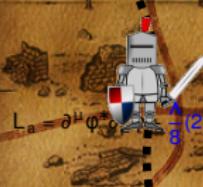
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$$L_a = \partial^\mu \phi_a$$

$$\frac{\Lambda}{8}$$

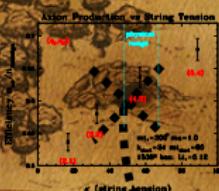
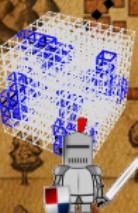
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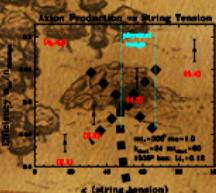






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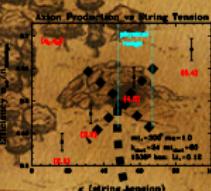
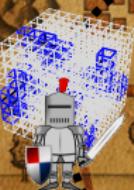
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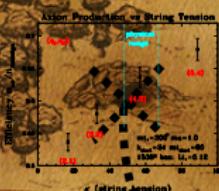
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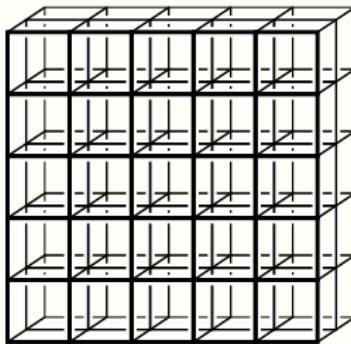
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Axion Cosmology solving on Lattice

You found a tool to calculate the EQM

- ▶ Put \mathcal{L}_a as classical field theory on **lattice**
- ▶ Simulation starts **after Inflation**
- ▶ ϕ is a complex field
- ▶ θ is chosen randomly
- ▶ Topological defects will arise
- ▶ Simulate with **Hubble drag**
- ▶ Count axions at the end



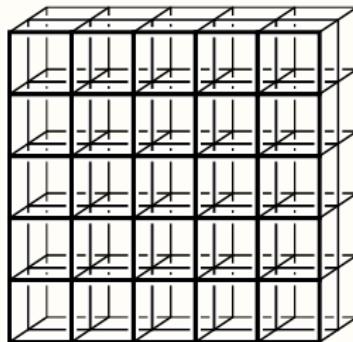
Solve it by hand

Use a computer

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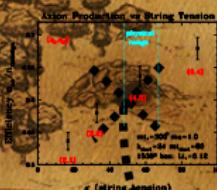
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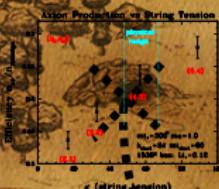
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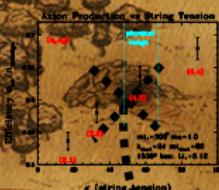
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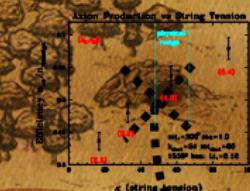
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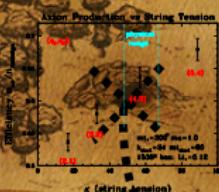
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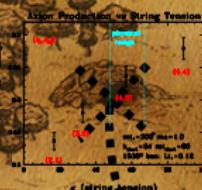
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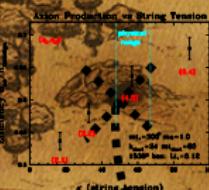
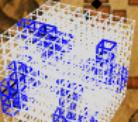
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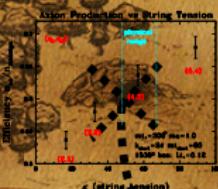
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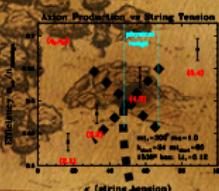
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$$L_a = \partial^\mu \phi^* \partial_\mu \phi + \frac{1}{8} (2\phi^* \phi - f_a^2)^2 + X(T)$$

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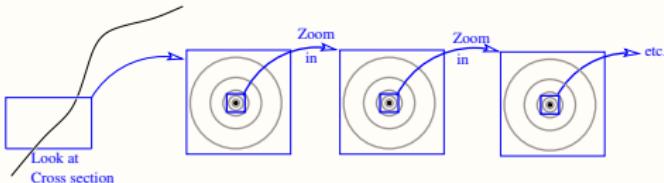


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Physical range not achievable!

$$E_{\text{str}} = \iiint r dr dz d\phi (\nabla \phi^* \nabla \phi \simeq f_a^2 / 2r^2) \simeq \pi l f_a^2 \int_{\sim f_a^{-1}}^{\sim H^{-1}} \frac{r dr}{r^2}$$



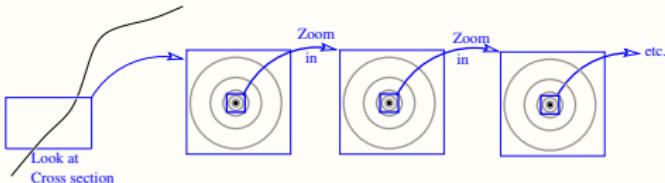
- ▶ **Log-large** string tension $T_{\text{str}} = \pi f_a^2 \ln(f_a/H)$
- ▶ $f_a/H \simeq 10^{30}$
- ▶ Equal energy in each x2 scale
- ▶ Scale range is $10^{30} \rightarrow \kappa = \ln(10^{30}) \simeq 70$, **achievable is**
 $\kappa = 6$

Find new methodes

Quit job

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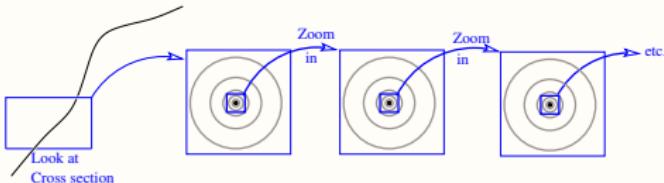
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[Find new methodes](#)

[Quit job](#)

Physical range not achievable!

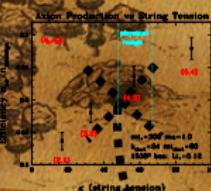
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Find new methodes

Quit job



$$L_a = \partial^\mu \phi^* \partial_\mu \phi - \frac{\lambda}{8} (2\phi^* \phi - f_c^2)^2 + \chi(T) \text{Re}\phi$$

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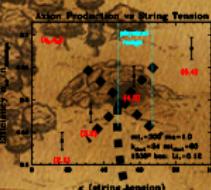
f_c

ω_1

ω_2

e

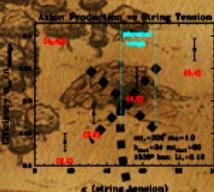
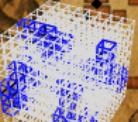
T



$$L_a = \partial^\mu \phi^* \partial_\mu \phi - \frac{\lambda}{8} (2\phi^* \phi - f^2)^2 + \chi(T) \text{Re}\phi$$

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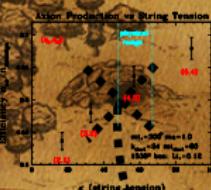
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$\frac{\lambda}{8}$

$(2\phi^* \phi - f_c^2)^2$

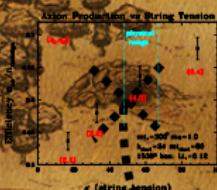
$+ \chi(T) \text{Re}\phi$



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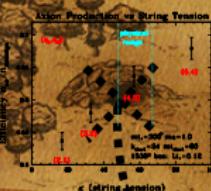




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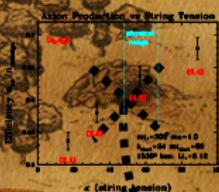




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$$L_a = \partial^\mu \varphi^* \partial_\mu \varphi - \frac{\Lambda}{8} (2\varphi^* \varphi - f_a^2)^2 + X(T)$$

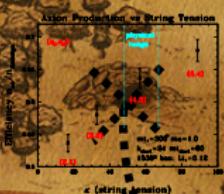
$$L(\phi_1, \phi_2, A_\mu) = \frac{1}{4} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2$$

$$+ \frac{\lambda}{8} [(2\varphi_1^* \varphi_1 - f^2) + (1$$

$$+ |(\partial_\mu - iq_1 e A_\mu) \Phi_1|^2$$

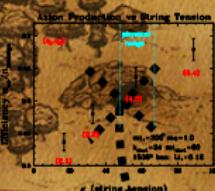
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$$L_a = \partial^\mu \phi^+ \partial_\mu \phi^- - \frac{\lambda}{8} (2\phi^* \phi - f_a^2)^2 + \chi(T)$$

$$L(\phi_1, \phi_2, A_\mu) = \frac{1}{4} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2 + \frac{\lambda}{8} [((2\phi_1^2 - 1)^2 + (2\phi_2^2 - 1)^2) + |(\partial_\mu - iq_1 e A_\mu)\phi_1|^2 + |(\partial_\mu - iq_2 e A_\mu)\phi_2|^2]$$



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Use a effective theory

Found an effective Lagrangian using two complex scalars and A_μ

$$\begin{aligned}\mathcal{L}(\phi_1, \phi_2, A_\mu) = & \frac{1}{4}(\partial_\mu A_\nu - \partial_\nu A_\mu)^2 \\ & + \frac{\lambda}{8}[(2\phi_1^* \phi_1 - f^2) + (2\phi_2^* \phi_2 - f^2)] \\ & + |(\partial_\mu - iq_1 e A_\mu)\phi_1|^2 + |(\partial_\mu - iq_2 e A_\mu)\phi_2|^2\end{aligned}$$

- ▶ Pick $q_1 \neq q_2$
- ▶ $T \simeq \pi f_a^2 (\kappa_{\text{eff}} + \kappa)$ with $\kappa_{\text{eff}} \simeq 2(q_1^2 + q_2^2)$
- ▶ $q_1 \theta_1 + q_2 \theta_2$ gauged, $q_2 \theta_1 - q_1 \theta_2$ global

Use new methode

Do nothing

Use a effective theory

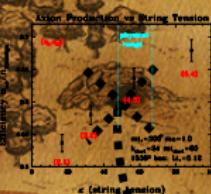
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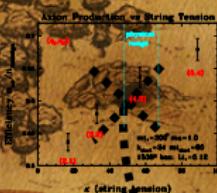


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$$\frac{\lambda}{8}$$



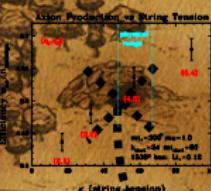
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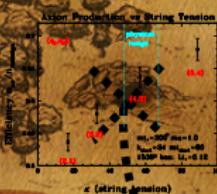




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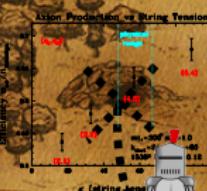
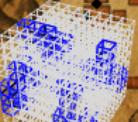




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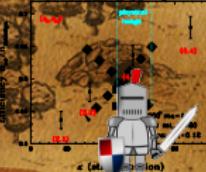
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Axion Production vs String Tension

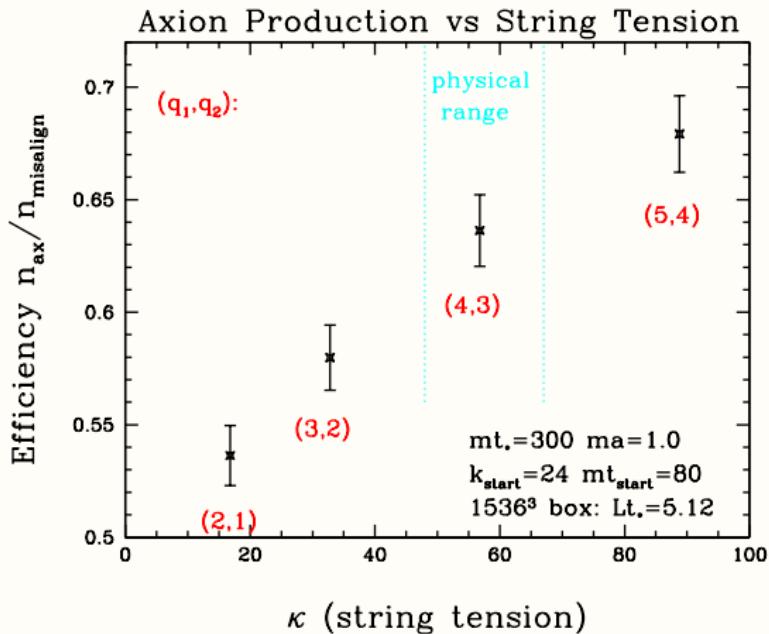


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Congratulation you did it!



- ▶ Axion production changes **only a little**
- ▶ $f_a = (2.21 \pm 0.29) \times 10^{11} \text{ GeV}$
- ▶ $m_a = 26.2 \pm 3.4 \mu\text{eV}$

Conclusion

- ▶ Axion can explain the **CP-problem** and the **dark matter**
- ▶ In early Universe the Axion dynamics are **string defects**
- ▶ With the **two-field-model** we get the string defect physics right
- ▶ We find **Axion mass** $m_a = 26.2 \pm 3.4 \mu\text{eV}$
 - ▶ Assuming Axions make all DM
 - ▶ θ is chosen randomly

Thank you for your Attention!