

Axion couplings to EW gauge bosons

Work in preparation

Pablo Quílez Lasanta

In collaboration with: *G. Alonso-Álvarez, M.B. Gavela.*

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Instituto de
Física
Teórica
UAM-CSIC



elusi**v**es in**v**isiblesPlus

The axion solution

Strong CP problem:

Why is it so small?



$$\mathcal{L} \supset \bar{\theta} \frac{\alpha_s}{8\pi} G\tilde{G}$$

$$\bar{\theta} \lesssim 10^{-10}$$

→ If $\bar{\theta}$ were a scalar field, its vev would be zero

$$\bar{\theta} \frac{\alpha_s}{8\pi} G\tilde{G} \longrightarrow \left(\bar{\theta} - \frac{a}{f_a} \right) \frac{\alpha_s}{8\pi} G\tilde{G}$$

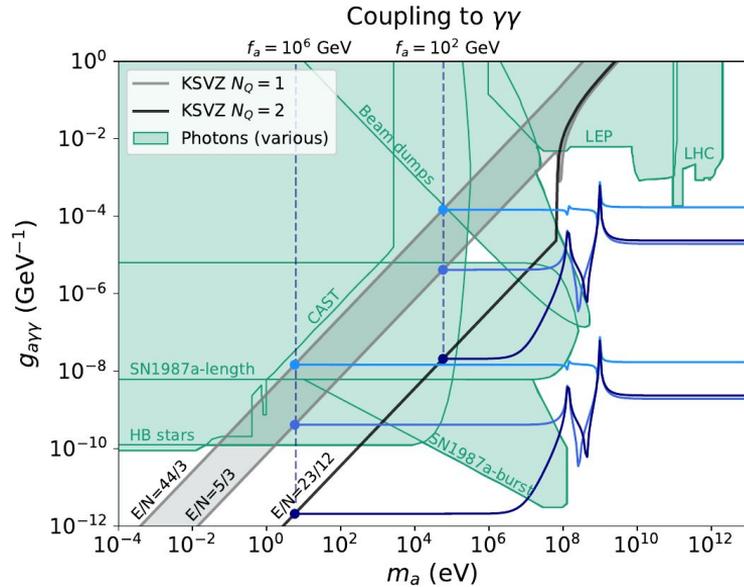
[Peccei+Quinn 77]

[Weinberg 78]

[Wilczek 78]

Axion phenomenology: photons

The most studied phenomenology of the axion: coupling to photons



$$\mathcal{L} \supset \frac{1}{4} g_{a\gamma\gamma} a F \tilde{F}$$

$$g_{a\gamma\gamma} \propto \frac{1}{f_a} \implies g_{a\gamma\gamma} \propto m_a$$

$$g_{a\gamma\gamma} = -\frac{1}{2\pi f_a} \alpha_{\text{em}} \left(\frac{E}{N} - 1.92(4) \right)$$

Motivation

Can axion couplings to W and Z -bosons help covering the axion parameter space?

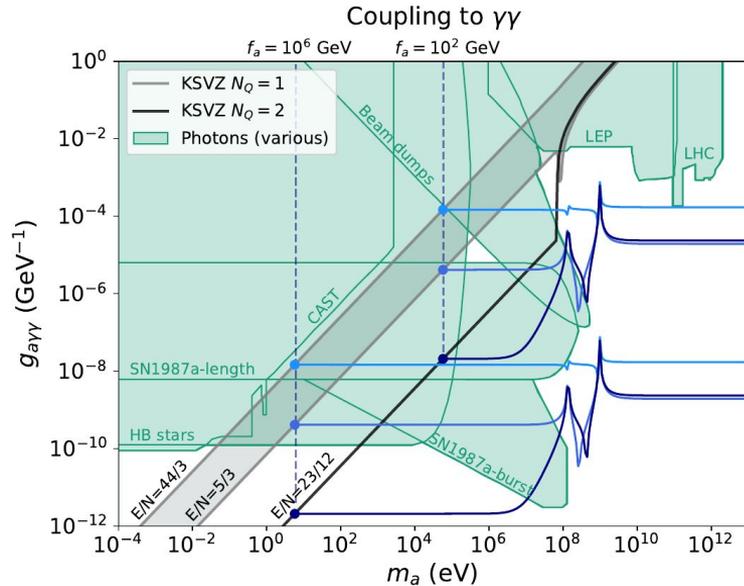
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Can axion couplings to W and Z -bosons help covering the axion parameter space?

Are couplings to photons enough?

Axion phenomenology: photons

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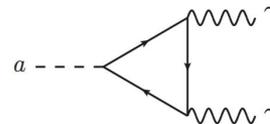
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Model dependent

Model independent



Where does it come from?

The axion mass matrix

There are two pseudoscalars that couple to the anomaly: the axion and the η :

$$\frac{\alpha}{8\pi} \left(2\frac{\eta_0}{f_\pi} - \frac{a}{f_a} \right) \tilde{G}G \longrightarrow \frac{1}{2} \Lambda_{QCD}^4 \left(2\frac{\eta_0}{f_\pi} + \frac{a}{f_a} \right)^2$$

$$M_{\{\pi_3, \eta_0, a\}}^2 = 4 \begin{pmatrix} B_0(m_u + m_d) & B_0(m_u - m_d) & 0 \\ B_0(m_u - m_d) & 4K/f_\pi + B_0(m_u + m_d) & 2K/(f_\pi f_a) \\ 0 & 2K/(f_\pi f_a) & K/f_a^2 \end{pmatrix}$$

The physical axion is a (model-independent) combination of the pion and the eta:

$$a_{phys} \simeq \hat{a} - \frac{f_\pi}{2f_a} \frac{m_d - m_u}{m_u + m_d} \pi_3 - \frac{f_\pi}{2f_a} \eta_0$$

Axion couplings bellow confinement

$$\mathcal{L} \supset \frac{\alpha_s}{8\pi} \frac{\hat{a}}{f_a} G\tilde{G} + \frac{1}{4} g_{aWW}^0 \hat{a} W\tilde{W} + \frac{1}{4} g_{aZZ}^0 \hat{a} Z\tilde{Z} + \frac{1}{4} g_{a\gamma\gamma}^0 \hat{a} F\tilde{F} + \frac{1}{4} g_{a\gamma Z}^0 \hat{a} F\tilde{Z}$$

$$g_{aXX} = g_{aXX}^0 + \theta_{a\pi} g_{\pi XX} + \theta_{a\eta'} g_{\eta' XX}$$

$$g_{a\gamma\gamma} = -\frac{1}{2\pi f_a} \alpha_{\text{em}} \left(\frac{E}{N} - \frac{2m_u + 4m_d}{3m_u + m_d} \right),$$

$$g_{aWW} = -\frac{1}{2\pi f_a} \frac{\alpha_{\text{em}}}{s_w^2} \left(\frac{L}{N} - \frac{3}{4} \right),$$

$$g_{aZZ} = -\frac{1}{2\pi f_a} \frac{\alpha_{\text{em}}}{s_w^2 c_w^2} \left(\frac{Z}{N} - \frac{11s_w^4 + 9c_w^4}{12} - \frac{s_w^2 (s_w^2 - c_w^2)}{2} \frac{m_d - m_u}{m_u + m_d} \right),$$

$$g_{a\gamma Z} = -\frac{1}{2\pi f_a} \frac{\alpha_{\text{em}}}{s_w c_w} \left(\frac{2K}{N} - \frac{9c_w^2 - 11s_w^2}{6} - \frac{1}{2} (c_w^2 - 3s_w^2) \frac{m_d - m_u}{m_u + m_d} \right).$$

~ 1.92

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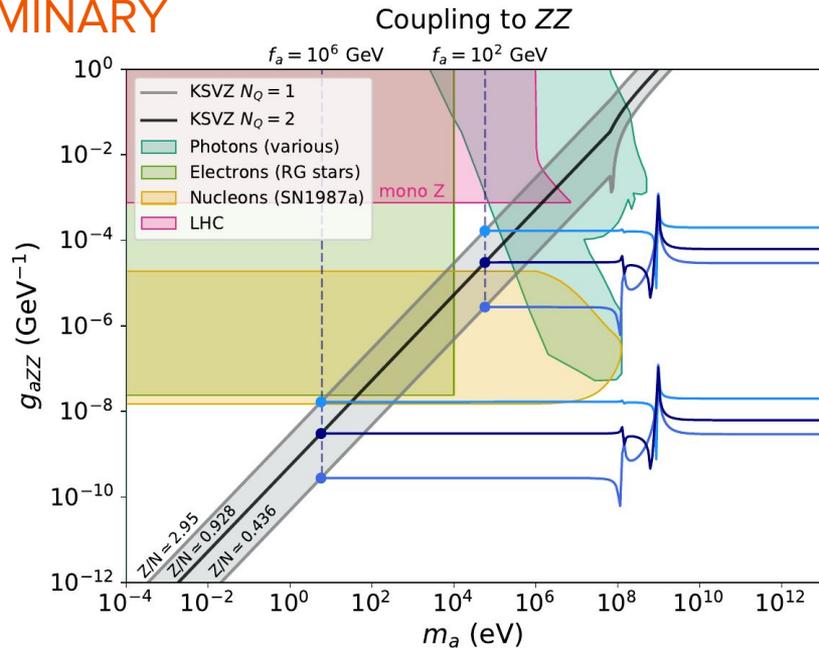
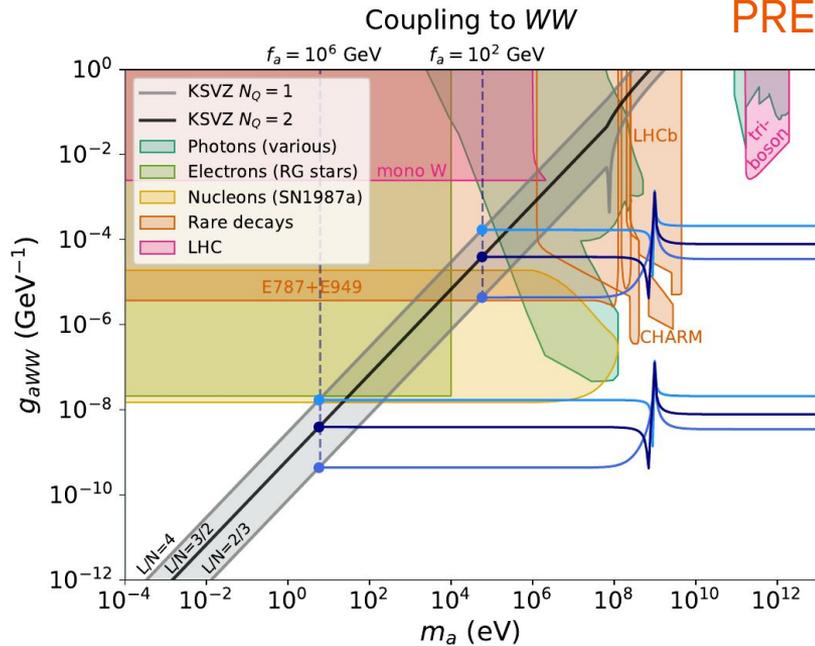
$$g_{a\gamma Z} = -\frac{1}{2\pi f_a} \frac{\alpha_{\text{em}}}{s_w c_w} \left(\frac{2K}{N} - \frac{9c_w^2 - 11s_w^2}{6} - \frac{1}{2}(c_w^2 - 3s_w^2) \frac{m_d - m_u}{m_u + m_d} \right).$$

Model independent
(present for $m_a < \Lambda_{\text{QCD}}$)

Model dependent

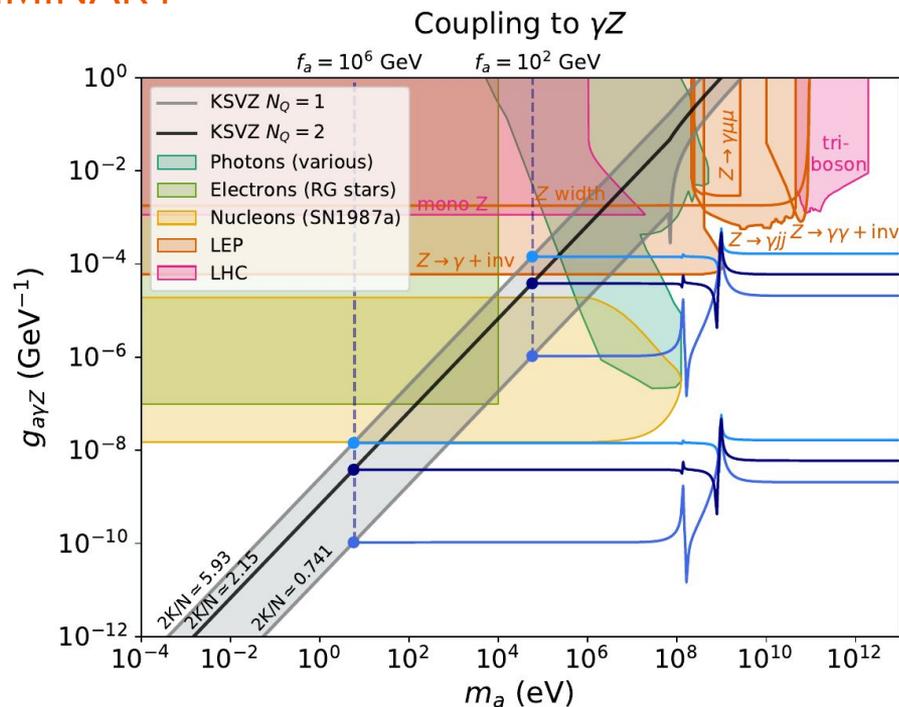
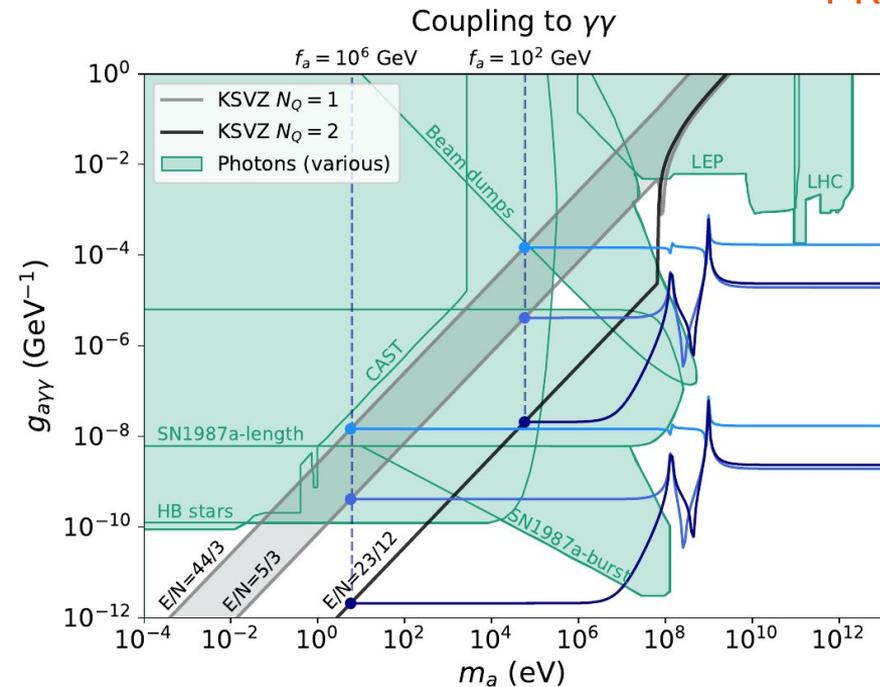
Axion phenomenology: g_{aWW} , g_{aZZ} and $g_{a\gamma Z}$

PRELIMINARY



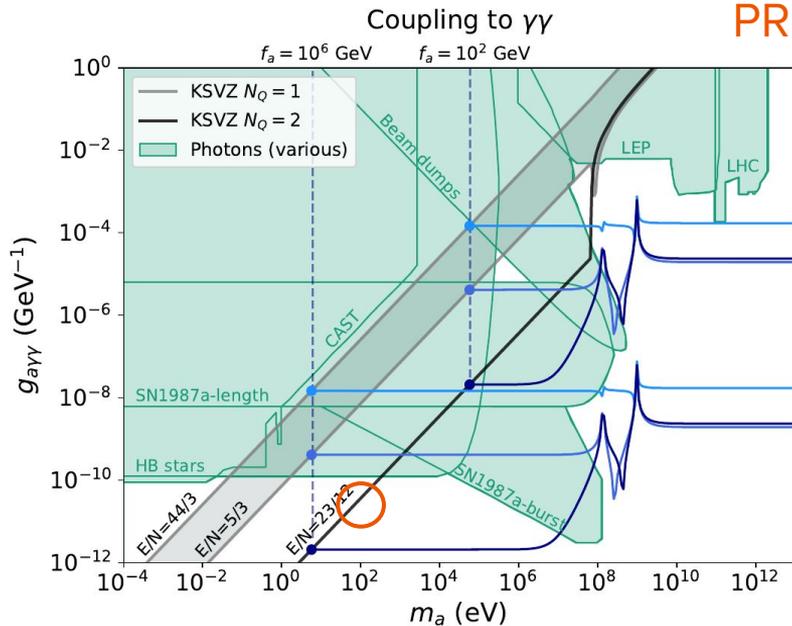
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PRELIMINARY

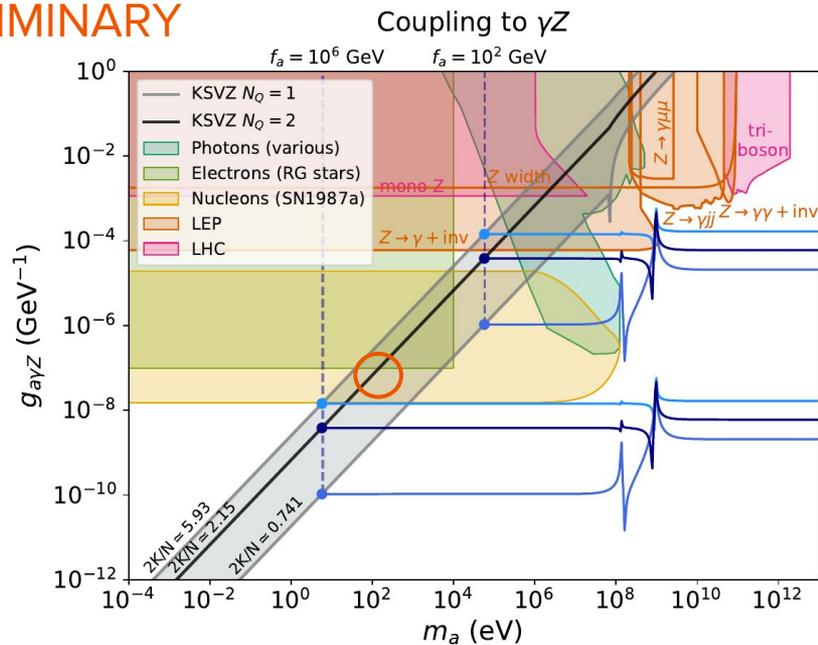


Applications: photophobic axion

PRELIMINARY



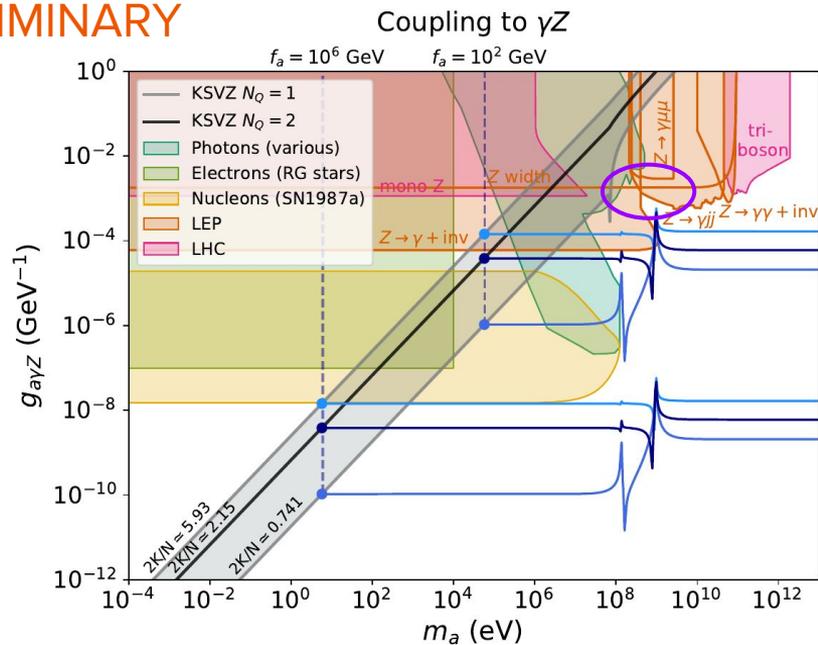
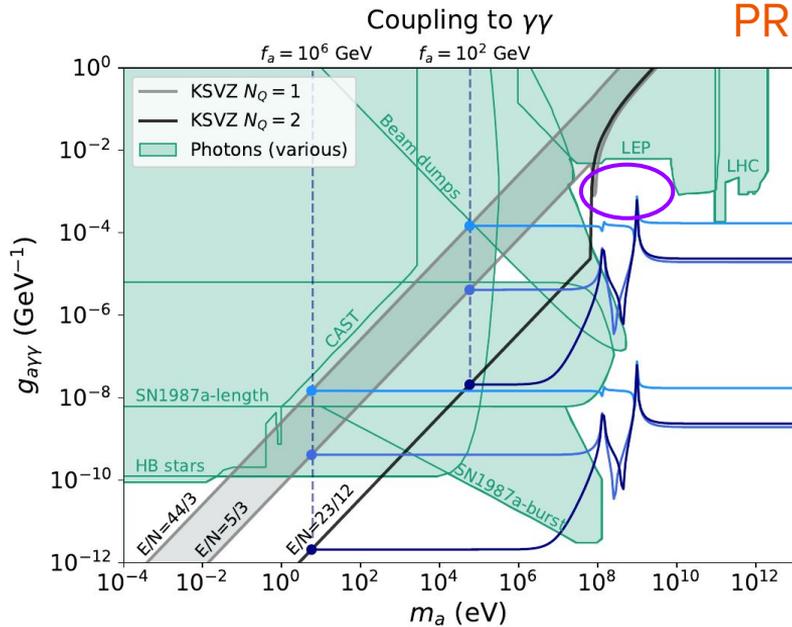
$$g_{a\gamma\gamma} = -\frac{1}{2\pi f_a} \alpha_{\text{em}} \left(\frac{E}{N} - \frac{2}{3} \frac{m_u + 4m_d}{m_u + m_d} \right)$$



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Applications: heavy axion with LEP

PRELIMINARY



Conclusions

- We have computed the model-independent contribution of the axion couplings to the heavy EW bosons for both, the standard and heavy axion.
- We have explored the phenomenology associated to these couplings.

Photophobic and heavy axion ($m_a \sim 0.1-100$ GeV) can be better probed through their EW boson couplings.

Thank you!

Back up I: heavy axion

$$\delta\mathcal{L}_a = \frac{1}{2} M^2 \hat{a}^2$$

$$\theta_{a\pi} \simeq -\frac{f_\pi}{2f_a} \frac{m_d - m_u}{m_u + m_d} \frac{1}{1 - \frac{M^2}{m_\pi^2}} \frac{1}{1 - \frac{M^2}{m_{\eta'}^2}}, \quad \theta_{A\eta'} \simeq -\frac{f_\pi}{2f_a} \frac{1}{1 - \frac{M^2}{m_{\eta'}^2}}$$

$$g_{aWW} = -\frac{1}{2\pi f_a} \frac{\alpha_{\text{em}}}{s_w^2} \left(\frac{L}{N} - \frac{3}{4} \frac{1}{1 - \left(\frac{M}{m_{\eta'}}\right)^2} \right),$$

$$g_{a\gamma\gamma} = -\frac{1}{2\pi f_a} \alpha_{\text{em}} \left[\frac{E}{N} - \frac{1}{1 - \left(\frac{M}{m_{\eta'}}\right)^2} \left(\frac{5}{3} + \frac{m_d - m_u}{m_u + m_d} \frac{1}{1 - \left(\frac{M}{m_\pi}\right)^2} \right) \right],$$

$$g_{aZZ} = -\frac{1}{2\pi f_a} \frac{\alpha_{\text{em}}}{s_w^2 c_w^2} \left[\frac{Z}{N} - \frac{1}{1 - \left(\frac{M}{m_{\eta'}}\right)^2} \left(\frac{11s_w^4 + 9c_w^4}{12} - \frac{s_w^2(s_w^2 - c_w^2)}{2} \frac{m_d - m_u}{m_u + m_d} \frac{1}{1 - \left(\frac{M}{m_\pi}\right)^2} \right) \right],$$

$$g_{a\gamma Z} = -\frac{1}{2\pi f_a} \frac{\alpha_{\text{em}}}{s_w c_w} \left[\frac{2K}{N} - \frac{1}{1 - \left(\frac{M}{m_{\eta'}}\right)^2} \left(\frac{9c_w^2 - 11s_w^2}{6} - \frac{1}{2}(c_w^2 - 3s_w^2) \frac{m_d - m_u}{m_u + m_d} \frac{1}{1 - \left(\frac{M}{m_\pi}\right)^2} \right) \right]$$

Back up I: RGE in the phenomenology

Effective couplings to fermions appear at one loop. Following [63], we get

$$\begin{aligned}
 \frac{c_f^{\text{eff}}}{f_a} &= \frac{c_f(\Lambda)}{f_a} - \frac{3}{4} \left(\frac{\alpha_L}{4\pi} \frac{3}{4} g_{aWW} + \frac{\alpha_Y}{4\pi} (Y_{fL}^2 + Y_{fR}^2) g_{aBB} \right) \log \frac{\Lambda^2}{m_W^2} - \frac{3}{2} Q_f^2 \frac{\alpha_{\text{em}}}{4\pi} g_{a\gamma\gamma} \log \frac{m_W^2}{m_f^2} \\
 &= \frac{c_f(\Lambda)}{f_a} - \frac{3}{2} \frac{\alpha_{\text{em}}}{4\pi} Q_f^2 g_{a\gamma\gamma} \log \frac{\Lambda^2}{m_f^2} - \frac{9}{16} \frac{\alpha_L}{4\pi} g_{aWW} \log \frac{\Lambda^2}{m_W^2} \\
 &\quad - \frac{3}{4} \frac{\alpha_Z}{4\pi} \left(\frac{3}{4} c_W^4 + (Y_{fL}^2 + Y_{fR}^2) s_W^4 \right) g_{aZZ} \log \frac{\Lambda^2}{m_W^2} \\
 &\quad - \frac{3}{4} \frac{\alpha_{\gamma Z}}{4\pi} \left(\frac{3}{4} c_W^2 - (Y_{fL}^2 + Y_{fR}^2) s_W^2 \right) g_{a\gamma Z} \log \frac{\Lambda^2}{m_W^2}.
 \end{aligned}$$

The photon coupling also receives 1-loop corrections in the presence of fermion or gauge boson couplings. The contribution can be expressed as

$$g_{a\gamma\gamma}^{\text{eff}} = g_{a\gamma\gamma}(\Lambda) + \sum_f N_C^f Q_f^2 \frac{\alpha_e m}{\pi} \frac{2c_f}{f_a} B_1 \tau_f + 2 \frac{\alpha_e m}{\pi} g_{aWW} B_2(\tau_W), \quad (3.12)$$