

Spontaneous CP violation and a new Higgs doublet below 500 GeV

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Why new physics? (via direct/indirect search)

- **Naturalness:** A long-standing guiding principle.
- Field theory solution to the hierarchy puzzle
= new physics not much heavier than TeV.
- A benchmark example: SUSY.
- Is naturalness **physically satisfying**?

Motivation and Introduction

SUSY is a decoupling theory

Question: **Can** the LHC rule out supersymmetry?

Answer: No. Supersymmetry is an example (one of many!) of a decoupling theory; the more you raise the masses of the new particles, the better it agrees with the Standard Model.

From S. Martin
pre-SUSY 2021

Naturalness require:

$$\Lambda_{\text{SUSY}} < 500 \text{ GeV}, 1 \text{ TeV}, 3 \text{ TeV}, 10 \text{ TeV}, \dots ?$$

Not too heavy is a vague statement.

Motivation for NP other than naturalness?

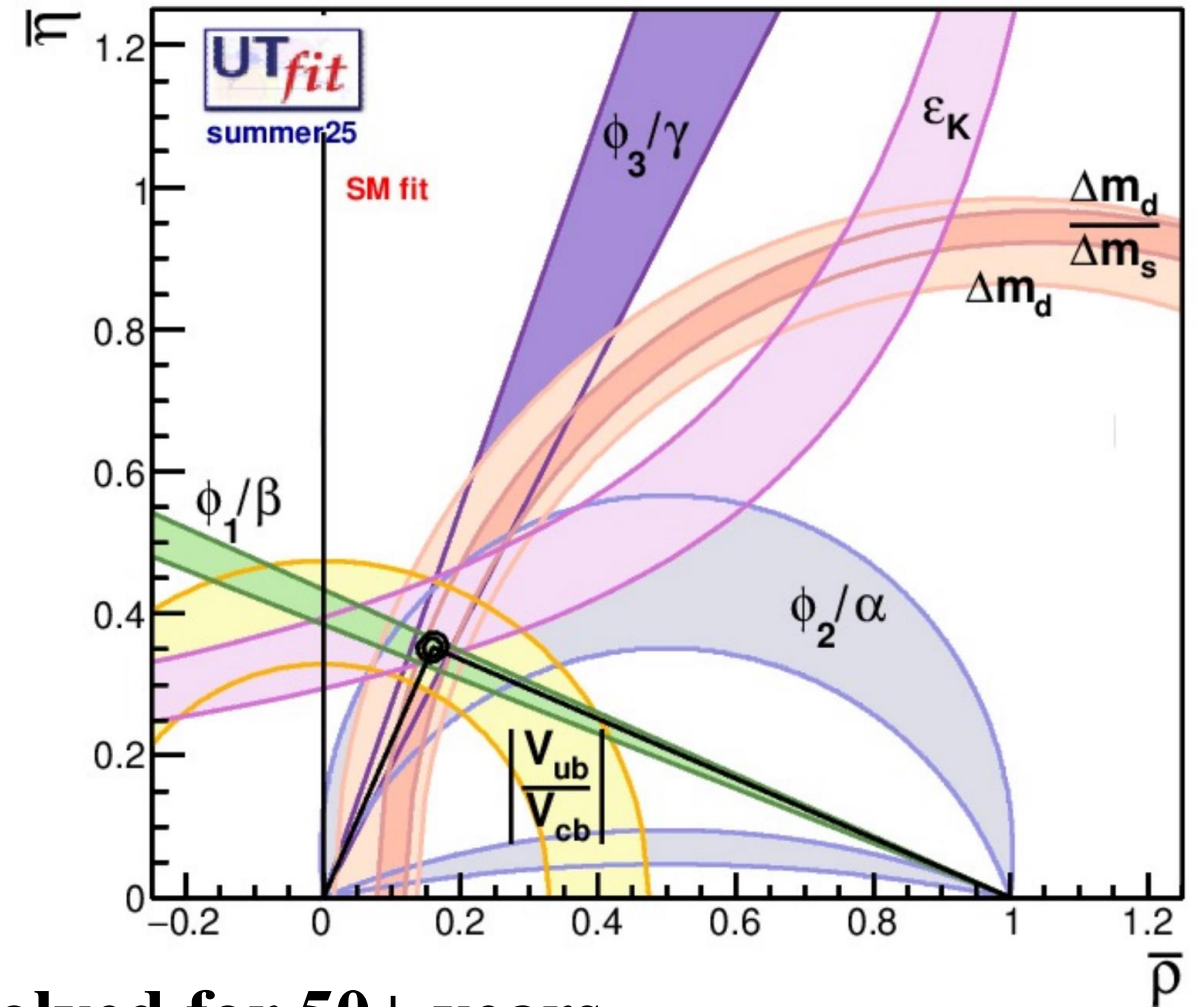
Spontaneous CP violation

Standard origin of CPV: KM phase

Kobayashi, Maskawa. '72

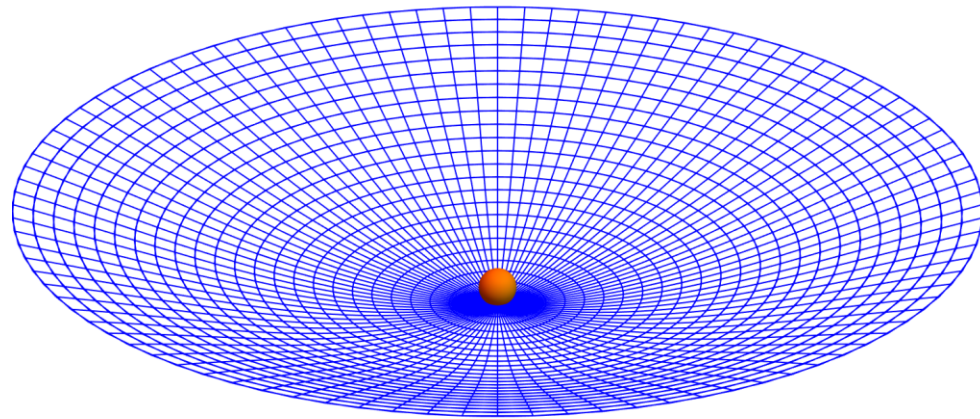
However:

- CP can serve as a fundamental symmetry.
- Vacuum solution leads to CPV.
- KM phase may arise from complex VEVs. T. D. Lee '73.

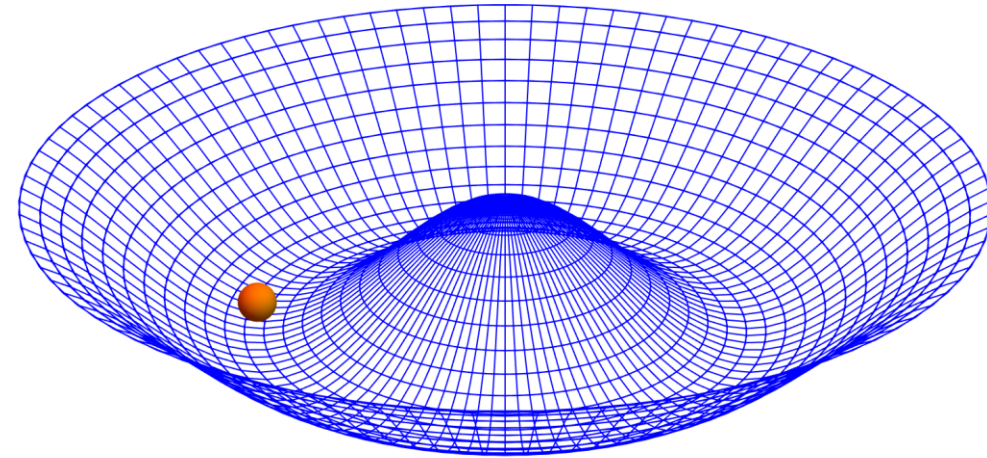


Explicit or Spontaneous CPV? Unsolved for 50+ years.

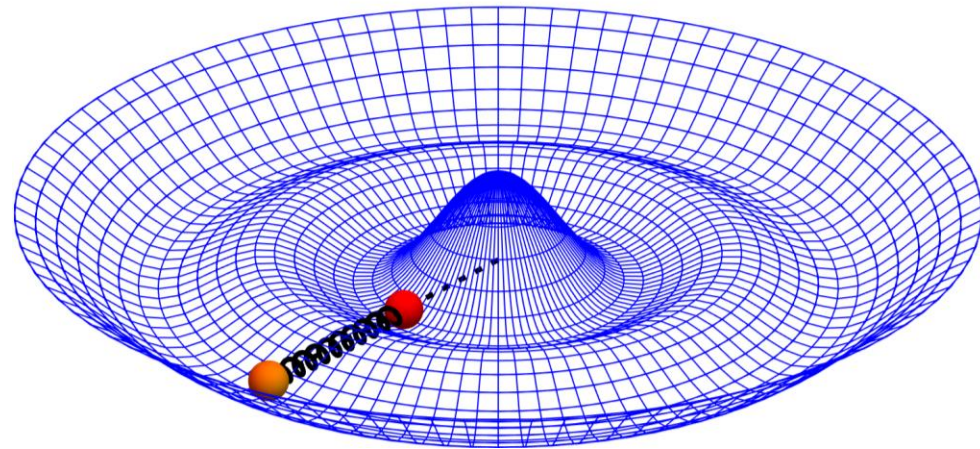
Spontaneous CP violation



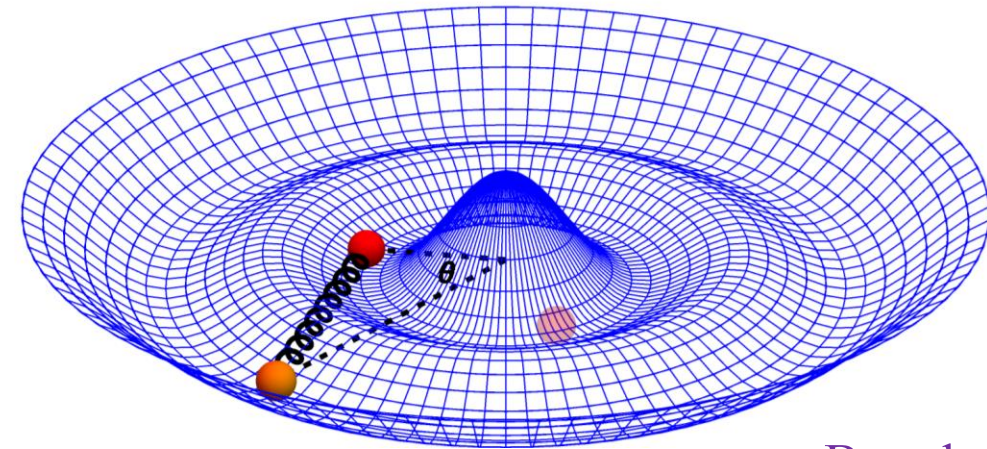
(a) No SSB



(b) SM SSB: Mexican Hat



(c) Unstable vacuum



(d) Spontaneous CPV

Based on the planar
diagrams in Lee '73

Spontaneous CP violation

The initial proposal of T. D. Lee '73: 2HDM

Two degenerate EW vacua v_i, v_i^*

→ **Two light Higgs Doublets $\lesssim 545$ GeV.**

→ **SM-like fine-tuning once again.**

Mohapatra, Senjanovic. '83;
Nebot, Botella, Branco. '19;
Nierste, Tabet, Ziegler, '20;
Miró, Nebot, Queiroz. '24.

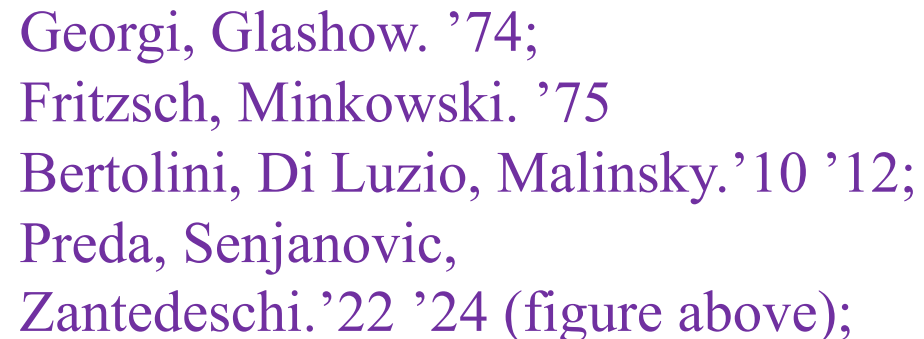
$$m_h^2 + m_H^2 + m_A^2 \lesssim M_{\text{LQT}}^2 = (700 \sim 800 \text{ GeV})^2.$$

No decoupling limit
Even if $\delta_{CP} \ll 1$!

A good guiding principle for NP? Still needs:

- To ensure no SCPV at higher scale.
- To constrain the couplings.

- Babu, Mohapatra. '89;
Foot, Lew, Volkas. '93;
Bressi, et al. '11;
Herms, Ruhdorfer. '24.



The theory: Minimal $SO(10) \times CP$

$$\begin{aligned}
 SO(10) \times CP & \xrightarrow[M_{GUT}]{\langle(1,1,1,0)\rangle \in 45_H} SU(3)_C \times SU(2)_L \times SU(2)_R \times U(1)_{B-L} \times CP \\
 & \xrightarrow[M_R]{\langle(1,1,3,1)\rangle \in 126_H} SU(3)_C \times SU(2)_L \times U(1)_Y \times CP \\
 & \xrightarrow[M_W]{\langle(1,2,2,0)\rangle \in 126_H, 10_H} SU(3)_c \times U(1)_{EM}.
 \end{aligned}$$

$$45_H \rightarrow 45_H, \quad 126_H \rightarrow \overline{126}_H, \quad 10_H \rightarrow 10_H^*$$

- Only real couplings in the limit of CP invariance.
- 45_H real, 126_H phase unphysical: **No high-scale SCPV**

The theory: Minimal $\text{SO}(10) \times \text{CP}$

Yukawa Sector: constrained by $\text{SO}(10)$:

$$-\mathcal{L}_Y = Y_{10} 16_F 10_H 16_F + \tilde{Y}_{10} 16_F 10_H^* 16_F \\ + Y_{126} 16_F \overline{126}_H 16_F + \text{h.c.}$$

Successful
numerical fit:
Patel '23.

$Y_{10}, \tilde{Y}_{10}, Y_{126}$ are **real, symmetric**, linear combinations of M_U, M_D, M_E .

$$M_D = D^* m_D D^\dagger, \quad M_U = U^* m_U U^\dagger, \quad M_E = E^* m_E E^\dagger; \\ V_{\text{CKM}} = U^\dagger D, \quad V_E = E^\dagger D.$$

The theory: Minimal $\text{SO}(10) \times \text{CP}$

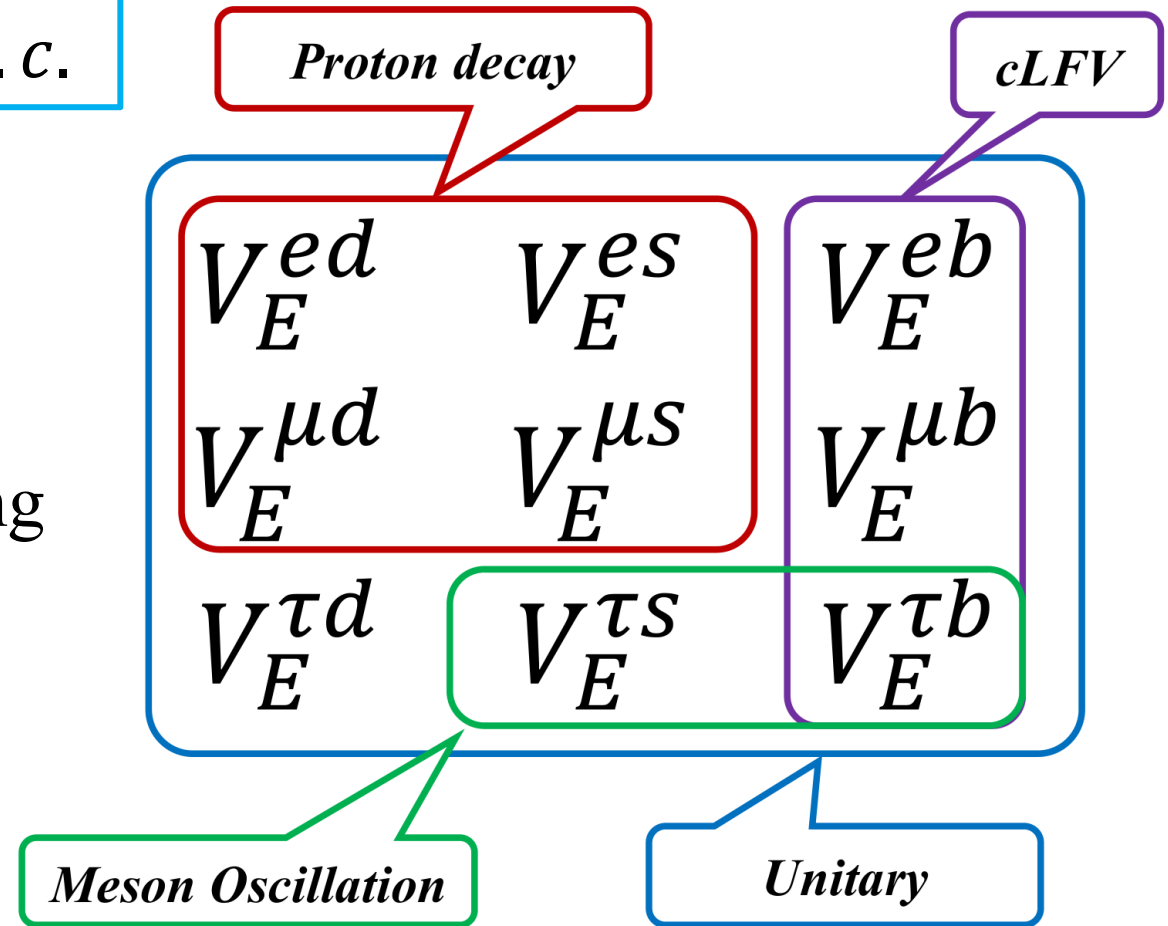
The low energy theory: Flavor violating 2HDM

$$-\mathcal{L}_Y \supset Y_E \overline{L}_L H_2 e_R + Y_D \overline{Q}_L H_2 d_R + h.c.$$

- NMFV-like, absolute strength swamped by fitting.
- Neglecting 1,2 gen quark masses.
- Neglecting small 1-3, 2-3 gen mixing angles.

$$Y_E^{\ell\ell'} \propto V_E^{\ell b} V_E^{\ell' b}, \quad \ell \neq \ell';$$

$$Y_D^{qq'} \propto V_E^{\tau q} V_E^{\tau q'}, \quad q \neq q'.$$



Phenomenology: A Window to Check SO(10)

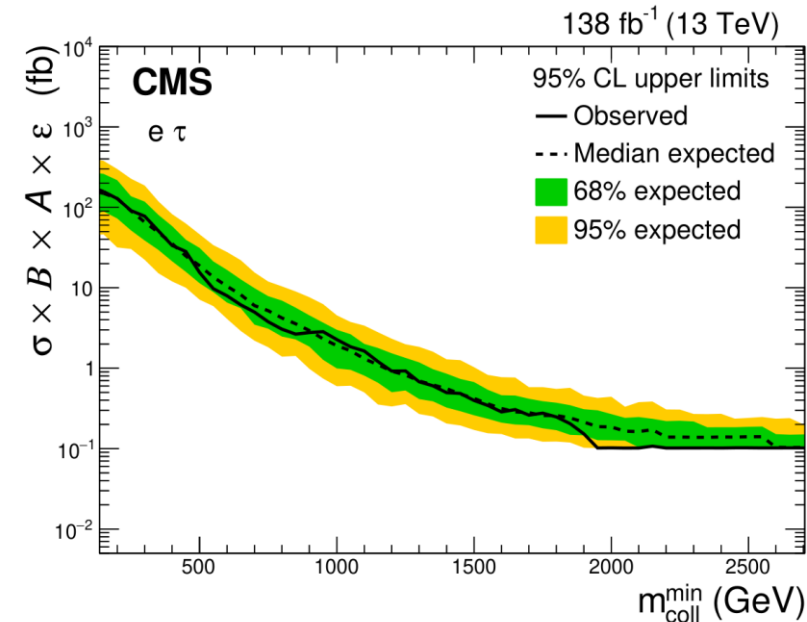
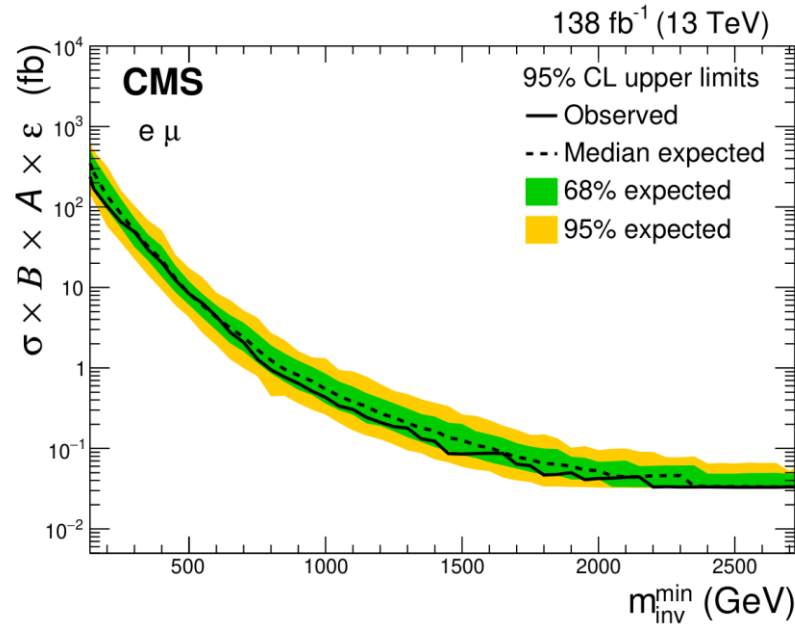
LFV decay of H and A (third column of V_E)

$$\sigma(pp \rightarrow H, A) \times \text{Br}(H, A \rightarrow \ell\ell') \propto |Y_E^{\ell\ell'}|^2 \propto |V_E^{\ell b} V_E^{\ell' b}|^2$$

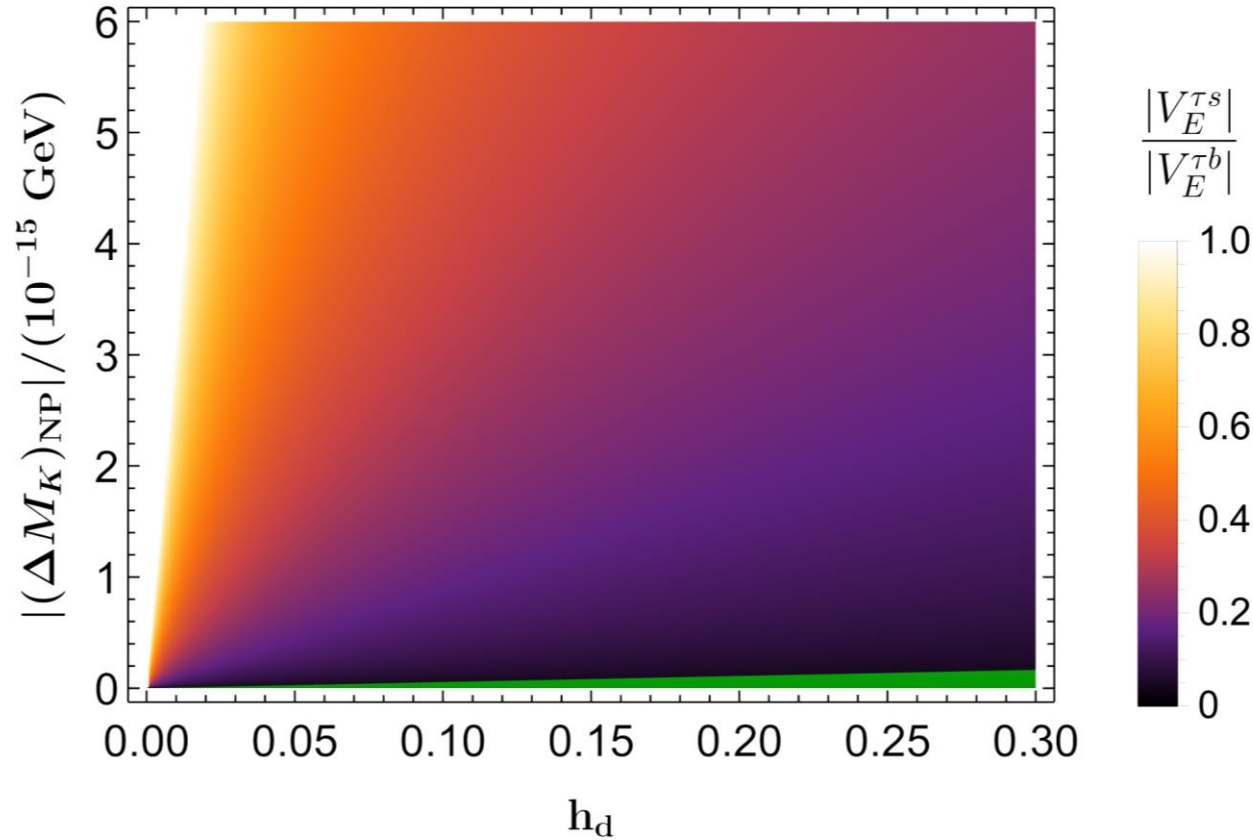
$$\frac{N_{e\mu}}{N_{\tau\mu}} = \frac{|V_E^{eb}|^2}{|V_E^{\tau b}|^2}, \quad \frac{N_{e\mu}}{N_{e\tau}} = \frac{|V_E^{\mu b}|^2}{|V_E^{\tau b}|^2}.$$

$N_{\ell\ell'}$, total # of excess evens,
normalized by detection
efficiency.

Limits:
CMS. 2205.06709



Neutral meson mixing (third row of V_E)



MFV only for the narrow green band.

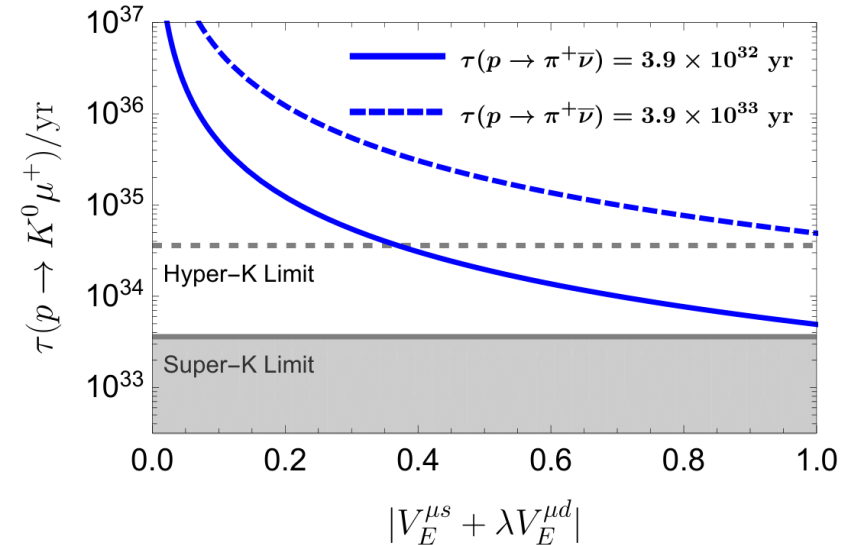
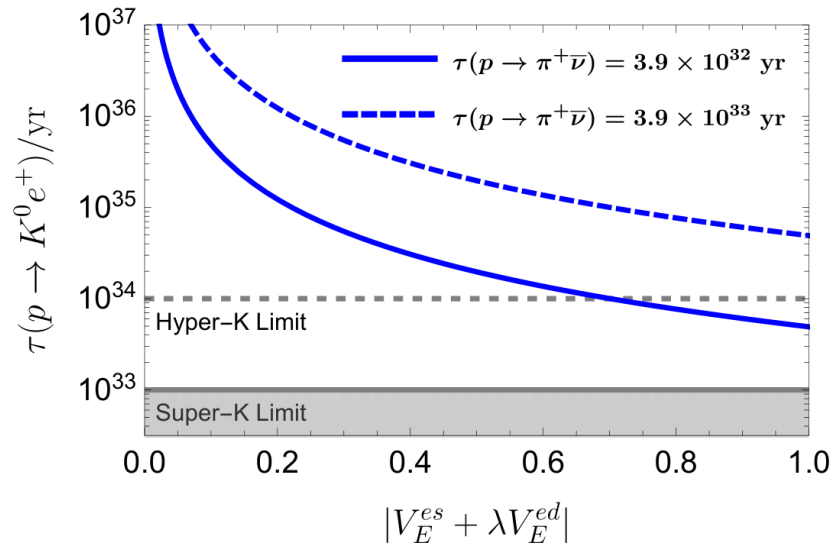
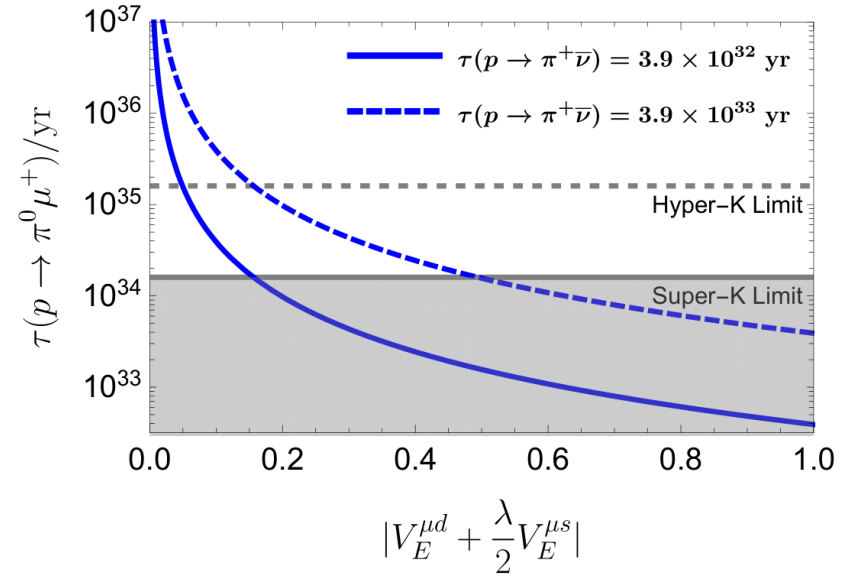
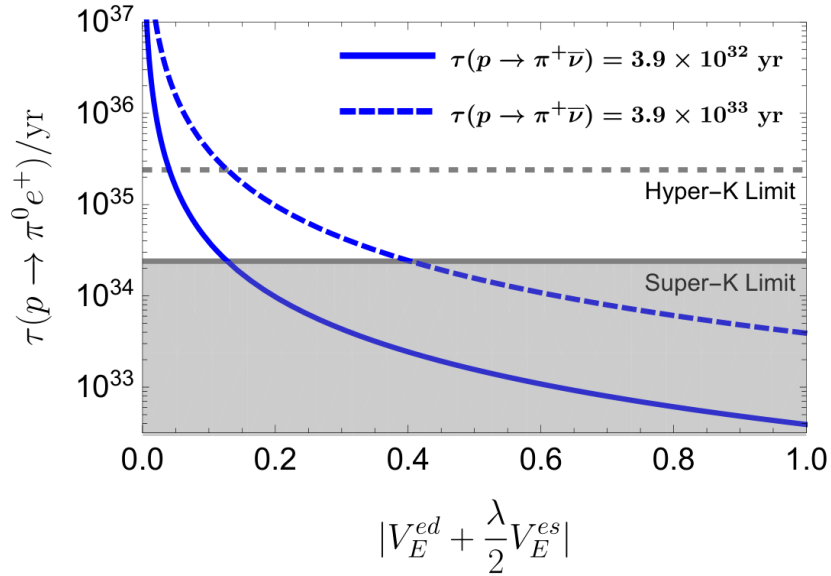
MFV:

$$\begin{aligned} \frac{|(\Delta M_K)_{\text{NP}}| m_K}{\langle K^0 | \bar{d}_L s_R \bar{d}_R s_L | K^0 \rangle} &\approx \lambda^4 \cdot \frac{2h_d |M_{12}^{d\text{SM}}| m_{B_d}}{\langle B_d^0 | \bar{b}_L d_R \bar{b}_R d_L | B_s^0 \rangle} \\ &\approx \lambda^6 \cdot \frac{2h_s |M_{12}^{s\text{SM}}| m_{B_s}}{\langle B_s^0 | \bar{b}_L s_R \bar{b}_R s_L | B_s^0 \rangle}. \end{aligned}$$

NMFV:

$$\begin{aligned} \frac{|(\Delta M_K)_{\text{NP}}|}{2h_d |M_{12}^{d\text{SM}}| \xi_B} &= \frac{|V_E^{\tau s}|^2}{|V_E^{\tau b}|^2}, \\ \xi_B &= \frac{m_B \langle K^0 | \bar{d}_L s_R \bar{d}_R s_L | K^0 \rangle}{m_K \langle B_d^0 | \bar{b}_L d_R \bar{b}_R d_L | B_d^0 \rangle}. \end{aligned}$$

Phenomenology: A Window to Check SO(10)



Combining π and K modes, λ goes away.

$$\begin{aligned} & \frac{2\Gamma(p \rightarrow \pi^0 e^+)}{\Gamma(p \rightarrow \pi^+ \bar{\nu})} - \frac{\Gamma(p \rightarrow K^0 e^+)}{\xi_K \Gamma(p \rightarrow K^+ \bar{\nu})} \\ &= 2|V_E^{ed}|^2 - |V_E^{es}|^2, \\ & \frac{2\Gamma(p \rightarrow \pi^0 \mu^+)}{\Gamma(p \rightarrow \pi^+ \bar{\nu})} - \frac{\Gamma(p \rightarrow K^0 \mu^+)}{\xi_K \Gamma(p \rightarrow K^+ \bar{\nu})} \\ &= 2|V_E^{\mu d}|^2 - |V_E^{\mu s}|^2. \end{aligned}$$

Conclusion and Outlook

Prediction: a concrete flavor **sum rule**:

$$\frac{2\Gamma(p \rightarrow \pi^0 \ell^+)}{\Gamma(p \rightarrow \pi^+ \bar{\nu})} - \frac{\Gamma(p \rightarrow K^0 \ell^+)}{\xi_K \Gamma(p \rightarrow K^+ \bar{\nu})}$$

$$= \left(\frac{3|(\Delta M_K)_{\text{NP}}|}{2\mathbf{h}_d |M_{12}^{d\text{SM}}| \xi_B} - \frac{N_{e\mu}}{N_{\tau\mu}} - \frac{N_{e\mu}}{N_{e\tau}} + 1 \right) \left(\frac{N_{e\mu}}{N_{\tau\mu}} + \frac{N_{e\mu}}{N_{e\tau}} + 1 \right)^{-1}$$

- Waiting for **Hyper-K**, **HL-LHC**, and more **lattice QCD** results.
- Hopefully, a hint for SO(10) in near future.
- What we need most? **Patience!**

Thanks

Part of the scalar potential

$$\phi: 45_H \quad \Sigma: 126_H$$

Bertolini, Di Luzio, Malinsky. '12

$$\langle (1, 1, 1, 0) \rangle \equiv \omega_{BL}, \quad \langle (1, 1, 3, 0) \rangle \equiv \omega_R,$$

$$\langle (1, 1, 3, +1) \rangle \equiv \sigma,$$

$$\begin{aligned} \langle V \rangle = & -\mu^2 (3\omega_{BL}^2 + 2\omega_R^2) + a_0 (12\omega_R^2\omega_{BL}^2 + 9\omega_{BL}^4 + 4\omega_R^4) \\ & + \frac{a_2}{2} (3\omega_{BL}^4 + 2\omega_R^4) - 2\nu^2 |\sigma|^2 + 4\lambda_0 |\sigma|^4 \\ & + 2\tau (3\omega_{BL} + 2\omega_R) |\sigma|^2 + 2\alpha (3\omega_{BL}^2 + 2\omega_R^2) |\sigma|^2 \\ & - 4\beta'_4 (6\omega_R\omega_{BL} + 3\omega_{BL}^2 + \omega_R^2) |\sigma|^2. \end{aligned} \quad (B1)$$

$$V = V_{45} + V_{126} + V_{\text{mix}}, \quad (1)$$

where

$$V_{45} = -\frac{\mu^2}{2} (\phi\phi)_0 + \frac{a_0}{4} (\phi\phi)_0 (\phi\phi)_0 + \frac{a_2}{4} (\phi\phi)_2 (\phi\phi)_2, \quad (2)$$

$$V_{126} = -\frac{\nu^2}{5!} (\Sigma\Sigma^*)_0 \quad (3)$$

$$\begin{aligned} & + \frac{\lambda_0}{(5!)^2} (\Sigma\Sigma^*)_0 (\Sigma\Sigma^*)_0 + \frac{\lambda_2}{(4!)^2} (\Sigma\Sigma^*)_2 (\Sigma\Sigma^*)_2 \\ & + \frac{\lambda_4}{(3!)^2 (2!)^2} (\Sigma\Sigma^*)_4 (\Sigma\Sigma^*)_4 + \frac{\lambda'_4}{(3!)^2} (\Sigma\Sigma^*)_{4'} (\Sigma\Sigma^*)_{4'} \\ & + \frac{\eta_2}{(4!)^2} (\Sigma\Sigma)_2 (\Sigma\Sigma)_2 + \frac{\eta_2^*}{(4!)^2} (\Sigma^*\Sigma^*)_2 (\Sigma^*\Sigma^*)_2, \end{aligned}$$

$$\begin{aligned} V_{\text{mix}} = & \frac{i\tau}{4!} (\phi)_2 (\Sigma\Sigma^*)_2 + \frac{\alpha}{2 \cdot 5!} (\phi\phi)_0 (\Sigma\Sigma^*)_0 \\ & + \frac{\beta_4}{4 \cdot 3!} (\phi\phi)_4 (\Sigma\Sigma^*)_4 + \frac{\beta'_4}{3!} (\phi\phi)_{4'} (\Sigma\Sigma^*)_{4'} \\ & + \frac{\gamma_2}{4!} (\phi\phi)_2 (\Sigma\Sigma)_2 + \frac{\gamma_2^*}{4!} (\phi\phi)_2 (\Sigma^*\Sigma^*)_2. \end{aligned} \quad (4)$$

The theory: Minimal $\text{SO}(10) \times \text{CP}$

The low energy theory: Flavor violating 2HDM

$$\begin{aligned}
-\mathcal{L}_{\Phi\bar{F}F} \supset & \left(\frac{m_E}{v} + \epsilon Y_E^{\ell\ell'}\right) h \bar{\ell}_L \ell'_R + \left(\frac{m_D}{v} + \epsilon Y_D^{qq'}\right) \bar{d}_L^q d_R^{q'} h + \left(\frac{m_U}{v} + \epsilon Y_U^{qq'}\right) \bar{u}_L^q u_R^{q'} h \\
& + \mathcal{Y}_E^{\ell\ell'} (H + iA) \bar{\ell}_L \ell'_R + \mathcal{Y}_D^{qq'} (H + iA) \bar{d}_L^q d_R^{q'} + \mathcal{Y}_U^{qq'} (H + iA) \bar{u}_L^q u_R^{q'} + \text{h.c.}
\end{aligned}$$

$$\begin{aligned}
Y_E &= \mathcal{C}_{EE} \frac{m_E}{v} + \mathcal{C}_{ED} V_E^* \frac{m_D}{v} V_E^\dagger + \mathcal{C}_{EU} V_E^* V_{\text{CKM}}^T \frac{m_U}{v} V_{\text{CKM}} V_E^\dagger, \\
Y_D &= \mathcal{C}_{DE} V_E^T \frac{m_E}{v} V_E + \mathcal{C}_{DD} \frac{m_D}{v} + \mathcal{C}_{DU} V_{\text{CKM}}^T \frac{m_U}{v} V_{\text{CKM}}, \\
Y_U &= \mathcal{C}_{UE} V_{\text{CKM}}^* V_E^T \frac{m_E}{v} V_E V_{\text{CKM}}^\dagger + \mathcal{C}_{UD} V_{\text{CKM}}^T \frac{m_D}{v} V_{\text{CKM}} + \mathcal{C}_{UU} \frac{m_U}{v},
\end{aligned}$$

where $\mathcal{Y}_F = e^{i\alpha_c} Y_F$, $\ell, \ell' = e, \mu, \tau$, $q, q' = d, s, b$ or u, c, t

- Flavor structure: **only V_E unknown.**
- $\mathcal{C}_{FF'}$: free, not predicted.

Phenomenology: Is the Theory Safe?

In which sense **FCNC** is **not ruled out**?

$$\begin{pmatrix} \mathcal{C}_{EE} & \mathcal{C}_{ED} & \mathcal{C}_{EU} \\ \mathcal{C}_{DE} & \mathcal{C}_{DD} & \mathcal{C}_{DU} \\ \mathcal{C}_{UE} & \mathcal{C}_{UD} & \mathcal{C}_{UU} \end{pmatrix} = \frac{v}{u} \begin{pmatrix} u_1 & -u_2 & -3u_3 \\ u_1 & -u_2 & u_3 \\ u_2^* & u_1^* & u_4^* \end{pmatrix} \begin{pmatrix} v_{10}^d & -v_{10}^{u*} & -3v_{126}^d \\ v_{10}^d & -v_{10}^{u*} & v_{126}^d \\ v_{10}^u & v_{10}^{d*} & v_{126}^u \end{pmatrix}^{-1}$$

- Naively, all $\mathcal{C}_{FF'}$ at $o(1)$. Strictly, not predicted.
- $\mathcal{C}_{EU} = \mathcal{C}_{DU}$. No other correlations.
- B_s mixing constrains \mathcal{C}_{DU} .
- **MFV**: B_d, K mixing less constrained.
- $\Delta F = 1$ processes: suppressed by **loop factor** or additional **chirality flipping**.

$$\Phi_i = \begin{pmatrix} \phi_i^+ \\ (v_i + \rho_i + i\eta_i)/\sqrt{2} \end{pmatrix},$$

$$u_i = v_i^* - \left(\sum_{j=1}^4 v_j^{*2} \right) \frac{v_i}{v}, \quad u = \sum_{i=1}^4 |u_i|^2.$$

$$|\mathcal{C}_{DU}| \lesssim \left(\frac{v}{m_t} \cdot \frac{1}{|V_{\text{CKM}}^{ts} V_{\text{CKM}}^{tb}|} \right) \times \frac{m_H/\sqrt{2}}{10^3 \text{ TeV}} \\ \approx 0.013 \times \frac{m_H}{500 \text{ GeV}}$$

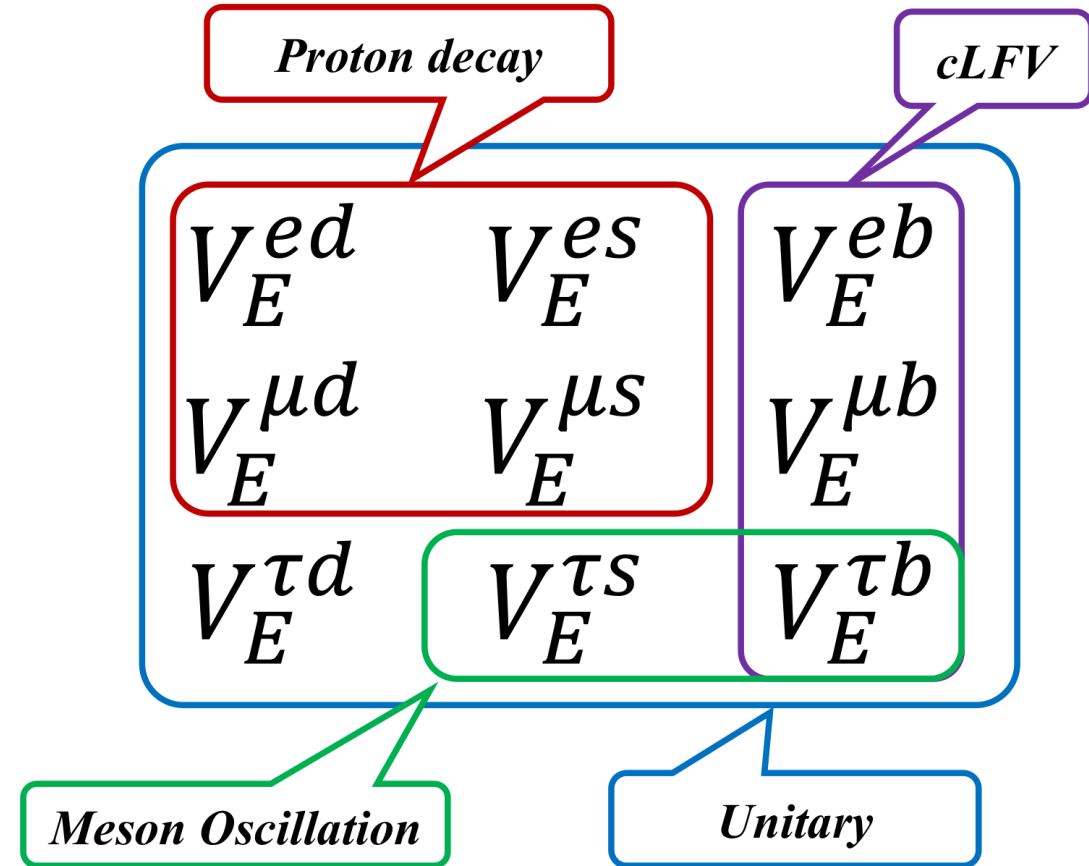
Understand the flavor structure:

- Charged Leptons: **NMFV**
- Down-type Quarks: **MFV** or **NMF**
- Up-type quarks: **$o(\lambda)$ corrections**.
 $\lambda \sim 0.2$, the Cabibbo angle

$$Y_E^{\ell\ell'} \propto (V_E^{\ell b} V_E^{\ell' b})^* + o(\lambda^2)(V_E^{\ell b} + V_E^{\ell' b})^*, \quad \ell \neq \ell',$$

$$Y_D^{qq'} \propto V_E^{\tau q} V_E^{\tau q'} + \frac{\mathcal{C}_{DU} m_t}{\mathcal{C}_{DE} m_\tau} V_{\text{CKM}}^{tq} V_{\text{CKM}}^{tq'}, \quad q \neq q',$$

$$Y_U^{qq'} \propto V_E^{\tau q} V_E^{\tau q'} + \frac{\mathcal{C}_{UD} m_b}{\mathcal{C}_{UE} m_\tau} V_{\text{CKM}}^{tq} V_{\text{CKM}}^{tq'} + o(\lambda), \quad q \neq q',$$



Additional Slides

Neutral meson mixing (third row of V_E)

$$\begin{aligned}
H_{\text{NP}}^q &= -\frac{1}{2m_H^2} \left(Y_D^{bq} \bar{b}_L q_R + Y_D^{qb*} \bar{b}_R q_L \right)^2 - \frac{1}{2m_A^2} \left(iY_D^{bq} \bar{b}_L q_R - iY_D^{qb*} \bar{b}_R q_L \right)^2 \\
&\approx -\frac{2|Y_D^{bq}|^2}{m_H^2} \bar{b}_L q_R \bar{b}_R q_L, \quad q = d, s.
\end{aligned}$$

$\bar{b}_L q_R \bar{b}_L q_R$ and $\bar{b}_R q_L \bar{b}_R q_L$ vanish in the limit $m_H = m_A$

Observable	Current limit	Future sensitivity
$\mathbf{h_d}$	0.26	0.049 (0.038)
$\mathbf{h_s}$	0.12	0.044 (0.031)
$ (\Delta M_K)_{\text{NP}} $	$5.2 \times 10^{-15} \text{ GeV}$	$0.2 \times 10^{-15} \text{ GeV}$

$$\frac{\langle B_q^0 | H_{\text{NP}}^q | \bar{B}_q^0 \rangle}{\langle B_q^0 | H_{\text{SM}}^q | \bar{B}_q^0 \rangle} = \mathbf{h_q} e^{i\sigma_q}$$

$$(\Delta M_K)_{\text{NP}} = \frac{1}{m_K} \langle K^0 | H_{\text{NP}}^K | \bar{K}^0 \rangle$$

CKMfit, 2006.04824; KTEV, 1011.0127;

Bai, et al. '14; Wang '23.

Phenomenology: A Window to Check SO(10)

Proton decay (2×2 top-left submatrix of V_E)

Symmetric Yukawa couplings simplifies a lot.

*Perez '04;
Nath, Perez '07.*

$$\frac{\Gamma(p \rightarrow \pi^+ \bar{\nu})}{\Gamma(p \rightarrow K^+ \bar{\nu})} = \frac{4(1 - m_K^2/m_p^2)^{-2} \langle \pi^+ | (du)_{Rd_L} | p \rangle^2}{\langle K^+ | (us)_{Rd_L} | p \rangle^2 + \lambda^2 \langle K^+ | (ud)_{Rs_L} | p \rangle^2}$$

$$\frac{\Gamma(p \rightarrow \pi^0 e^+)}{\Gamma(p \rightarrow \pi^+ \bar{\nu})} = |V_E^{ed} + \frac{\lambda}{2} V_E^{es}|^2, \quad \frac{\Gamma(p \rightarrow \pi^0 \mu^+)}{\Gamma(p \rightarrow \pi^+ \bar{\nu})} = |V_E^{\mu d} + \frac{\lambda}{2} V_E^{\mu s}|^2,$$

$$\frac{\Gamma(p \rightarrow K^0 e^+)}{\xi_K \Gamma(p \rightarrow K^+ \bar{\nu})} = |V_E^{es} + \lambda V_E^{ed}|^2, \quad \frac{\Gamma(p \rightarrow K^0 \mu^+)}{\xi_K \Gamma(p \rightarrow K^+ \bar{\nu})} = |V_E^{\mu s} + \lambda V_E^{\mu d}|^2,$$

$$\text{with } \xi_K = \frac{2 \langle K^0 | (us)_{Ru_L} | p \rangle^2}{\langle K^+ | (us)_{Rd_L} | p \rangle^2 + \lambda^2 \langle K^+ | (ud)_{Rs_L} | p \rangle^2} \approx 6.4.$$

Decay Mode	$\ell = e^+$	$\ell = \mu^+$	$\ell = \bar{\nu}$
$p \rightarrow \pi \ell$	$> 2.4 \times 10^{34} \text{ yr}$	$> 1.6 \times 10^{34} \text{ yr}$	$> 3.9 \times 10^{32} \text{ yr}$
$p \rightarrow K \ell$	$> 1.0 \times 10^{33} \text{ yr}$	$> 3.6 \times 10^{33} \text{ yr}$	$> 5.9 \times 10^{33} \text{ yr}$

RGE towards M_I changes:

The IR theory (M_{EW}): SM \rightarrow 2HDM

Just estimate an upper limit here.
Long way to the complete GUT
Yukawa texture.

$$(Y_E^{\tau\ell} - Y_E^{\ell\tau})|_{M_{EW}} = (Y_E^{\tau\ell} - Y_E^{\ell\tau})|_{M_I}, (Y_U^{tq'} - Y_U^{q't})|_{M_{EW}} = (Y_U^{tq'} - Y_U^{q't})|_{M_I},$$

$$(Y_D^{bq} - Y_D^{qb})|_{M_{EW}} = (Y_D^{bq} - Y_D^{qb})|_{M_I} + Y_D^{bq} \times \frac{y_t^2}{16\pi^2} \log(M_I/M_{EW}) \left(\frac{3}{2} |C_{UU}|^2 - \frac{1}{2} \right),$$

$$(M_D^{\tau\ell} - M_D^{\ell\tau})|_{M_{EW}} = (M_D^{\tau\ell} - M_D^{\ell\tau})|_{M_I}, (M_D^{tq'} - M_D^{q't})|_{M_{EW}} = (M_D^{tq'} - M_D^{q't})|_{M_I},$$

$$(M_D^{bq} - M_D^{qb})|_{M_{EW}} = (M_D^{bq} - M_D^{qb})|_{M_I} + (m_b V_E^{\tau b} V_E^{\tau q}) \times \frac{y_t^2}{16\pi^2} \log(M_I/M_{EW}) \left(C_{UU} C_{DE} \frac{2m_\tau}{m_b} \right).$$

$$\frac{y_t^2}{16\pi^2} \log(M_I/M_{EW}) \approx 3\%.$$

Additional Slides

More on V_E :

$$V_E^T (V_{\text{PMNS}}^* m_{\nu_L}^{\text{diag}} V_{\text{PMNS}}^\dagger - k_1 m_E^{\text{diag}}) V_E = -k_1 m_D^{\text{diag}} + k_2 (V_{\text{CKM}}^T m_U^{\text{diag}} V_{\text{CKM}}) + k_3 M_T,$$

$$k_1 = \frac{8(v_4^*)^2}{v_3 \langle \Delta_R^0 \rangle} - \frac{\langle \Delta_L \rangle}{8v_3}, \quad k_2 = \frac{16v_4^*}{\langle \Delta_R^0 \rangle}, \quad k_3 = \frac{8v_3}{\langle \Delta_R^0 \rangle}.$$

$$M_T = V_{\text{CKM}}^T m_U^{\text{diag}} V_{\text{CKM}} (V_E^T m_E^{\text{diag}} V_E - m_D^{\text{diag}})^{-1} V_{\text{CKM}}^T m_U^{\text{diag}} V_{\text{CKM}}.$$

M_T is ‘diagonal’ when $m_t \rightarrow \infty$ and $V_{\text{CKM}} \rightarrow \mathbf{1}$

- $k_1 m_\tau \gg m_{\nu_L} \sim 0.1 \text{ eV}$, $V_E \sim \mathbf{1}$.
- $k_1 m_\tau \lesssim m_{\nu_L} \sim 0.1 \text{ eV}$, $V_E \sim V_{\text{PMNS}}$.