# Quark-mass effects in Higgs-boson production in gluon fusion

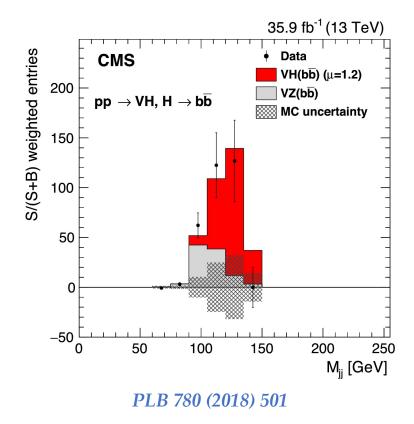
Project A1a

M. Czakon, R. Harlander

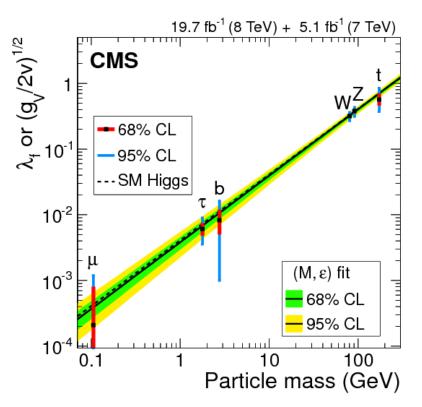
#### Plan:

- Status
- Some open problems in Higgs-boson production description:
  - Total inclusive cross section
  - Differential cross section
  - Kinematical limits
  - Resummation and bottom Yukawa coupling
- Semi-numerical methods
- Conclusions

#### Status experiment



 $5.1 \text{ fb}^{-1} (7 \text{ TeV}) + 19.7 \text{ fb}^{-1} (8 \text{ TeV}) + 35.9 \text{ fb}^{-1} (13 \text{ TeV})$  Observed **CMS**  $--\pm 1\sigma$  (stat  $\oplus$  syst) <u>±1</u>σ (syst) tīH(WW\*) — ±2σ (stat ⊕ syst)  $t\bar{t}H(ZZ^*)$  $t\bar{t}H(\gamma\gamma)$  $t\bar{t}H(\tau^+\tau^-)$  $t\bar{t}H(b\bar{b})$ 7+8 TeV 13 TeV Combined  $\mu_{t\bar{t}H}^3$ 



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# Status theory

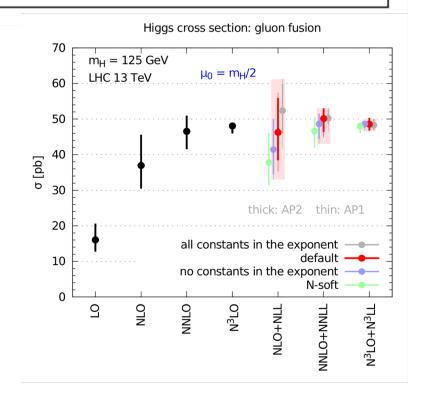
Handbook of LHC Higgs cross sections:

4. Deciphering the nature of the Higgs sector Report of the LHC Higgs Cross Section Working Group `16

LHC @13 TeV

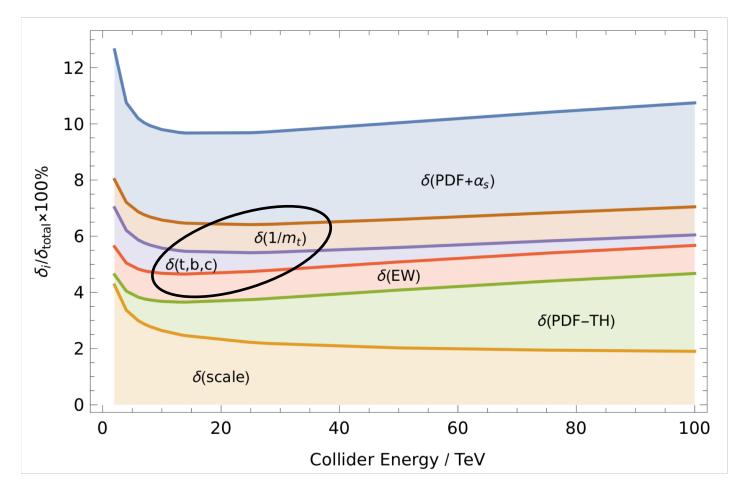
$$\sigma = 48.58 \, \mathrm{pb}_{-3.27 \, \mathrm{pb} \, (-6.72\%)}^{+2.22 \, \mathrm{pb} \, (+4.56\%)}$$
 (theory)  $\pm 1.56 \, \mathrm{pb} \, (3.20\%)$  (PDF+ $\alpha_s$ ).

$$48.58 \, \mathrm{pb} = 16.00 \, \mathrm{pb} \qquad (+32.9\%) \qquad (\mathrm{LO}, \mathrm{rEFT}) \\ + 20.84 \, \mathrm{pb} \qquad (+42.9\%) \qquad (\mathrm{NLO}, \mathrm{rEFT}) \\ - 2.05 \, \mathrm{pb} \qquad (-4.2\%) \qquad ((t, b, c), \, \mathrm{exact} \, \mathrm{NLO}) \\ + 9.56 \, \mathrm{pb} \qquad (+19.7\%) \qquad (\mathrm{NNLO}, \mathrm{rEFT}) \\ + 0.34 \, \mathrm{pb} \qquad (+0.7\%) \qquad (\mathrm{NNLO}, 1/m_t) \\ + 2.40 \, \mathrm{pb} \qquad (+4.9\%) \qquad (\mathrm{EW}, \, \mathrm{QCD-EW}) \\ + 1.49 \, \mathrm{pb} \qquad (+3.1\%) \qquad (\mathrm{N}^3 \mathrm{LO}, \, \mathrm{rEFT})$$



## Status theory

Higgs Physics at the HL-LHC and HE-LHC Report from Working Group 2 on the Physics of the HL-LHC, and Perspectives at the HE-LHC `19

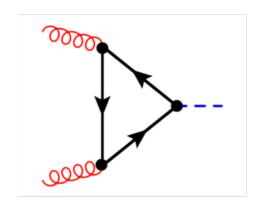


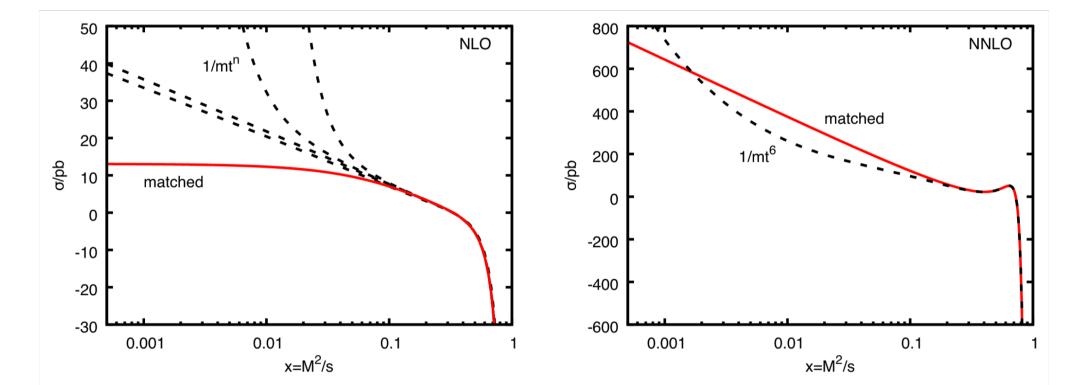
- Missing higher-order effects of QCD corrections beyond N3LO ( $\delta$ (scale)).
- Missing higher-order effects of electroweak and mixed QCD-electroweak corrections at and beyond O( $\alpha_s$   $\alpha$ ) ( $\delta$ (EW)).
- Effects due to finite quark masses neglected in QCD corrections beyond NLO ( $\delta(t,b,c)$  and  $\delta(1/mt)$ )
- Mismatch in the perturbative order of the parton distribution functions (PDF) evaluated at NNLO and the perturbative QCD cross sections evaluated at N3LO ( $\delta$ (PDF-TH)).

# Top quark mass & inclusive cross section gg -> H + X

#### Exploiting the large mass expansion

RH, Ozeren, `09 RH, Neumann, Ozeren, M. Wiesemann, `12 RH, Mantler, S. Marzani, K. Ozeren, `10 Marzani, Ball, Del Duca, Forte, Vicini `08 Pak, Rogal, Steinhauser `10

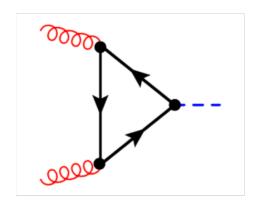




# Top quark mass & inclusive cross section gg -> H + X

#### ...and the high-energy limit

RH, Ozeren, `09 RH, Neumann, Ozeren, M. Wiesemann, `12 RH, Mantler, S. Marzani, K. Ozeren, `10 Marzani, Ball, Del Duca, Forte, Vicini `08 Pak, Rogal, Steinhauser `10



$$\Delta_{gg}(N, \tau, \mu_{\rm F}) = h(0, \tau, \gamma_s, \gamma_s) R^2(\gamma_s) \left(\frac{M_H^2}{\mu_{\rm F}^2}\right)^{2\gamma_s}$$

$$h(N,\tau,M_1,M_2) = M_1 M_2 \int_0^1 d\zeta \zeta^{N-1} \int_0^\infty d\xi_1 \xi_1^{M_1-1} \int_0^\infty d\xi_2 \xi_2^{M_2-1} \int_0^{2\pi} \frac{d\varphi}{2\pi}$$
$$M_H^2 \sigma^{\text{off}}(\zeta,\tau,\xi_1,\xi_2,\varphi),$$

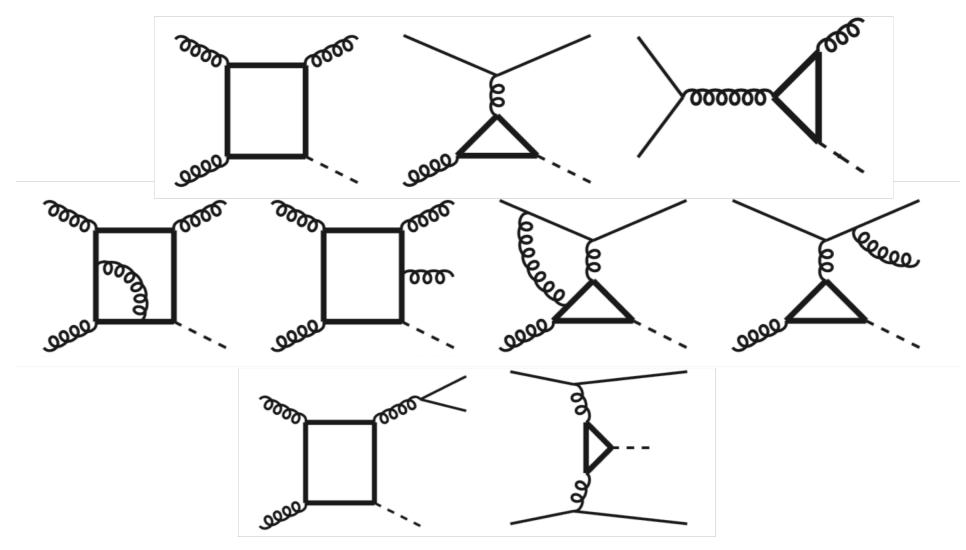
$$R = 1 + \mathcal{O}\left(\left(\frac{\alpha_s}{N}\right)^3\right)$$

$$\xi_i = \frac{|\mathbf{k}_i|^2}{M_H^2}, \quad \zeta = \frac{M_H^2}{2(k_1 \cdot k_2 - \mathbf{k}_1 \cdot \mathbf{k}_2)}$$

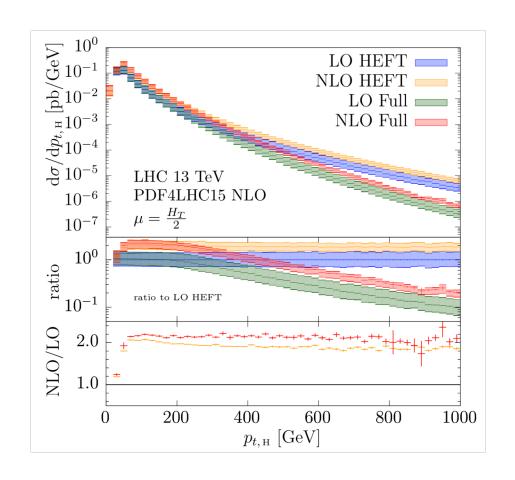
$$M_1 = M_2 = \gamma_s \left(\frac{\alpha_s}{N}\right) \qquad \qquad \gamma_s \left(\frac{\alpha_s}{N}\right) = \sum_{k=1}^{\infty} c_k \left(\frac{C_A \alpha_s}{\pi N}\right)^k, \quad c_k = 1, 0, 0, 2\zeta(3), \dots$$

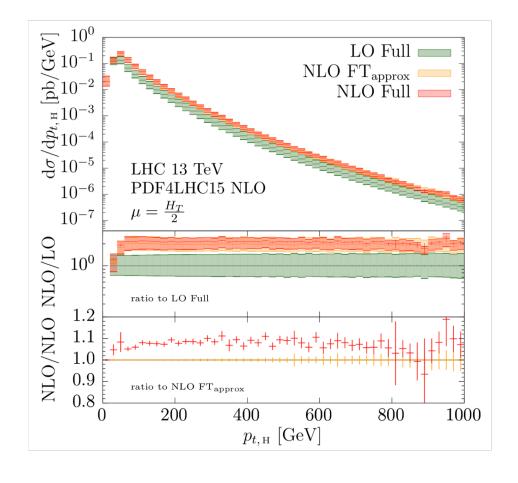
- Extend beyond leading high-energy behavior?
- How about threshold behavior?

Complicated two-loop amplitudes

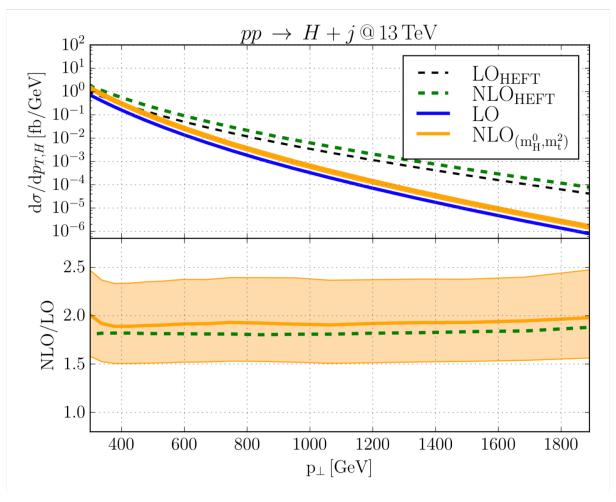


• Exact numerical result using sector decomposition for two-loop amplitudes Jones, Kerner, Luisioni `18



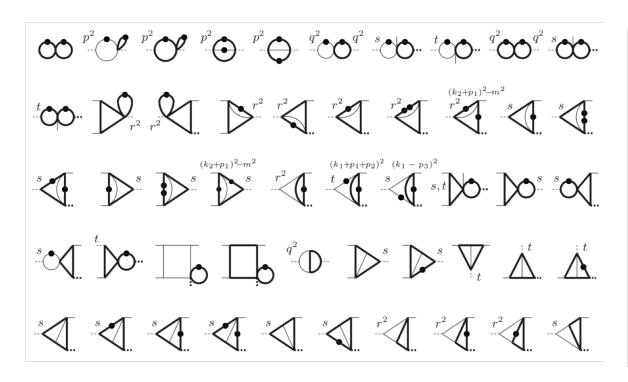


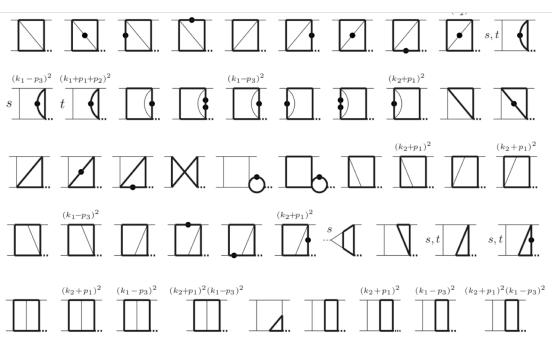
 Results using full two-loop amplitudes for nearly massless top quarks and exact real radiation Lindert, Kudashkin, Melnikov, Wever `18



Exact two-loop amplitudes in the planar approximation

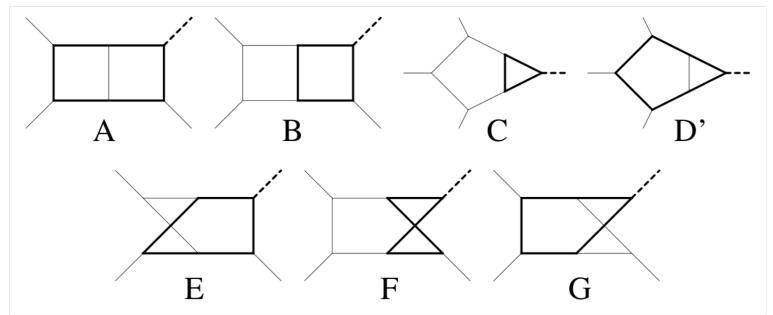
Bonciani, Del Duca, Frellesvig, Henn, Moriello, Smirnov `16





Exact two-loop amplitudes and some non-planars

Bonciani, Del Duca, Frellesvig, Henn, Moriello, Smirnov `18



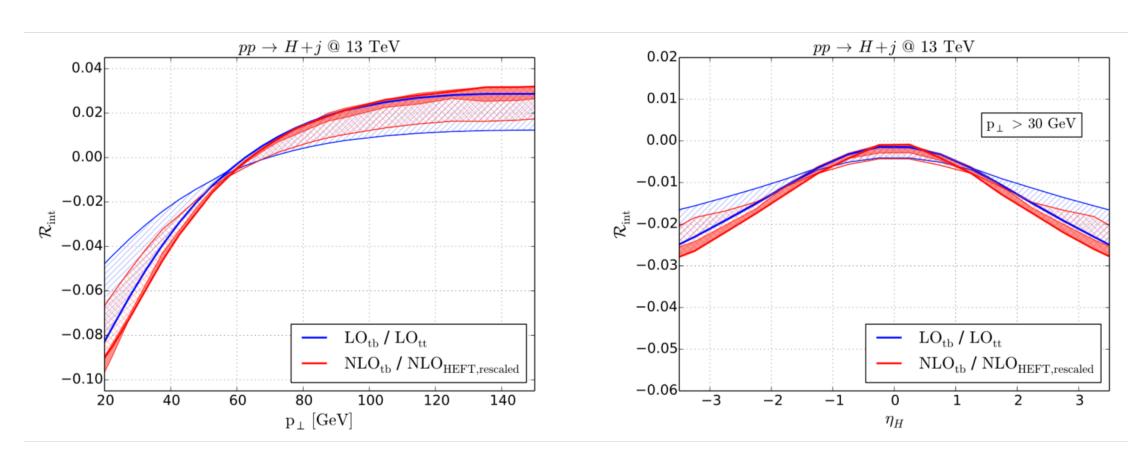
Elliptic polylogs

$$A_{66}|_{\text{max cut}} \propto \int \frac{dz}{s\sqrt{P_{4;1}(z)}}$$

$$A_{66}|_{\text{max cut}} \propto \int \frac{dz}{s_3 \sqrt{P_{4:1}(z)}}$$
  $P_{4;1} \equiv \left( (m_H^2 + z)^2 - 4m_H^2 m_t^2 \right) \left( 4m_t^2 t u/s + (t+z)^2 \right)$ 

# Top-bottom-loop interference for Higgs + jet

Lindert, Melnikov, Tancredi, Wever `17



Integrate to obtain the effect on the inclusive cross section?

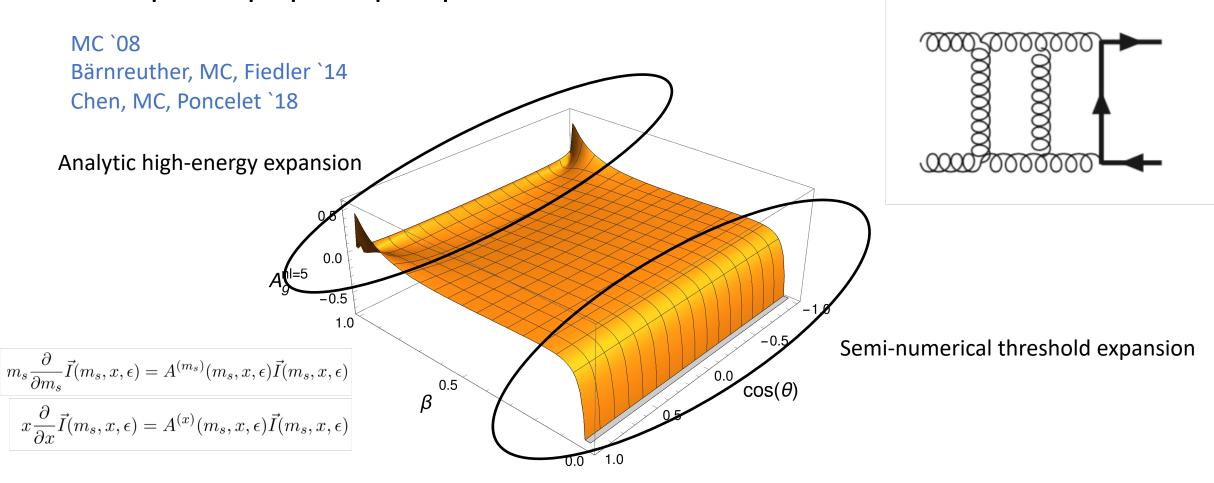
## Cross sections with mass effects: summary

 3-loop 3-points amplitudes and 2-loop 4-point amplitudes allow to evaluate the missing mass effects at NNLO by means of a direct Monte Carlo integration using an NNLO subtraction scheme, e.g. STRIPPER MC `10, MC and Heymes `14

Do we have to wait for the complete analytic results?

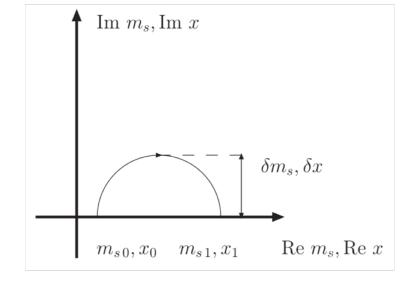
Numerical solution of differential equations

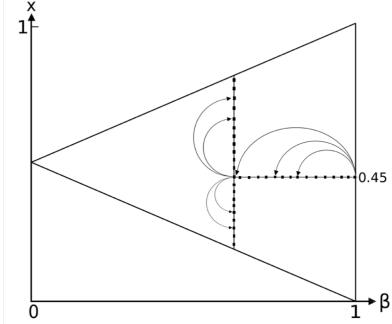
• Example: top-quark pair production in hadron collisions



- Numerical solution of differential equations
- Example: top-quark pair production in hadron collisions

MC `08 Bärnreuther, MC, Fiedler `14 Chen, MC, Poncelet `18



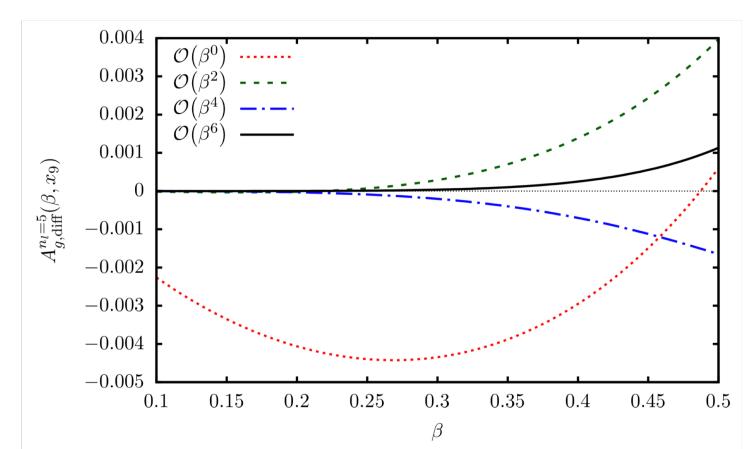


$$m_s \frac{\partial}{\partial m_s} \vec{I}(m_s, x, \epsilon) = A^{(m_s)}(m_s, x, \epsilon) \vec{I}(m_s, x, \epsilon)$$
$$x \frac{\partial}{\partial x} \vec{I}(m_s, x, \epsilon) = A^{(x)}(m_s, x, \epsilon) \vec{I}(m_s, x, \epsilon)$$

- Numerical solution of differential equations
- Example: top-quark pair production in hadron collisions

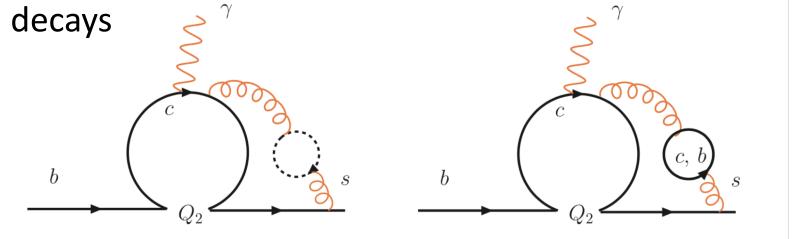
MC `08 Bärnreuther, MC, Fiedler `14 Chen, MC, Poncelet `18

$$m_s \frac{\partial}{\partial m_s} \vec{I}(m_s, x, \epsilon) = A^{(m_s)}(m_s, x, \epsilon) \vec{I}(m_s, x, \epsilon)$$
$$x \frac{\partial}{\partial x} \vec{I}(m_s, x, \epsilon) = A^{(x)}(m_s, x, \epsilon) \vec{I}(m_s, x, \epsilon)$$

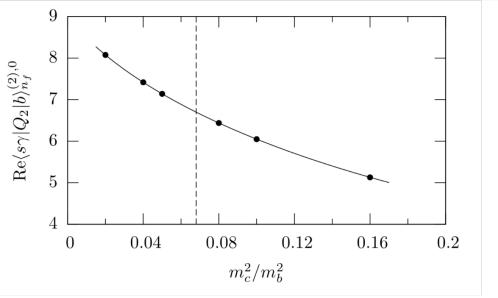


Numerical solution of differential equations

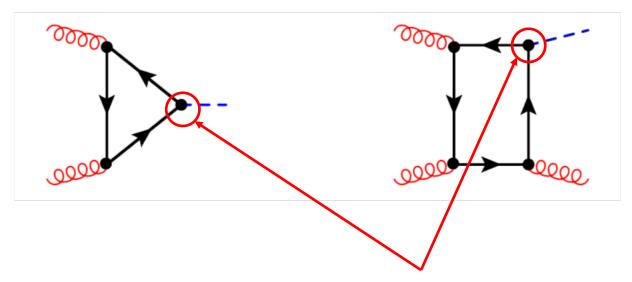
• Example: Rare B-meson decays



Boughezal, MC, Schutzmeier `07 MC, Fiedler, Huber, Misiak, Schutzmeier, Steinhauser `15 Misiak et al `15



### One last question...



How to renormalize the bottom Yukawa coupling?

Hints from subleading Sudakov resummation? Melnikov, Penin `16

#### Conclusions

- In principle, one can obtain both total and differential cross sections using purely numerical methods
- However, there is potential for physical insight by analytic methods in high-energy, threshold and other regimes
- Progress expected through synergy between the two approaches

#### Conclusions

#### Links to other projects:

 A1b - Higgs boson physics with higher order QCD corrections within the Higgs Effective Theory

 A3b - Precision predictions for Higgs boson properties as a probe for New Physics

- B1a Production of colour-singlet final states through N3LO QCD
- B1b Precision top-quark physics at the LHC

#### Methodologically related to:

Subtraction methods
Massive virtual amplitudes
Evaluation of master integrals
A1a, B1a
A3b, B1b
C1a, C1b

#### Relevant input to Higgs-boson physics analyses:

• Bottom-mass effects in BSM theories A3a, A3b

• Mass effects in Higgs-boson production

A1b

