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Theoretical Advances in Parton Showers

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at the

TRR 257 Kick-off Meeting

Karlsruhe | 18 March 2019

Indispensable input
for experiments & phenomenology.

Realistic, fully detailed description
spanning orders of magnitude in
relevant energy scales.

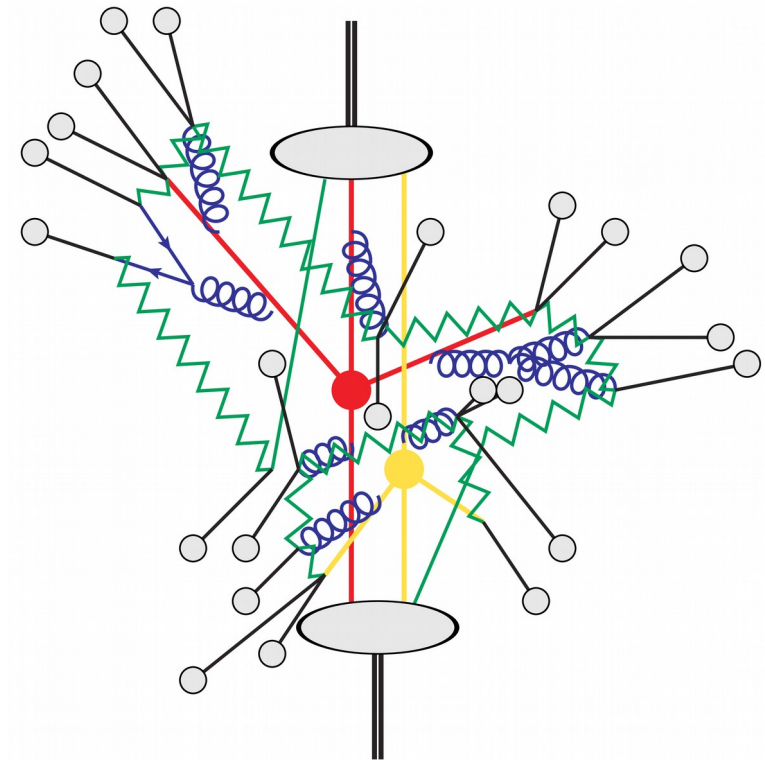
Factorization dictates work flow.

Hard partonic scattering

Jet evolution

Multiple interactions

Hadronization

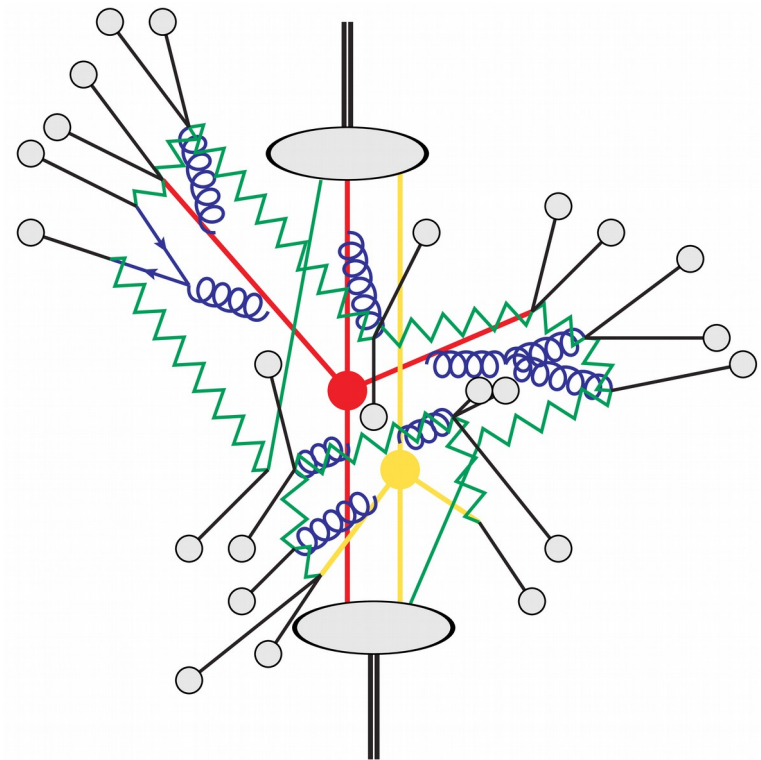


$$d\sigma \sim d\sigma_{\text{hard}}(Q) \times \text{PS}(Q \rightarrow \mu) \times \text{Had}(\mu \rightarrow \Lambda) \times \dots$$

Precision collider phenomenology demands **reliable simulations** with the highest level of theoretical control.

Open questions in **existing** shower algorithms and **future development**:

- Formal accuracy?
- Relation to analytic resummation?
- Relation to effective field theories?
- New formulation, methods and algorithms?



$$d\sigma \sim d\sigma_{\text{hard}}(Q) \times \text{PS}(Q \rightarrow \mu) \times \text{Had}(\mu \rightarrow \Lambda) \times \dots$$

Parton shower algorithms

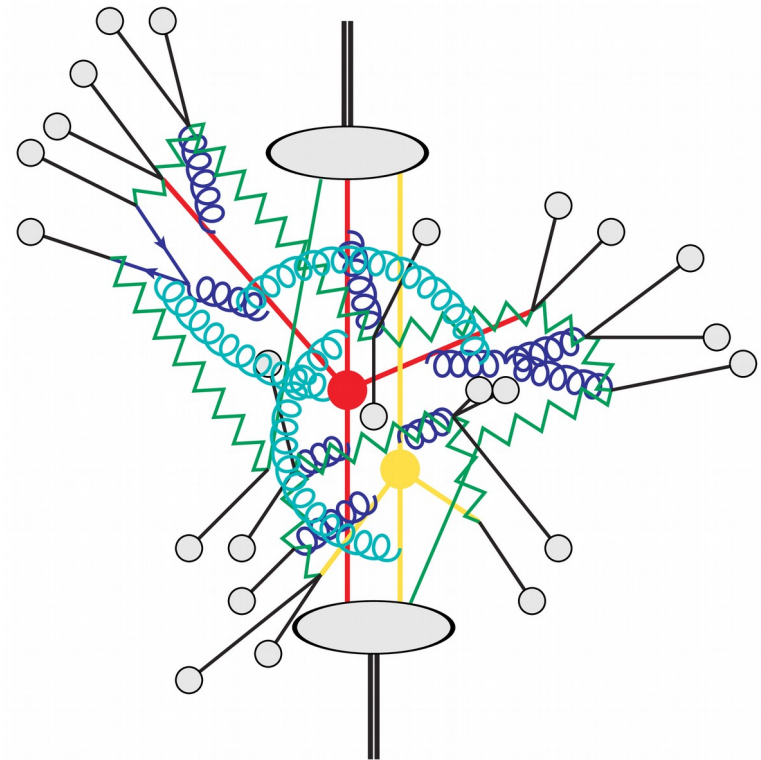
Lack a systematic expansion,
obstruct NNLO for the hard process.

Hadronization models

Lack constraints from perturbative
evolution: Hiding perturbative effects?

Rethink foundations of parton showers:

Systematic picture including virtual
corrections and quantum mechanical
interference.



$$d\sigma \sim d\sigma_{\text{hard}}(Q) \times \text{PS}(Q \rightarrow \mu) \times \text{Had}(\mu \rightarrow \Lambda) \times \dots$$

Forshaw + De Angelis, Holguin, ...



Plätzer + Ruffa, ...



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CVolver

RWTHAACHEN
UNIVERSITY

Czakov + NN, ...



H7

A highly **ambitious program:**

Need to combine **diverse expertise** from different networks to gain momentum.



Gieseke + Löschner, Simpson-Dore, ...

Seek a framework to systematically address and improve **a new kind of parton shower algorithm**, not relying on ad-hoc constructions, treating colour and spin exactly as far as possible.

Gain most detailed analytic control over the algorithm to make **decisive statements about logarithmic accuracy** for a certain class of observables.

Extend the algorithm **beyond the customary leading colour limit**, include spin correlations and prepare to include higher orders in the shower algorithm itself.

Determine the relevant ingredients to have **resummation beyond (N)LL** in reach.

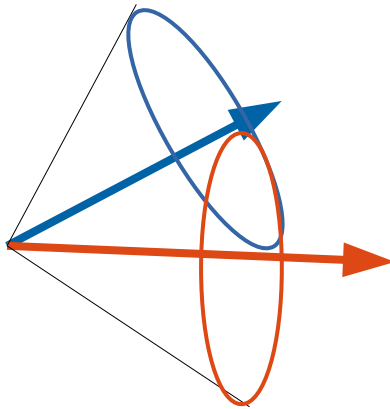
Facilitate the **matching to NNLO** fixed-order perturbative calculations.

Contribute to the development of the **CVolver** package which targets new evolution algorithms, and include into the multi-purpose event generator **Herwig 7**.

Coherent branching algorithms essential to direct QCD resummation of global event shapes, and to designing parton shower algorithms.

[Catani, Marchesini, Webber]

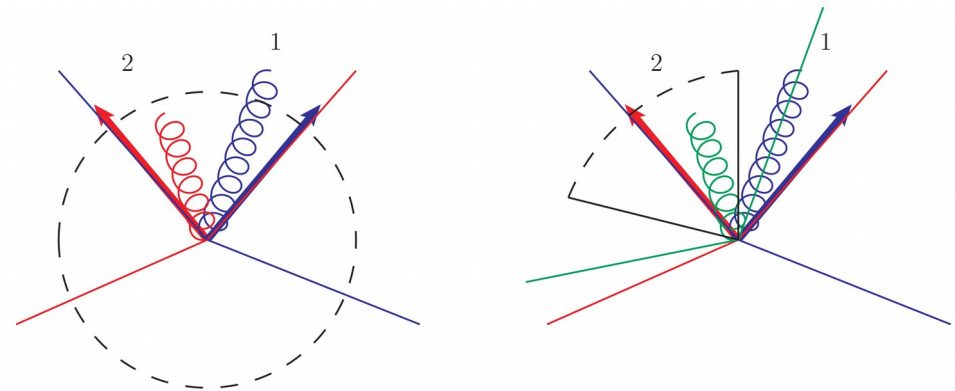
[Gieseke, Stephens, Webber]



$$\begin{aligned}
 & \text{Diagram: } Q^2 \text{ with a wavy line and } q_L \text{ label} = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} \\
 & \xrightarrow{q_L^2 \ll Q^2} \text{Diagram 4} + \mathcal{O}(q_L^2/Q^2) \\
 & \quad \quad \quad \uparrow \quad \quad \quad \uparrow \\
 & \quad \quad \text{gauge invariant} \quad \text{decomposition}
 \end{aligned}$$

Large-angle soft effects included on average by a clever choice of ordering variable.

Non-global logarithms require dipole-type soft gluon evolution to take into account change in colour structure after each emission.



Systematic inclusion of collinear effects needs to be addressed.

[Dasgupta, Salam; Banfi, Marchesini, Smye]
[Angeles, DeAngelis, Forshaw, Plätzer, Seymour]

[Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044]

[see also Nagy, Soper]

Unified framework requires evolution at the amplitude level as most general basis:

$$\sigma = \sum_n \int \text{Tr} [\mathbf{A}_n(\mu)] u(p_1, \dots, p_n) d\phi_n$$

$$\mathbf{A}_n(\mu) = |\mathcal{M}_n(\mu)\rangle \langle \mathcal{M}_n(\mu)|$$

Evolved **density operator**

Observable

Phase space

General expression of (partonic) cross section including all multiplicities and **virtual corrections and colour mixing** in all orders perturbation theory.

$$|\mathcal{M}_n(\mu)\rangle = \mathbf{Z}^{-1}(\mu, \epsilon) |\tilde{\mathcal{M}}_n(\epsilon)\rangle$$

[Sterman, Becher & Neubert, ...]

[Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044]

Parton shower picture encoded in recursive definition including **emission** and **virtual evolution** operators

$$\mathbf{A}_n(E) = \mathbf{V}(E, E_n) \mathbf{D}_n \mathbf{A}_{n-1}(E_n) \mathbf{D}_n^\dagger \mathbf{V}^\dagger(E, E_n) \theta(E - E_n)$$

Corresponds to tower of evolution equations for each partonic multiplicity.

Questions to be addressed include:

- Origin and generalizations of the evolution towards full parton showers, constraints for certain classes of observables
↔ **systematic expansion**
- Role and impact of the ordering variable, relation to virtual corrections
- Anomalous dimensions at higher orders
- Kinematic mappings and factorization properties of amplitudes

[Angeles, Forshaw, Seymour]

Soft evolution at leading order established, even beyond for NGL [Weigert, Caron-Huot]

$$\mathbf{A}_n(E) = \mathbf{V}(E, E_n) \mathbf{D}_n \mathbf{A}_{n-1}(E_n) \mathbf{D}_n^\dagger \mathbf{V}^\dagger(E, E_n) \theta(E - E_n)$$

$$\mathbf{V}_n(E, Q) = \text{P exp} \left(- \int_E^Q \frac{dq}{q} \mathbf{\Gamma}_n(q) \right) \quad \mathbf{D}_n = \sum_{i=1}^{n-1} \frac{p_i \cdot \epsilon^*(p_n)}{p_i \cdot p_n} \mathbf{T}_i$$

$$\mathbf{\Gamma}_n = \frac{\alpha_s}{\pi} \sum_{i < j} \int d\Omega \frac{n_i \cdot n_j}{n_i \cdot n \ n \cdot n_j} (-\mathbf{T}_i \cdot \mathbf{T}_j)$$

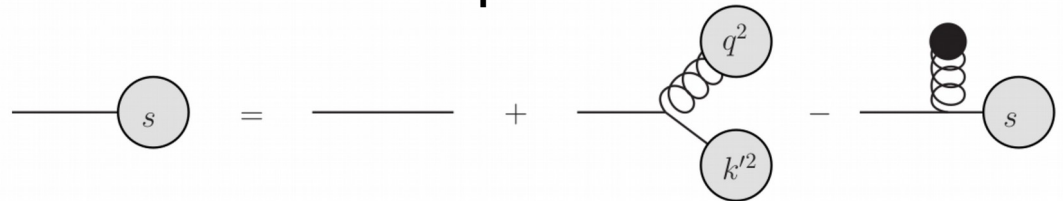
Next steps to be undertaken are:

- Systematic inclusion of soft- and hard-collinear effects
 - ↔ relate algorithm to coherent branching, access spin correlations
- Implement momentum conservation and clarify role of kinematic mappings
- Push to next-to-leading order in the soft sector, and beyond
- Smooth coverage of phase space and algorithmic construction desirable
- Make contact with NNLO subtraction

Analytic results are vital to judge on the accuracy of the algorithms: Link the general algorithm to analytic results focusing on both **direct QCD** as well as **effective field theory** setups. Not at all impossible!

Coherent branching calculation of jet mass with massive quarks

$$J(s, Q^2) = \delta(s) + \int_0^{Q^2} \frac{d\tilde{q}^2}{\tilde{q}^2} \int_0^1 dz P_{qq}[\alpha_s(z(1-z)\tilde{q}), z] \\ \times \left[\int_0^\infty dk'^2 \int_0^\infty dq^2 \delta\left(s - \frac{k'^2}{z} - \frac{q^2}{1-z} - z(1-z)\tilde{q}^2\right) J(k'^2, z^2\tilde{q}^2) J_g(q^2, (1-z)^2\tilde{q}^2) \right. \\ \left. - J(s, \tilde{q}^2) \right],$$



[Hoang, Plätzer, Samitz – JHEP 10 (2018) 200]

Link soft gluon evolution algorithm to effective theory for jets

$$E \frac{\partial \mathbf{G}_n(E)}{\partial E} = -\mathbf{\Gamma} \mathbf{G}_n(E) - \mathbf{G}_n(E) \mathbf{\Gamma}^\dagger + \mathbf{D}_n \mathbf{G}_{n-1}(E) \mathbf{D}_n^\dagger u(E, \hat{p}_n)$$

[Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044]

Derive how (non-)global resummations are recovered from colour-space evolution.

Fully **integrate new evolution inside the Herwig 7** event generator and perform extensive comparison against collider data.

- Role of colour reconnection models needs to be clarified.
- Availability of fixed-order calculations at amplitude level required.

Perform true **matching to NNLO calculations.**

- Link shower kinematics and kernels to subtraction schemes
- Devise matching subtractions from expanding new evolution to second order

Loads of preliminary studies and 'tools' underway which ensure that the **project gets moving very quickly**:

- Density operator evolution in colour matrix element corrections in Herwig

[Plätzer, Sjö Dahl – JHEP 1207 (2012) 042]

[Plätzer, Sjö Dahl, Thoren – JHEP 11 (2018) 009]

- Soft gluon evolution algorithm and systematic expansion around large-N

[Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044]

[De Angelis, Forshaw, Plätzer – soon]

- Spin correlation algorithm in Herwig

[Webster, Richardson – arXiv:1807.01955]

- MC algorithms for non-probabilistic evolution

[Plätzer, Sjö Dahl – EPJ Plus 127 (2012) 26]

- Kinematics and factorization of double real emissions

[Gieseke, Plätzer, Simpson – in progress]

- First steps towards collinear contributions

[Forshaw, Holguin, Plätzer – soon]

- Touching ground with colour reconnection models

[Gieseke, Kirchgaesser, Plätzer, Siodmok – JHEP 11 (2018) 149]

- NNLO cross sections available within STRIPPER framework

[Czakon – PLB 693 (2010) 259]

[Czakon, Heymes – NPB 890 (2014) 152]

Thank you!

Evolution equations for factorizing observables $u(p_1, \dots, p_n) = u(E_1, \hat{p}_1) \cdots u(E_n, \hat{p}_n)$
integrating over the intermediate scales:

$$E \frac{\partial \mathbf{G}_n(E)}{\partial E} = -\mathbf{\Gamma} \mathbf{G}_n(E) - \mathbf{G}_n(E) \mathbf{\Gamma}^\dagger + \mathbf{D}_n \mathbf{G}_{n-1}(E) \mathbf{D}_n^\dagger u(E, \hat{p}_n)$$

Recursive observables: $u_n(p_1, \dots, p_n) = u(p_n, \{p_1, \dots, p_{n-1}\}) u_{n-1}(p_1, \dots, p_{n-1})$

Genuine non-global case: $u(k, \{q\}) = \Theta_{\text{out}}(k) + \Theta_{\text{in}}(k) u_{\text{in}}(k, \{q\})$

Can identify global and non-global contributions simply by splitting anomalous dimension into `out' and `in' contributions.

Reproduce all available literature.

[Dasgupta, Salam – Phys.Lett. B512 (2001) 323]

[Forshaw, Kyrieleis, Seymour – JHEP 0608 (2006) 059]

[Weigert – Nucl.Phys. B685 (2004) 321]

[Caron-Huot – JHEP 1803 (2018) 036]

[Becher, Neubert, Rothen, Shao – JHEP 1611 (2016) 019]

Includes a **re-derivation of the BMS equation** at leading-N, being able to calculate subleading-N corrections to it.

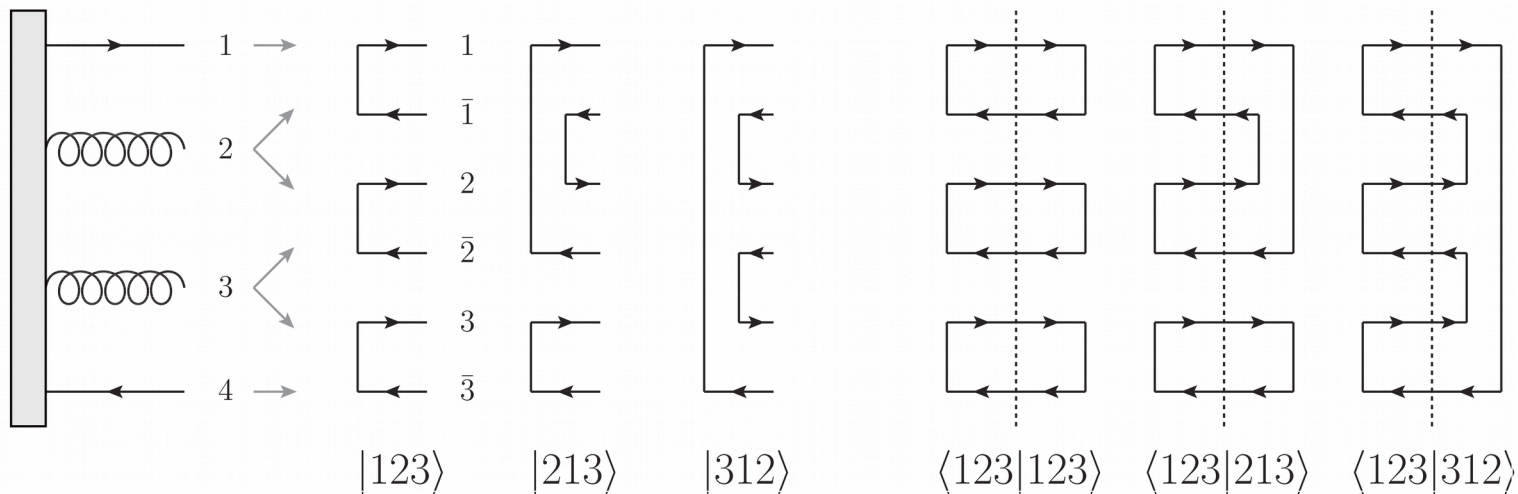
[Banfi, Marchesini, Smye – JHEP 0208 (2002) 006]

[Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044]

Express amplitudes in combinations of fundamental/anti-fundamental indices:

$$|\mathcal{M}\rangle = \sum_{\sigma} \mathcal{M}_{\sigma} |\sigma\rangle$$

$$|\sigma\rangle = \left| \begin{array}{ccc} 1 & \cdots & n \\ \sigma(1) & \cdots & \sigma(n) \end{array} \right\rangle = \delta_{\bar{\alpha}_{\sigma(1)}^{\alpha_1}} \cdots \delta_{\bar{\alpha}_{\sigma(n)}^{\alpha_n}}$$



Non-orthogonal basis:

$$1 = \sum_{\sigma} |\sigma\rangle [\sigma| \quad [\sigma|\tau\rangle = \langle\sigma|\tau\rangle = \delta_{\tau\sigma} \quad \text{Tr}[\mathbf{A}] = \sum_{\tau, \sigma} [\tau|\mathbf{A}|\sigma] \langle\sigma|\tau\rangle$$

Also overcomplete ... but computationally very handy: It's all about permutations.

[Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044]

Evolution operator in colour flow basis:

$$\begin{aligned}
 [\tau|e^{\Gamma}|\sigma\rangle &= \sum_{k=0}^{\infty} \frac{(-1)^k}{N^k} \sum_{l=0}^{\infty} \frac{(-\rho)^l}{l!} \sum_{\sigma_0, \dots, \sigma_{l-k}} \delta_{\tau\sigma_0} \delta_{\sigma_{l-k}\sigma} \prod_{\alpha=0}^{l-k-1} \Sigma_{\sigma_{\alpha}\sigma_{\alpha+1}} R(\{\sigma_0, \dots, \sigma_{l-k}\}) \\
 &= \delta_{\tau\sigma} \left(e^{-N\Gamma_{\sigma}} + e^{-N\Gamma_{\sigma}} \frac{\rho}{N} \right) - \frac{1}{N} \frac{e^{-N\Gamma_{\tau}} - e^{-N\Gamma_{\sigma}}}{\Gamma_{\tau} - \Gamma_{\sigma}} \Sigma_{\tau\sigma} + \text{NNLC}
 \end{aligned}$$

[Plätzer – EPJ C 74 (2014) 2907]

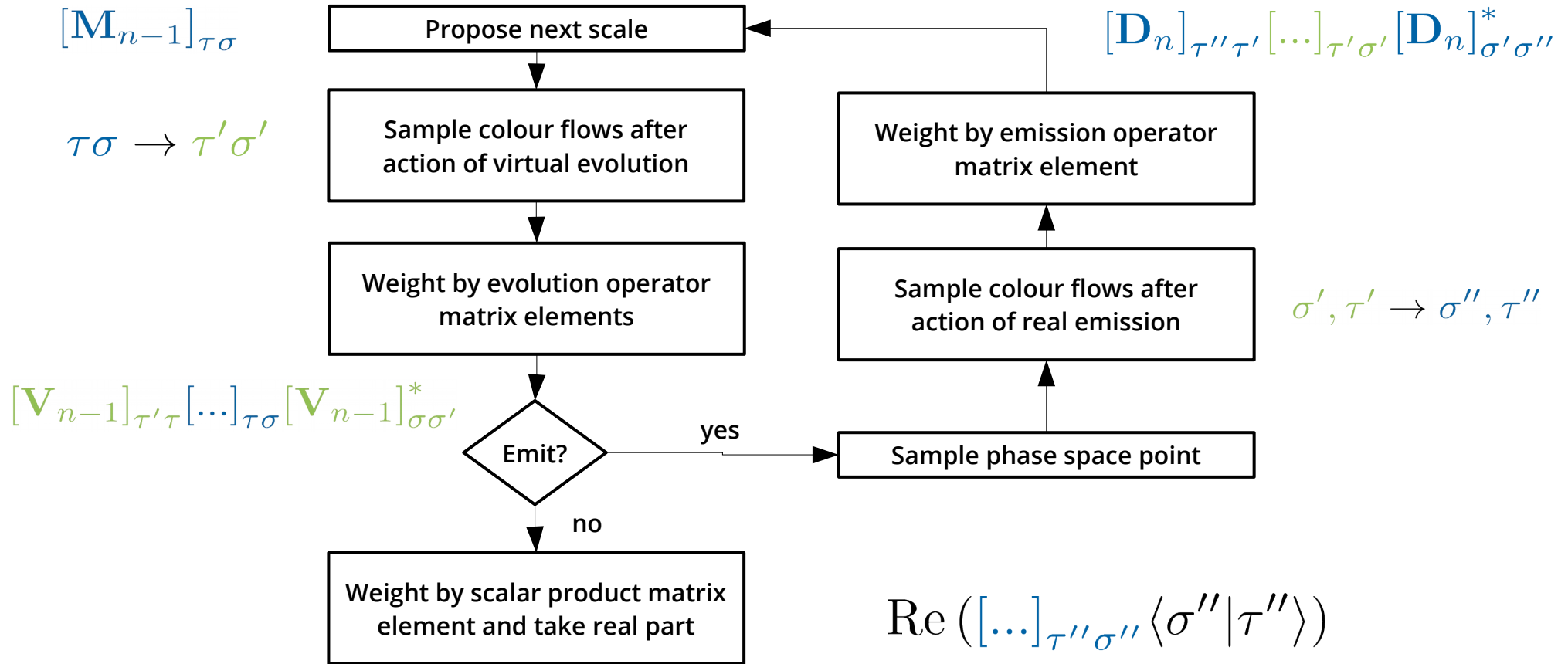
Sum terms enhanced by $\alpha_s N$ to all orders, insert perturbations in $1/N$.

Take into account real emission contributions and the final suppression by the scalar product matrix element.

				virtuals	reals
N^3			Γ^3	$(0 \text{ flips}) \times 1 \times (\alpha_s N)^n$	$(t[\dots]t _{0 \text{ flips}})^{r-1} t[\dots]t _{2 \text{ flips}} \times 1$ $(t[\dots]t _{0 \text{ flips}})^{r-1} t[\dots]s _{1 \text{ flip}} \times N^{-1}$ $(t[\dots]t _{0 \text{ flips}})^{r-1} s[\dots]s _{0 \text{ flips}} \times N^{-2}$
N^2		Γ^2	$\Sigma\Gamma^2$	$(1 \text{ flip}) \times \alpha_s \times (\alpha_s N)^n$	$(t[\dots]t _{0 \text{ flips}})^r$ $(t[\dots]t _{0 \text{ flips}})^{r-1} t[\dots]s _{1 \text{ flip}} \times N^{-1}$
N^1	Γ	$\Sigma\Gamma$	$\rho\Gamma^2$	$(0 \text{ flips}) \times \alpha_s N^{-1} \times (\alpha_s N)^n$	$(t[\dots]t _{0 \text{ flips}})^r$
N^0	1	Σ	$\rho\Sigma\Gamma$	$(0 \text{ flips}) \times \alpha_s^2 \times (\alpha_s N)^n$ $(2 \text{ flips}) \times \alpha_s^3 \times (\alpha_s N)^n$	$(t[\dots]t _{0 \text{ flips}})^r$ $(t[\dots]t _{0 \text{ flips}})^{r-1} t[\dots]t _{2 \text{ flips}}$
N^{-1}		$\rho 1$	$\rho\Sigma$		
N^{-2}			$\rho^2 1$		
N^{-3}					
	α_s^0	α_s^1	α_s^2	α_s^3	

[De Angelis, Forshaw, Plätzer – arXiv:1902.xxxxx, first presented at PSR '18 and others]

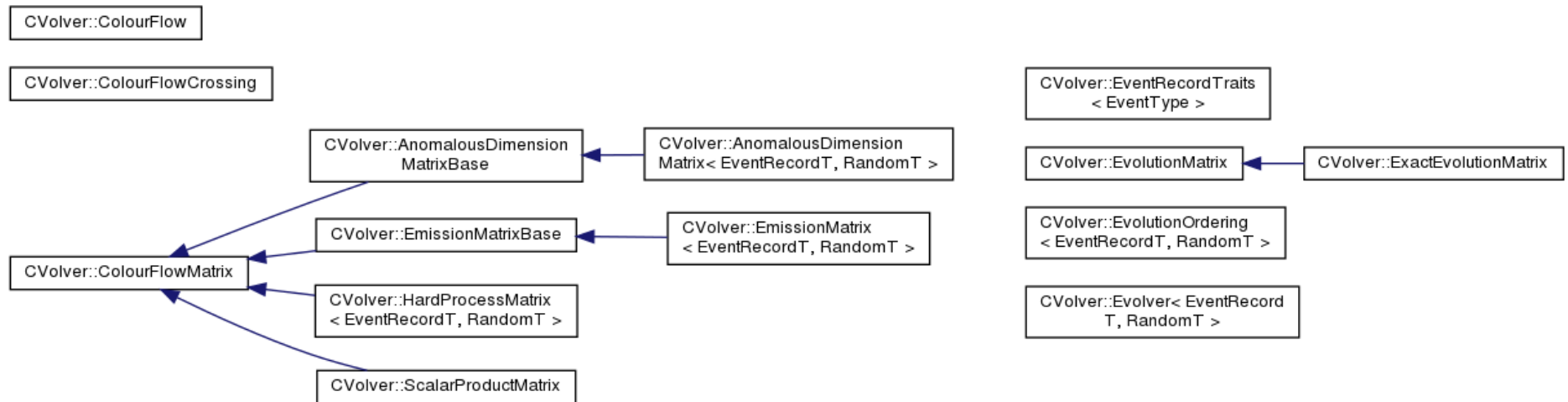
$$\mathbf{A}_n(E) = \mathbf{V}(E, E_n) \mathbf{D}_n \mathbf{A}_{n-1}(E_n) \mathbf{D}_n^\dagger \mathbf{V}^\dagger(E, E_n) \theta(E - E_n)$$



[De Angelis, Forshaw, Plätzer – arXiv:1902.xxxxx, first presented at PSR '18 and others]

A framework to solve multi-differential evolution equations in colour space.
Concise, simple, and light-weight code structure.

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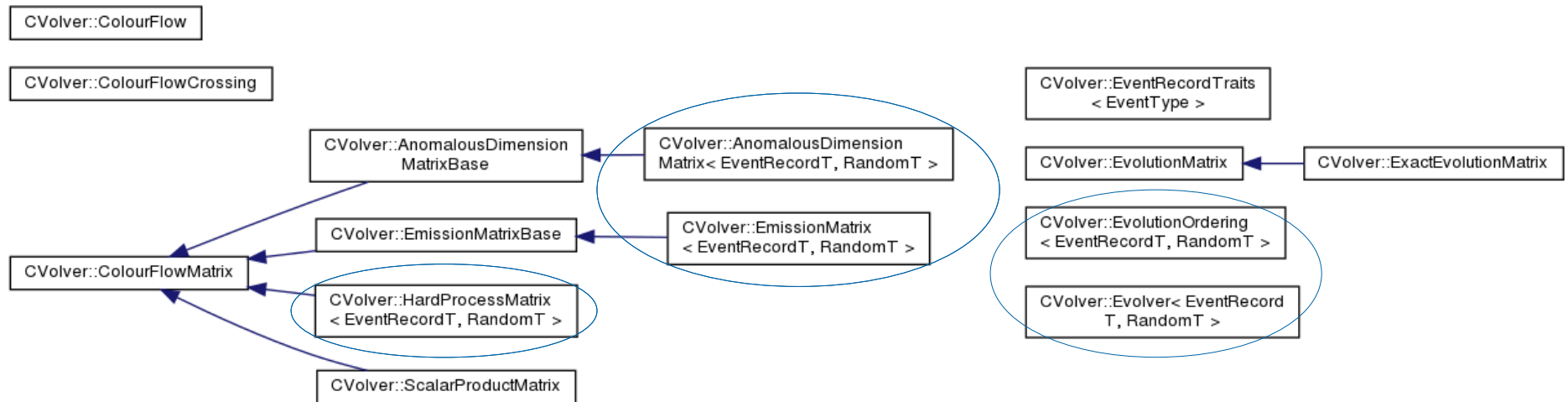
Dedicated Monte Carlo algorithms to sample colour structures.

Plugin approach can accommodate anything from (N)GLs to full parton showers.

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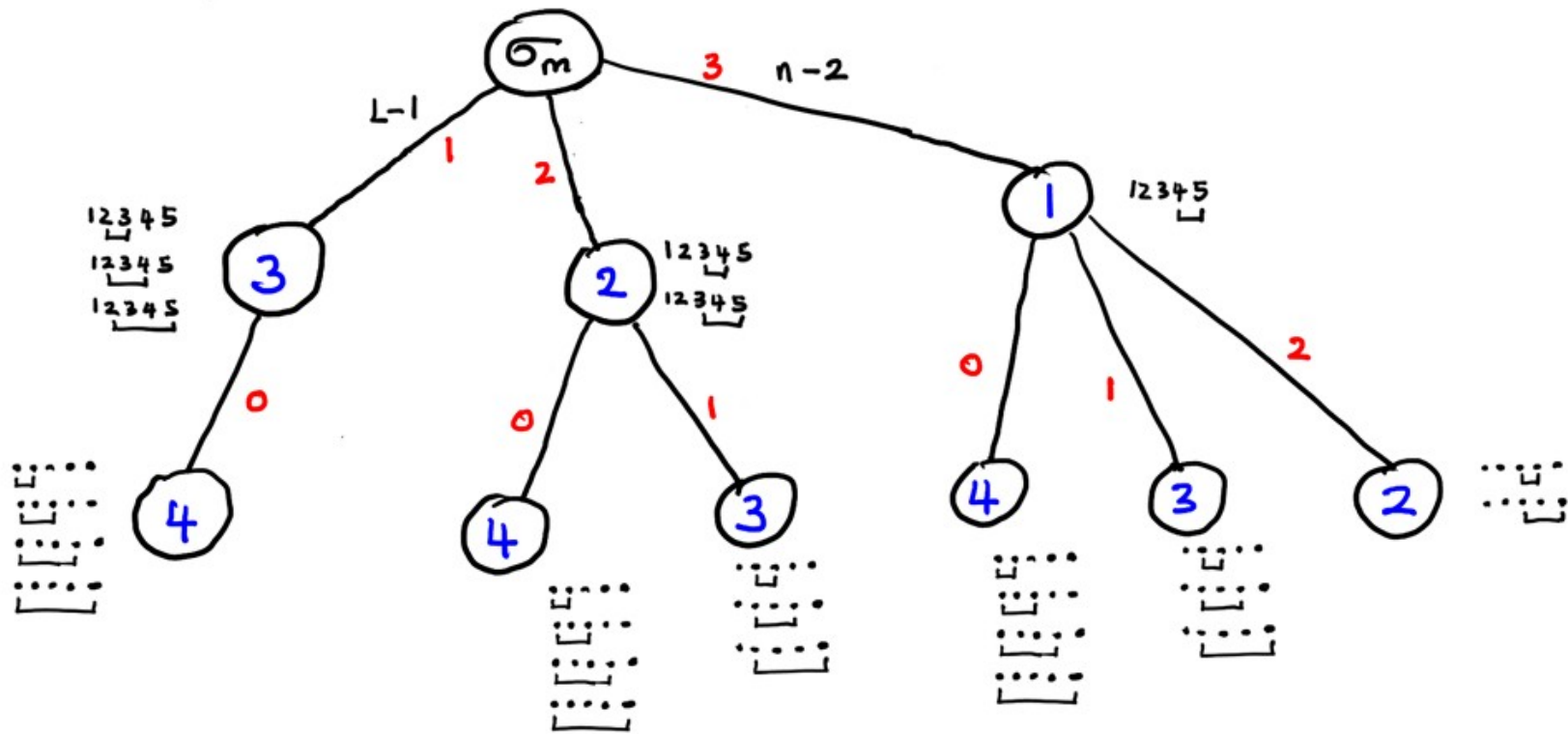
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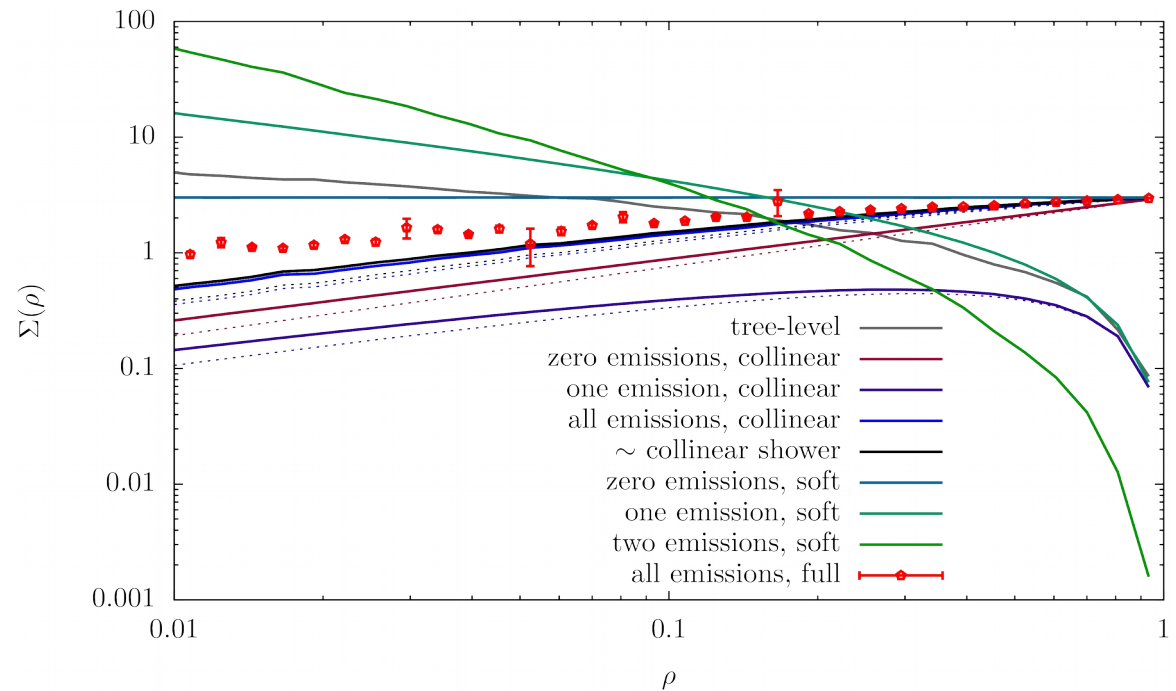
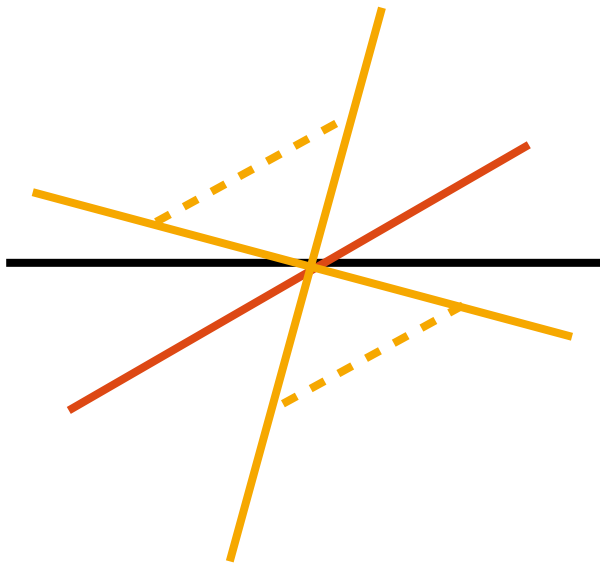
Importance sampling in colour space rules: $\#(\tau, \tau') \sim 1/N^{\#(\tau, \tau')}$

Enumerate and address permutations with fixed cycle length:



[De Angelis, Forshaw, Plätzer – arXiv:1902.xxxxx, first presented at PSR '18 and others]

Code is differential for a large class of (non-global) observables.
Example: Cone-dijet veto cross section.



1/N breakdowns possible, scales up to several 10s of emissions for $d=2$.

The cluster model is based on a single colour flow after the shower has stopped.

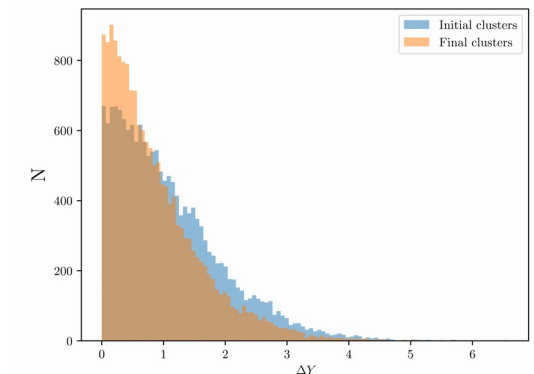
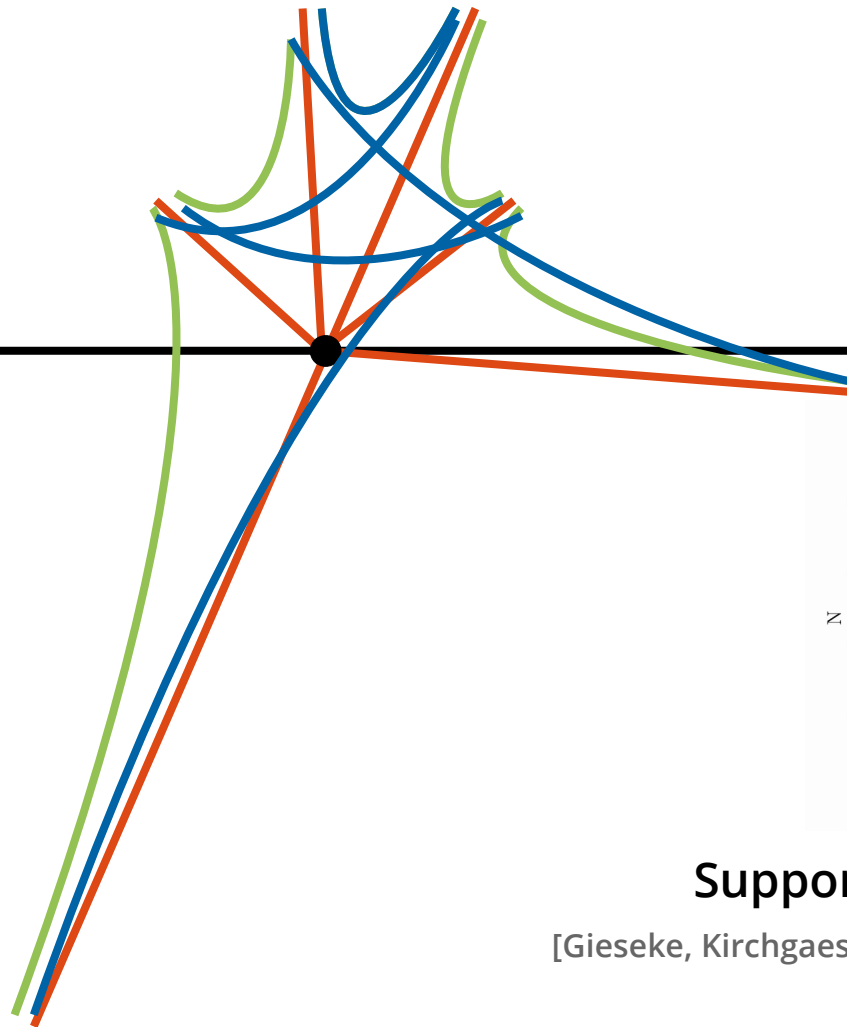
Essentially no evolution is considered, just decays.

View as an evolving amplitude, driven by a single initial colour flow:

$$|\mathcal{M}\rangle = e^{\Gamma} |\text{clusters}\rangle$$

$$P_{\text{reco}} \sim |\langle \text{clusters}' | \mathcal{M} \rangle|^2$$

[Gieseke, Kirchgaesser, Plätzer, Siodmok – JHEP 11 (2018) 149]



Supports geometric models

[Gieseke, Kirchgaesser, Plätzer – EPJ C 78 (2018) 99]