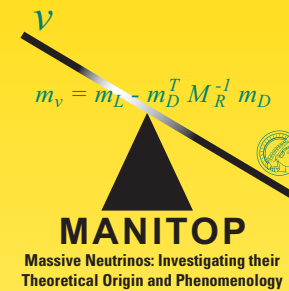


Neutrinoeigenschaften



WERNER RODEJOHANN
AT-PHYSIK IN D
30/09/14



**Astroteilchenphysik
in Deutschland**
Status und Perspektiven

30.9. - 1.10.2014 am KIT Karlsruher Institut für Technologie
Campus Süd, Tuilla-Hörsaal

Programmkomitee Karsten Danzmann Roland Diehl Dieter Horns Josef Jochum Karl-Heinz Kampert Lothar Oberauer Elisa Resconi Günter Sigl Christian Weinheimer (chair)	Lokales Organisationskomitee Johannes Blümer Klaus Ebel Andreas Haungs Tim Hüne Marie-Christine Kautmann Bianca Kethauer
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<https://indico.scc.kit.edu/indico/event/ATP2014>

HELMHOLTZ-GESELLSCHAFT
Bundeshilfen für Bildung und Forschung
Deutsche Forschungsgemeinschaft DFG



Neutrinoeigenschaften in Kürze

- massiv
- mischen
- wahrscheinlich Majorana

Standardmodel

addiere m_ν zu

Spezies	#	Σ
Quarks	10	10
Leptonen	3	13
Ladung	3	16
Higgs	2	18
starke CP	1	19

Standardmodel

addiere m_ν

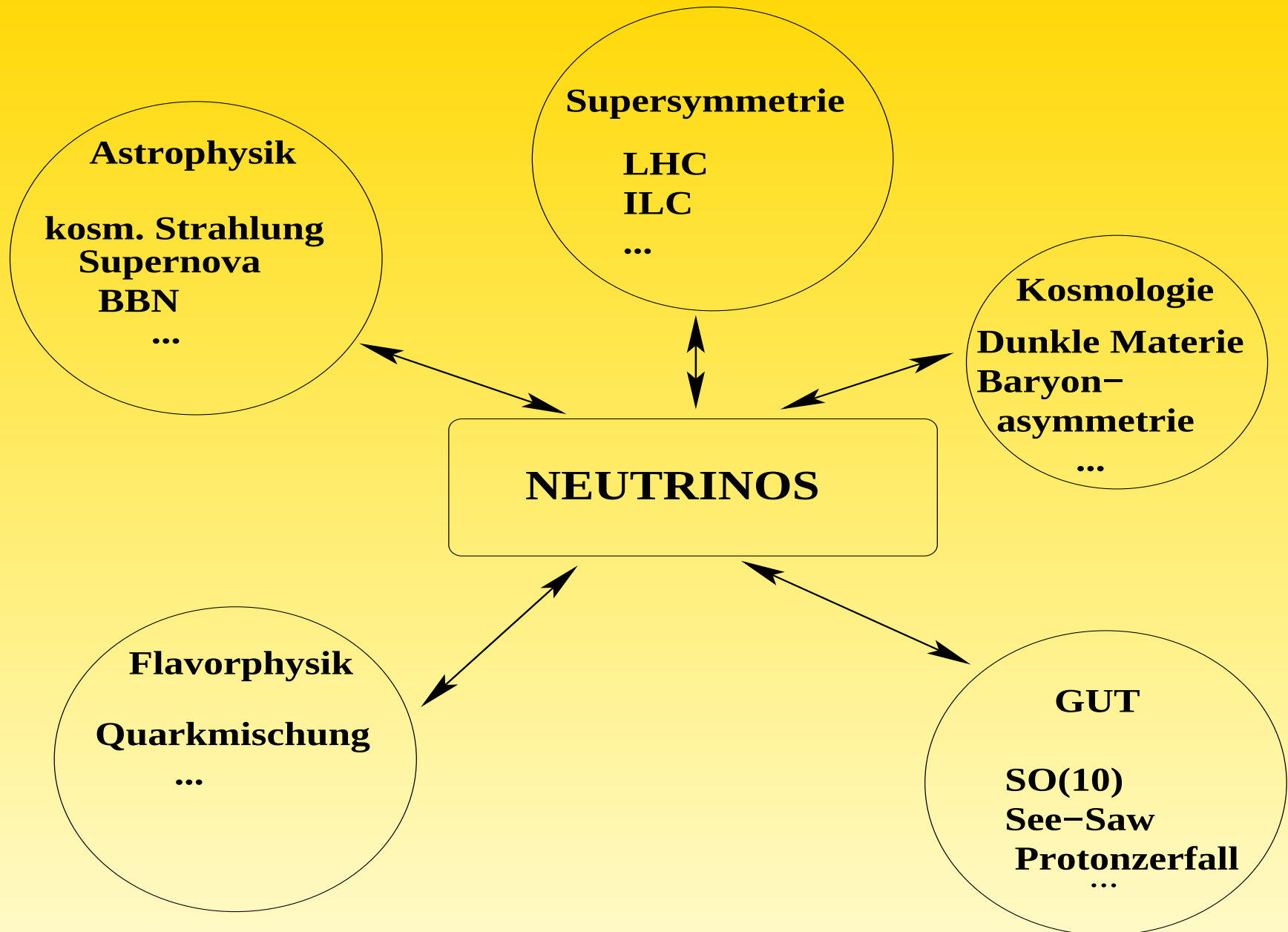
Spezies	#	Σ		Spezies	#	Σ
Quarks	10	10		Quarks	10	10
Leptonen	3	13	→	Leptonen	12 (10)	22 (20)
Ladung	3	16		Ladung	3	25 (23)
Higgs	2	18		Higgs	2	27 (25)
starke CP	1	19		starke CP	1	28 (26)

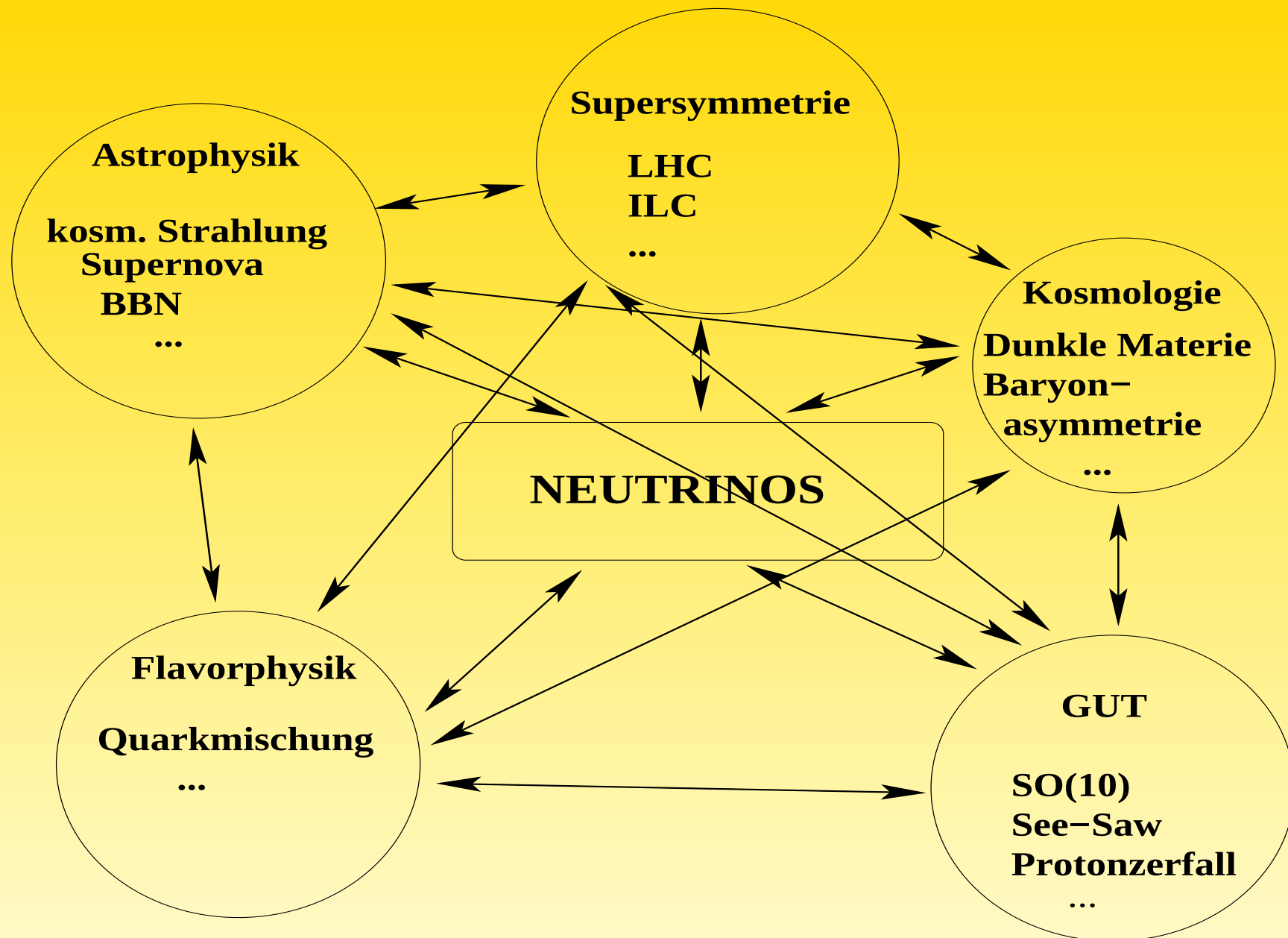
Standardmodel*

addiere m_ν

Spezies	#	Σ		Spezies	#	Σ
Quarks	10	10		Quarks	10	10
Leptonen	3	13	→	Leptonen	12 (10)	22 (20)
Ladung	3	16		Ladung	3	25 (23)
Higgs	2	18		Higgs	2	27 (25)
starke CP	1	19		starke CP	1	28 (26)

- plus neue Energieskala
- plus neue SM Darstellung
- plus neue Konzepte
- plus...





Neutrino parameter bei niedriger Energie

$$\mathcal{L} = \frac{1}{2} \nu^T m_\nu \nu \quad \text{mit} \quad m_\nu = U \text{diag}(m_1, m_2, m_3) U^T$$

und PMNS Matrix

$$U = \begin{pmatrix} c_{12} c_{13} & s_{12} c_{13} & s_{13} e^{-i\delta} \\ -s_{12} c_{23} - c_{12} s_{23} s_{13} e^{i\delta} & c_{12} c_{23} - s_{12} s_{23} s_{13} e^{i\delta} & s_{23} c_{13} \\ s_{12} s_{23} - c_{12} c_{23} s_{13} e^{i\delta} & -c_{12} s_{23} - s_{12} c_{23} s_{13} e^{i\delta} & c_{23} c_{13} \end{pmatrix} P$$

mit $P = \text{diag}(e^{i\alpha}, e^{i\beta}, 1)$ ($\leftrightarrow$ Majorana, Leptonzahlverletzung)

$\Rightarrow$ 3 Winkel, 3 Phasen, 3 Massen

3 Aufgaben

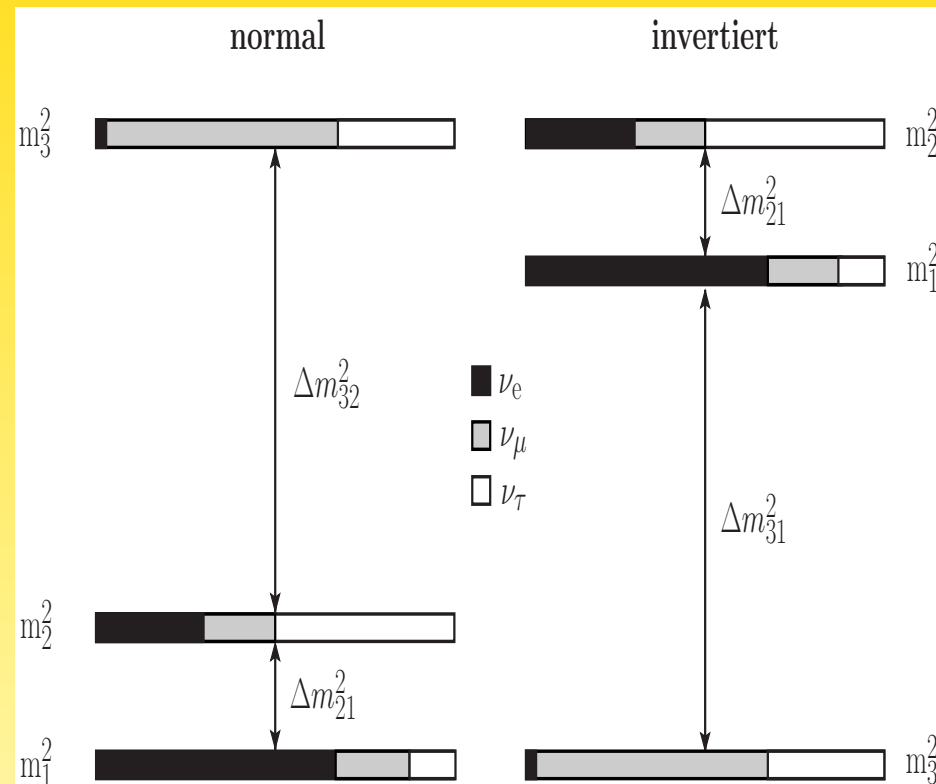
- Parameter bestimmen
- Parameter erklären
- Überprüfen ob minimale Beschreibung ("3 Majorana-Neutrino Paradigma") korrekt ist

Status 2014

9 physikalische Parameter in m_ν

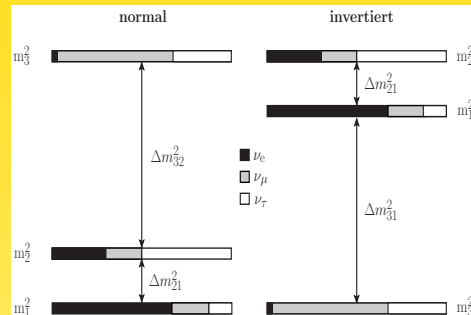
- θ_{12} und $m_2^2 - m_1^2$
- θ_{23} und $|m_3^2 - m_2^2|$
- θ_{13}
- m_1, m_2, m_3
- $\text{sgn}(m_3^2 - m_2^2)$
- Dirac-Phase δ
- Majorana-Phasen α und β

Massenskala und -ordnung



Skala und normal/invertiert fundamentaler Input für Modellbauer!

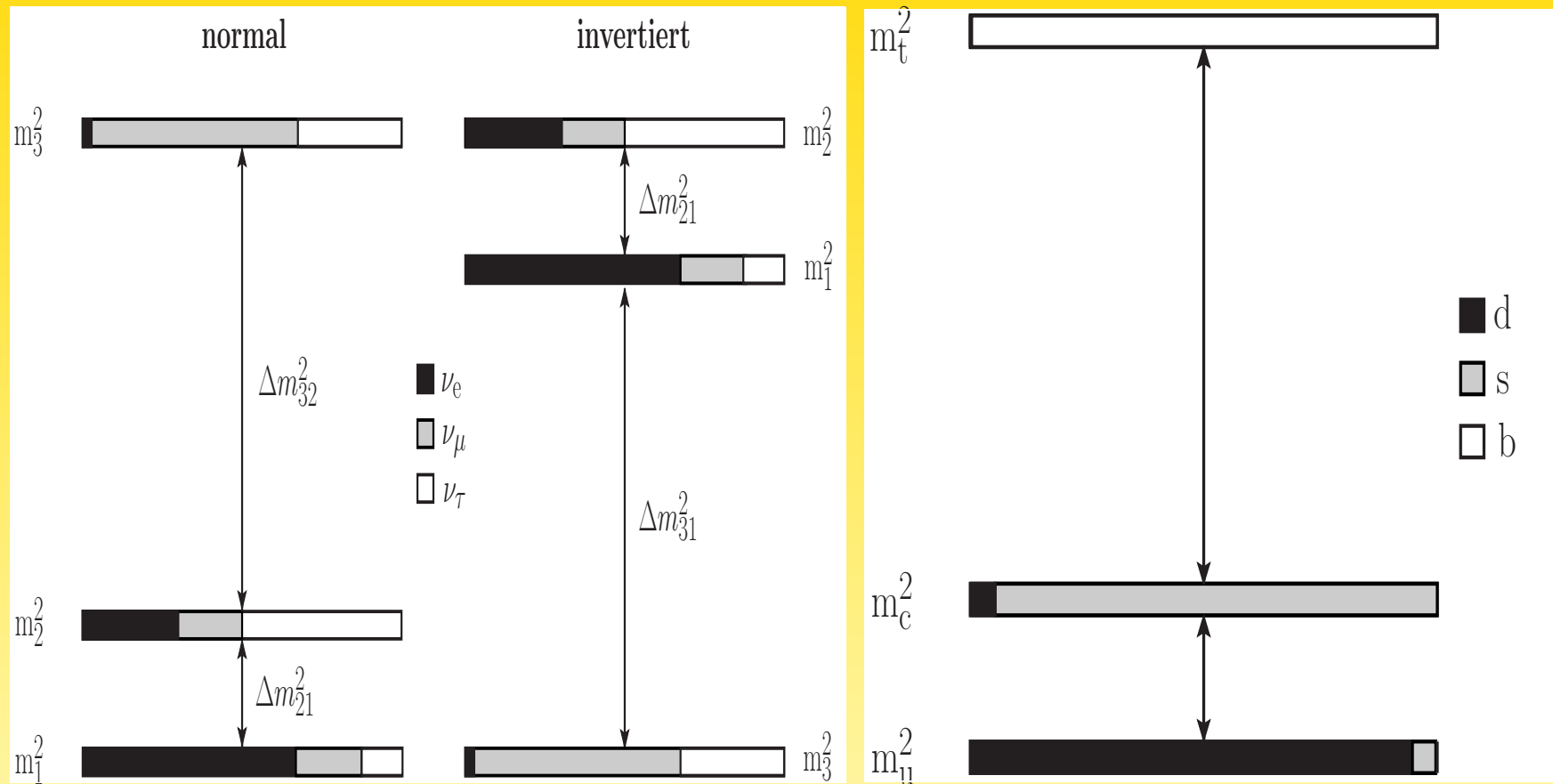
Massenskala und -ordnung



$$(m_\nu)_{\text{NH}} \sim \begin{pmatrix} \epsilon^2 & \epsilon & \epsilon \\ \epsilon & 1 & 1 \\ \epsilon & 1 & 1 \end{pmatrix} \quad \text{versus} \quad (m_\nu)_{\text{IH}} \sim \begin{pmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix}$$

Allgemein: $m_\nu \propto 1/\Lambda_{\text{NP}}$

Skala und normal/invertiert fundamentaler Input für Modellbauer!



Warum so verschieden?? $\leftrightarrow$ Flavoursymmetrien!

Neutrinomasse

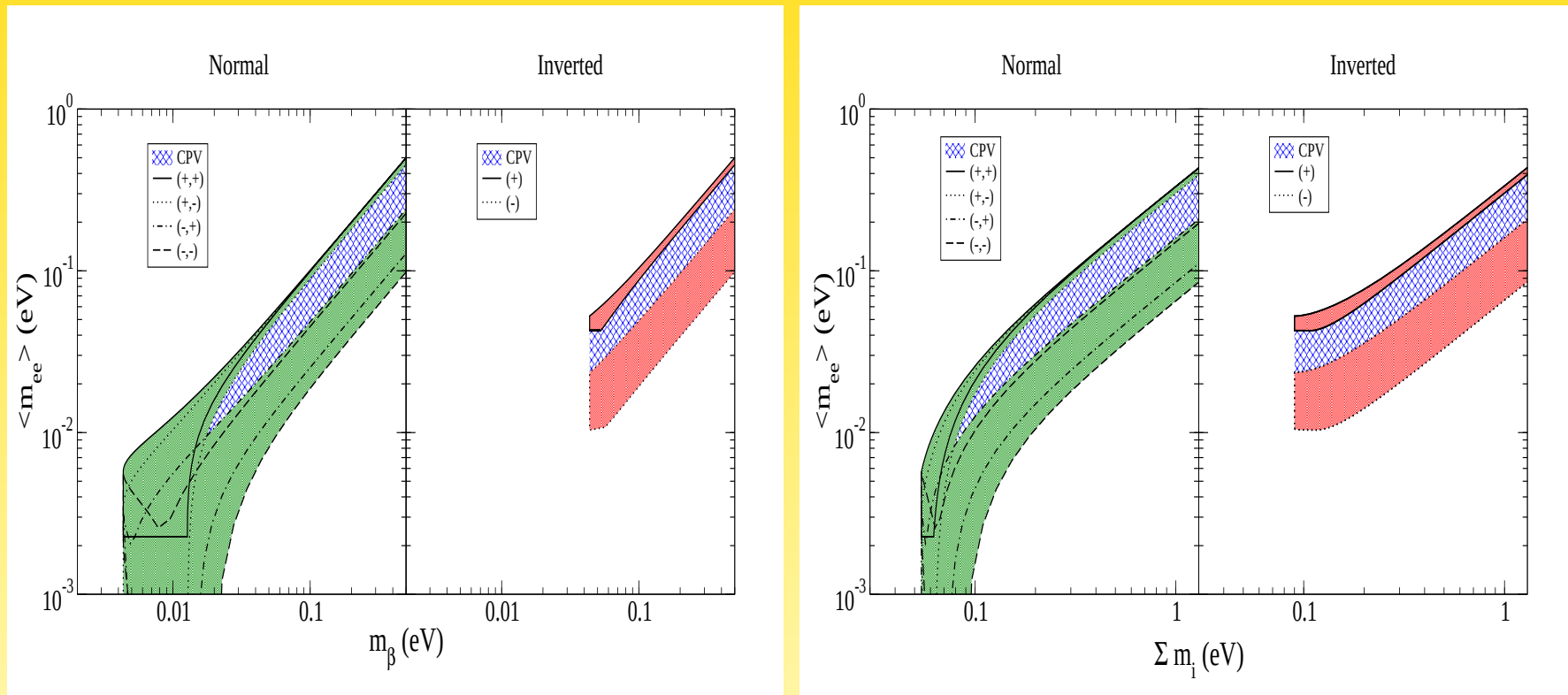
$$m(\text{schwerste}) \gtrsim \sqrt{|m_3^2 - m_1^2|} \simeq 0.05 \text{ eV}$$

3 **komplementäre** Methoden zur Messung:

Methode	Observable	geg. [eV]	nah/fern [eV]	pro	con
Kurie	$\sqrt{\sum U_{ei} ^2 m_i^2}$	2.3	0.2/0.1	model-unabh.; theo. sauber	final?; schwächstes
Kosmo.	$\sum m_i$	0.7	0.3/0.05	bestes; NH/IH	Systematik; model-abh.
$0\nu\beta\beta$	$ \sum U_{ei}^2 m_i $	0.3	0.1/0.05	fundamental; NH/IH	model-abh.; theo. schmutzig

Alle 3 unverzichtbar!

Komplementarität!



$$|m_{ee}| = |U_{ei}^2 m_i|, \quad m_\beta = \sqrt{|U_{ei}|^2 m_i^2} \quad \text{und} \quad \Sigma = \sum m_i$$

Massenskala

$$\begin{pmatrix} u \\ d \end{pmatrix} \text{ mit } m_u \simeq m_d \quad \text{versus} \quad \begin{pmatrix} \nu_e \\ e \end{pmatrix} \text{ mit } m_\nu \simeq 10^{-6} m_e$$

warum m_ν winzig? $\Rightarrow$ Seesaw-Mechanismus



Ursprung kleiner Neutrinomassen: Seesaw-Mechanismus

Einfachstes Konzept^a führt drei neue Aspekte ein:

- fermionische Singlets $N_R \sim (1, 0)$
- neue Energieskala $M_R (\propto 1/m_\nu)$
- Leptonzahlverletzung



$$m_\nu = m_D^2/M_R = m_{\text{SM}} \epsilon \text{ mit } \epsilon = m_{\text{SM}}/M_R$$

^aVorhersage von $SO(10)$ Modellen!

Tests des Seesaw?

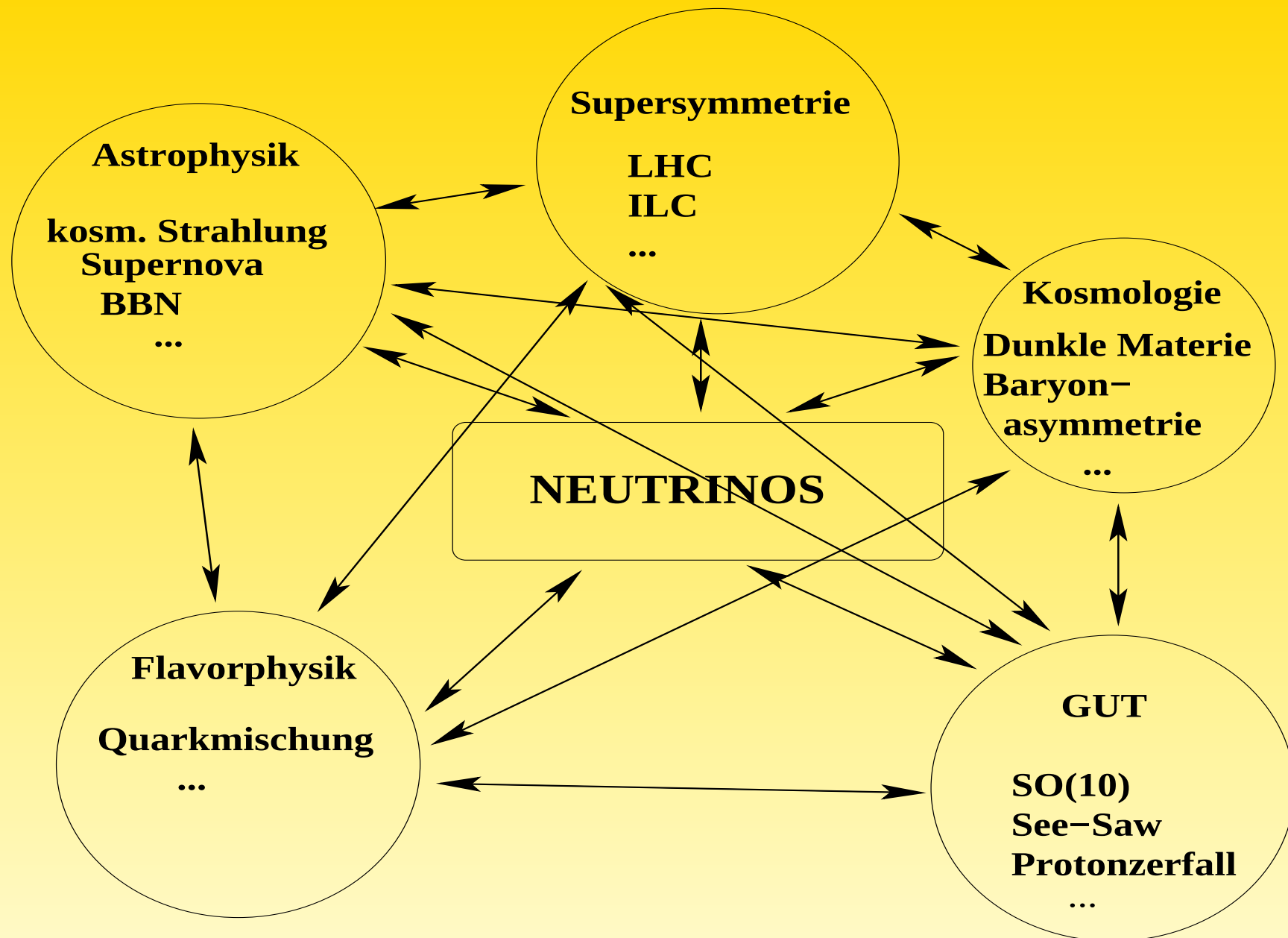
naiv, $M_R = m_D^2 / m_\nu = v^2 / \sqrt{\Delta m_A^2} \sim 10^{15} \text{ GeV} \dots$

Seesaw-Portal: $\bar{L} \Phi N_R \longrightarrow m_D:$

Vertex aus Lepton-Doublet L , Higgs Φ , Singlet N_R !

- $N_R \rightarrow L \Phi$: Leptogenese
- $L_\alpha \rightarrow N_R \Phi \rightarrow L_\beta$: Leptonflavorverletzung
- $\Phi \rightarrow N_R L \rightarrow \Phi$: Vakuumstabilität

(plus neutrinoloser doppelter Betazerfall)



Pfade zu Neutrinomassen

	Inhalt	Quantenzahl Messenger	$\mathcal{L}$	m_ν	Skala
“SM” (Dirac-Masse)	RH ν	$N_R \sim (1, 0)$	$h \overline{N_R} \Phi L$	$h v$	$h = \mathcal{O}(10^{-12})$
“effektiv” (dim 5 Operator)	neue Skala + LZV	–	$h \overline{L^c} \Phi \Phi L$	$\frac{h v^2}{\Lambda}$	$\Lambda = 10^{14} \text{ GeV}$
“direkt” (type II seesaw)	Higgs-Triplet + LZV	$\Delta \sim (3, 2)$	$h \overline{L^c} \Delta L + \mu \Phi \Phi \Delta$	$h v_T$	$\Lambda = \frac{1}{h \mu} M_\Delta^2$
“indirekt 1” (type I seesaw)	RH ν + LZV	$N_R \sim (1, 0)$	$h \overline{N_R} \Phi L + \overline{N_R} M_R N_R^c$	$\frac{(h v)^2}{M_R}$	$\Lambda = \frac{1}{h} M_R$
“indirekt 2” (type III seesaw)	Fermion-Triplets + LZV	$\Sigma \sim (3, 0)$	$h \overline{\Sigma} L \Phi + \text{Tr} \overline{\Sigma} M_\Sigma \Sigma$	$\frac{(h v)^2}{M_\Sigma}$	$\Lambda = \frac{1}{h} M_\Sigma$

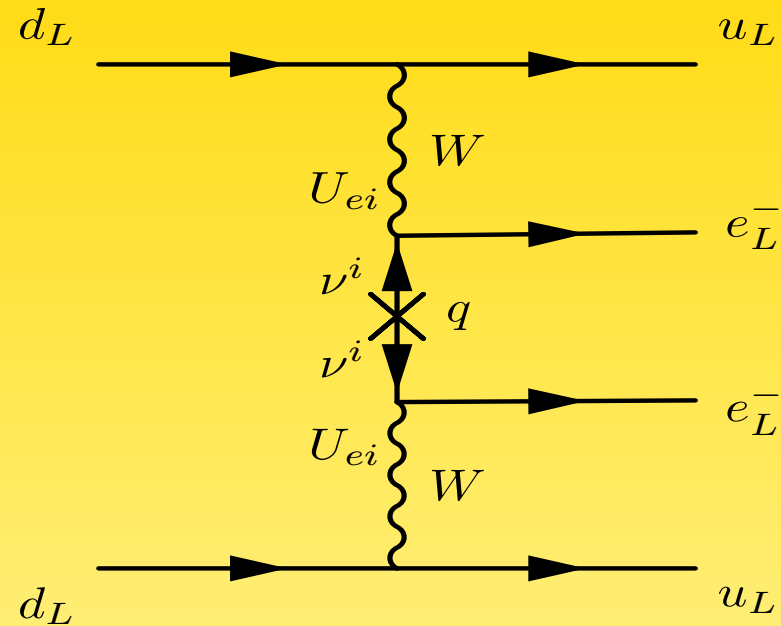
plus Seesaw-Varianten (linear, double, inverse, ...)

plus radiative Mechanismen

plus zusätzliche Raumdimensionen

plusplusplus

Konsequenz fast aller Mechanismen: Leptonzahlverletzung



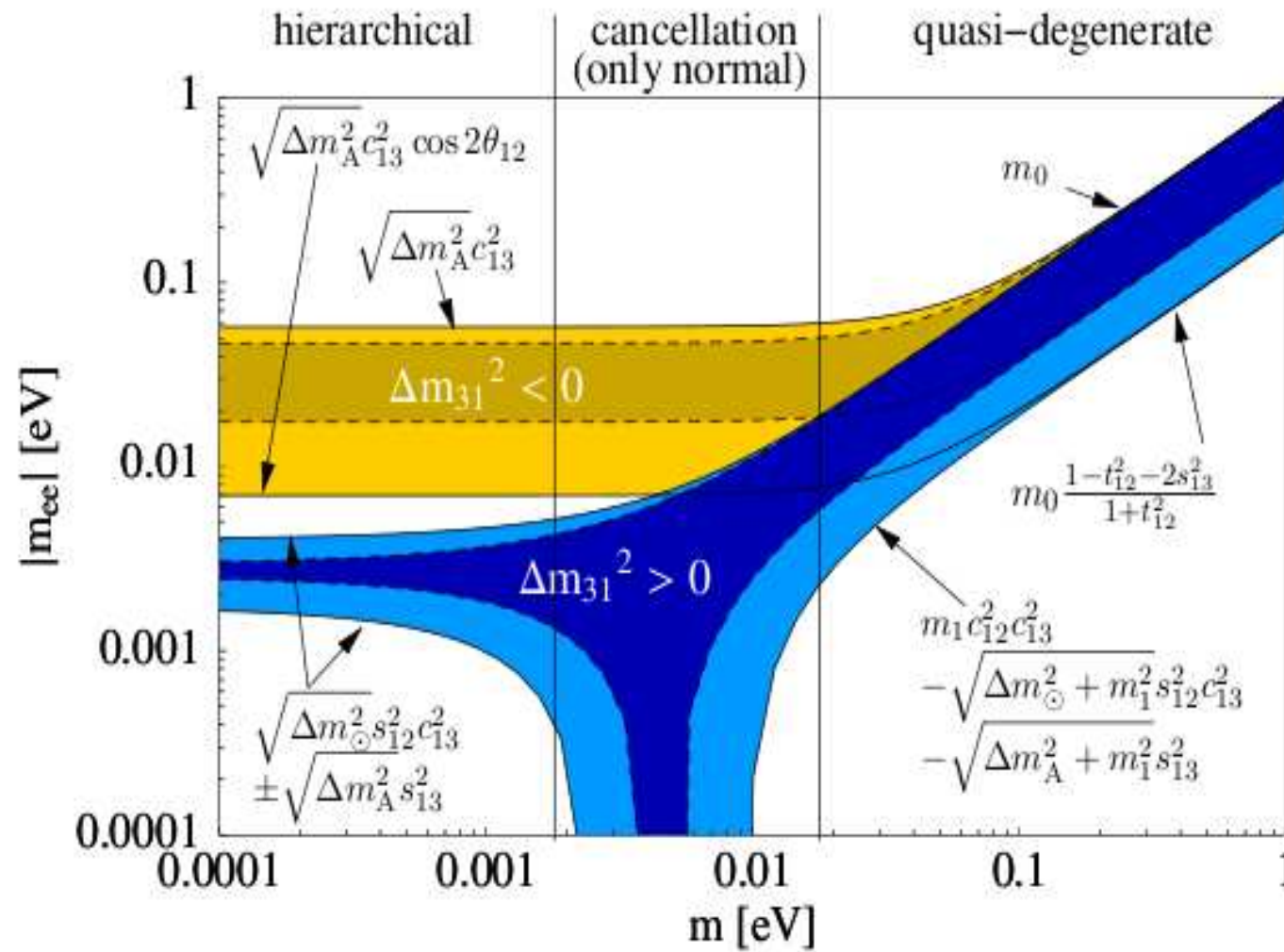
Neutrinoloser doppelter Betazerfall $(A, Z) \rightarrow (A, Z + 2) + 2 e^-$

$$|m_{ee}| \equiv \left| \sum U_{ei}^2 m_i \right| = \left| |U_{e1}|^2 m_1 + |U_{e2}|^2 m_2 e^{2i\alpha} + |U_{e3}|^2 m_3 e^{2i\beta} \right|$$

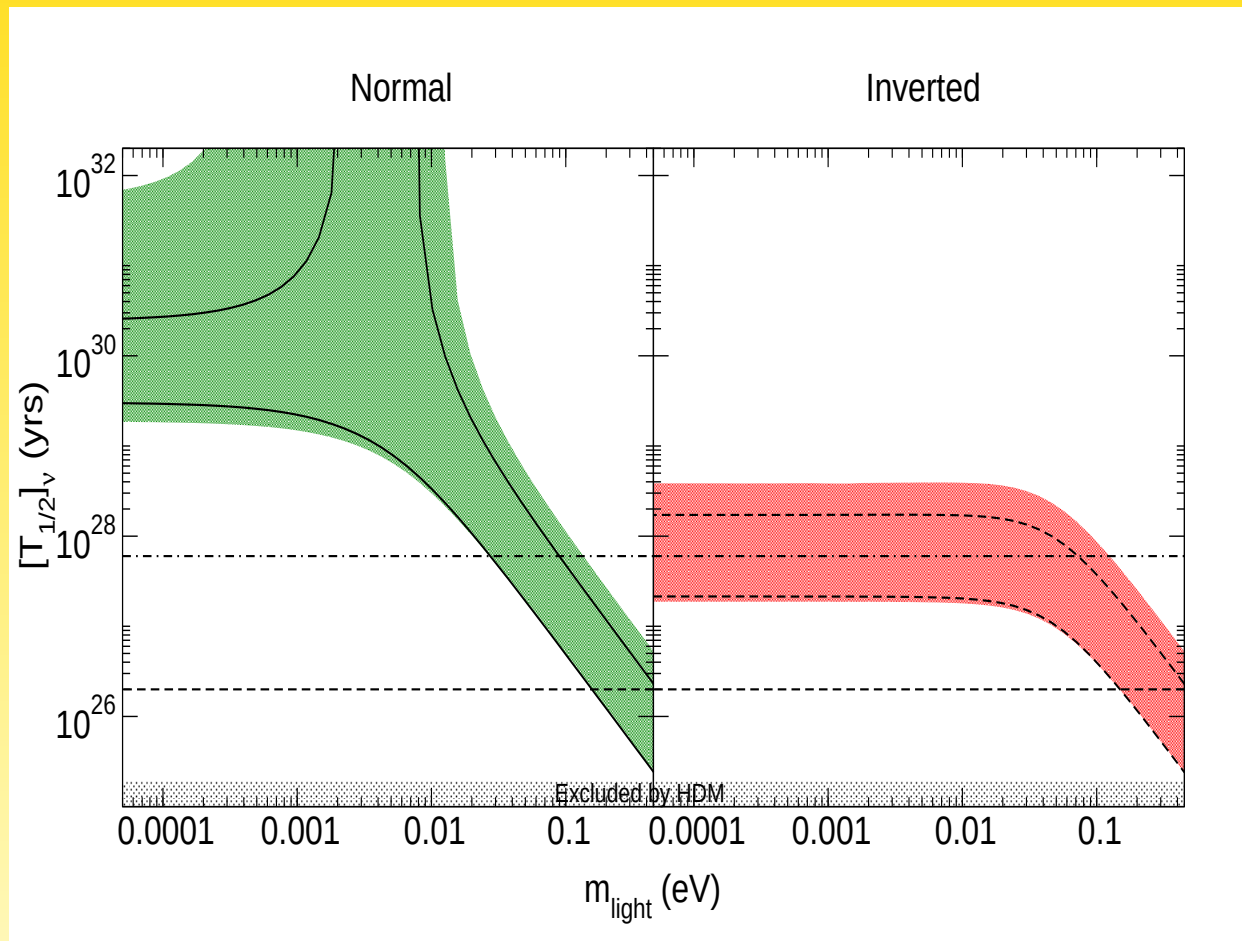
$$= f(\theta_{12}, |U_{e3}|, m_i, \text{sgn}(\Delta m_A^2), \alpha, \beta)$$

7 von 9 Parametern!

Statt des üblichen Plots...



...Lebensdauer vs. kleinste Masse

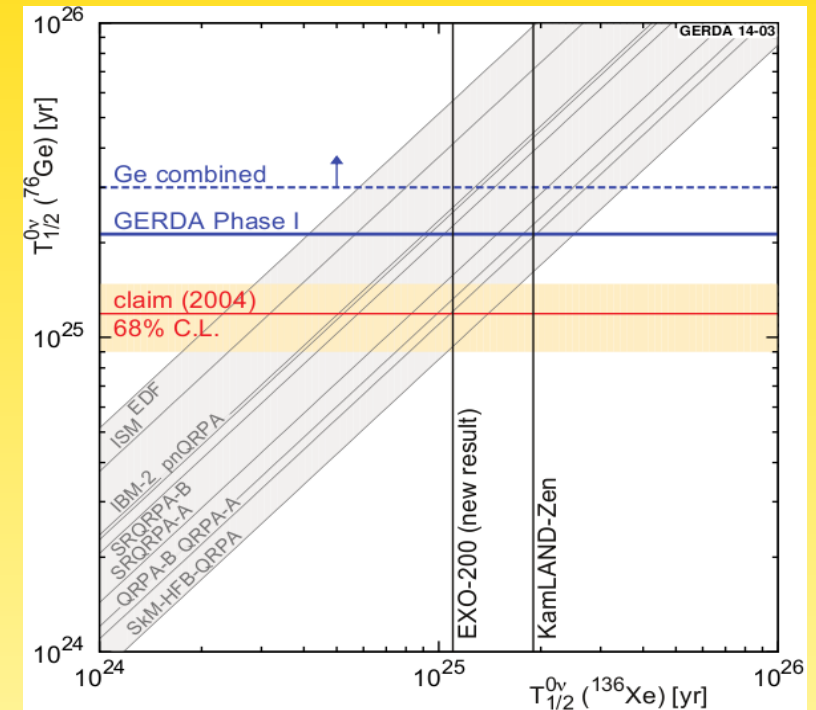


Laufende, bald startende Experimente

HDM-Limit von 2001 erst 2012/13 verbessert...

Name	Isotop	Quelle = Detektor; kalorimetrisch mit			Quelle $\neq$ Detektor Topologie
		ΔE hoch	ΔE niedr.	Topologie	
 AMoRE	^{100}Mo	✓	—	—	—
CANDLES	^{48}Ca	—	✓	—	—
 COBRA	^{116}Cd (und ^{130}Te)	—	—	✓	—
CUORE	^{130}Te	✓	—	—	—
DCBA/MTD	^{82}Se / ^{150}Nd	—	—	—	✓
 EXO	^{136}Xe	—	—	✓	—
 GERDA	^{76}Ge	✓	—	—	—
IHE	^{82}Se / ^{100}Mo / ^{116}Cd / ^{130}Te	✓	—	—	—
KamLAND-Zen	^{136}Xe	—	✓	—	—
LUCIFER	^{82}Se / ^{100}Mo / ^{130}Te	✓	—	—	—
 LUMINEU	^{100}Mo	✓	—	—	—
MAJORANA	^{76}Ge	✓	—	—	—
MOON	^{82}Se / ^{100}Mo / ^{150}Nd	—	—	—	✓
NEXT	^{136}Xe	—	—	✓	—
 SNO+	^{130}Te	—	✓	—	—
SuperNEMO	^{82}Se / ^{150}Nd	—	—	—	✓
XMASS	^{136}Xe	—	✓	—	—

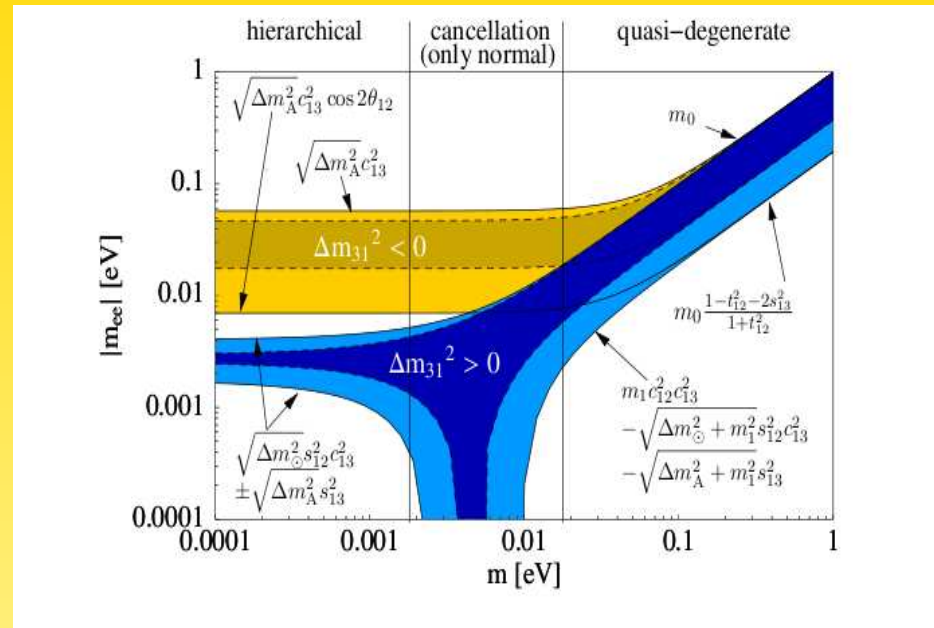
NME	^{76}Ge		^{136}Xe	
	GERDA	comb	KLZ	comb
EDF(U)	0.32	0.27	0.13	—
ISM(U)	0.52	0.44	0.24	—
IBM-2	0.27	0.23	0.16	—
pnQRPA(U)	0.28	0.24	0.17	—
SRQRPA-B	0.25	0.21	0.15	—
SRQRPA-A	0.31	0.26	0.23	—
QRPA-A	0.28	0.24	0.25	—
SkM-HFB-QRPA	0.29	0.24	0.28	—



GERDA, Schwingenheuer

Bhupal Dev, Goswami, Mitra,
W.R., Phys. Rev. **D88**

Invertierte Ordnung



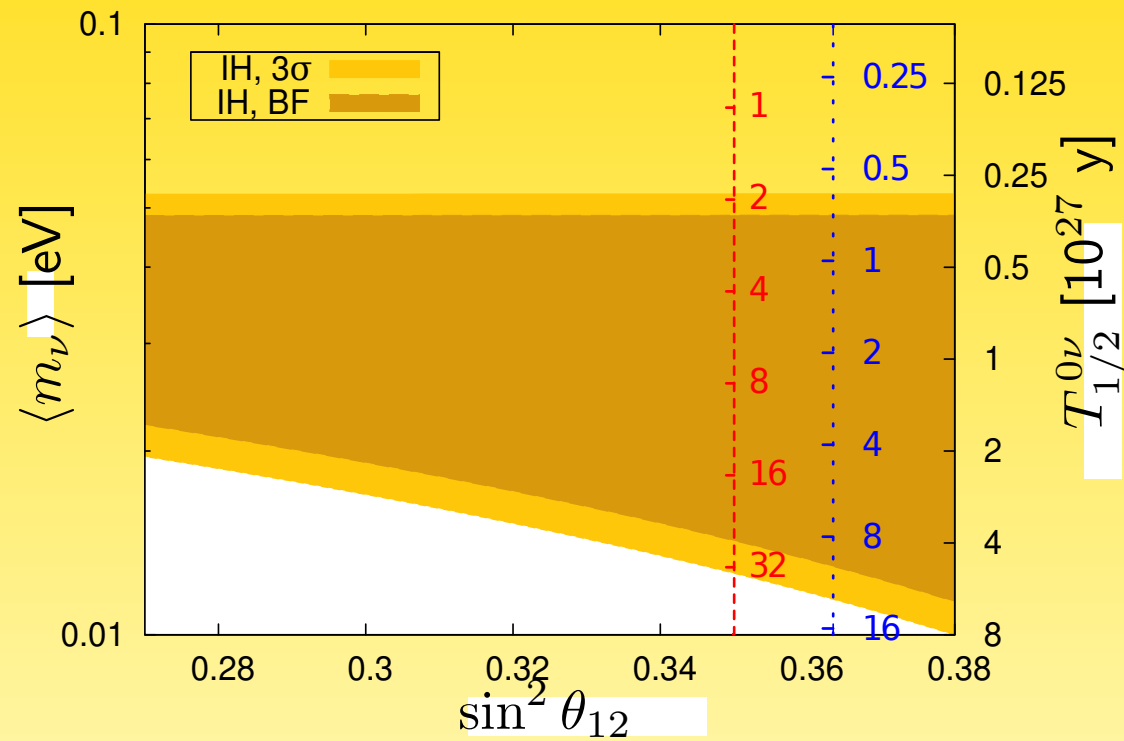
wichtige Skala:

$$|m_{ee}|_{\min}^{\text{IH}} \simeq c_{13}^2 \sqrt{\Delta m_A^2} \cos 2\theta_{12}$$

kennen wir $|m_{ee}|_{\min}^{\text{IH}}$ gut genug?

Invertierte Ordnung

$m_3 = 0.001 \text{ eV}$



Dueck, W.R., Zuber, PRD **83**

Invertierte Ordnung ausschließen

$$\begin{aligned} |m_{ee}|_{\min}^{\text{IH}} &= (1 - |U_{e3}|^2) \sqrt{|\Delta m_A^2|} (1 - 2 \sin^2 \theta_{12}) \\ &= (0.01 \dots 0.02) \text{ eV} \end{aligned}$$

$\Rightarrow$ 3σ -Bereich von $\sin^2 \theta_{12}$ gibt Faktor ~ 2 Unsicherheit für $|m_{ee}|_{\min}^{\text{IH}}$

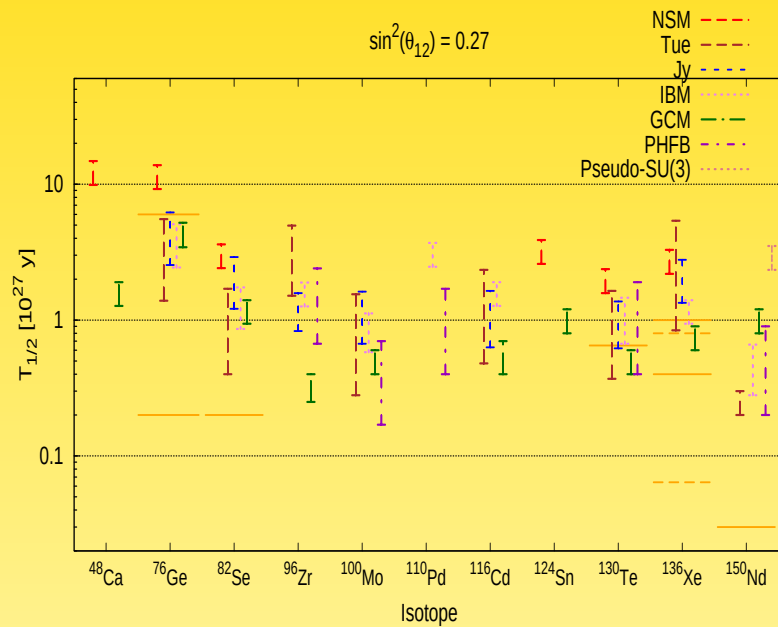
$$(T_{1/2}^{0\nu})^{-1} \propto |m_{ee}|^2 \propto \sqrt{\frac{M t}{B \Delta E}}$$

$\Rightarrow$ Faktor ~ 16 für $M \times t \times B \times \Delta E$

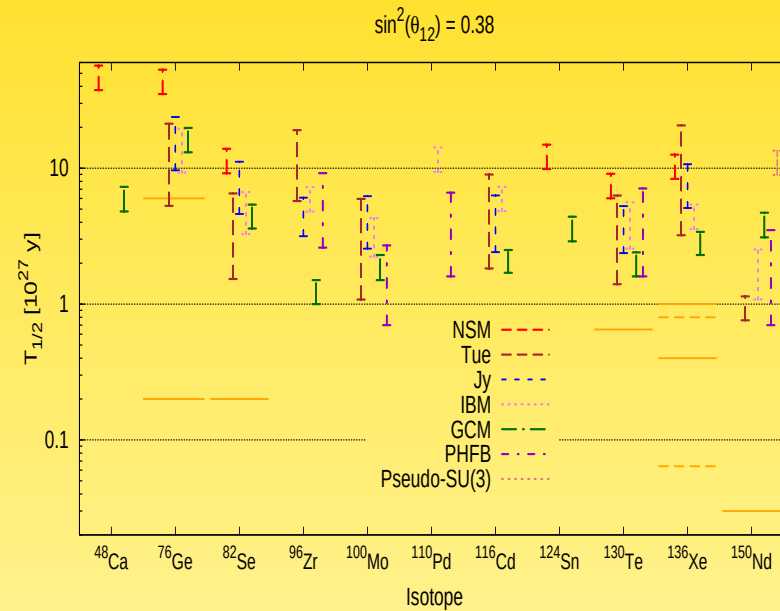
$\Rightarrow$ Präzisionsmessung von θ_{12} nötig (e.g. JUNO)

Dueck, W.R., Zuber, PRD **83**

Invertierte Ordnung ausschließen



$$\sin^2 \theta_{12} = 0.27$$



$$\sin^2 \theta_{12} = 0.38$$

Matrixelemente **und** $\theta_{12}!!$

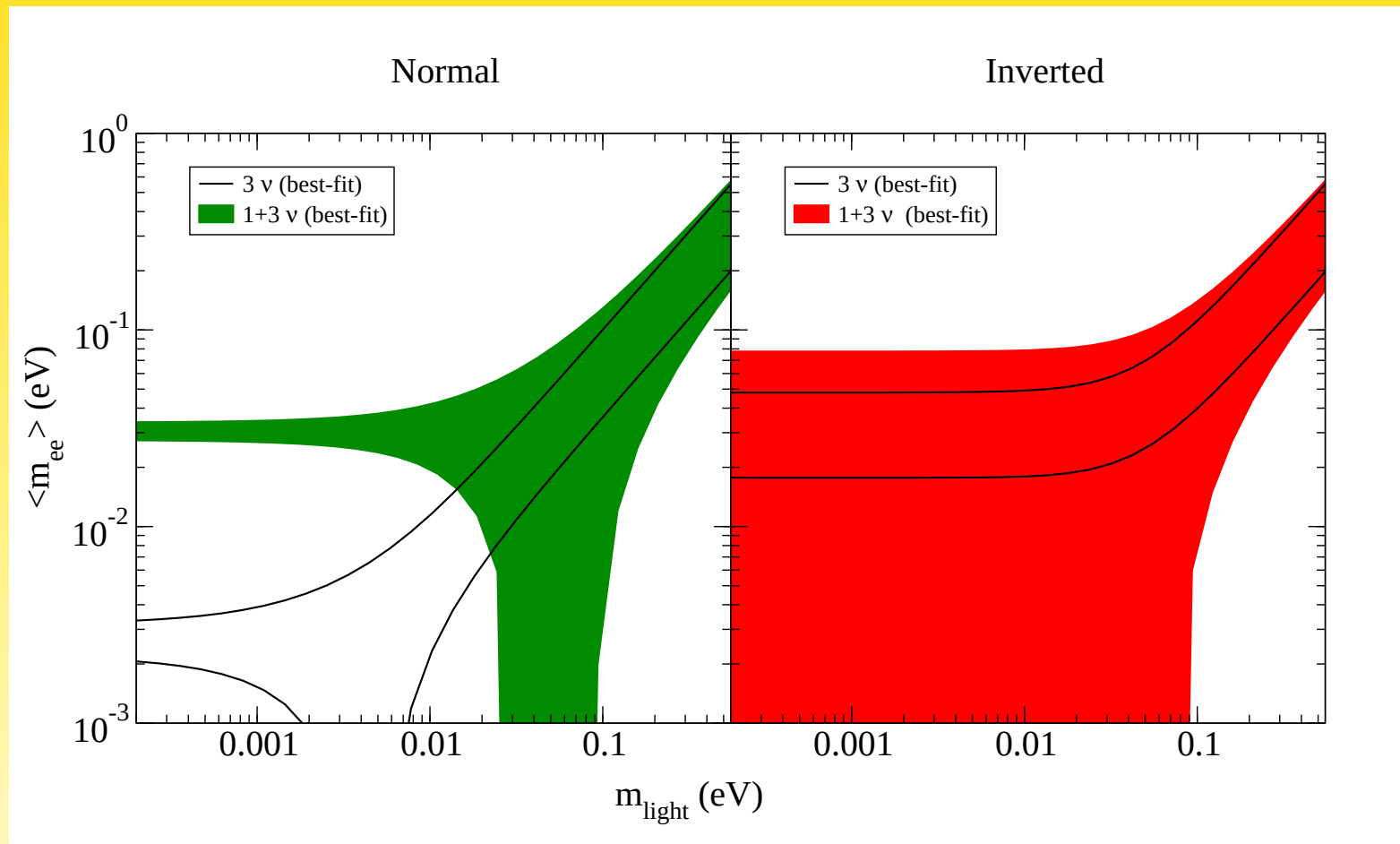
Vorhersagen von Grand Unified Theories

Modell	Fit	$ m_{ee} $ [meV]	m_0 [meV]	M_3 [GeV]	χ^2
$10_H + \overline{126}_H$	NH	0.49	2.40	3.6×10^{12}	23.0
$10_H + \overline{126}_H + SS$	NH	0.44	6.83	1.1×10^{12}	3.29
$10_H + \overline{126}_H + 120_H$	NH	2.87	1.54	9.9×10^{14}	11.2
$10_H + \overline{126}_H + 120_H + SS$	NH	0.78	3.17	4.2×10^{13}	6.9×10^{-6}
$10_H + \overline{126}_H + 120_H$	IH	35.52	30.2	1.1×10^{13}	13.3
$10_H + \overline{126}_H + 120_H + SS$	IH	24.22	12.0	1.2×10^{13}	0.6

Dueck, W.R., JHEP **1309**

Falls sterile Neutrinos...

$|m_{ee}| = 0$ für IH und $|m_{ee}| \neq 0$ für NH!



Barry, W.R., Zhang, JHEP 1107

Interpretation des Doppelbetazerfalls

- **Standardinterpretation:**

Neutrinoloser doppelter Betazerfall wird von leichten massiven Majorana-Neutrinos erzeugt

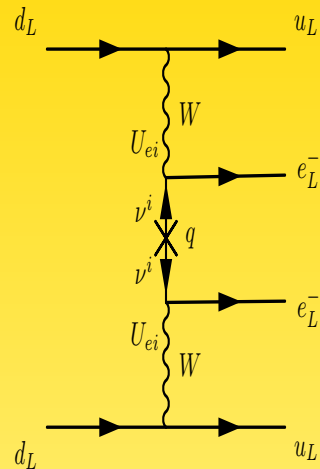
- **Nicht-Standardinterpretationen:**

Neutrinoloser doppelter Betazerfall wird nicht von leichten massiven Majorana-Neutrinos erzeugt

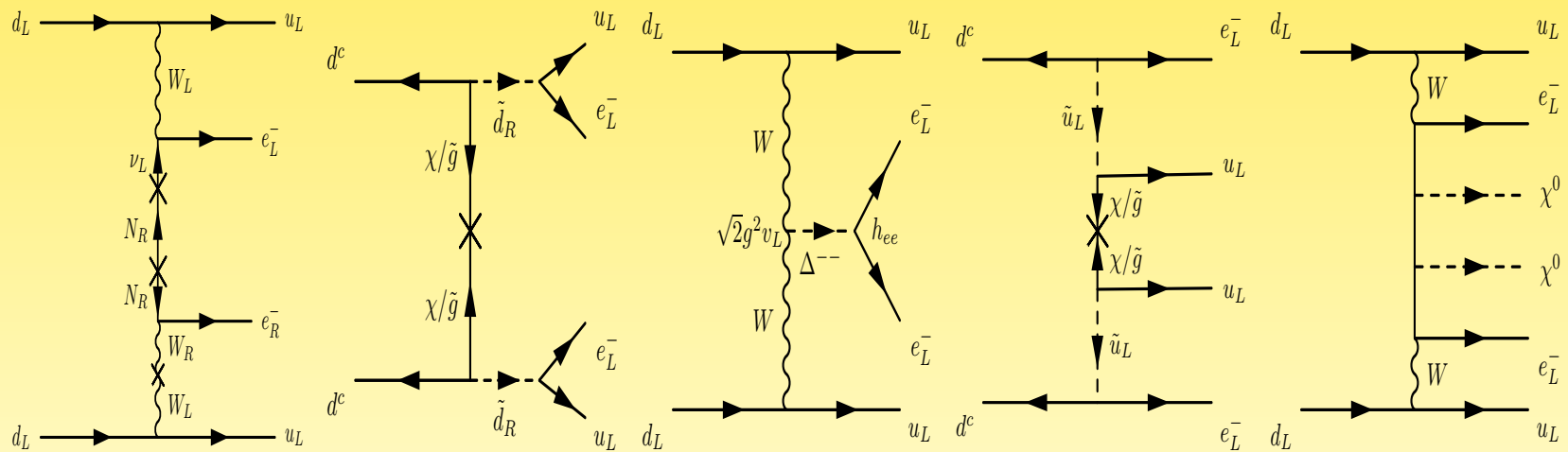
W.R., Int. J. Mod. Phys. **E20**, 1833-1930 (2011)

$0\nu\beta\beta$ -Experimente sind keine reinen Neutrinoexperimente!

- **Standardinterpretation:**



- **Nicht-Standardinterpretationen:**

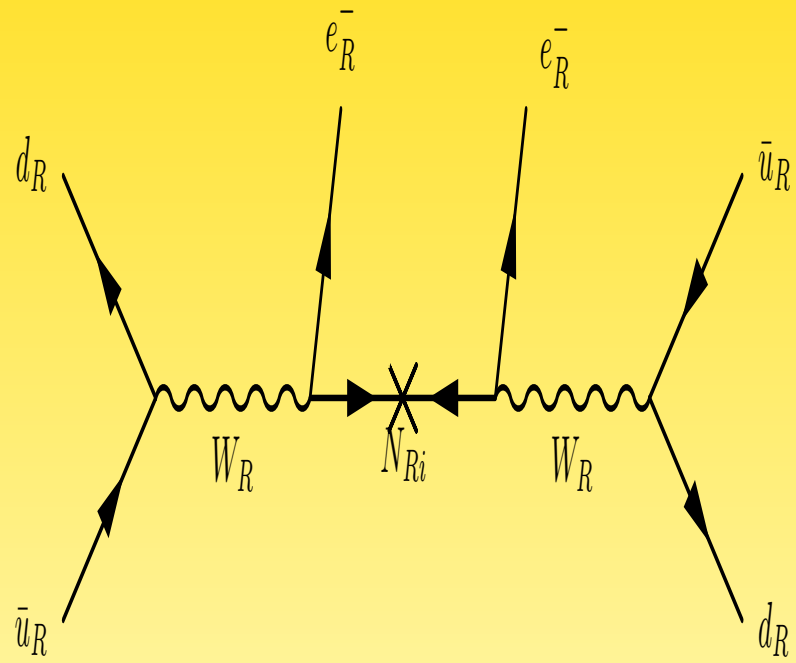
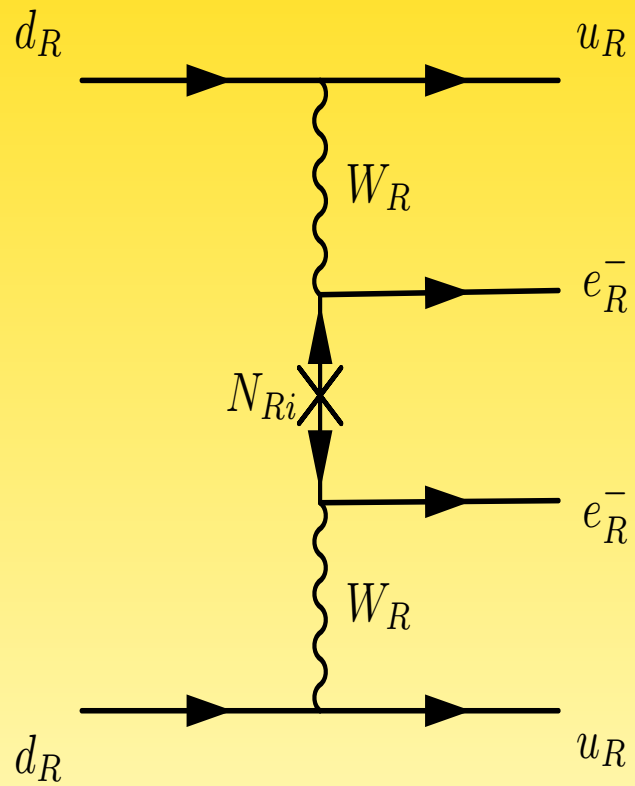


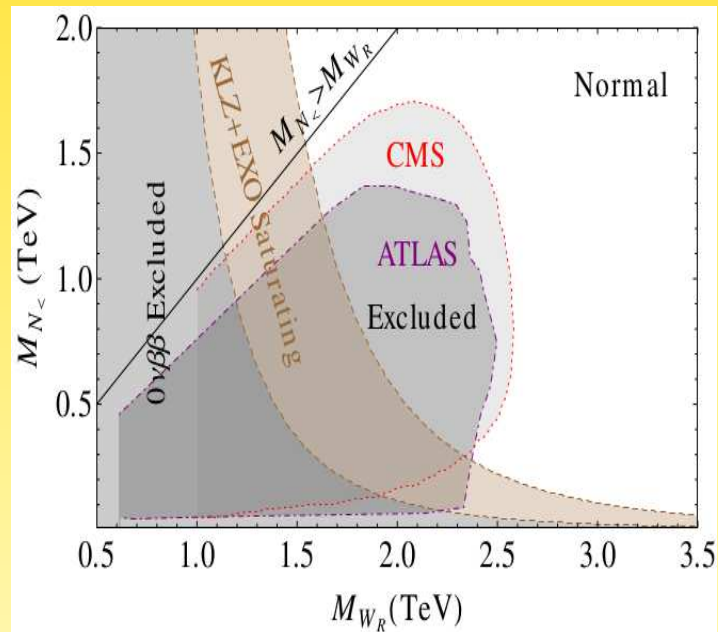
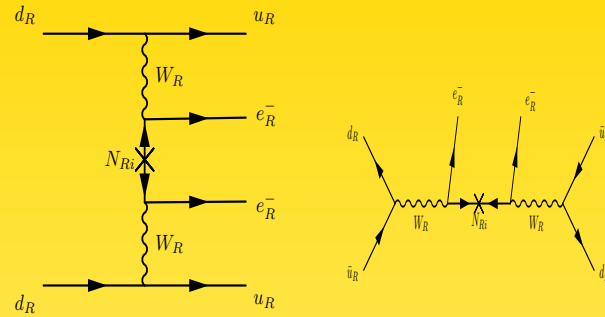
Nützlicher Zufall: $T_{1/2}^{0\nu}(1 \text{ eV}) = T_{1/2}^{0\nu}(1 \text{ TeV})$

- RPV Supersymmetrie
- links-rechts symmetrische Modelle
- schwere Neutrinos
- Color-octets
- Leptoquarks
- effektive Operatoren
- Extra Dimensionen
- ...

⇒ Lösung des inversen Problems...

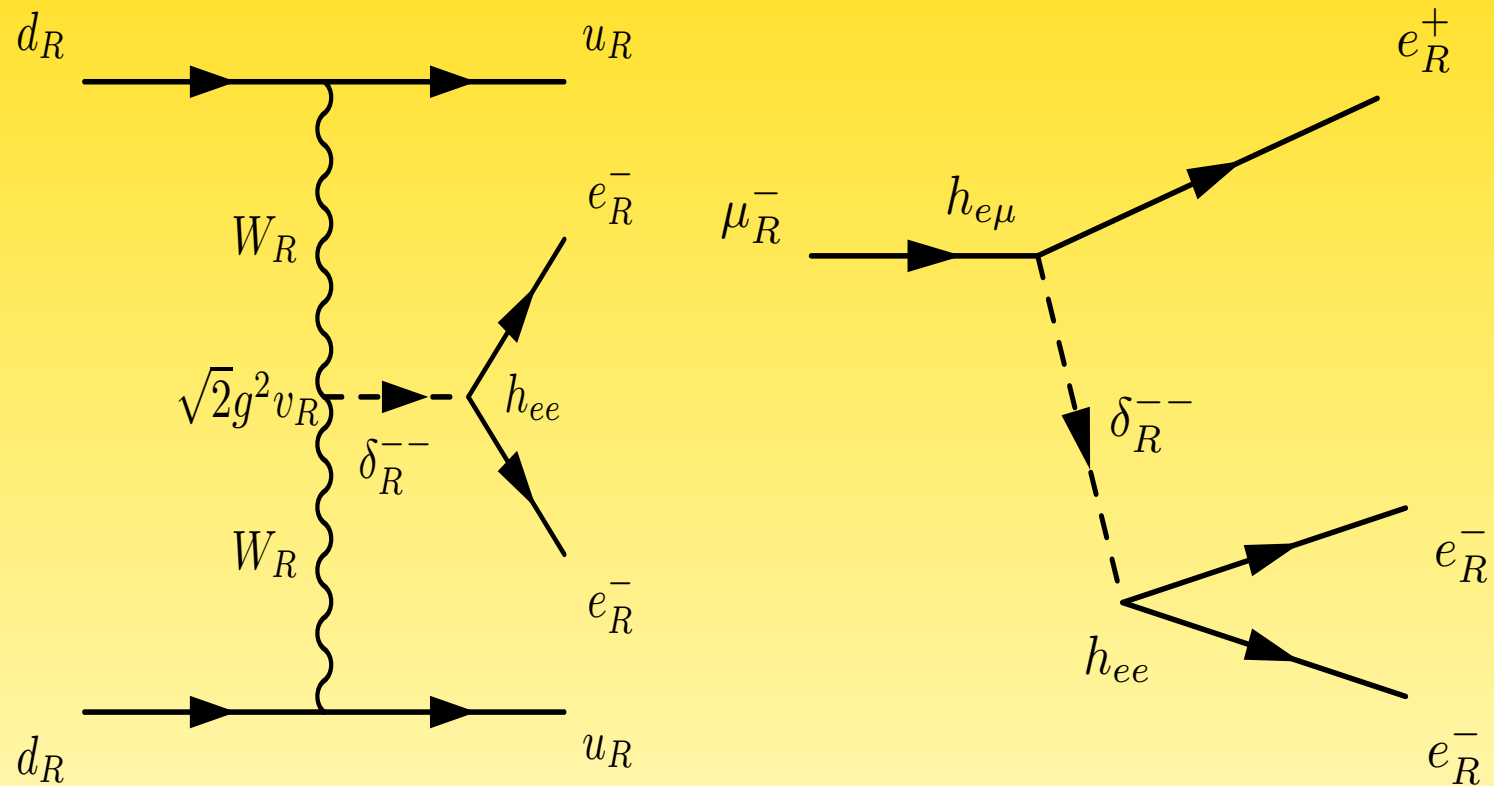
LHV vs. $0\nu\beta\beta$



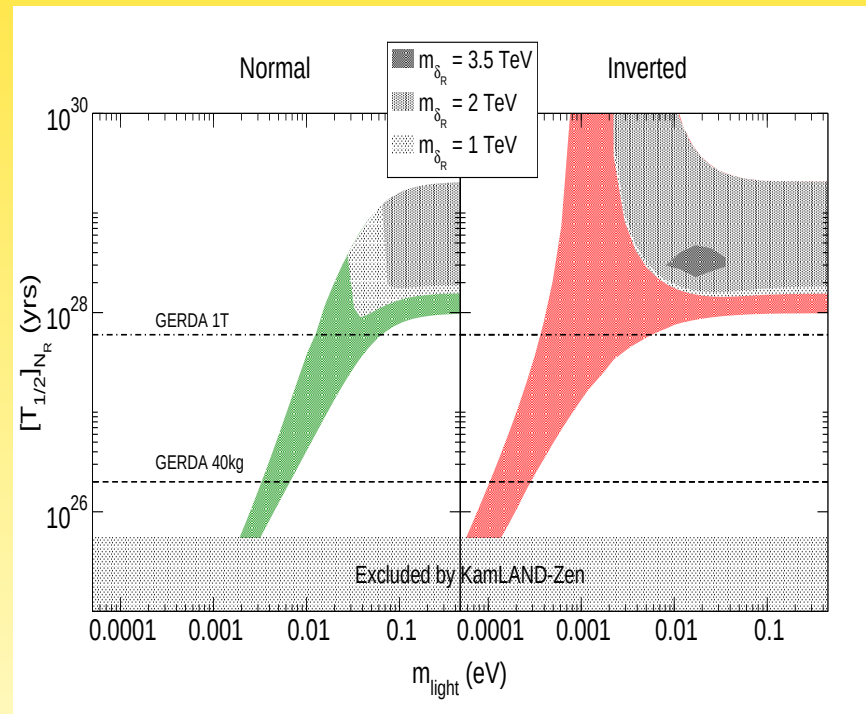
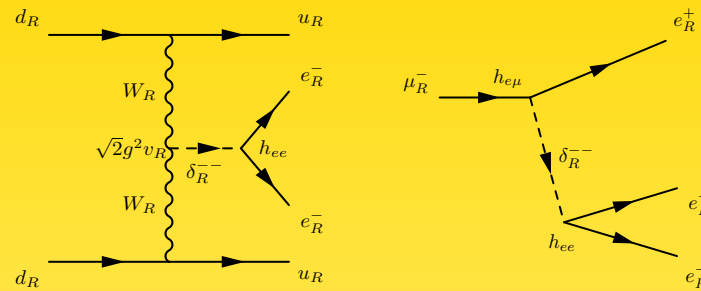


Bhupal Dev, Goswami, Mitra, W.R., Phys. Rev. **D88**

Leptonflavorverletzung

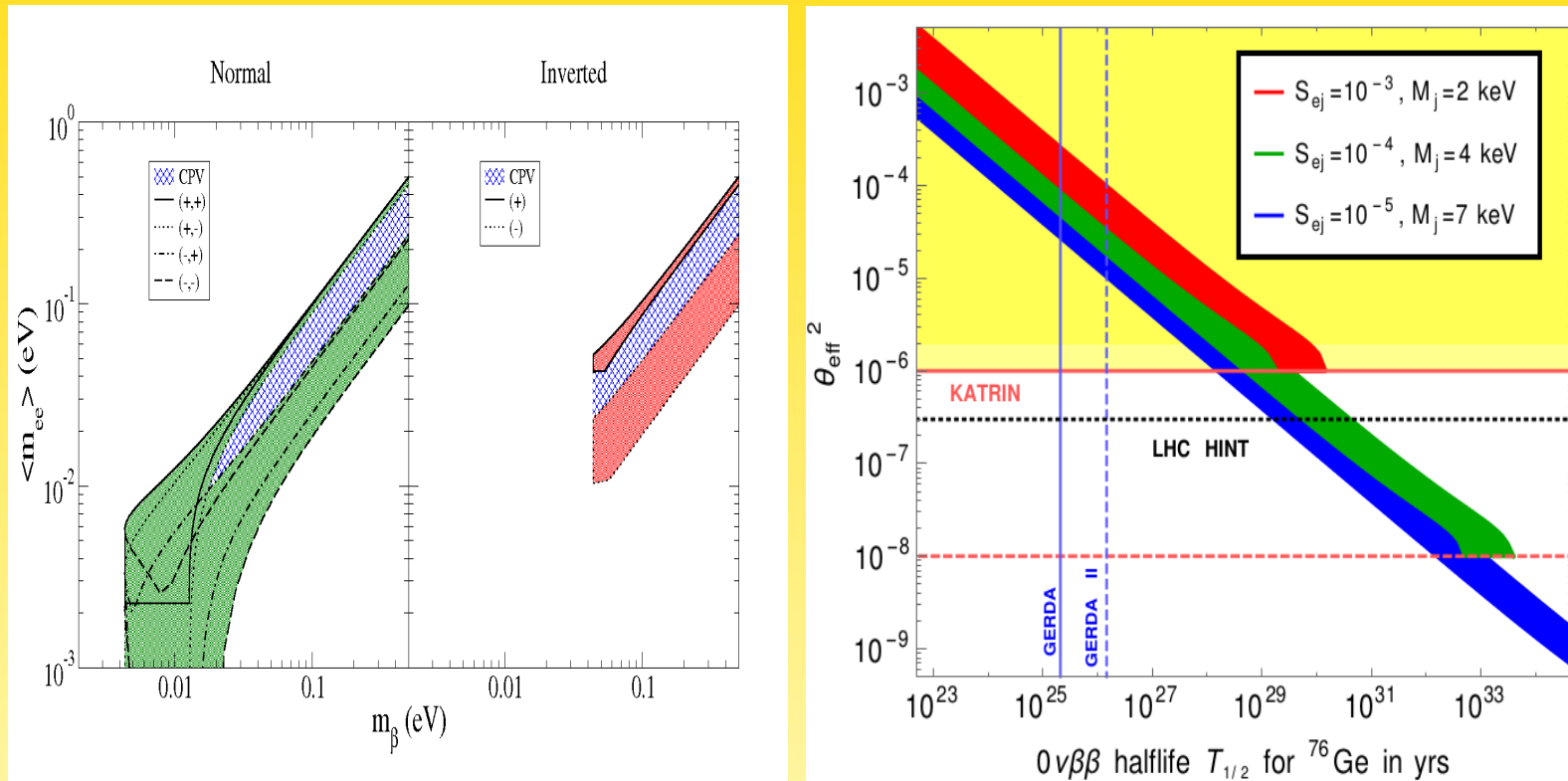


Leptonflavorverletzung



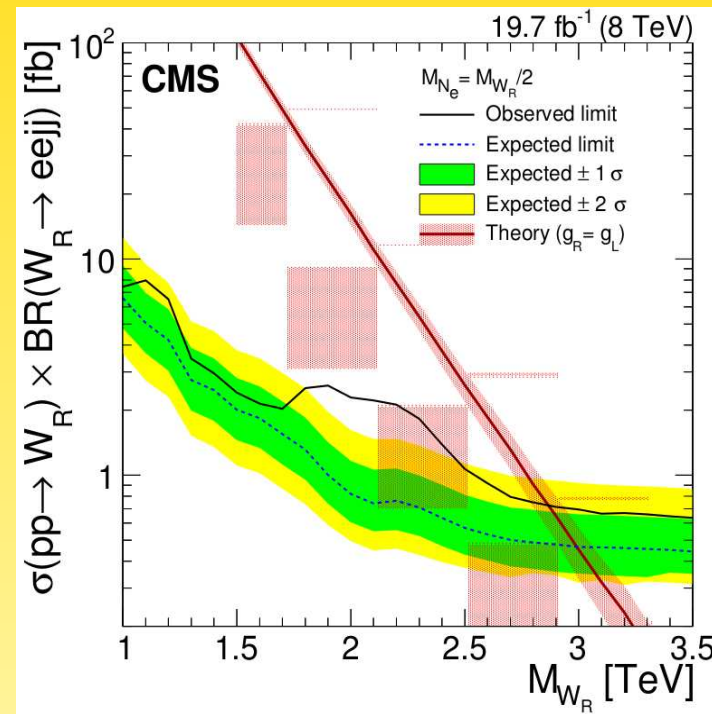
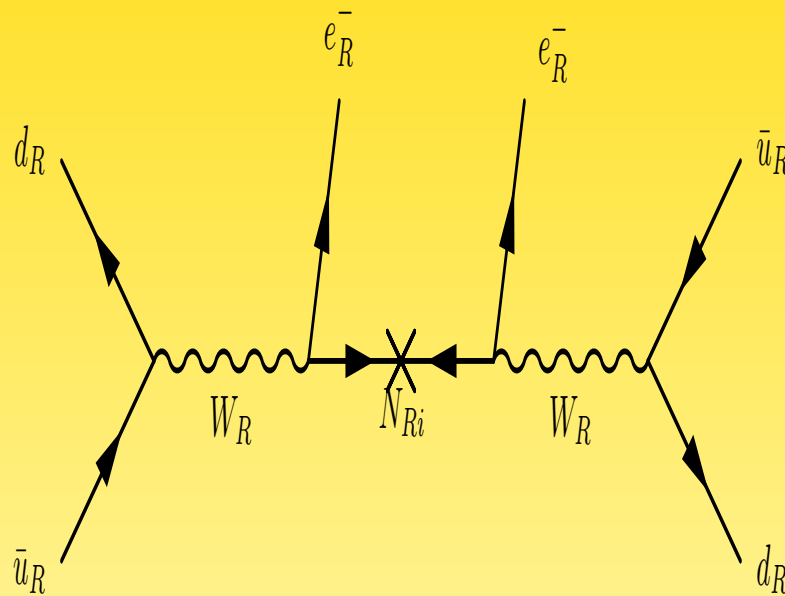
Barry, W.R., JHEP 1309

Auch hier Zusammenhang zu (modifizierter) KATRIN?



Barry, Heeck, W.R., 1404.5955

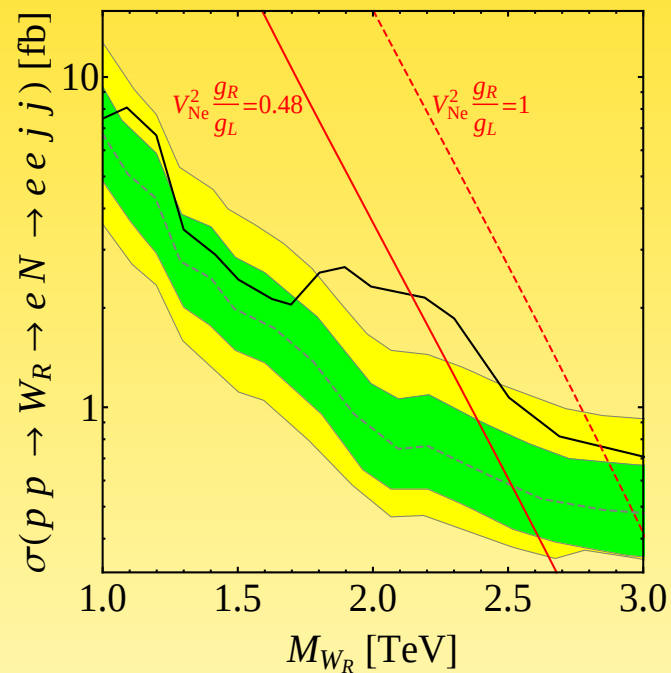
Neues aus dem Sommerloch: ein LHC-Signal in $eejj$? (1407.3683)



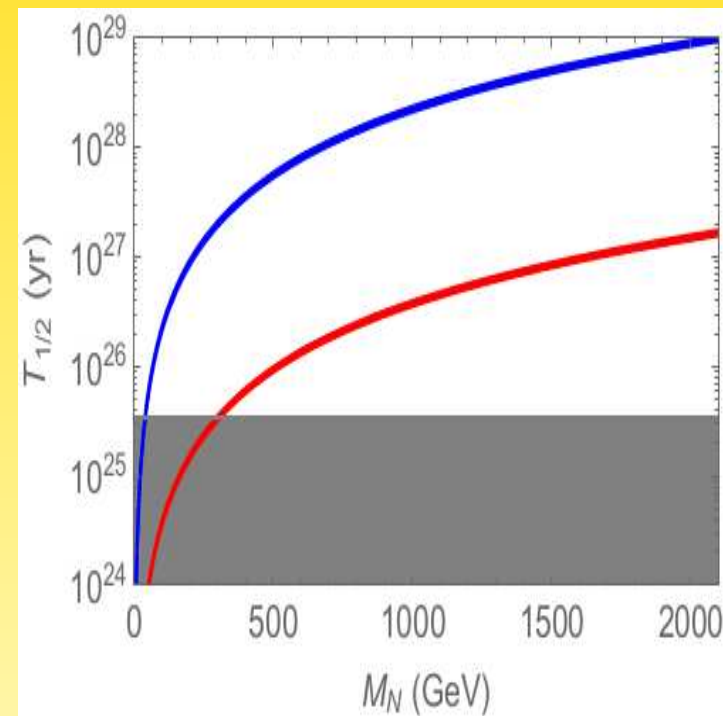
lokale 2.8σ bei $M_{W_R} = 2.1 \text{ TeV}$, nur im ee -Kanal

Interpretation I: LRSM

Möglich mit modifizierter links-rechts Symmetrie, $g_R = 0.6 g_L$ (Skalarfelder der $SU(2)_L$ haben verschiedene Masse als Skalarfelder der $SU(2)_R$)

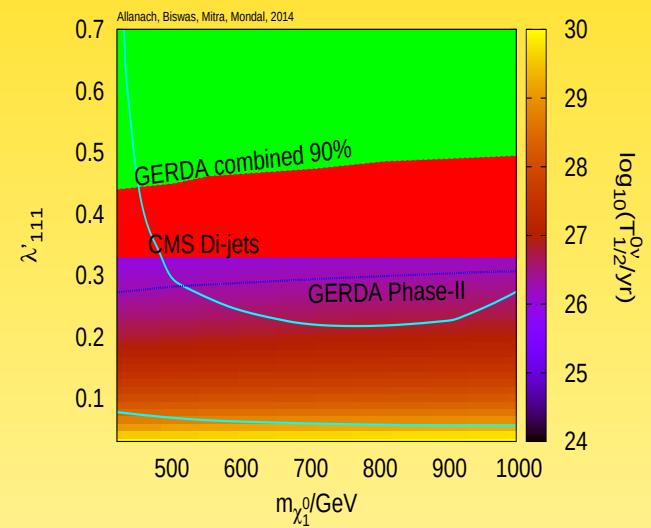
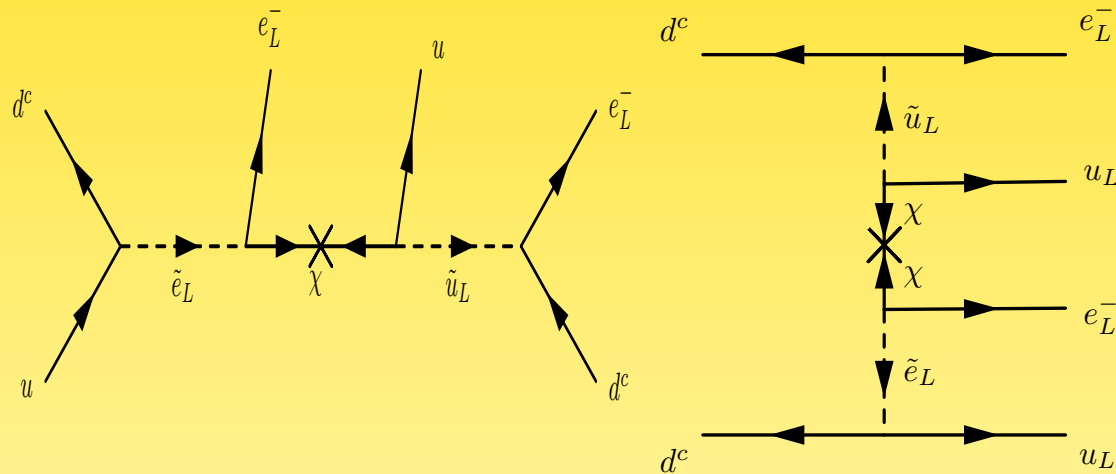


Deppisch *et al.*, 1407.5384



Bhupal Dev, W.R.

Interpretation II: RPV



Allanach *et al.*, 1408.5439

Zusammenfassung

Chi l'ha visto ?



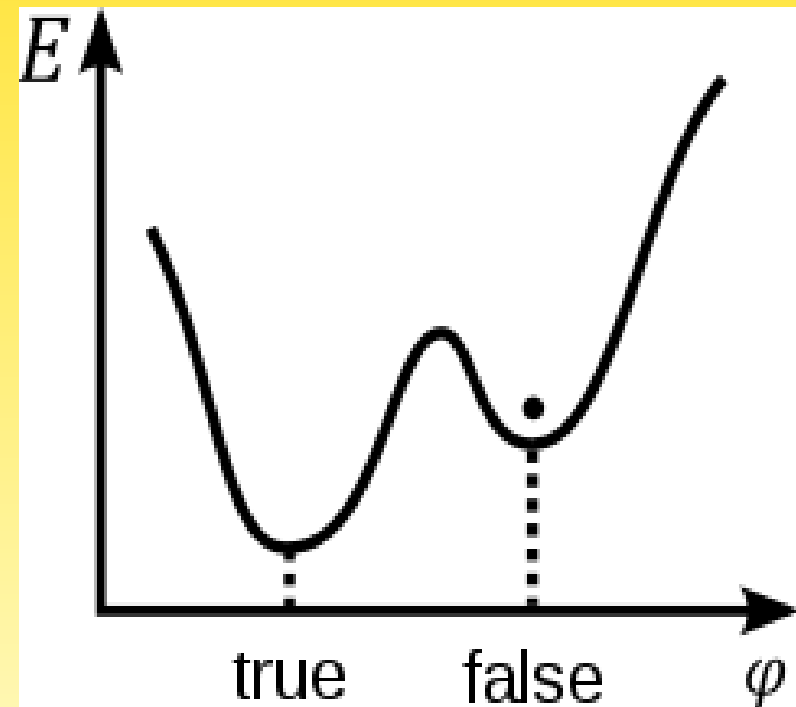
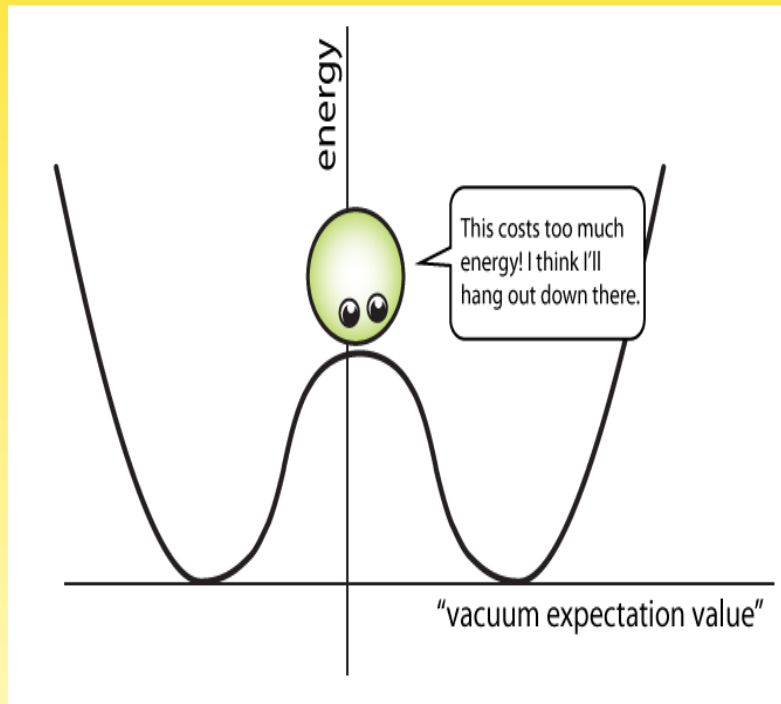
Ettore Majorana, ordinario di fisica teorica all'Università di Napoli, è misteriosamente scomparso dagli ultimi di marzo. Di anni 31, alto metri 1,70, snello, con capelli neri, occhi scuri, una lunga cicatrice sul dorso di una mano. Chi ne sapesse qualcosa è pregato di scrivere al R. P. E. Marianecchi, Viale Regina Margherita 66 - Roma.

Phänomenologie schwerer Singlets: Higgs

Higgs-Potenzial ist flach $\Rightarrow$ Korrekturen gefährlich

$$\dot{\lambda} \propto -24 \operatorname{Tr} (Y_u^\dagger Y_u)^2 \propto m_t^4 \Rightarrow \text{Vakuumstabilität}$$

(Holthausen, Lim, Lindner; Bezrukov *et al.*; Degrandi *et al.*; Masina)



(Vakuum vermutlich metastabil)

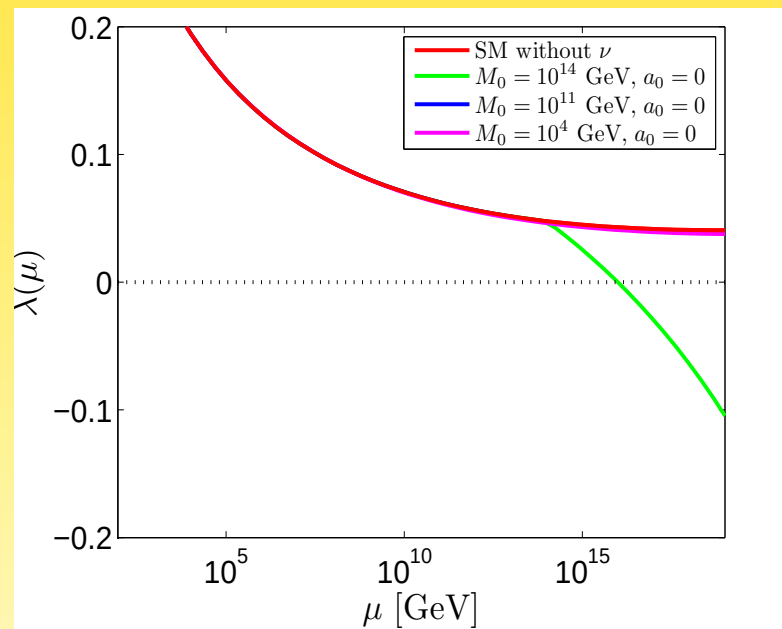
Phänomenologie schwerer Singlets: Higgs

Dirac-Beitrag $\bar{L} \Phi N_R$ zu λ :

$$\dot{\lambda} \propto -8 m_D^4$$

Casas *et al.*; Strumia *et al.*; W.R., Zhang

Zerstört Vakuumstabilität



(auch für TeV-Neutrinos die an Beschleunigern produzierbar sind)

Sterile Neutrinos

immer noch viele Hinweise auf (sub-)eV sterile Neutrinos (LSND, MiniBooNE, Gallium, Reaktor, etc.)

Falls diese sterile Neutrinos existieren...

- ...ändern sich die Massenobservablen und ihre Korrelationen
- ...müssen wir Flavorsymmetriemodelle überdenken
- ...müssen wir Seesaw-Modelle überdenken
- ...könnte Standard-Kosmologie falsch sein
- ...etc.

teilweise ebenso für keV sterile Neutrinos ($\leftrightarrow$ WDM)

Warum sollten wir Leptonzahlerhaltung testen?

- L und B zufällig im SM erhalten
- effektive Theorie: $\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_{\text{LZV}} + \frac{1}{\Lambda^2} \mathcal{L}_{\text{LFV, BZV}} + \dots$
- Baryogenese: B verletzt
- B, L oft in GUTs zusammenhängend
- GUTs haben Seesaw und Majorana-Neutrinos
- (chirale Anomalien: $\partial_\mu J_{B,L}^\mu = c G_{\mu\nu} \tilde{G}^{\mu\nu} \neq 0$ mit $J_\mu^B = \sum \bar{q}_i \gamma_\mu q_i$ und $J_\mu^L = \sum \bar{\ell}_i \gamma_\mu \ell_i$)

$\Rightarrow$ Leptonzahlverletzung genau so wichtig wie Baryonzahlverletzung
($0\nu\beta\beta$ ist viel mehr als ein Neutrinoexperiment)

Vorhersagen von Grand Unified Theories

Yukawa-Struktur von $SO(10)$ Modellen hängt von Higgs-Darstellungen ab

$$10_H (\leftrightarrow H), \overline{126}_H (\leftrightarrow F), 120_H (\leftrightarrow G)$$

Gibt Relationen zwischen Massenmatrizen:

$$m_{\text{up}} \propto r(H + sF + it_u G)$$

$$m_{\text{down}} \propto H + F + iG$$

$$m_D \propto r(H - 3sF + it_D G)$$

$$m_\ell \propto H - 3F + it_l G$$

$$M_R \propto r_R^{-1} F$$

$\leftrightarrow$ Fit mit RG, Higgs, θ_{13}

Dueck, W.R., JHEP **1309**

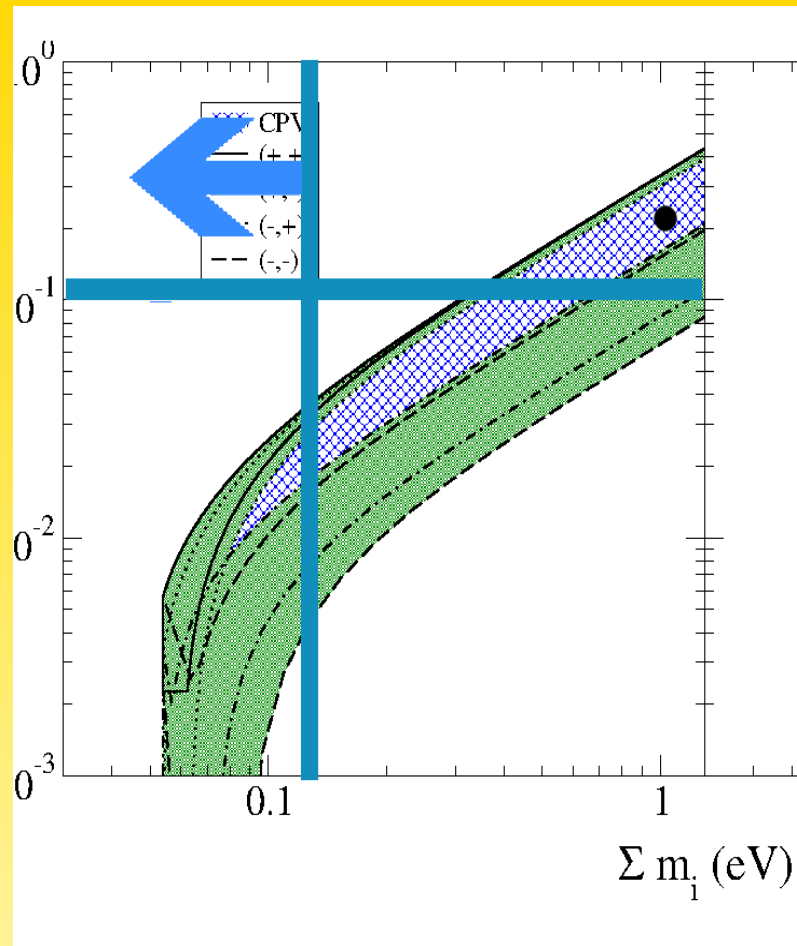
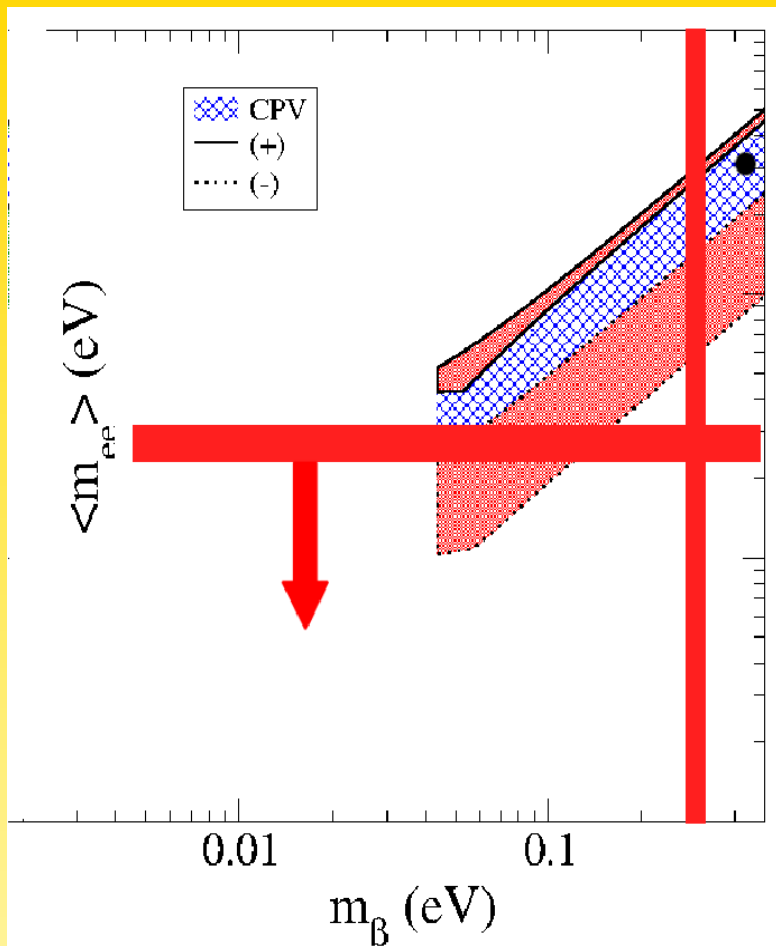
PMNS-Matrix:

$$|U| = \begin{pmatrix} 0.801 \dots 0.845 & 0.514 \dots 0.580 & 0.137 \dots 0.158 \\ 0.225 \dots 0.517 & 0.441 \dots 0.699 & 0.614 \dots 0.793 \\ 0.246 \dots 0.529 & 0.464 \dots 0.713 & 0.590 \dots 0.776 \end{pmatrix}$$

CKM-Matrix:

$$|V| = \begin{pmatrix} 0.97427 \pm 0.00015 & 0.22534 \pm 0.00065 & 0.00351^{+0.00015}_{-0.00014} \\ 0.22520 \pm 0.00065 & 0.97344^{+0.00016}_{-0.00016} & 0.0412^{+0.0011}_{-0.0005} \\ 0.00867^{+0.00029}_{-0.00031} & 0.0404^{+0.0011}_{-0.0005} & 0.999146^{+0.000021}_{-0.000046} \end{pmatrix}$$

Warum so verschieden?? $\leftrightarrow$ Flavoursymmetrien!



CP-Verletzung!

Dirac-Neutrinos!

Etwas anderes erzeugt $0\nu\beta\beta$!

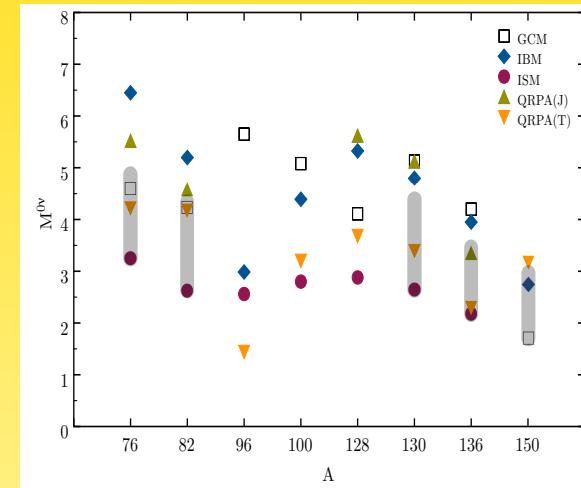
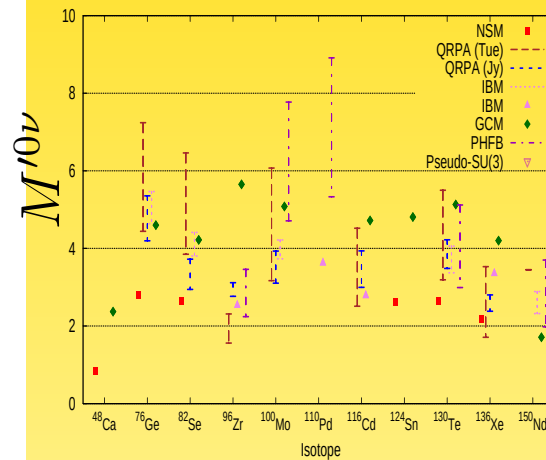
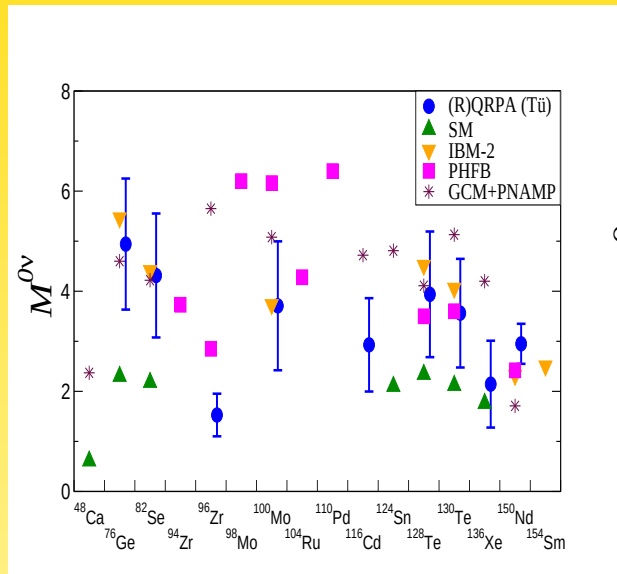
Komplementarität!

- sind in Ära der Komplementarität:
 - Kosmologie-Limits schließen aus dass leichte Neutrinos $0\nu\beta\beta$ -Limits erfüllen
 - $0\nu\beta\beta$ -Limits schließen aus dass leichte Neutrinos Mainz/Troitsk-Limit erfüllen
- interessante Möglichkeiten falls Inkonsistenzen auftauchen...

Komplementarität und Nicht-Standardphysik

- Kosmologie $\neq \Lambda$ CDM
 - Pan-STARRS1 Survey (1310.3828): $\omega_{\text{DE}} = -1.142^{+0.076}_{-0.065}$:
 Σ -Limit schwächer
 - BICEP-2 vs. Planck: $N_{\text{eff}} = 4$ bevorzugt, m_s wird kleiner (1403.7028, 1403.8049, 1404.1794)
 - σ_8 von Planck vs. lokales σ_8 : $\Sigma = 0.5 \pm 0.2$ eV (1403.4852)
- keine Neutrinos in $0\nu\beta\beta$
- KATRIN: RH Ströme, extra Dimensionen, tachyonische Neutrinos,...

Von Lebensdauer zu Teilchenphysik: Kernmatrixelemente



Xe-Limit ist besser als Ge-Limit wenn:

$$T_{\text{Xe}} > T_{\text{Ge}} \frac{G_{\text{Ge}}}{G_{\text{Xe}}} \left| \frac{\mathcal{M}_{\text{Ge}}}{\mathcal{M}_{\text{Xe}}} \right|^2 \text{ yrs}$$

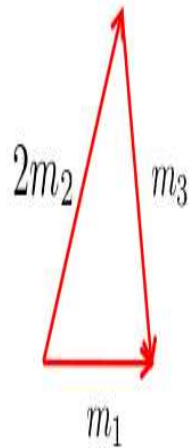
The Zoo (of A_4 models)

Type	L_i	ℓ_i^c	ν_i^c	Δ	References
A1				-	[1–14] [15] [#]
A2	$\underline{3}$	$\underline{1}, \underline{1}', \underline{1}''$	-	$\underline{1}, \underline{1}', \underline{1}'', \underline{3}$	[16–18]
A3				$\underline{1}, \underline{3}$	[19]
B1	$\underline{3}$	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{3}$	-	[4, 20–27] [#] [28–30]* [31–45]
B2				$\underline{1}, \underline{3}$	[46] [#]
C1				-	[2, 47, 48]
C2	$\underline{3}$	$\underline{3}$	-	$\underline{1}$	[49, 50] [51] [#]
C3				$\underline{1}, \underline{3}$	[52]
C4				$\underline{1}, \underline{1}', \underline{1}'', \underline{3}$	[53]
D1				-	[54, 55] [#] [56, 57]* [58]
D2	$\underline{3}$	$\underline{3}$	$\underline{3}$	$\underline{1}$	[59] [60]*
D3				$\underline{1}'$	[61]*
D4				$\underline{1}', \underline{3}$	[62]*
E	$\underline{3}$	$\underline{3}$	$\underline{1}, \underline{1}', \underline{1}''$	-	[63, 64]
F	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{3}$	$\underline{3}$	$\underline{1}$ or $\underline{1}'$	[65]
G	$\underline{3}$	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{1}, \underline{1}', \underline{1}''$	-	[66]
H	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	-	-	[67]
I	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	$\underline{1}, \underline{1}, \underline{1}$	-	[68]*
J	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	$\underline{3}$	-	[12, 39, 69, 70]
K	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	$\underline{1}, \underline{1}$	$\underline{1}$	[71]*
L	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	$\underline{1}$	-	[72]*
M	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{1}, \underline{1}'', \underline{1}'$	$\underline{3}, \underline{1}$	-	[73, 74]
N	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{1}, \underline{1}'', \underline{1}'$	$\underline{3}, \underline{1}', \underline{1}''$	-	[75]

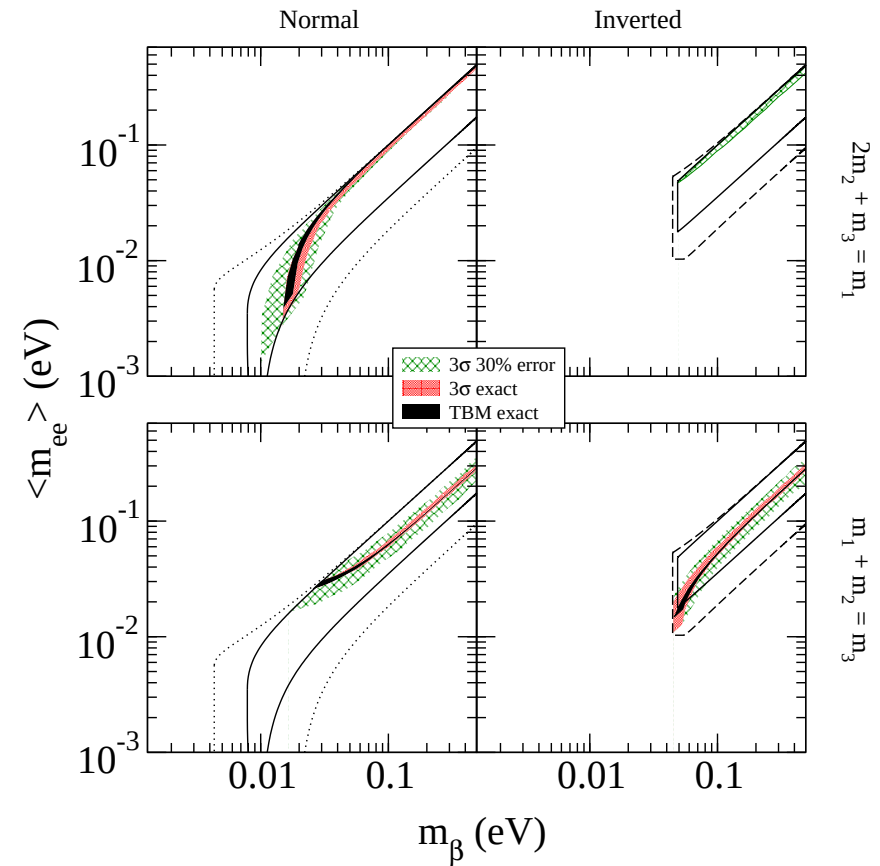
Barry, W.R., PRD **81**, 093002 (2010)

all give the same mixing scheme...

Neutrino mass sum-rules



Sum-rule	Flavour symmetry
$2m_2 + m_3 = m_1$	$A_4, T', (S_4)$
$m_1 + m_2 = m_3$	$S_4, (A_4)$
$\frac{2}{m_2} + \frac{1}{m_3} = \frac{1}{m_1}$	A_4, T'
$\frac{1}{m_1} + \frac{1}{m_2} = \frac{1}{m_3}$	S_4



constrains masses and Majorana phases

Barry, W.R., NPB **842**

Seesaw-Formalismus

$$\mathcal{L} = \frac{1}{2}(\bar{\nu}_L, \bar{N}_R^c) \begin{pmatrix} 0 & m_D \\ m_D^T & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ N_R \end{pmatrix}$$

6×6 Matrix, diagonalisiert mit

$$\mathcal{U}_\nu \simeq \begin{pmatrix} 1 - \frac{1}{2}BB^\dagger & B \\ -B^\dagger & 1 - \frac{1}{2}B^\dagger B \end{pmatrix} \begin{pmatrix} U & 0 \\ 0 & V_R \end{pmatrix} \quad \text{mit } B = m_D M_R^{-1}$$

leichte Neutrinos (mischen nicht-unitär):

$$m_\nu = -m_D M_R^{-1} m_D^T = U \text{diag}(m_1, m_2, m_3) U^T$$

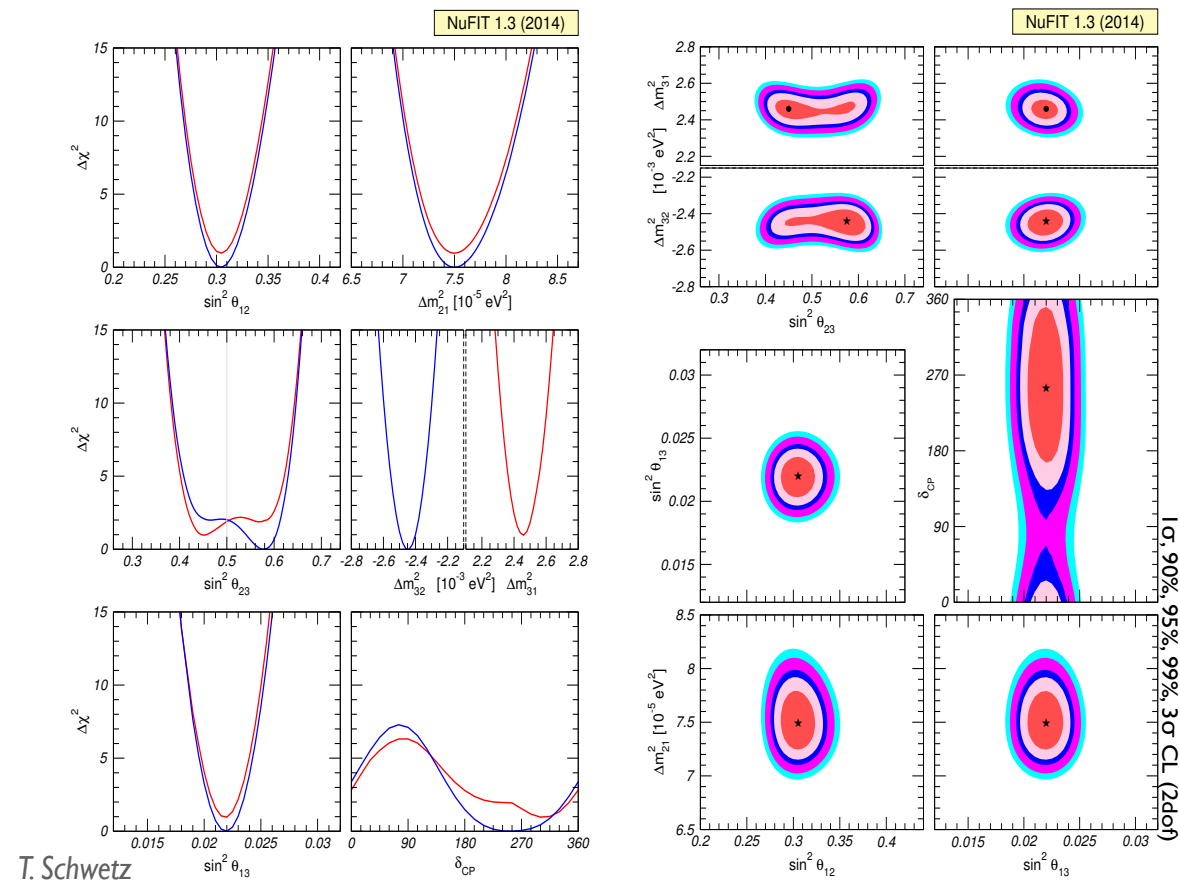
schwere Neutrinos:

$$M_R = V_R \text{diag}(M_1, M_2, M_3) V_R^T$$

(N_R sind *sterile Neutrinos*)

Status globaler Fits (Maltoni *et al.*, post Nu2014)

3-flavour global fit to oscillation data

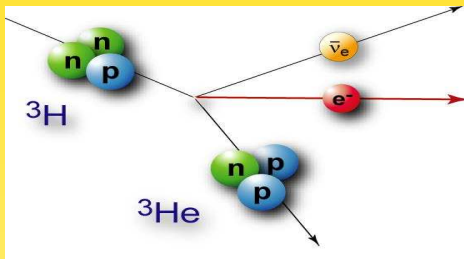


Ergebnisse der Fits

		<i>Fogli et al.</i>	<i>Maltoni et al.</i>	<i>Valle et al.</i>
$\sin^2 \theta_{23}$	nor	0.437	0.452	0.567
	inv	0.455	0.579	0.573
$\delta/^\circ$	nor	250	306	266
	inv	236	254	266

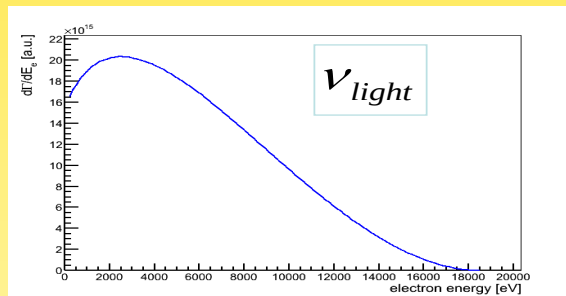
- Oktant von θ_{23} : T2K misst signifikante $P_{\mu e} \propto \sin^2 \theta_{13} \sin^2 \theta_{23} + \text{Materie}, \delta$, aber stark von Details atm. nus abhängig (sub-GeV e -Exzess)
- $\delta \simeq \frac{3}{2}\pi$: Korrelation von θ_{13} aus LBL und KL; 1σ -Bereich ist $\sin \delta \in [-1, 0]$
- CP Verletzung ist bei $\lesssim 1.5\sigma$, korreliert mit θ_{23}
- Keine signifikante Präferenz bei Normal vs. Invertiert

Imprint of keV neutrinos on β -spectrum

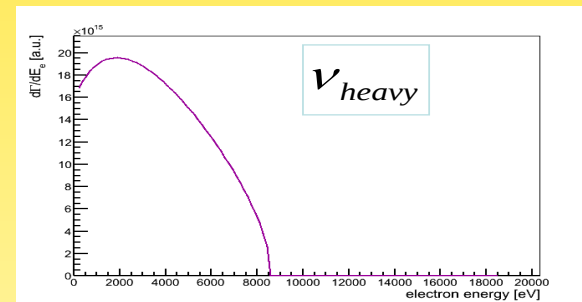


$$\begin{pmatrix} \nu_e \\ \nu_s \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_{light} \\ \nu_{heavy} \end{pmatrix}$$

$\cos^2(\theta)$



+ $\sin^2(\theta)$



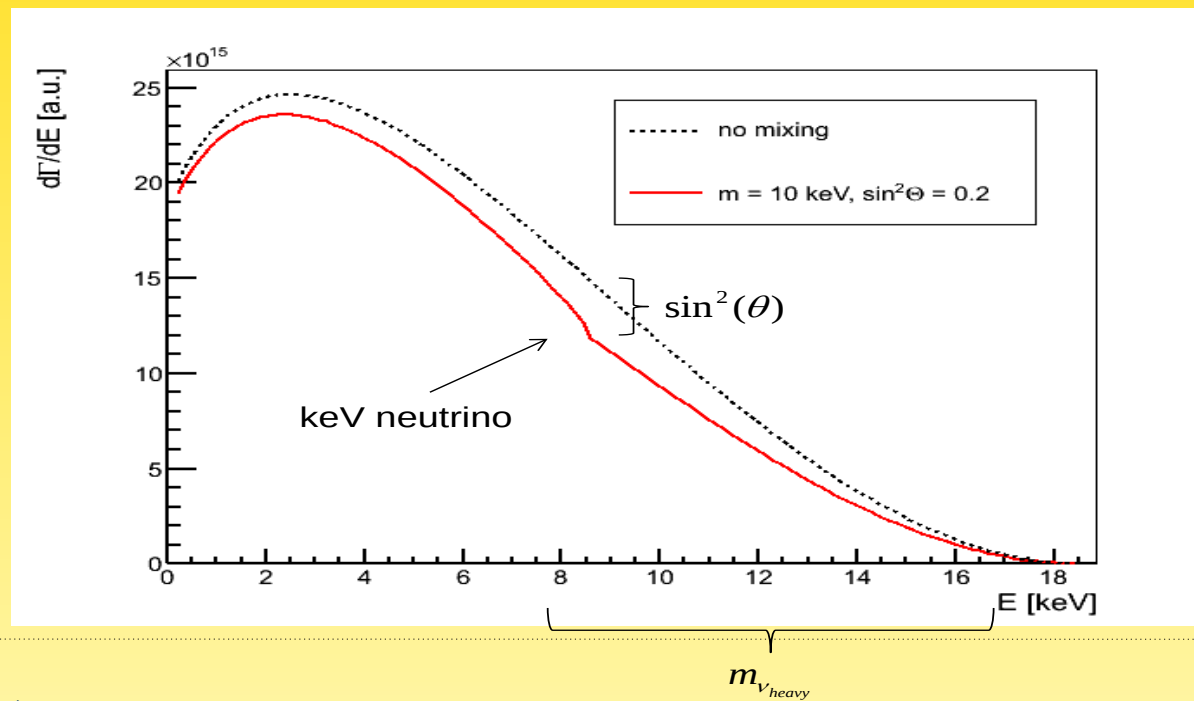
Susanne Mertens

12



Mertens *et al.*, 1409.0920

Imprint of keV neutrinos on β -spectrum



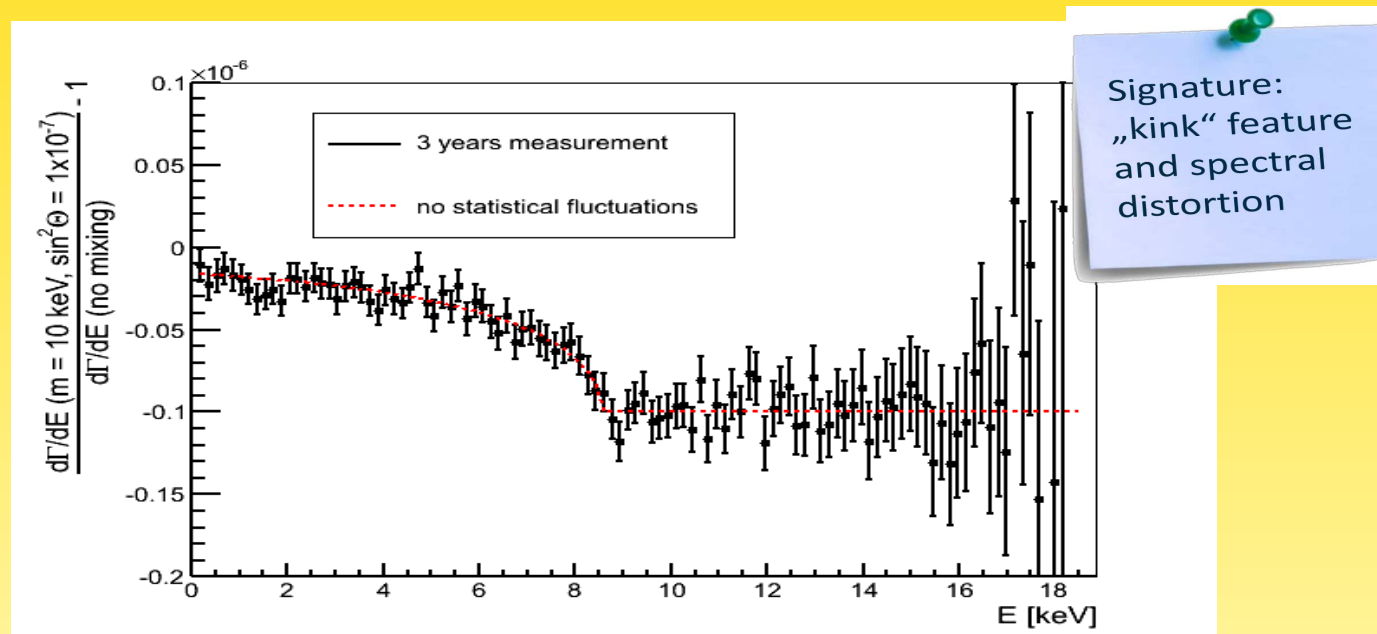
Susanne Mertens

13



Mertens *et al.*, 1409.0920

Imprint of keV neutrinos on β -spectrum

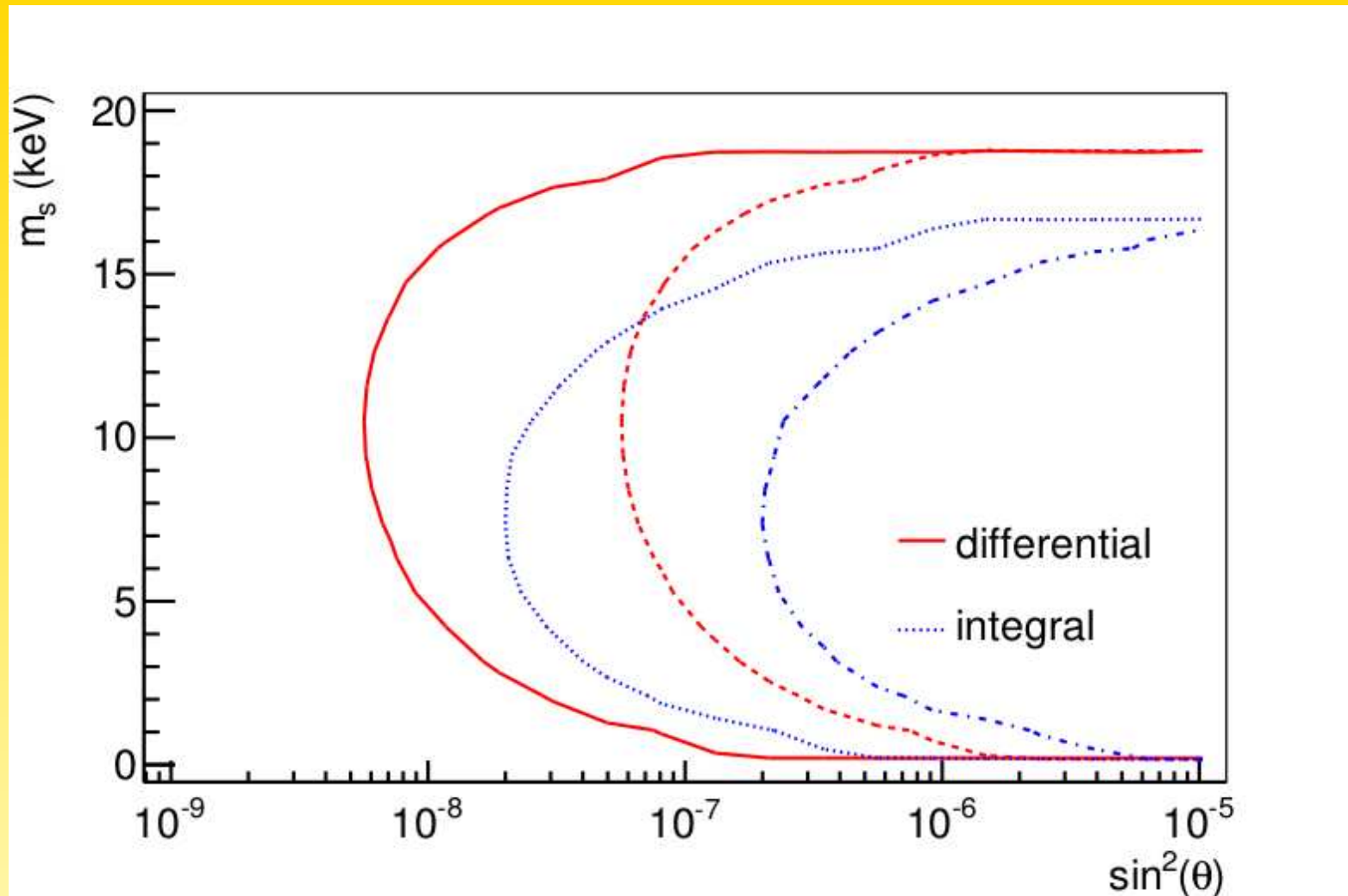


Susanne Mertens

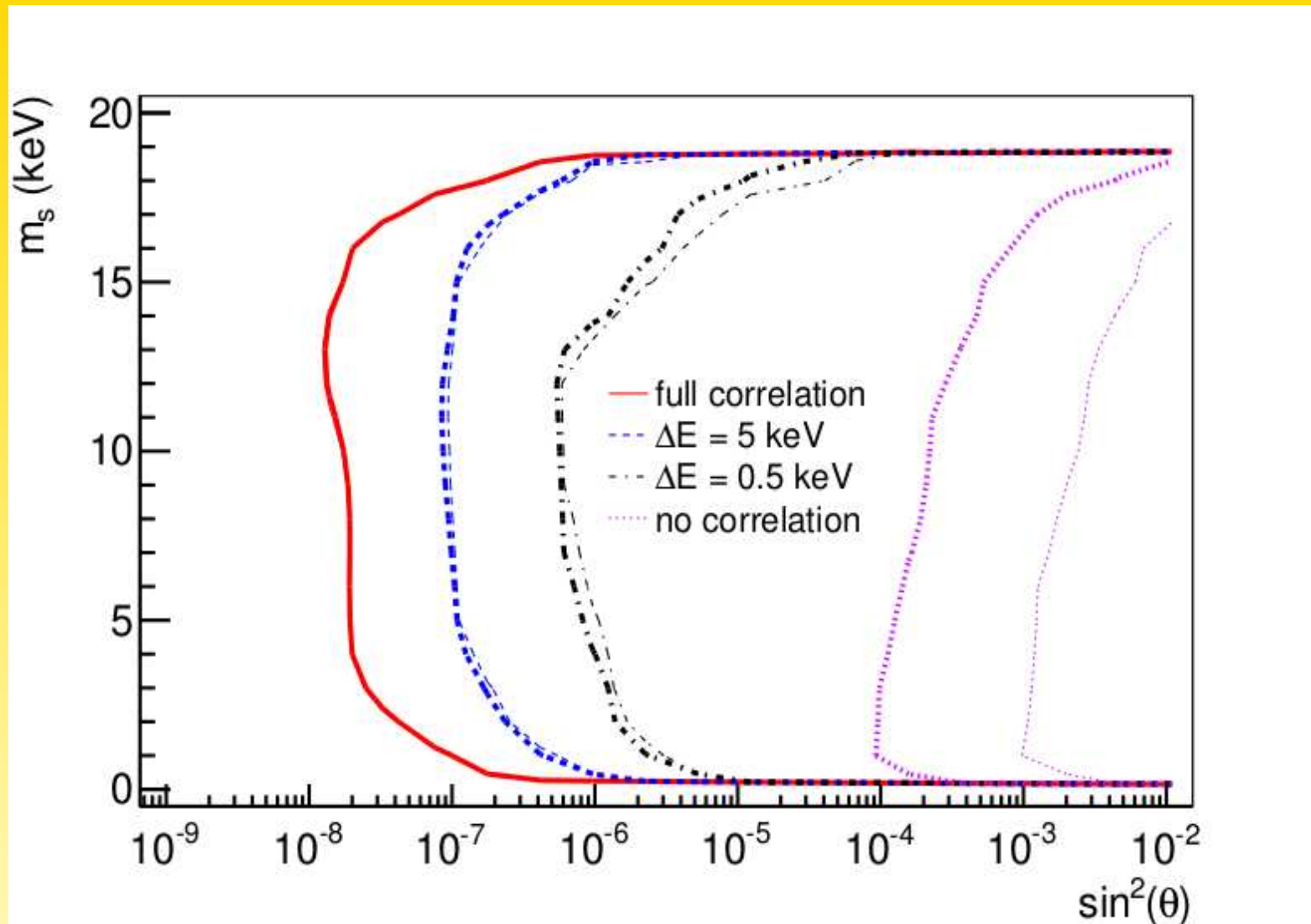
15



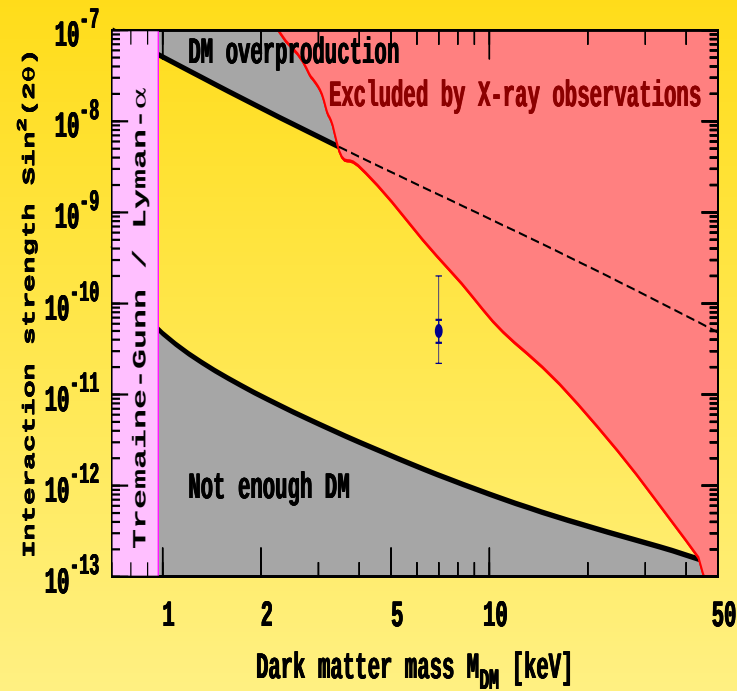
Mertens *et al.*, 1409.0920



(i) energy resolving detector (differential) or (ii) counting detector (integral) or
(iii) time-of-flight

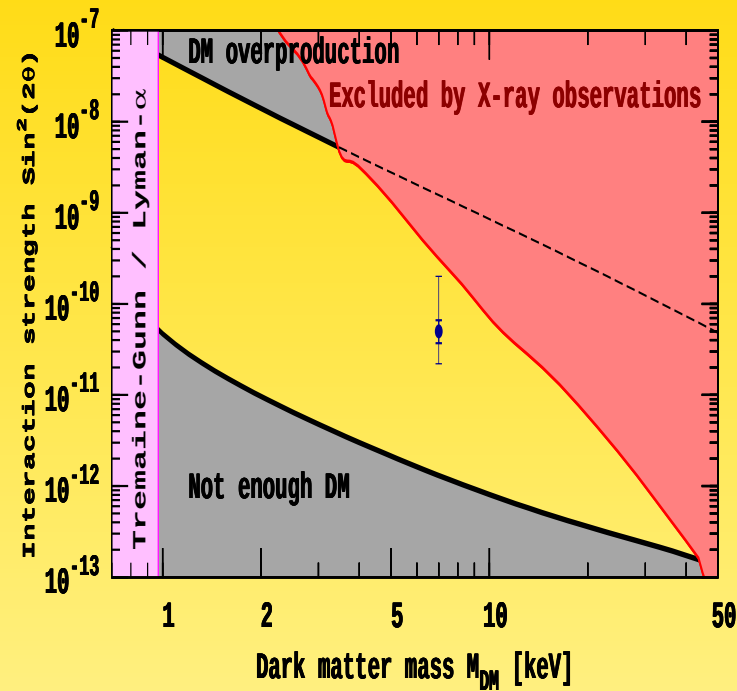


$\Rightarrow$ mixing down to 10^{-7} in reach!?



⇒ mixing down to 10^{-7} in reach!?

hard for KATRIN to see something...

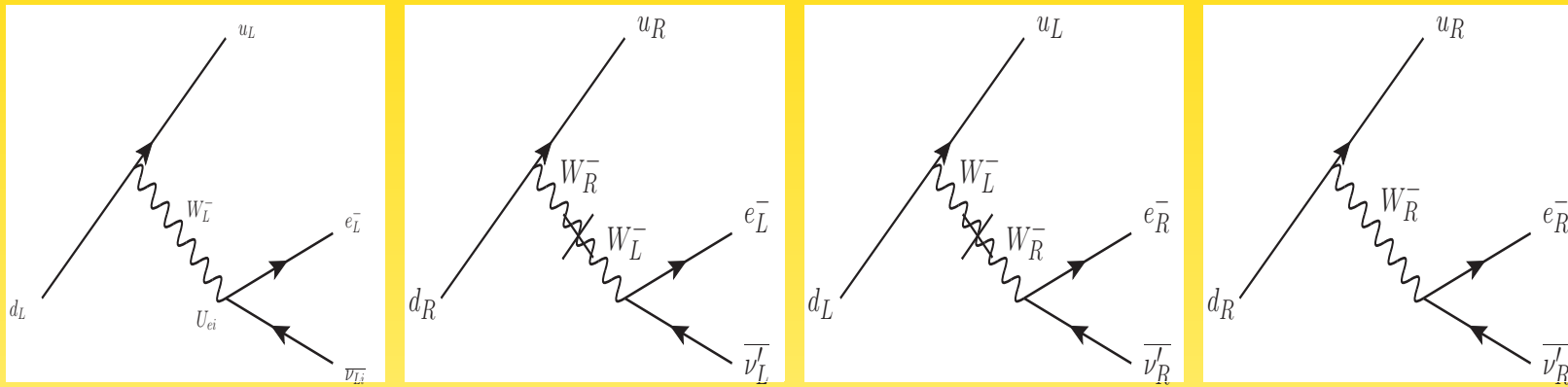


⇒ mixing down to 10^{-7} in reach!?

hard for KATRIN to see something...

...can increase signal with additional interactions, e.g. right-handed currents

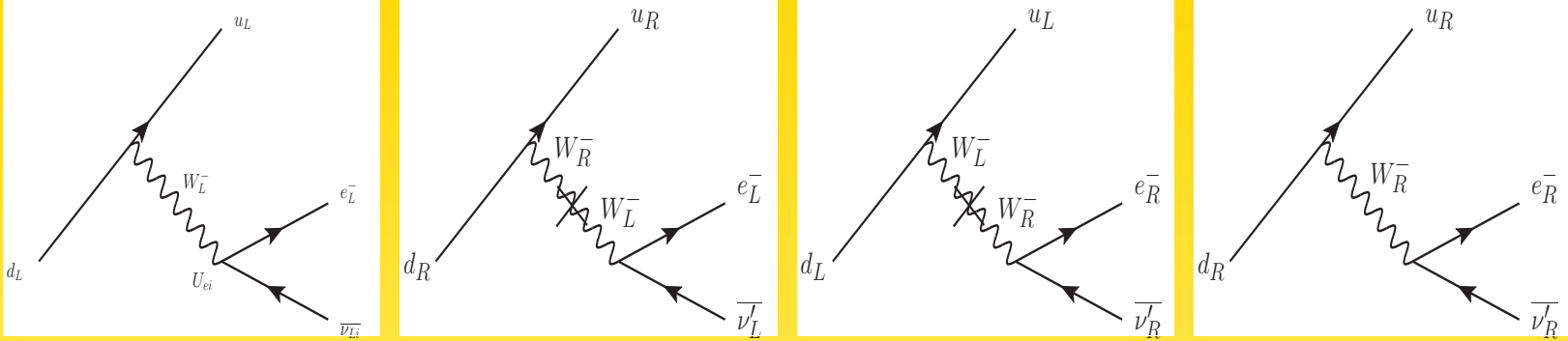
KATRIN and right-handed currents



- left-handed contribution
- right-handed contribution
- interference contribution (Bonn *et al.*)

Neutrino masses up to $m = 18.6$ keV testable

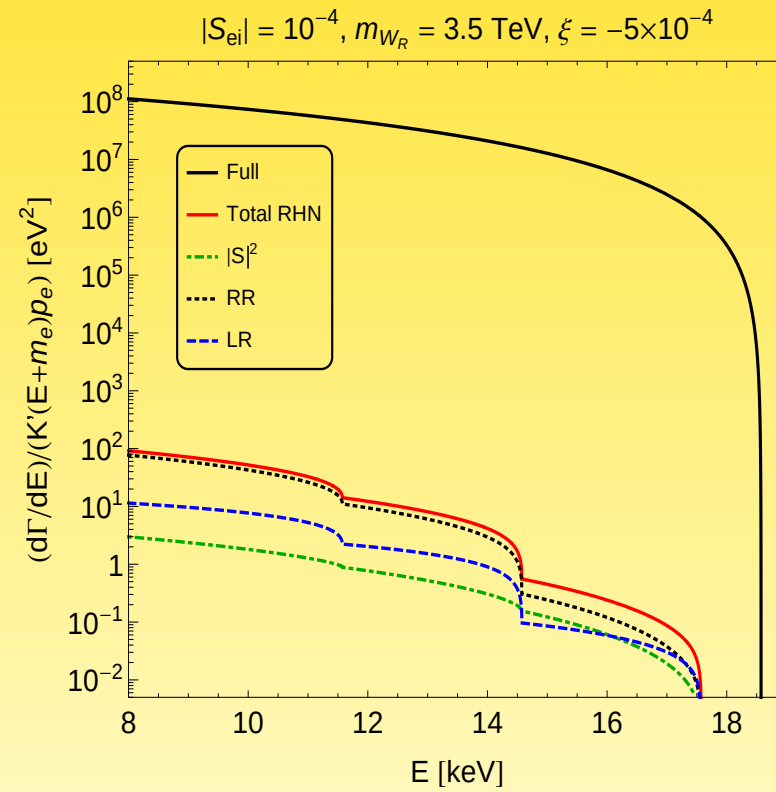
(Mertens *et al.*; deVega *et al.*)



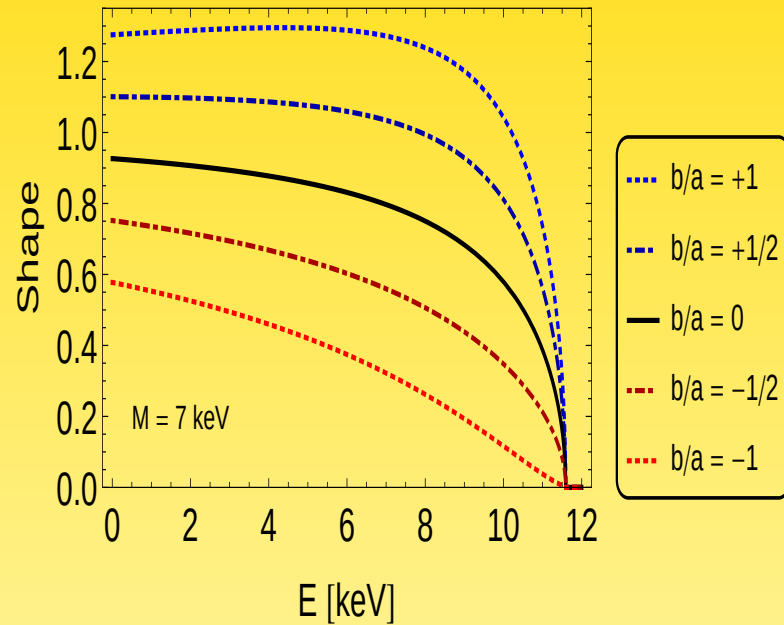
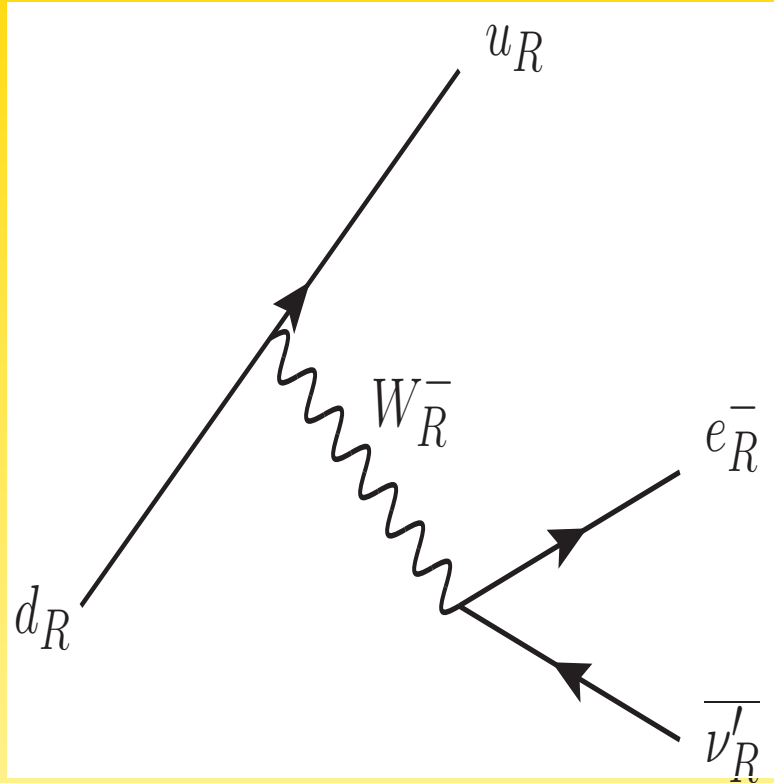
$$\begin{aligned}
\left(\frac{d\Gamma}{dE}\right)_{LL} &= K'(E + m_e)p_e X [1 + 2C \tan \xi] \\
&\times \left[|U_{ei}|^2 \sqrt{X^2 - m_i^2} \Theta(X - m_i) + |S_{ei}|^2 \sqrt{X^2 - M_i^2} \Theta(X - M_i) \right] \\
\left(\frac{d\Gamma}{dE}\right)_{RR} &\simeq K'(E + m_e)p_e X \left[\frac{m_{W_L}^4}{m_{W_R}^4} + \tan^2 \xi + 2C \frac{m_{W_L}^2}{m_{W_R}^2} \tan \xi \right] \\
&\times |V_{ei}|^2 \sqrt{X^2 - M_i^2} \Theta(X - M_i) \\
\left(\frac{d\Gamma}{dE}\right)_{LR} &= -2K' m_e p_e \text{Re} \left\{ \left[\left(\frac{m_{W_L}}{m_{W_R}} \right)^2 + C \tan \xi \right] \right. \\
&\times \left. \left[U_{ei} T_{ei} m_i \sqrt{X^2 - m_i^2} \Theta(X - m_i) + S_{ei} V_{ei} M_i \sqrt{X^2 - M_i^2} \Theta(X - M_i) \right] \right\} \\
&\text{with } X = E_0 - E
\end{aligned}$$

in total one has

$$\frac{d\Gamma}{dE} = \left(\frac{d\Gamma}{dE} \right)_{LL} + \left(\frac{d\Gamma}{dE} \right)_{RR} + \left(\frac{d\Gamma}{dE} \right)_{LR}$$



Focus for simplicity on



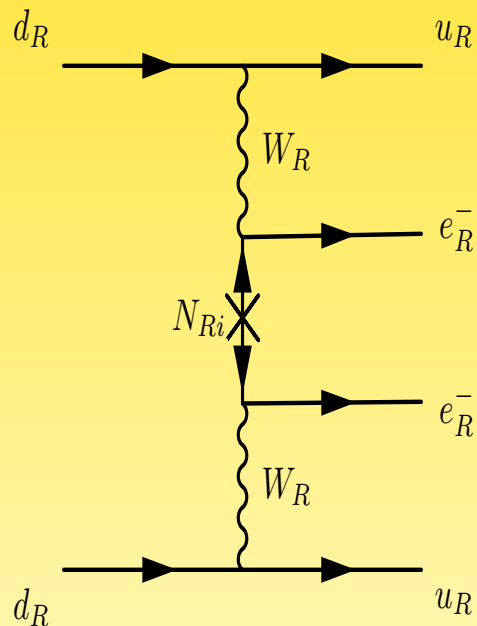
and assume keV neutrino with RH current

$$\frac{d\Gamma/dE - (d\Gamma/dE)_{\text{std}}}{(d\Gamma/dE)_{\text{std}}} \simeq \left(a + b \frac{M}{E_0 - E} \right) \sqrt{1 - \frac{M^2}{(E_0 - E)^2}} \Theta(E_0 - E - M)$$

neglect interference term b :

$$\theta_{\text{eff}}^2 \simeq |S_{ej}|^2 + 1.1 \times 10^{-6} |V_{ej}|^2 \left(\frac{2.5 \text{ TeV}}{m_{W_R}} \right)^4$$

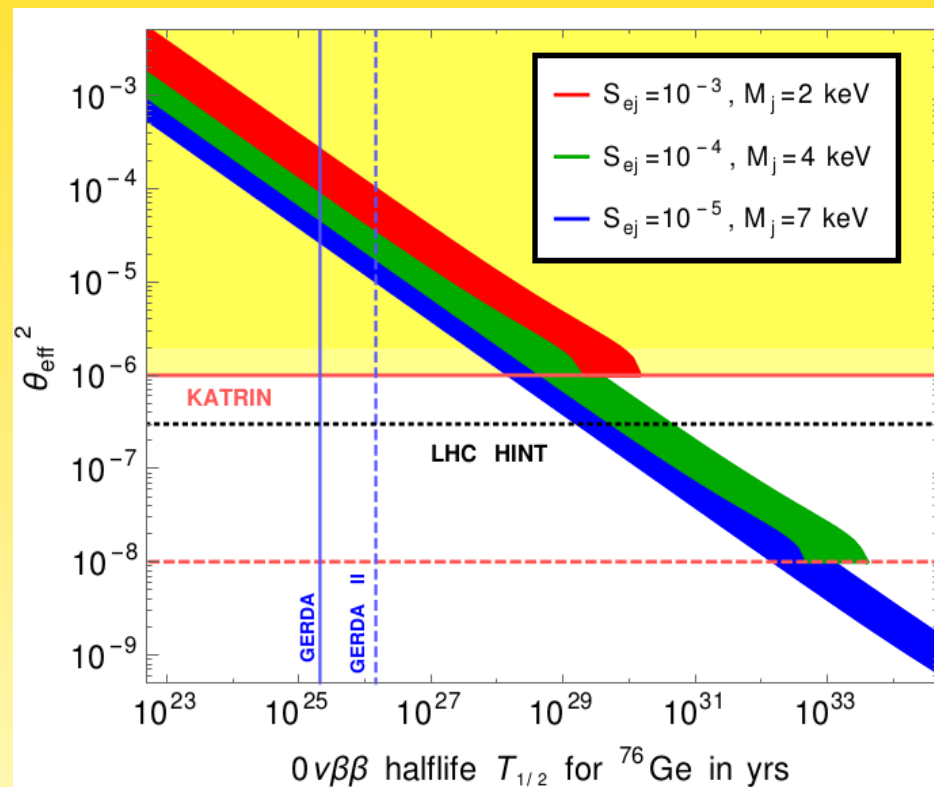
and note that M does $0\nu\beta\beta$ with amplitude $\propto |V_{ej}|^2 (m_W/m_{W_R})^4 M/q^2$



$\Rightarrow$ connection to $0\nu\beta\beta$ constraints!

connection to $0\nu\beta\beta$ constraints:

$$\theta_{\text{eff}}^2 = |S_{ej}|^2 + \frac{m_e}{M_j} \left[|\mathcal{M}_\nu^{0\nu}|^{-2} (G_{01}^{0\nu})^{-1} \left(T_{1/2}^{0\nu}\right)^{-1} - |S_{ej}^2 M_j / m_e|^2 \right]^{\frac{1}{2}}$$



Barry, Heeck, W.R., 1404.5955

How the additional interactions save the day

- double beta decay without RHC: $\theta^2 M = 7 \times 10^{-10} \text{ keV} = 70 \mu\text{eV}$
- double beta decay with RHC: $(m_{W_L}/m_{W_R})^4 |V_{ei}|^2 M = 8 \text{ meV}$
- decay: $\frac{\Gamma_{\text{RHC}}(N_j \rightarrow \bar{\nu}\gamma)}{\Gamma_{\text{SM}}(N_j \rightarrow \nu\gamma)} \simeq \frac{m_{W_L}^4 |S_{ei}|^2}{m_{W_R}^4 |T_{ei}|^2} \simeq \frac{m_{W_L}^4}{m_{W_R}^4}$
- beta decay: $\theta_{\text{eff}}^2 \simeq |S_{ej}|^2 + 1.1 \times 10^{-6} |V_{ej}|^2 \left(\frac{2.5 \text{ TeV}}{m_{W_R}} \right)^4 > |S_{ej}|^2$

Implizieren Dirac-Neutrinos Leptonzahlerhaltung?

Model basierend auf $B - L$, gebrochen um 4 Einheiten

$\Rightarrow$ Neutrinos sind Dirac, $\Delta L = 2$ verboten, aber $\Delta L = 4$ erlaubt. . .

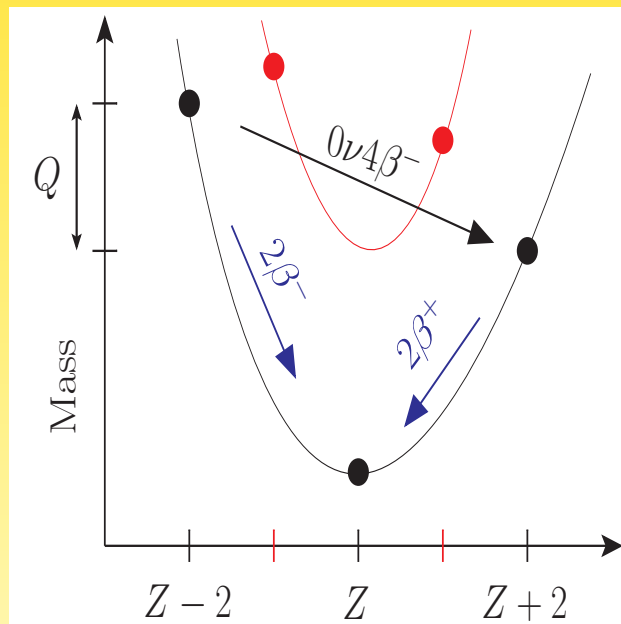
Heeck, W.R., EPL 103

Implizieren Dirac-Neutrinos Leptonzahlerhaltung?

Model basierend auf $B - L$, gebrochen um 4 Einheiten

⇒ Neutrinos sind Dirac, $\Delta L = 2$ verboten, aber $\Delta L = 4$ erlaubt...

⇒ Observable: neutrinoloser vierfacher Betazerfall $(A, Z) \rightarrow (A, Z + 4) + 4e^-$



	$Q_{0\nu 4\beta}$	Zerfälle	NA
${}^{96}_{40}\text{Zr} \rightarrow {}^{96}_{44}\text{Ru}$	0.629	$\tau_{1/2}^{2\nu 2\beta} \simeq 2 \times 10^{19}$	2.8
${}^{136}_{54}\text{Xe} \rightarrow {}^{136}_{58}\text{Ce}$	0.044	$\tau_{1/2}^{2\nu 2\beta} \simeq 2 \times 10^{21}$	8.9
${}^{150}_{60}\text{Nd} \rightarrow {}^{150}_{64}\text{Gd}$	2.079	$\tau_{1/2}^{2\nu 2\beta} \simeq 7 \times 10^{18}$	5.6

Heeck, W.R., EPL 103