

Precision observables for new physics searches

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Status of new physics searches today

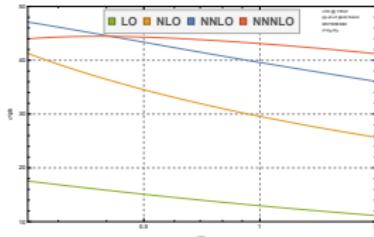
- No direct evidence for new physics at the LHC (so far)
- Use available data to maximum extent possible
- Use precision calculations to search for / constrain new physics

Goals of this talk

- Show progress in precision calculations for the LHC
- Highlight examples with connection to BSM searches
- Discuss future challenges for higher-order calculations

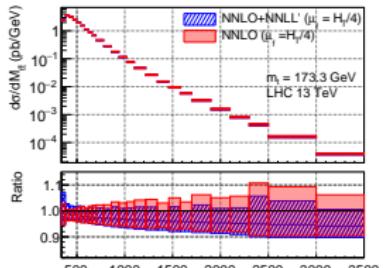
Precision calculations for the LHC

$gg \rightarrow h$



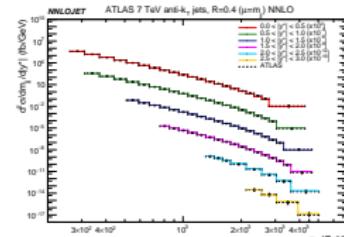
[Anastasiou, Duhr, Dulat, Herzog, Mistlberger '15]

$pp \rightarrow t\bar{t}$



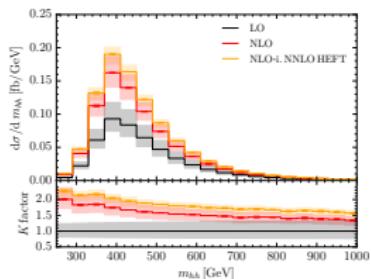
[Czakon, Ferroglio, Heymes, Mitov, Pecjak, Scott, Wang, Yang '18]

$pp \rightarrow jj$



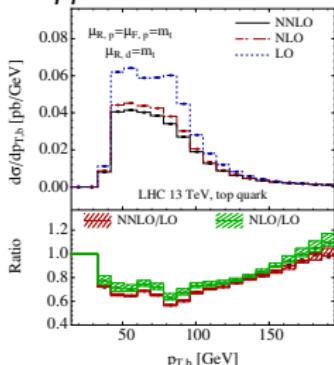
[Currie, Gehrmann-De Ridder, Gehrmann, Glover, Huss, Pires '17]

$pp \rightarrow hh$



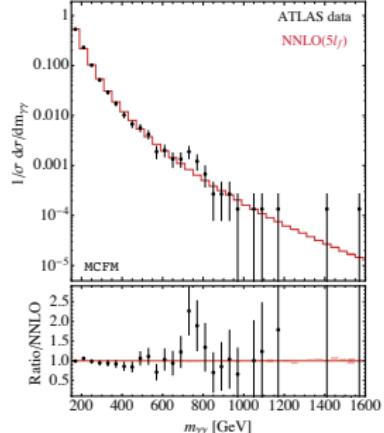
[Borwka, Greiner, Heinrich, Jones, Kerner, Schlenk, Zirke '16]

$pp \rightarrow tX \rightarrow b\ell\nu X$



[Berger, Gao, Yuan, Zhu '16]

$pp \rightarrow \gamma\gamma$

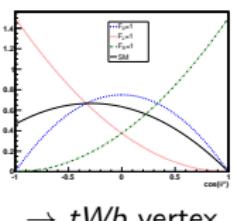


[Campbell, Ellis, Li, Williams '16]

New physics searches using precision calculations

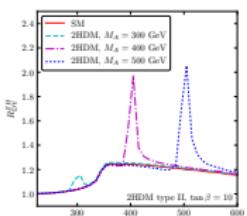
- Precision calculations become increasingly important for new physics searches
- New physics searches become possible that would not be possible otherwise

W helicity fractions



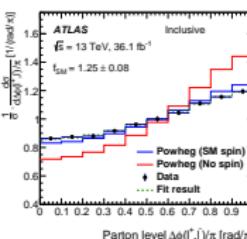
→ tWb vertex

ZH/WH



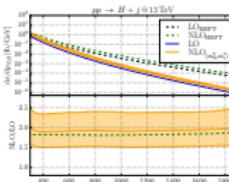
→ BSM in ZH

$t\bar{t}$ spin correlations



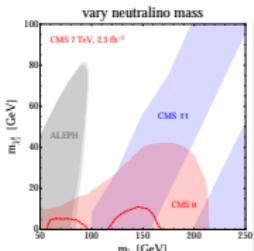
→ stealth stops & CMDM operator

Higgs @ high p_T



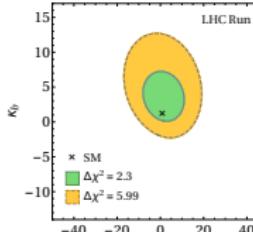
→ ggH vertex:
 y_t vs. c_g

$t\bar{t}$ cross-section



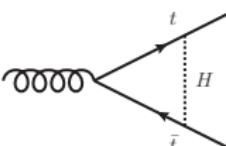
→ stealth stops

Higgs @ low p_T



→ light Yukawas

$m_{t\bar{t}}$ distributions

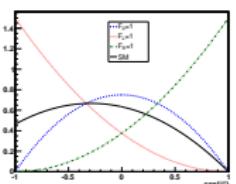


→ top Yukawa

New physics searches using precision calculations

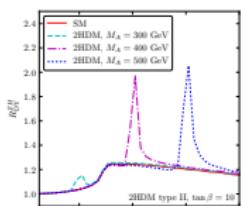
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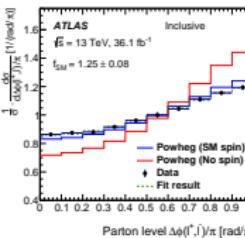
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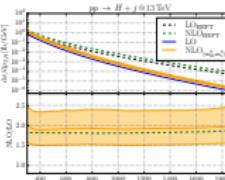
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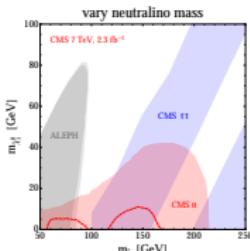
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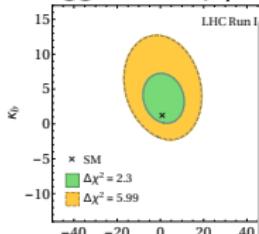
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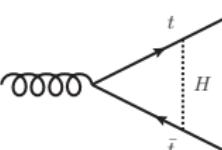
→ stealth stops

Higgs @ low p_T



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$m_{t\bar{t}}$ distributions



→ top Yukawa

Top decay: W helicity fractions

Observable

- Top decay: Differential distribution of charged lepton/ d -type quark

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta^*} = \frac{3}{8} (1 - \cos \theta^*)^2 F_L + \frac{3}{4} (\sin \theta^*)^2 F_0 + \frac{3}{8} (1 + \cos \theta^*)^2 F_R$$

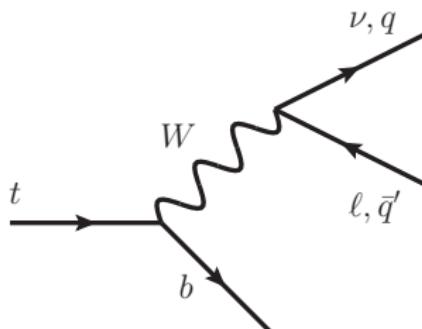
- $\cos \theta^*$: Angle between charged lepton/ d -type quark and reversed direction of b quark in W rest frame

W helicity fractions

- W is massive vector boson
 \Rightarrow 3 polarisations: L, R & longitudinal

$$F_{L,R,0} = \frac{\Gamma_{L,R,0}}{\Gamma}$$

- Measures production mode of W bosons in top decays

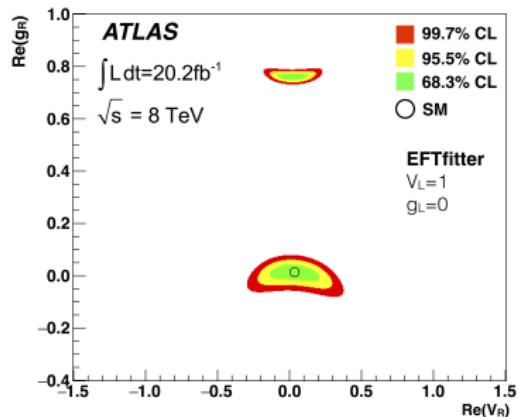


Relevance for BSM searches

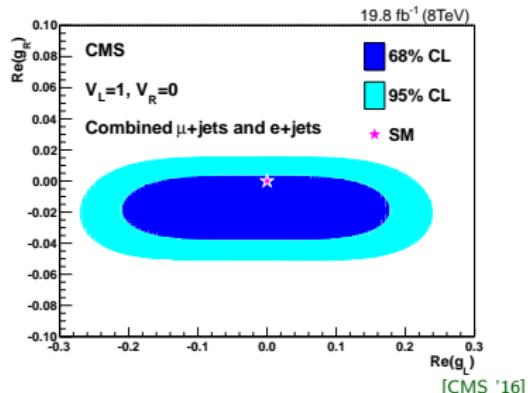
- Sensitive to structure of tWb vertex, e.g.,

$$\begin{aligned}\mathcal{L}_{tWb} = & -\frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (V_L P_L + V_R P_R) t W_\mu^- \\ & - \frac{g}{\sqrt{2}} \bar{b} \frac{i \sigma^{\mu\nu} q_\nu}{m_W} (g_L P_L + g_R P_R) t W_\mu^- + \text{h.c.}\end{aligned}$$

- Longitudinal W polarisation related to Goldstone bosons
⇒ sensitive to modified Higgs sector
- Used by ATLAS and CMS to set limits on anomalous couplings



[ATLAS '16]

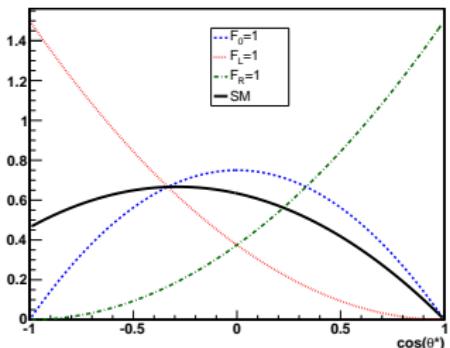


[CMS '16]

Experimental situation

Use angular distribution of W decay products to measure helicity of W boson

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta^*} = \frac{3}{8}(1 - \cos \theta^*)^2 F_L + \frac{3}{4}(\sin \theta^*)^2 F_0 + \frac{3}{8}(1 + \cos \theta^*)^2 F_R$$



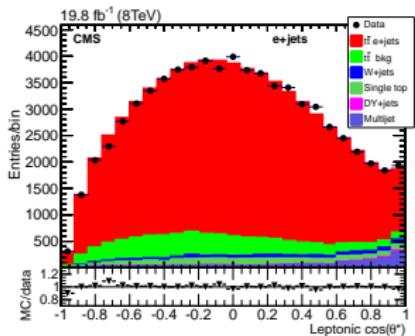
[CMS '16]

	F_0	F_L	F_R
ATLAS [ATLAS '17]	0.709 ± 0.019	0.299 ± 0.015	-0.008 ± 0.014
NNLO QCD + NLO EW	0.687 ± 0.005	0.311 ± 0.005	0.0017 ± 0.0001
[Czarnecki, Körner, Piclum '10]			

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[CMS '16]

Measurement

- $t\bar{t}$ production in lepton+jets channel
- Reconstruct $t\bar{t}$ system
⇒ associate b jet correctly
- Extrapolate fiducial to inclusive measurement
- Fit distribution to extract $F_{L,R,0}$

	F_0	F_L	F_R
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Precision predictions

W helicity fractions

- Measured very precisely (percent level accuracy)
 - Perturbative corrections converge well [Czarnecki, Körner, Piclum '10]
e.g., $F_L = 0.3032 \cdot (1_{\text{LO}} + 2.15\%_{\text{NLO}} + 0.70\%_{\text{NNLO}})$
 - Question: How close is what is calculated to what is being measured?
 - Decay in isolation:
Impossible to compute an observable following a set-up similar to experimental one
 - No realistic selection criteria, no contamination from the production, no combinatorics
- ⇒ Experiments have to extrapolate from fiducial to inclusive measurement

Recent developments

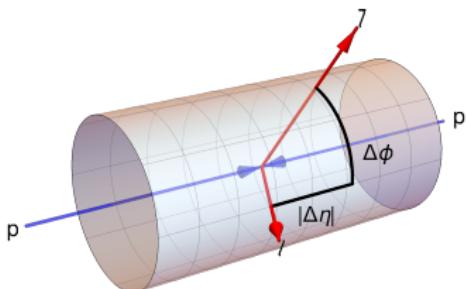
- Progress in fixed-order calculations:
⇒ $t\bar{t}$ production with decays at NNLO QCD
- More realistic description of final states

→ Moves interface between theory and experiment closer to experimental side

Spin correlations in $t\bar{t}$ production

Observable

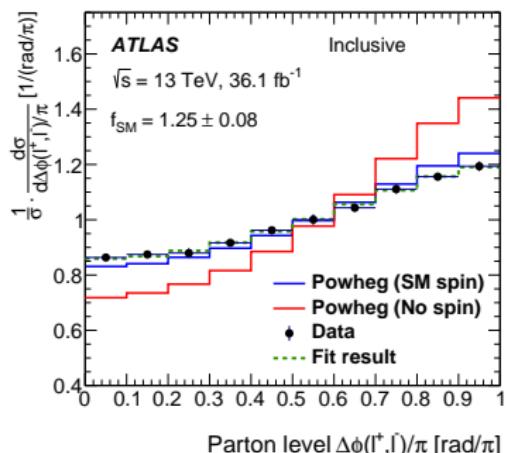
- $t\bar{t}$ production in dilepton channel
- Tops and anti-tops from $t\bar{t}$ essentially unpolarised, but spins highly correlated
- Differential distributions of lepton kinematics in lab frame
 - Lepton azimuthal opening angle $\Delta\phi(\ell, \bar{\ell})$
 - Lepton rapidity difference $|\Delta\eta(\ell, \bar{\ell})|$



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Competing effects

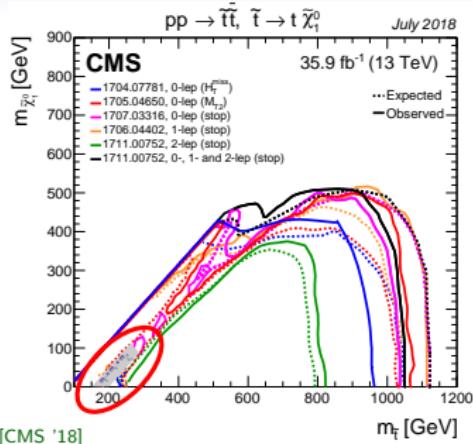
- Boosts of top quarks enhance anti-parallel leptons
- Spin correlations counteract and enhance parallel leptons

[ATLAS '19]

New physics searches with $t\bar{t}$ spin correlations

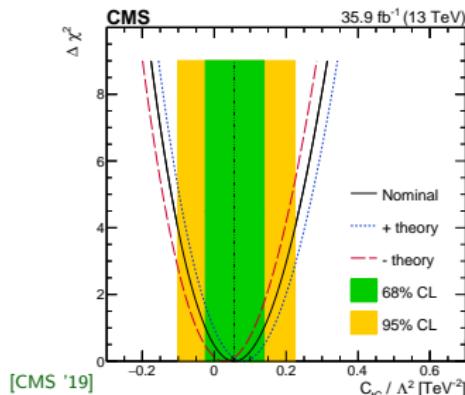
Example: Stealth stops

- Top squarks can be light in the MSSM
- Light stops have mostly been excluded by direct searches
- But if $m_{\tilde{t}} \approx m_t$: Difficult to distinguish from SM tops \rightarrow “Stealth stops”
- Squarks are scalar particles
 \Rightarrow Use spin correlations to search for stealth tops



Example: Top quark CMDM operator

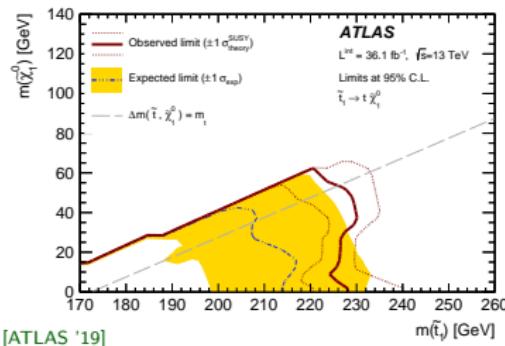
- Chromomagnetic dipole moment operator
- $$O_{tG} = y_t g_s (\bar{Q} \sigma^{\mu\nu} T^a t) \tilde{\phi} G^a_\mu$$
- Modifies $t\bar{t}$ spin correlations
 - Can be constrained using these observables



New physics searches with $t\bar{t}$ spin correlations

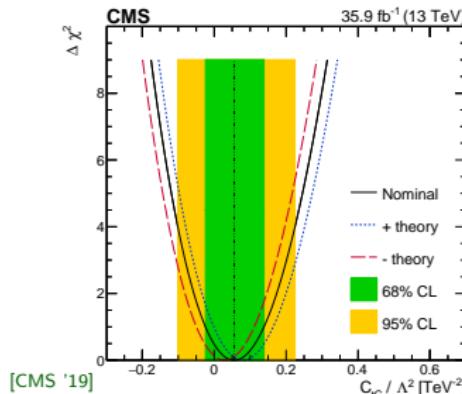
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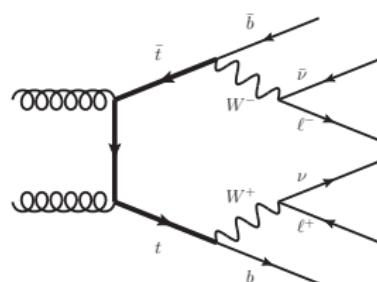


NNLO predictions for $t\bar{t}$ spin correlations

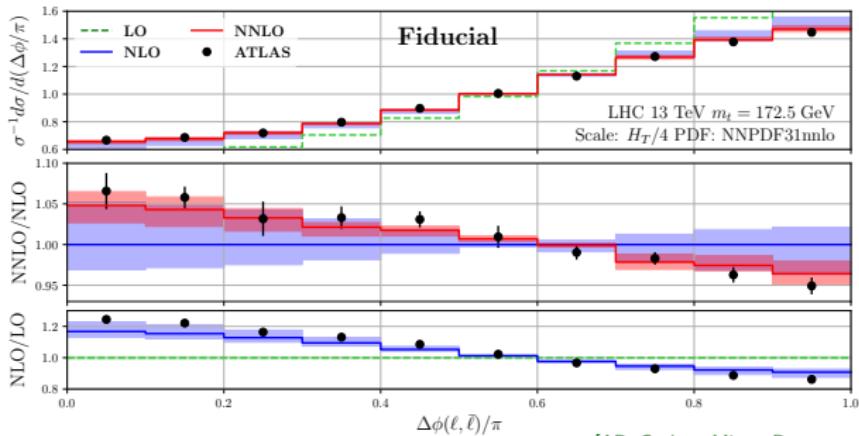
Why has this not been possible before?

- Complicated process with complex final state
⇒ Use narrow width approximation (NWA)
to factorise production and decay
- Virtual corrections: Needs polarised 2-loop ME
 - Decay: Available since ~ 2008 [Bonciani, Ferroglio '08]
[Asatrian, Greub, Pecjak '08] [Beneke, Huber, Li '09]
 - Production: Recently completed
[Chen, Czakon, Poncelet '18]
- Real radiation: Needs efficient subtraction
for many IR limits
⇒ Extension of sector improved residue
subtraction scheme for NWA

→ Payoff: NNLO QCD predictions for realistic final states



Dilepton $\Delta\phi(\ell, \bar{\ell})$ distributions at NNLO



- Agreement between data and theory improves at NNLO
- Higher-order corrections (\rightarrow additional radiation) can mimic spin-correlations
 \rightarrow Precision calculations are important to model observables accurately

Higgs p_T distribution

Motivation

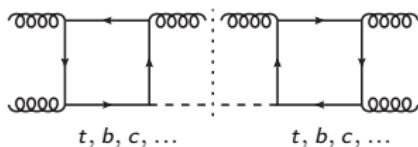
- Charm Yukawa coupling y_c is mostly unconstrained so far, $\kappa_c = \frac{y_c}{y_c^{\text{SM}}}$
- y_c is difficult to measure directly @ LHC

Idea: Detect/place limits on deviations of y_Q

Use p_T spectrum of Higgs at $m_c \ll p_T \ll m_h$ [Bishara, Haisch, Monni, Re '17] [Soreq, Zhu, Zupan '16]

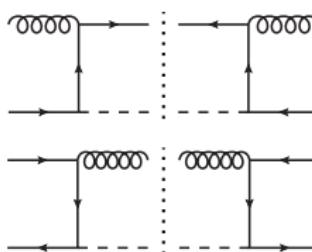
- $gg \rightarrow hg$:

$$d\sigma = d\sigma_{tt} + d\sigma_{tQ} + d\sigma_{QQ} + \dots$$



- $d\sigma_{tQ}$: $\kappa_Q \frac{m_Q^2}{m_h^2} \ln^2 \left(\frac{p_\perp^2}{m_Q^2} \right)$ terms
 \Rightarrow enhance b and c contributions at low p_T
- Most relevant for $\kappa_Q \lesssim 10$

- $gQ \rightarrow hQ$ and $Q\bar{Q} \rightarrow hg$:



- Scales like $\kappa_Q^2 \frac{m_Q^2}{m_h^2}$
- But suppressed by PDFs
- Most relevant for $\kappa_Q \gtrsim 10$

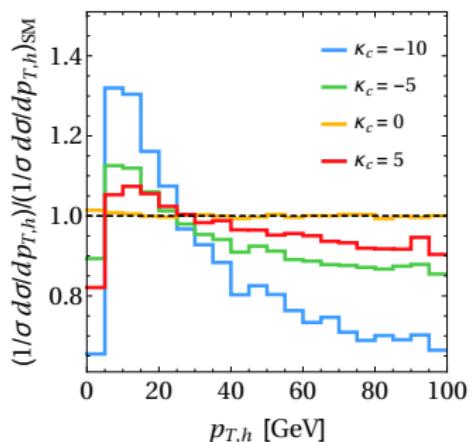
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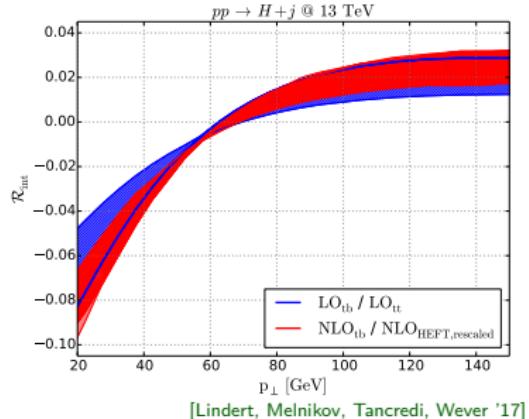
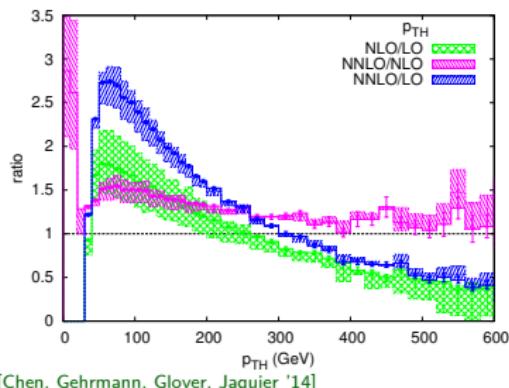


[Bishara, Haisch, Monni, Re '17]

- Higgs p_T spectrum receives large perturbative corrections
- Charm contribution is small compared to corrections from top loop

Precision predictions

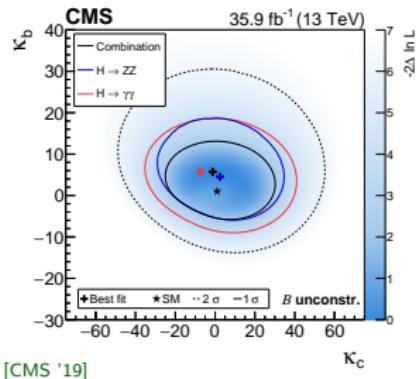
- Top-loop contributions known
 - to NNLO + N^3LL with $m_t \rightarrow \infty$ (HEFT)
 - to NLO with full m_t dependence
- Top-bottom & top-charm interference known
 - to LO with full m_Q
 - to NLO with $m_Q \rightarrow 0$
- Contributions from b and c quarks are $O(1 - 5\%)$ effects in the SM



→ Requires precision calculations since perturbative corrections are large

Experimental situation

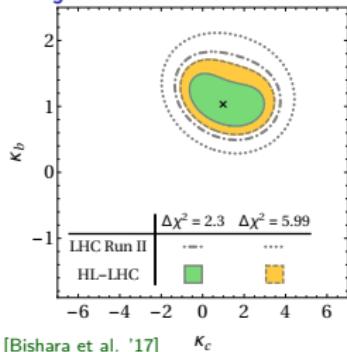
Current situation



[CMS '19]

- Based on NLO QCD calculation
- Corresponds to LO for top-charm interference
- Constrains $-33 < \kappa_c < 38$ at 95% CL

Projections for Run II and HL-LHC



[Bishara et al. '17]

- Assumptions on theoretical uncertainties for
 - Run II: $\pm 5\%$
 - HL-LHC: $\pm 2.5\%$
- Requires at least NLO top-charm interference
- Would constrain charm Yukawa coupling to
 - Run II: $-1.4 < \kappa_c < 3.8$
 - HL-LHC: $-0.6 < \kappa_c < 3.0$

Challenges for the future

Recent developments allow us to think about new challenging opportunities

- Mixed QCD-EW corrections to $Z + j$ at high p_T
 - Background to dark matter studies at LHC
 - Amplitudes with nearly massless EW gauge bosons
 - Progress from expansion techniques?
- $gg \rightarrow ZH$
 - Search for new physics in top loop
 - Evaluation of integrals with different masses
 - Progress from numerical techniques?
- tth production
 - Measure top Yukawa coupling precisely
 - Many masses/scales
 - challenging integral reductions
 - evaluation of very complicated integrals
- ttZ production
 - Probe top gauge couplings
 - Similar challenges as for tth

Conclusions

- Remarkable progress in higher-order corrections in recent years
- More realistic final states and fiducial calculations become possible
- Precision calculations enable us to make the most of (HL-)LHC data
- New physics searches profit from increased precision
→ New opportunities for measurements that would be impossible otherwise
- Communication with different communities (experiment, EFT, model building, mathematics, . . .) will be helpful to guide and inspire progress