

DM PAIR PRODUCTION IN SIMPLIFIED MODELS AND THE MSSM

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based on Eur.Phys.J. C79 (2019) no.5, 428

[C. Borschensky, GC, B. Jäger]

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Outline

- 1 Motivation
- 2 Simplified dark matter models
 - s-channel model
 - t-channel model
- 3 Comparison of MSSM and simplified models
- 4 Summary



The dark matter model space

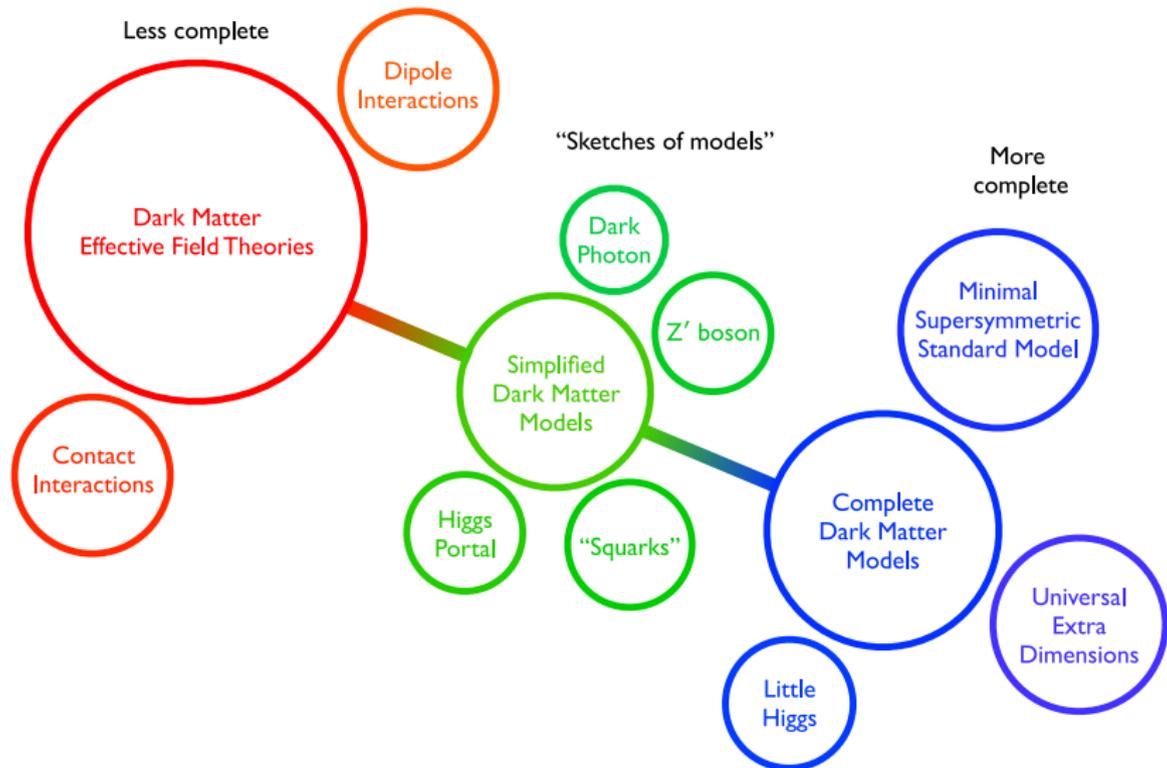


Figure taken from [1506.03116]



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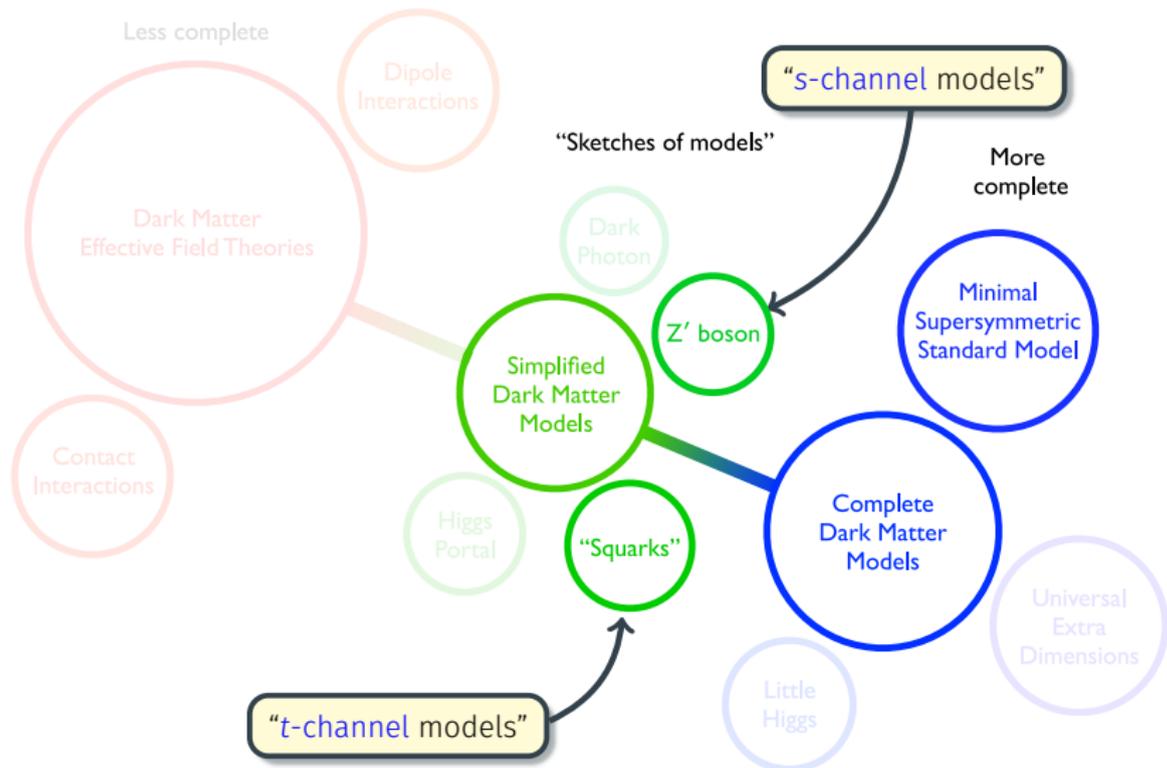


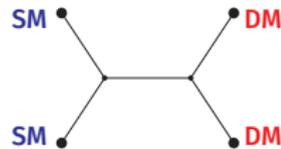
Figure taken from [1506.03116]



s-channel model

Assumption: DM is a singlet under $SU(3) \times SU(2) \times U(1)$ and ...

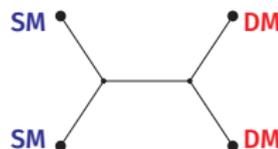
- ▶ ... consists of Dirac/Majorana fermions χ
- ▶ ... interacts with the SM via the topology:



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Possible Lagrangians with a scalar or vector mediator:

$$\mathcal{L}_S = g_X^S \bar{\chi} \chi S + \sum_q g_q^S \bar{q} q S$$

$$\mathcal{L}_P = i g_X^P \bar{\chi} \gamma_5 \chi P + \sum_q g_q^P \bar{q} \gamma_5 q P$$

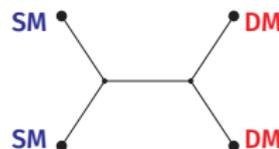
$$\mathcal{L}_V = \bar{\chi} \gamma_\mu \left[g_X^V - g_X^A \gamma_5 \right] \chi V^\mu + \sum_q \bar{q} \gamma_\mu \left[g_q^V - g_q^A \gamma_5 \right] q V^\mu$$



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In this talk

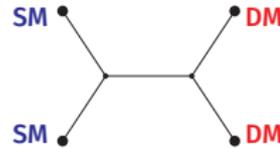
$$\mathcal{L}_V = \bar{\chi} \gamma_\mu [g_X^V - g_X^A \gamma_5] \chi V^\mu + \sum_q \bar{q} \gamma_\mu [g_q^V - g_q^A \gamma_5] q V^\mu$$



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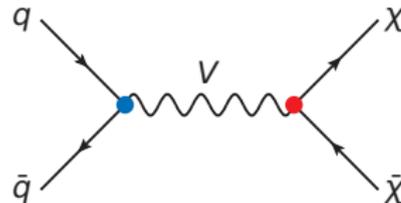
Interaction Lagrangian for a vector mediator [Dudas et al. '09]:

$$\mathcal{L}_V = \bar{\chi} \gamma_\mu [g_X^V - g_X^A \gamma_5] \chi V^\mu + \sum_q \bar{q} \gamma_\mu [g_q^V - g_q^A \gamma_5] q V^\mu$$

with q : quark fields, χ : DM field, V^μ : vector mediator field, g^V, g^A : vector and axialvector couplings

Properties

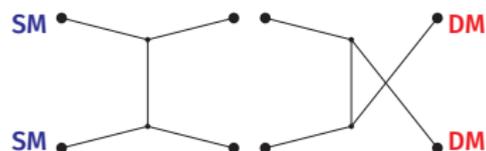
- ▶ V is uncoloured and massive (M_V)
- ▶ Added to SM by spontaneously broken $U(1)'$ symmetry to generate V mass
- ▶ Decays only into SM or DM pairs



t -channel model

Assumption: DM is a singlet under $SU(3) \times SU(2) \times U(1)$ and ...

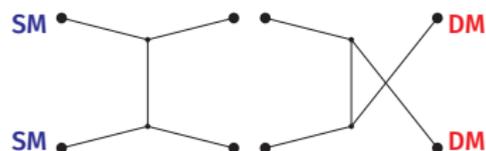
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Interaction Lagrangian for a coloured scalar mediator:

$$\mathcal{L}_{\tilde{Q}} = - \left[\lambda_{Q_L} \bar{\chi} P_L \tilde{Q}_L^\dagger \cdot Q + \lambda_{u_R} \tilde{Q}_{u_R}^* \bar{\chi} P_R u + \lambda_{d_R} \tilde{Q}_{d_R}^* \bar{\chi} P_R d + \text{h.c.} \right]$$

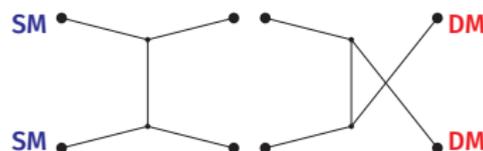
with $\tilde{Q}_L = \begin{pmatrix} \tilde{Q}_{u_L} \\ \tilde{Q}_{d_L} \end{pmatrix}$ an $SU(2) \times U(1)$ doublet



t -channel model

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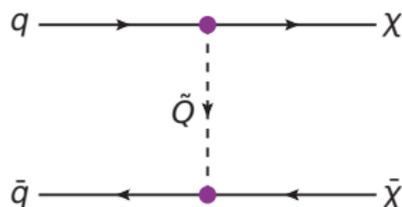
Interaction Lagrangian for a coloured scalar mediator [An et al. '14]:

$$\mathcal{L}_{\tilde{Q}} = - \left[\lambda_{Q_L} \left(\tilde{Q}_{u_L}^* \bar{\chi} P_L u + \tilde{Q}_{d_L}^* \bar{\chi} P_L d \right) + \lambda_{u_R} \tilde{Q}_{u_R}^* \bar{\chi} P_R u + \lambda_{d_R} \tilde{Q}_{d_R}^* \bar{\chi} P_R d + \text{h.c.} \right]$$

with u, d : up- and down-type quark fields, χ : DM field, $\tilde{Q}_{q_{L/R}}$: coloured scalar mediator fields, λ : DM-quark-squark Yukawa couplings, $P_{L/R}$: left- and right-handed chirality projectors

Properties

- ▶ \tilde{Q} are coloured and flavoured (#12)
- ▶ Heavier than χ so that the decay $\tilde{Q} \rightarrow q\chi$ is possible ($M_{\tilde{Q}} > m_\chi$)



Tools and numerical setup

Roadmap of the calculation:

- ▶ *Generate points in MSSM parameter space*

Spectrum generator: **SPheno 4.0.3** [Porod '03; Porod, Staub '12]

CMSSM [Adeel Ajaib, Gogoladze '17]	pMSSM10 [de Vries et al. '15]	
$M_0 \in [0, 10]$ TeV	$M_1 \in [-1, 1]$ TeV	$M_2 \in [0, 4]$ TeV
$m_{1/2} \in [0, 10]$ TeV	$M_3 \in [-4, 4]$ TeV	$m_{\tilde{q}_{1/2}} \in [0, 4]$ TeV
$A_0 \in [-3, 3] \times M_0$	$m_{\tilde{g}_3} \in [0, 4]$ TeV	$m_{\tilde{t}_1} \in [0, 2]$ TeV
$\tan \beta \in [2, 60]$	$M_A \in [0, 4]$ TeV	$A \in [-5, 5]$ TeV
$\text{sign } \mu > 0$	$\mu \in [-5, 5]$ TeV	$\tan \beta \in [1, 60]$

5000 points where $\tilde{\chi}_1^0$ is the LSP and the lightest Higgs mass satisfies $124 \text{ GeV} \leq m_h \leq 126 \text{ GeV}$

- ▶ *Fix parameters of s- and t-channel models*

Choose: $m_{\tilde{\chi}} = m_{\tilde{\chi}_1^0}$, $M_V = 1 \text{ TeV}$ and 10 TeV , $M_{\tilde{Q}} =$ average of $\tilde{u}_{L/R}$, $\tilde{d}_{L/R}$, $\tilde{c}_{L/R}$, $\tilde{s}_{L/R}$, $\tilde{b}_{1/2}$ masses,

$g_X^{V/A} = g_q^{V/A} = g = 0.5$, $\lambda_{Q_L} = \lambda_{U_R} = \lambda_{D_R} = \lambda = 1$

- ▶ *Calculate $pp \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0$ cross section in MSSM for each point*

POWHEG-BOX [Alioli, Nason, Oleari, Re '10] with weakino code [Baglio, Jäger, Kesenheimer '16]

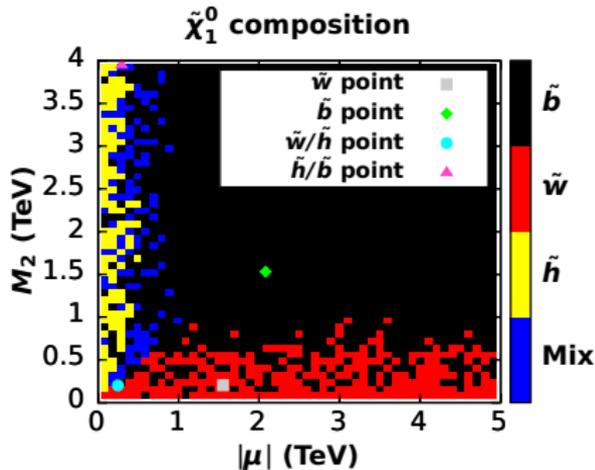
- ▶ *Calculate $pp \rightarrow \tilde{\chi} \tilde{\chi}$ cross section in SDMMs for each point*

POWHEG-BOX and for the t-channel model COLLIER-1.2 [Denner, Dittmaier, Hofer '17]

LHC at $\sqrt{s} = 13 \text{ TeV}$, PDFs used: PDF4LHC15 NLO MC PDFs [Butterworth et al. '16]



Parameter scan in the pMSSM10



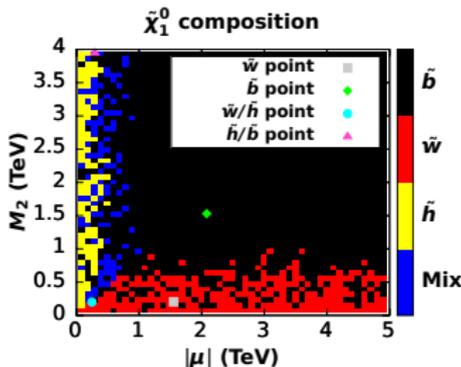
[C. Borschensky, GC, B. Jäger; Eur.Phys.J. C79 (2019) no.5, 428]

$\tilde{\chi}_1^0$ composition

- Distinguish between \tilde{b} ino, \tilde{w} ino, \tilde{h} iggsino
- Pure bino/wino: no Z exchange possible



Parameter scan in the pMSSM10

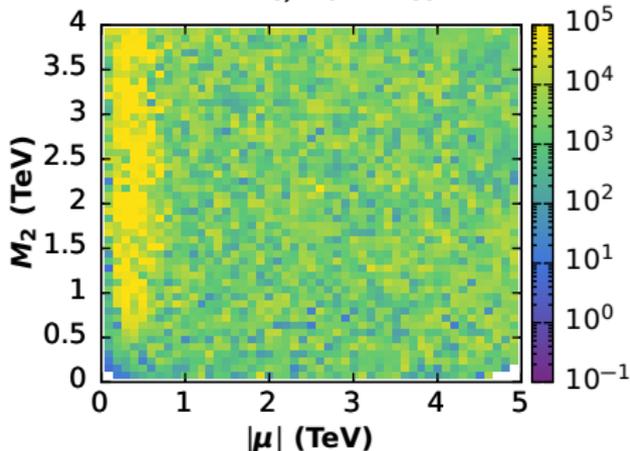


[C. Borschensky, GC, B. Jäger; Eur.Phys.J. C79 (2019) no.5, 428]

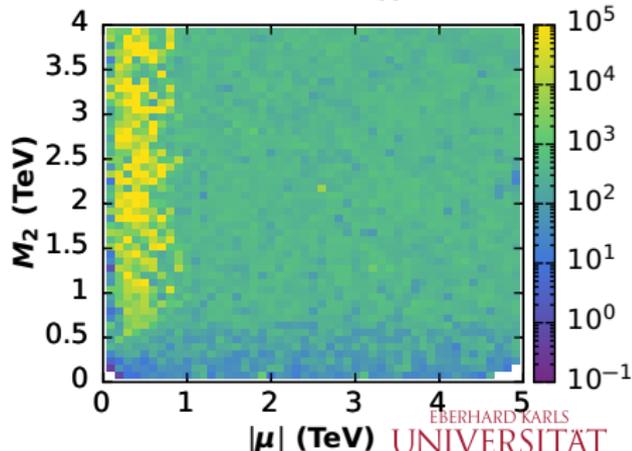
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- ▶ Distinguish between \tilde{b} ino, \tilde{w} ino, \tilde{h} iggsino
- ▶ Pure bino/wino: **no Z exchange** possible
- ▶ Higher s-channel mediator mass M_V smoothens the bino region a bit, **difficult to distinguish from t-channel model**

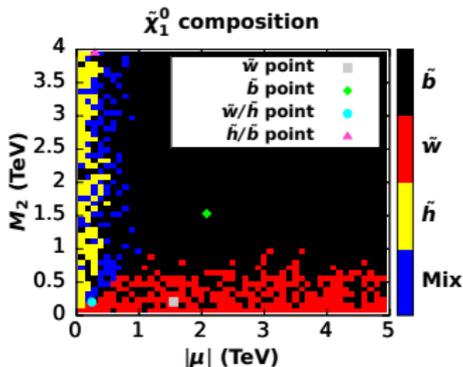
Ratio $\sigma_{s,1 \text{ TeV}}/\sigma_{\text{MSSM}}$



Ratio $\sigma_t/\sigma_{\text{MSSM}}$



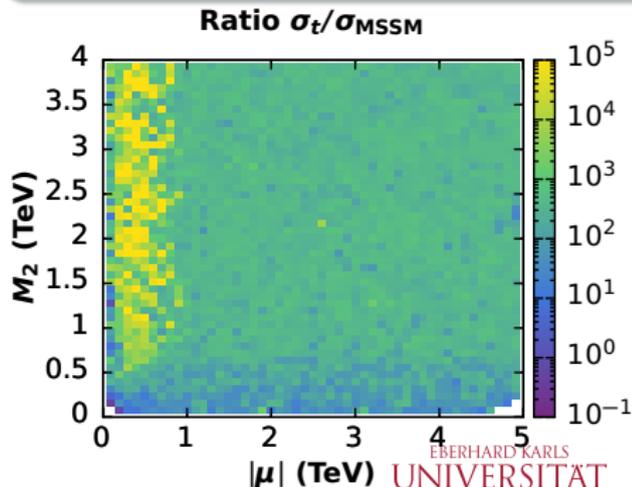
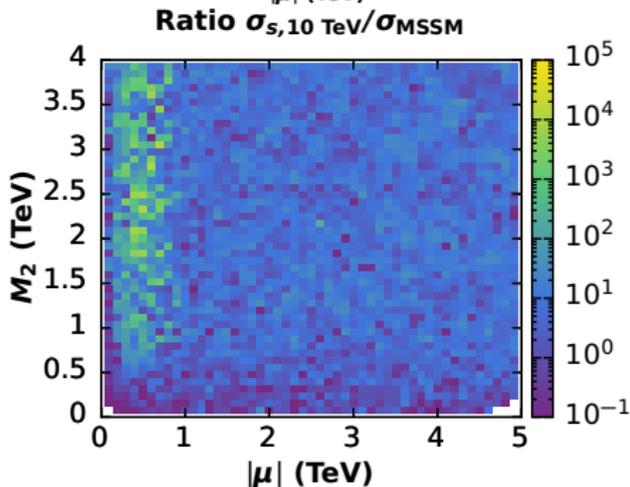
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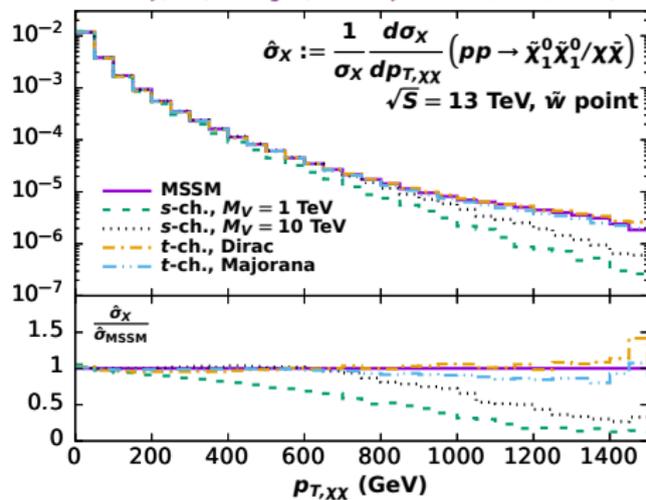
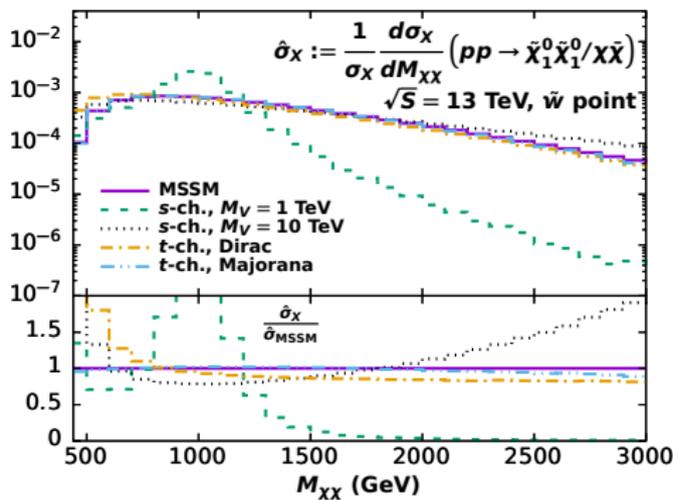
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Distributions: \tilde{w} point

[C. Borschensky, GC, B. Jäger; Eur.Phys.J. C79 (2019) no.5, 428]



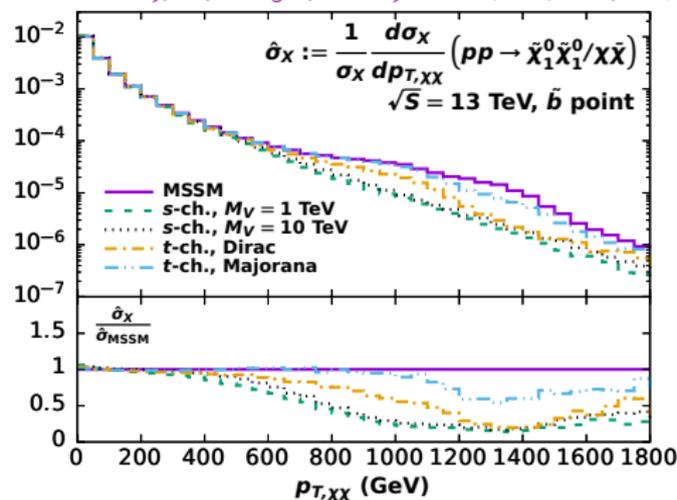
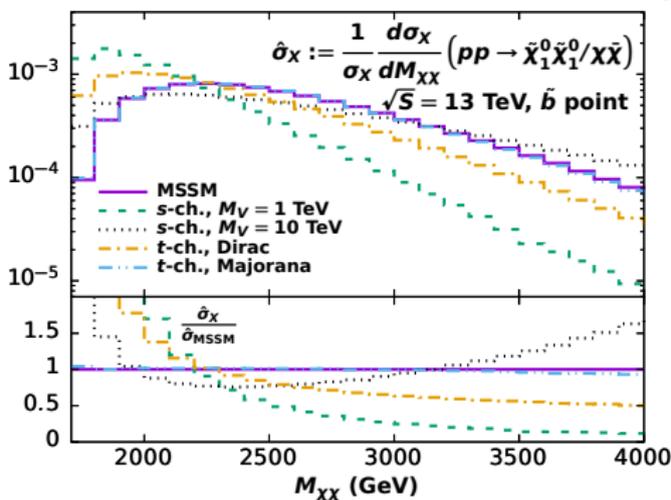
Analysis at NLO + parton shower accuracy

- ▶ $\tilde{\chi}_1^0$ mainly \tilde{w} ino, DM mass $\sim 220 \text{ GeV}$, squark masses $\sim 3 \text{ TeV}$
- ▶ t-channel very close to MSSM, agreement with s-channel ($M_V = 1 \text{ TeV}$) is worst



Distributions: \tilde{b} point

[C. Borschensky, GC, B. Jäger; Eur.Phys.J. C79 (2019) no.5, 428]



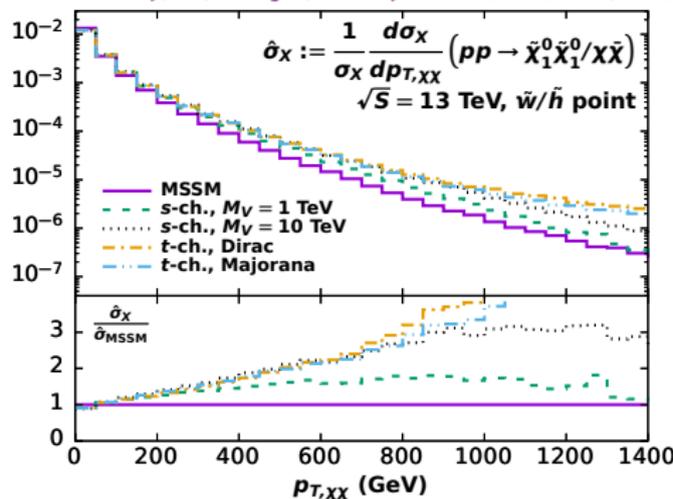
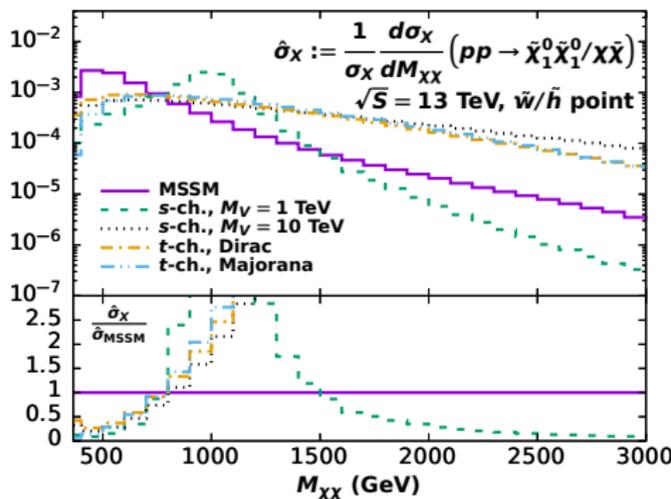
Analysis at NLO + parton shower accuracy

- ▶ $\tilde{\chi}_1^0$ mainly \tilde{b} ino, DM mass ~ 850 GeV, squark masses ~ 2.5 TeV
- ▶ t-channel close to MSSM, Majorana case better than Dirac
- ▶ Bump around $p_{T,XX} \approx 1.25$ TeV remnant of on-shell subtraction procedure



Distributions: \tilde{w}/\tilde{h} point

[C. Borschensky, GC, B. Jäger; Eur.Phys.J. C79 (2019) no.5, 428]



Analysis at NLO + parton shower accuracy

- ▶ $\tilde{\chi}_1^0$ mainly \tilde{w} ino- \tilde{h} iggsino, DM mass ~ 180 GeV, squark masses ~ 3.4 TeV
- ▶ No agreement between simplified models and MSSM for M_{XX} and $p_{T,XX}$ distributions



Summary and conclusions

Simplified models: studying DM with a minimal set of parameters



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Two specific models studied:

- ▶ s-channel model with a vector mediator, and t-channel model with coloured scalar mediators
- ▶ NLO QCD corrections including PS (PYTHIA 6) calculated for DM pair production at the LHC and implemented in the POWHEG-BOX framework



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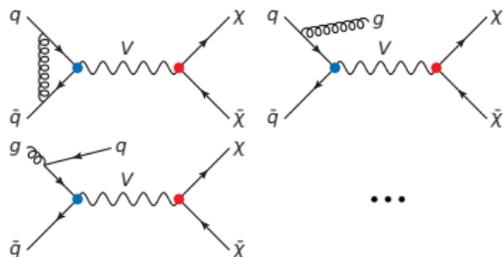
Comparison with $\tilde{\chi}_1^0$ pair-production in the MSSM:

- ▶ Simplified models can reproduce some MSSM features, in particular the t-channel model with only three parameters ($m_\chi, M_{\tilde{Q}}, \lambda$)
- ▶ However, poor agreement for our studied models in other regions
- ▶ Require more complex models, or simplified models better suited for description of some other non-SUSY theory?

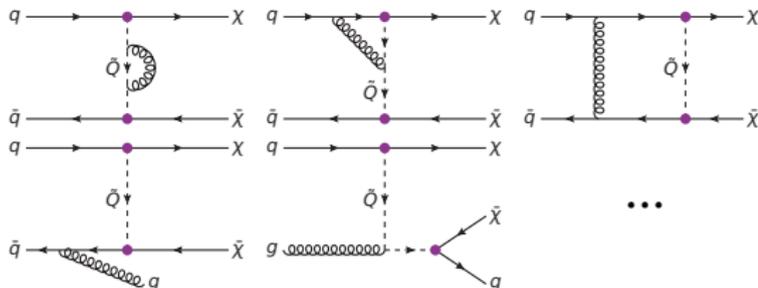


Backup :NLO QCD corrections

s-channel



t-channel

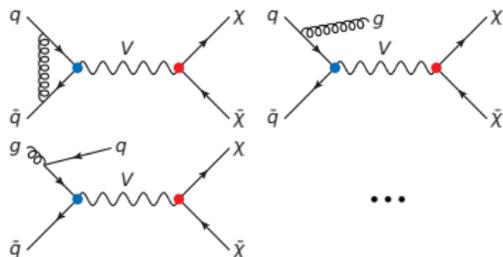


$$|\mathcal{M}_{\text{NLO}}|^2 = |\mathcal{M}_{\text{Born}}|^2 + 2 \text{Re}(\mathcal{M}_{\text{Born}} \mathcal{M}_{\text{virt}}^*) + |\mathcal{M}_{\text{real}}|^2 + \mathcal{O}(\alpha_s^2)$$

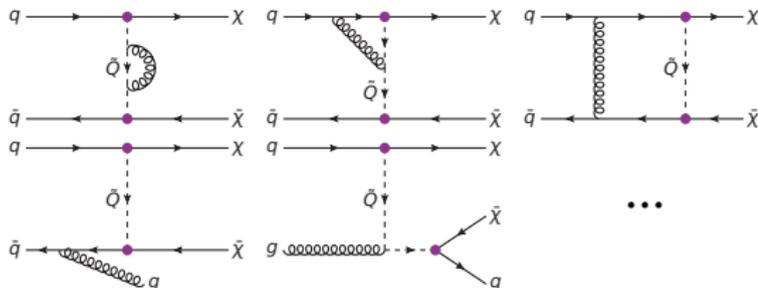


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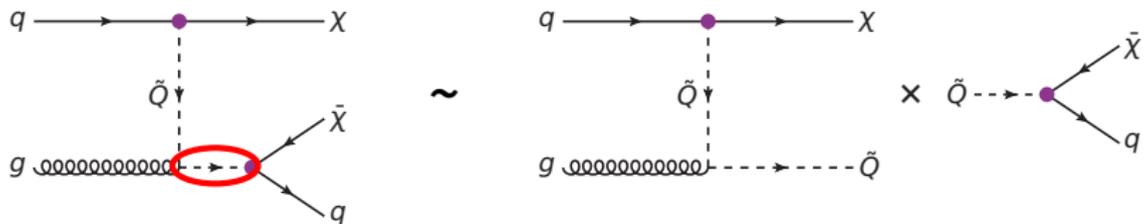
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- ▶ **Born, spin-/colour-correlated Born, and real amplitudes** calculated with **MadGraph 4** [Murayama et al., Stelzer et al., Alwall et al. '92-'07]
- ▶ **Virtual amplitudes** calculated with **FeynArts 3.9/FormCalc 9.4** [Hahn et al. '98-'16]



Backup: On-shell subtraction for the t -channel model

Subtlety in real corrections: intermediate \tilde{Q} becomes on-shell



On-shell subtraction (follow mainly [Baglio, Jäger, Kesenheimer '16-'17])

- ▶ Split up amplitude into **non-resonant (nr)** and **possibly resonant (r)** diagrams:

$$|\mathcal{M}_{\text{real}}|^2 = |\mathcal{M}_{\text{nr}}|^2 + 2 \operatorname{Re}(\mathcal{M}_{\text{nr}}\mathcal{M}_r^*) + |\mathcal{M}_r|^2$$

- ▶ Construct **local (i.e. for each phase space point) counterterm**

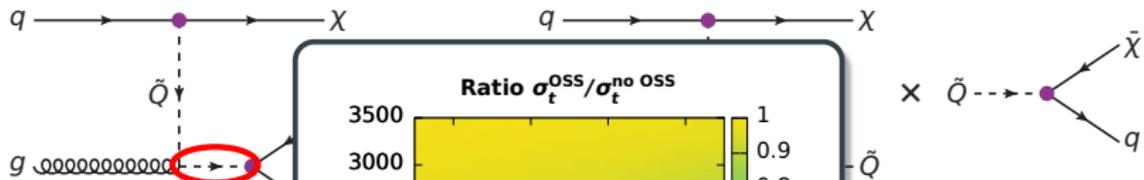
$$\text{CT} \propto \text{BW} \times |\mathcal{M}_r^{\text{OS}}|^2$$

with $\mathcal{M}_r^{\text{OS}}$ mapped to on-shell kinematics $p^2 = M_{\tilde{Q}}^2$, and where BW is a Breit-Wigner factor shaping the resonant propagator



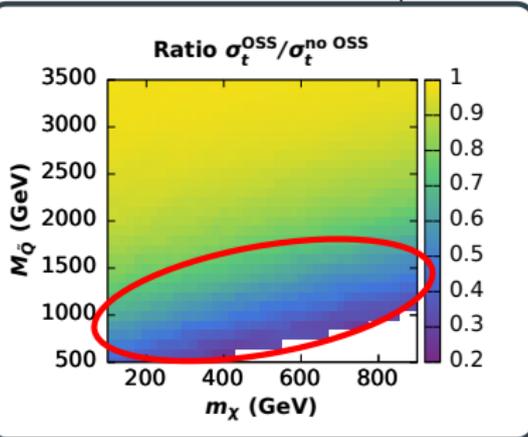
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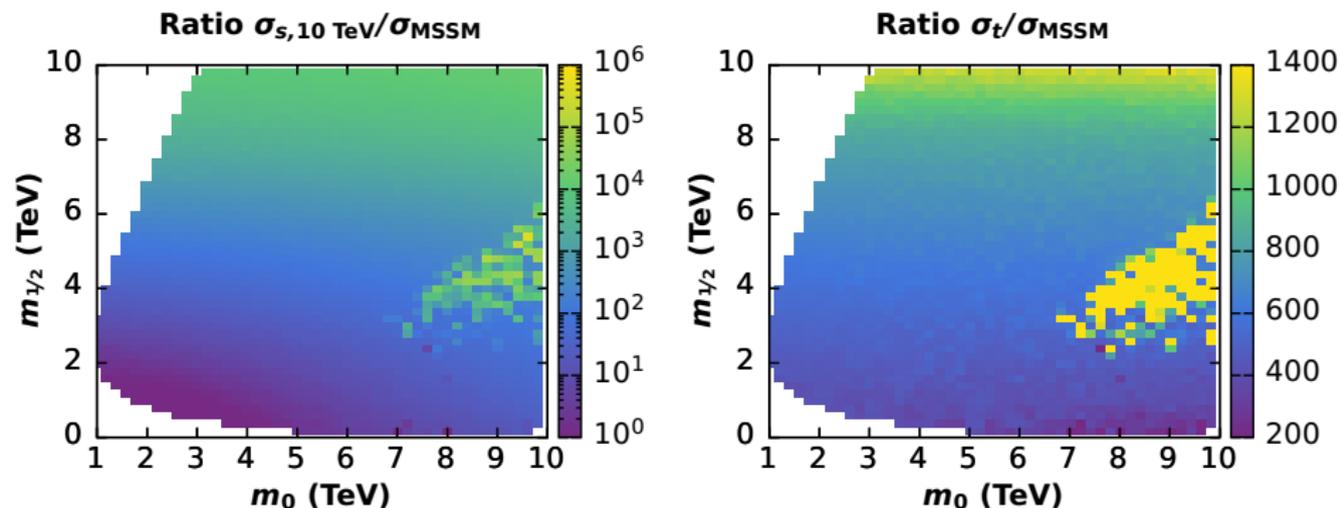
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Backup: Parameter scan in the CMSSM



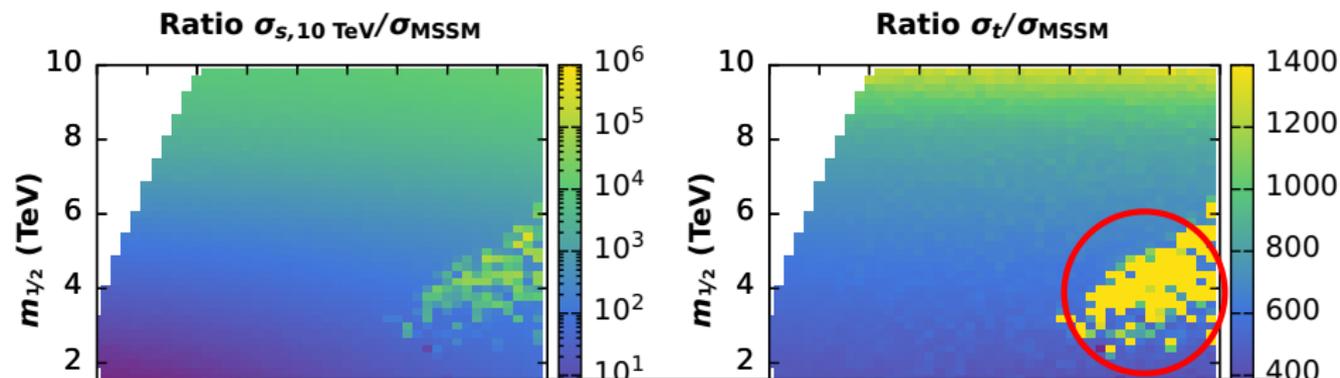
Ratios with respect to MSSM at NLO QCD

- ▶ Ratio involving s -channel model varies by several orders of magnitude
- ▶ Ratio involving t -channel model almost constant (\sim factor 3)
- ▶ Effect at high M_0 and intermediate $m_{1/2}$ due to “higgsino mixing matrix suppression” in MSSM $\Rightarrow \sigma_{\text{MSSM}}$ low

[GC, C. Borshensky, B. Jäger; 1812.08704]



Backup: Parameter scan in the CMSSM



In the MSSM:

$$\frac{ie}{2s_W c_W} \gamma^\mu \left((Z_N^{4i*} Z_N^{4j} - Z_N^{3i*} Z_N^{3j}) P_L - (Z_N^{4i} Z_N^{4j*} - Z_N^{3i} Z_N^{3j*}) P_R \right)$$

$\mathcal{M} \propto |N_{14}|^2 - |N_{13}|^2$

[Rosiek '95-'09]

Ratio

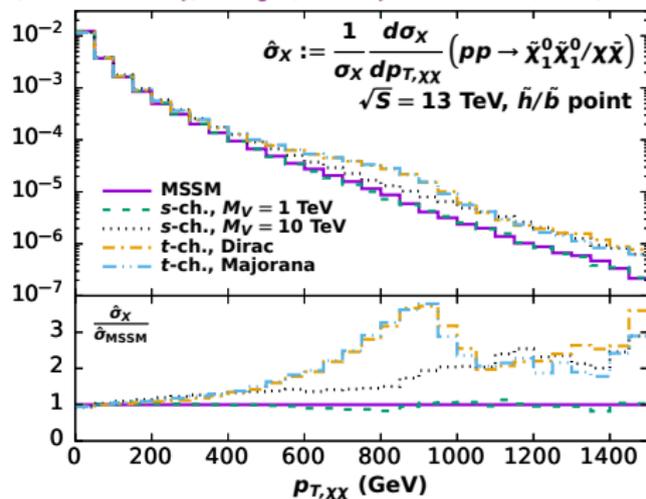
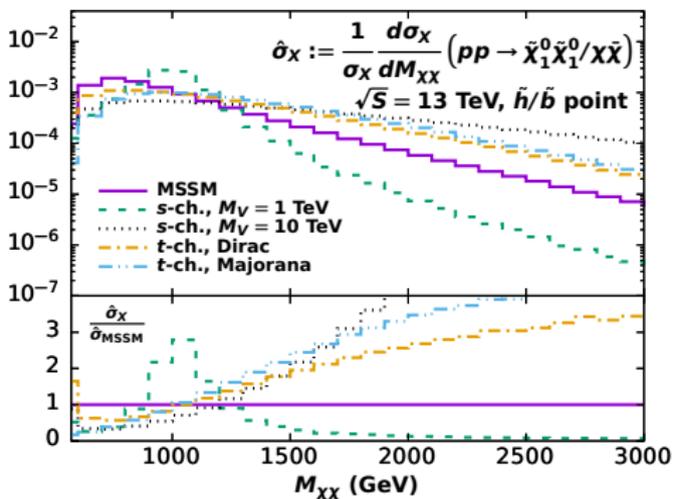
- ▶ Ratio involving s-channel model varies by several orders of magnitude
- ▶ Ratio involving t-channel model almost constant (~ factor 3)
- ▶ Effect at high M_0 and intermediate $m_{1/2}$ due to "higgsino mixing matrix suppression" in MSSM $\Rightarrow \sigma_{\text{MSSM}}$ low

[GC, C. Borshensky, B. Jäger; 1812.08704]



Backup: Distributions: \tilde{h}/\tilde{b} point

[GC, C. Borshensky, B. Jäger; Eur.Phys.J. C79 (2019) no.5, 428]



Analysis at NLO + parton shower accuracy

- ▶ $\tilde{\chi}_1^0$ mainly \tilde{h} iggsino-mix, DM mass ~ 290 GeV, squark masses ~ 1.8 TeV
- ▶ No agreement between simplified models and MSSM for M_{XX} distribution
- ▶ Good agreement with s-channel ($M_V = 1 \text{ TeV}$) for $p_{T,XX}$ distribution

