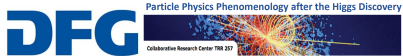


Project C2a

Hadronic Matrix Elements and Exclusive Semileptonic Decays

Thorsten Feldmann and Thomas Mannel

CRC Annual Meeting, 6.-8. October 2020



Work Areas:

- WA 1 Factorization and Light-Cone Distribution Amplitudes
- WA 2 Development of QCD Sum Rules and Related Methods
- WA 3 New Channels and Multi-Hadron Final states
- WA 4 Inclusive Rate from the Sum over Exclusive Channels

People:

- | | |
|-----------------------------|----------|
| • Jao Ji (PD) | C2a |
| • Nicolas Seitz (PhD) | C2a |
| • Marzia Bordone (PD) | C2b |
| • Robin Brüser (PD) | C2b |
| • Alex Khodjamirian (Prof.) | U Siegen |
| • Rusa Mandal (PD) | AvH |
| • Thorsten Feldmann | U Siegen |
| • Thomas Mannel | U Siegen |

Project C2a: Papers and Current Projects

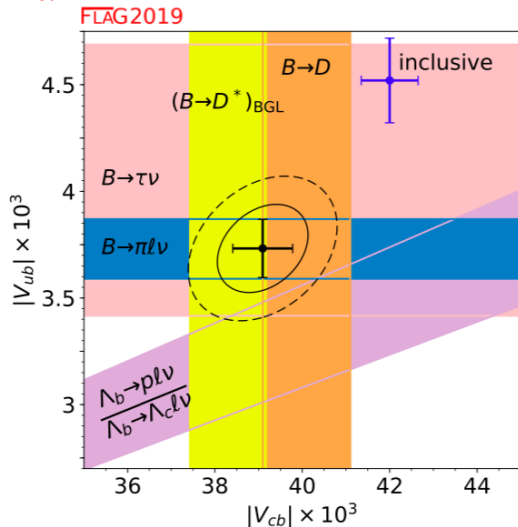
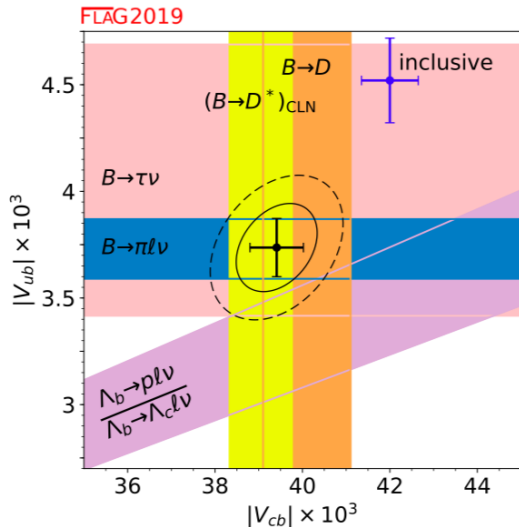
Papers:

- ➊ **M. Bordone**, N. Gubernari, D. van Dyk and M. Jung, “Heavy-Quark Expansion for $\bar{B}_s \rightarrow D_s^{(*)}$ Form Factors and Unitarity Bounds beyond the $SU(3)_F$ Limit,” arXiv:1912.09335 [hep-ph].
- ➋ **R. Brüser**, Z. L. Liu and M. Stahlhofen, “Three-loop soft function for heavy-to-light quark decays,” arXiv:1911.04494 [hep-ph].
- ➌ **M. Bordone**, M. Jung and D. van Dyk, “Theory determination of $\bar{B} \rightarrow D^{(*)} \ell^- \bar{\nu}$ form factors at $\mathcal{O}(1/m_c^2)$,” Eur. Phys. J. C **80** (2020) no.2, 74
- ➍ **A. Khodjamirian, R. Mandal and T. Mannel**, “Inverse moment of the B_s -meson distribution amplitude from QCD sum rule,” [arXiv:2008.03935 [hep-ph]], JHEP in print
- ➎ S. Cheng, **A. Khodjamirian** and A. V. Rusov, “The pion light-cone distribution amplitude from the pion electromagnetic form factor,” [arXiv:2007.05550 [hep-ph]].

Current Projects:

- NLO-improved parametrization of B-meson LCDA (**TF**, van Dyk)
- Sum rule for $B_c \rightarrow J/\psi \ell \nu$ (**TM**, **Bordone**, **Khodjamirian**)
- Sum rule for $B \rightarrow D^{**} \ell \bar{\nu}$ (**TM**, **Khodjamirian**, **Mandal**)

Context: Current tensions in the determination of V_{cb}



Scrutinize the exclusive determination of V_{cb}

- Form factors, including the $1/m$ expansion
- Other effects: QED etc.

Scrutinize the inclusive determination of V_{cb} (part of C1a)

- Contaminations from $b \rightarrow u$ and $b \rightarrow \tau$
- Duality violations

Here: Discussion of the form factors in $\bar{B} \rightarrow D^{(*)} \ell^- \bar{\nu}$

Reminder: From factors for $\bar{B} \rightarrow D^{(*)}$ transitions:

$$\langle D(k) | \bar{c} \gamma^\mu b | \bar{B}(p) \rangle = \left[(p+k)^\mu - \frac{M_B^2 - M_D^2}{q^2} q^\mu \right] f_+^{B \rightarrow D}(q^2) + \frac{M_B^2 - M_D^2}{q^2} q^\mu f_0^{B \rightarrow D}(q^2),$$

$$\langle D(k) | \bar{c} \sigma^{\mu\nu} b | \bar{B}(p) \rangle = \frac{2i}{M_B + M_D} (k^\mu p^\nu - p^\mu k^\nu) f_T(q^2, \mu),$$

$$\langle D^*(k, \eta) | \bar{c} \gamma^\mu b | \bar{B}(p) \rangle = -\epsilon^{\mu\nu\rho\sigma} \eta_\nu^*(k) p_\rho k_\sigma \frac{2V(q^2)}{M_B + M_{D^*}},$$

$$\begin{aligned} \langle D^*(k, \eta) | \bar{c} \gamma^\mu \gamma_5 b | \bar{B}(p) \rangle = i\eta_\nu^* \Big\{ & 2M_{D^*} A_0(q^2) \frac{q^\mu q^\nu}{q^2} + 16 \frac{M_B M_{D^*}^2}{\lambda} A_{12} \left[2p^\mu q^\nu - \frac{M_B^2 - M_{D^*}^2 + q^2}{q^2} q^\mu q^\nu \right] \\ & + (M_B + M_{D^*}) A_1(q^2) \left[g^{\mu\nu} + \frac{2(M_B^2 + M_{D^*}^2 - q^2)}{\lambda} q^\mu q^\nu - \frac{2(M_B^2 - M_{D^*}^2 - q^2)}{\lambda} p^\mu q^\nu \right] \Big\} \end{aligned}$$

$$\begin{aligned} \langle D^*(k, \eta) | \bar{c} \sigma^{\mu\nu} b | \bar{B}(p) \rangle = i\eta_\alpha^* \epsilon^{\mu\nu\rho\sigma} \Big\{ & - \left[\left((p+k)^\rho - \frac{M_B^2 - M_{D^*}^2}{q^2} q^\rho \right) g^{\alpha\sigma} + \frac{2}{q^2} p^\alpha p^\rho k^\sigma \right] T_1(q^2) \\ & - \left(\frac{2}{q^2} p^\alpha p^\rho k^\sigma - \frac{M_B^2 - M_{D^*}^2}{q^2} q^\rho g^{\alpha\sigma} \right) T_2(q^2) + \frac{2}{M_B^2 - M_{D^*}^2} p^\alpha p^\rho k^\sigma T_3(q^2) \Big\}. \end{aligned}$$

Analysis up to $1/m_b, 1/m_c^2, \alpha_s$

Theoretical Machinery

- Heavy Quark Limit
- Fixed Point QCD Sum Rules
- Light-Cone Sum rules

Heavy Quark Limit relations: $q = M_B v - M_{D^{(*)}} v'$, $w = vv'$

- Only a single form factor for $m_c, m_b \rightarrow \infty$
- Corrections are of order $1/m_c$ and $1/m_b$
- Some Corrections are forbidden due to Ademollo Gatto / Luke's theorem
- Calculable hard Gluon Corrections $\alpha_s(m_c m_b)$

Analysis up to $1/m_b, 1/m_c^2, \alpha_s$

Conformal Variable:

$$z = \frac{(\sqrt{w+1} - \sqrt{2})}{(\sqrt{w+1} + \sqrt{2})}$$

- BGL (Boyd, Grinstein, Lebed) parametrization

$$F(z) = \frac{1}{P_F(z)\phi_F(z)} \sum_{n=0}^{\infty} a_n z^n \quad \sum_n a_n^2 \leq 1$$

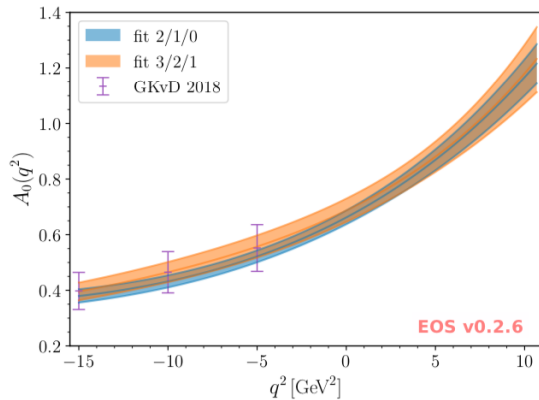
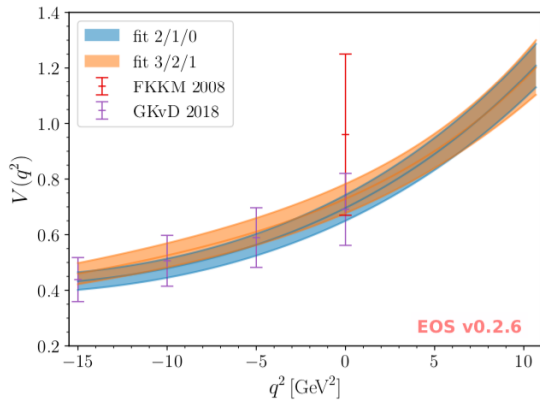
- CLN (Caprini, Lellouch, Neubert): One parameter form:

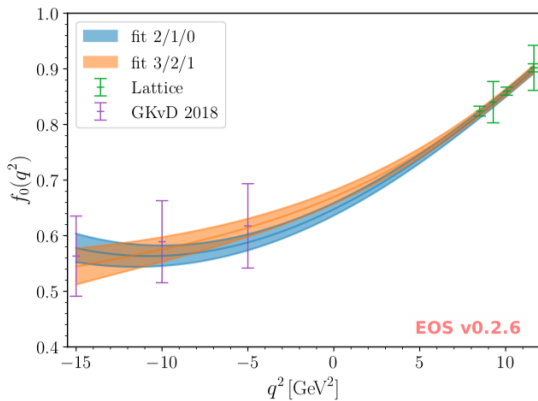
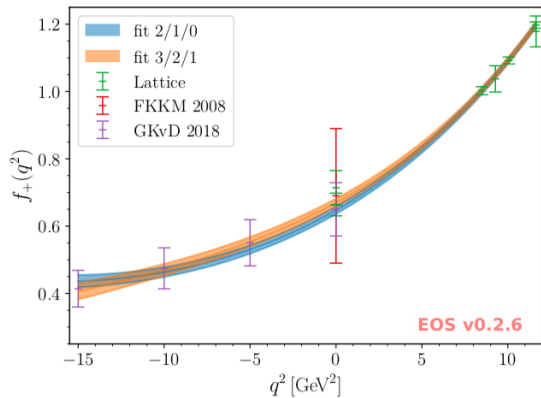
$$F(z) = F(0) [1 - 8\rho^2 z + (53\rho^2 - 15)z^2 + (231\rho^2 - 91)z^3]$$

Fit to theoretical predictions from

- Lattice data from most recent simulations from HPQCD and FNAL/MILC collaborations
- Subleading form factor contributions from three-point sum rules (Neubert, Ligeti Nir 93)
- Light-Cone Sum rules for the form factors (Gubernari, Kokulu, van Dyk 18)

Results





model	scenarios				
	2/1/0	3/2/1	3/2/1	3/2/1	3/2/1
exp. likelihood	—	—	2017	2018	all exp.
$\mathcal{B}(\bar{B}^0 \rightarrow D^+ \{e^-, \mu^-\} \bar{\nu})/ V_{cb} ^2$	12.99 ± 0.35	13.48 ± 0.37	—	—	13.56 ± 0.35
$\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \{e^-, \mu^-\} \bar{\nu})/ V_{cb} ^2$	32.33 ± 1.28	33.16 ± 2.15	31.74 ± 1.46	32.19 ± 1.03	32.00 ± 1.03
correlation	0.34	0.14	—	—	0.10
$ V_{cb} \times 10^3$ from $\bar{B} \rightarrow D \{e^-, \mu^-\} \bar{\nu}$	41.5 ± 1.2	40.7 ± 1.2	—	—	40.6 ± 1.1
$ V_{cb} \times 10^3$ from $\bar{B} \rightarrow D^* \{e^-, \mu^-\} \bar{\nu}$	39.8 ± 1.2	39.3 ± 1.7	40.1 ± 1.3	39.8 ± 1.0	40.0 ± 1.1
	—	—	(39.5 ± 1.9)	(39.0 ± 1.3)	—
$ V_{cb} \times 10^3$ combined incl. corr.	40.7 ± 1.0	40.2 ± 1.0	—	—	40.3 ± 0.8

$$V_{cb}^{\text{excl}} = (40.3 \pm 0.8) \times 10^{-3}$$

Calculation for non-strange B mesons: Brain, Ivanov, Korchemski 2004

Light Cone distribution for the B_s

$$\langle 0 | \bar{s}(tn) \not{n}[tn, 0] \gamma_5 h_v(0) | \bar{B}_s(v) \rangle = i F_{B_s}(\mu) \int_0^\infty dk e^{-ikt} \phi_+^{B_s}(k, \mu),$$

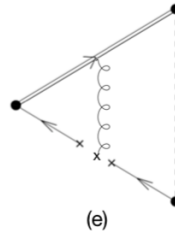
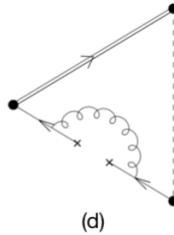
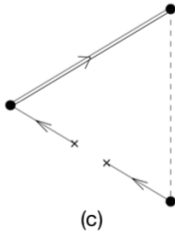
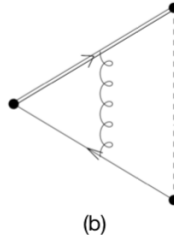
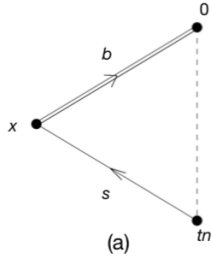
Inverse moment

$$\lambda_{B(s)}^{-1}(\mu) = \int_0^\infty \frac{dk}{k} \phi_+^{B(s)}(k, \mu),$$

Correlation function in HQET $\rightarrow$ sum rule

$$\begin{aligned} \mathcal{P}_{5s}(\omega, t) &= i \int d^4x e^{-i\omega v \cdot x} \langle 0 | T \{ \bar{s}(tn) \not{n}[tn, 0] \gamma_5 h_v(0) \bar{h}_v(x) i \gamma_5 s(x) \} | 0 \rangle \\ &= \frac{\langle 0 | \bar{s}(tn) \not{n}[tn, 0] \gamma_5 h_v(0) | B_s(v) \rangle \langle B_s(v) | \bar{h}_v i \gamma_5 s | 0 \rangle}{\bar{\Lambda}_s - \omega} + \dots = \frac{[F_{B_s}(\mu)]^2}{2(\bar{\Lambda}_s - \omega)} \int_0^\infty dk e^{-ikt} \phi_+^{B_s}(k, \mu) + \dots \end{aligned}$$

Partonic Calculation



Condensate Contributions are singular, need to be modeled by non-local condensates

$$\langle 0 | \bar{q}(x) [x, 0] q(0) | 0 \rangle = \langle \bar{q} q \rangle f_S(x^2) \quad f_S(x^2) = \int_0^\infty \tilde{f}_S(\nu) e^{\nu x^2} d\nu,$$

Local OPE yields

$$\tilde{f}_S(\nu) = \delta(\nu) - \frac{m_0^2}{16} \delta'(\nu) + \dots$$

Two models (Braun, Ivanov, Korchemsky)

- Model 1: $\tilde{f}_S(\nu) = \delta(\nu - m_0^2/16)$
- Model 2: $\tilde{f}_S(\nu) = \frac{\lambda^{p-2}}{\Gamma(p-2)} \nu^{1-p} e^{-\lambda/\nu}, \quad p = 3 + \frac{16\lambda}{m_0^2}$

$$\lambda_{B_s} = 438 \pm 150 \text{ MeV}, \quad \lambda_B = 383 \pm 153 \text{ MeV}$$

$$\frac{\lambda_{B_s}}{\lambda_B} = 1.19 \pm 0.14$$

Snapshot on some Work in Progress

Start from the correlation function

$$T(q, P) = i \int d^4x e^{iqx} \langle B_c(P) | T [\bar{b}(x) \Gamma_1 c(x) \bar{c}(0) \Gamma_2 c(0)] | 0 \rangle$$

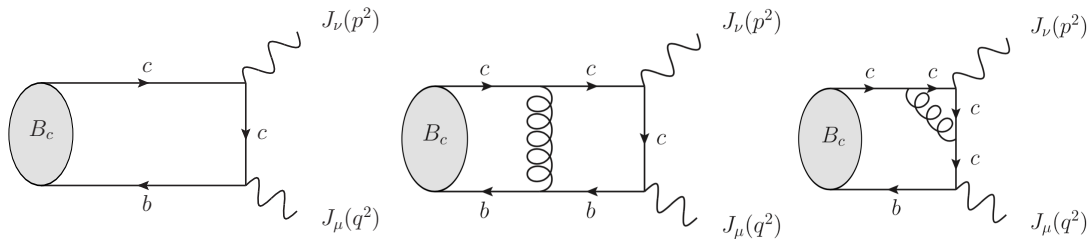
The dispersion relation in the variable p^2 has the spectral function

$$W_{\mu\nu}(q, P) = i \int d^4x e^{iqx} \langle B_c(P) | \bar{b}(x) \gamma_\mu c(x) \bar{c}(0) \gamma_\nu c(0) | 0 \rangle$$

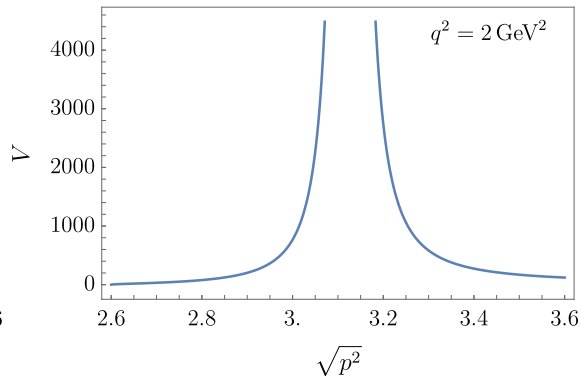
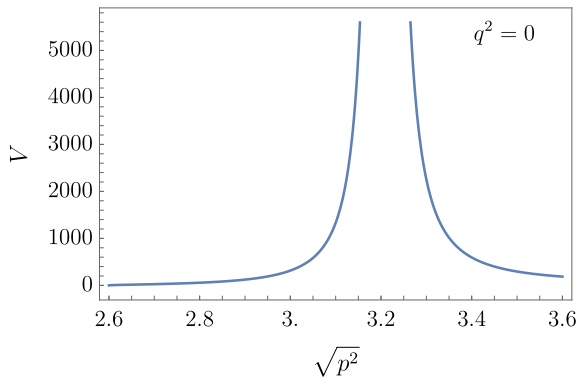
from which we get the sum rule

$$\begin{aligned} W_{\mu\nu}(q, P) &= \int d^4x e^{iqx} \sum_n \langle B_c(P) | \bar{b}(x) \gamma_\mu c(x) | n \rangle \langle n | \bar{c}(0) \gamma_\nu c(0) | 0 \rangle \\ &= \int \widetilde{dp} (2\pi)^4 \delta^4(P - p - q) \langle B_c(P) | \bar{b} \gamma_\mu c | J/\psi(p) \rangle \langle J/\psi(p) | \bar{c} \gamma_\nu c | 0 \rangle + \dots \end{aligned}$$

For heavy b and c quark we can compute perturbatively



Spectral function for $m_c = 1.3$ GeV (Preliminary!!):



To get the $B_c \rightarrow J/\psi$ Form Factors

- Fix the duality threshold
- Numerical Evaluation
- Study the q^2 dependence

WA1:

- generic parametrizations for higher-twist B -meson LCDAs
- implications for phenomenology, in particular for $B \rightarrow \gamma \ell \nu$
- applying the same ideas to shape functions for inclusive decays

[→ C2b]

[→ WA4 + C1a]

WA2:

- Study the new sum rules for B_c form factors
- Study the subleading terms in the new B_c sum rule

WA3:

- semileptonic processes with bottom baryons
- B_c phenomenology

WA4:

- study of non-resonant contributions in $b \rightarrow c$ and $b \rightarrow u$
- decays into excited hadrons