

Project B2b

Operator Analysis of New Physics in Top-Quark Observables

1910.03606

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And ongoing work with S. Bruggisser, D. Van Dyk, R. Schaefer, S. Westhoff

Presenter

PI



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Why the Top-Sector?

- Couples the strongest to the Higgs → Top-sector influences Higgs-sector and vice-versa → part of EW symmetry breaking?
- Many BSM models modify the Top-sector: Light Top partners in SUSY/Composite Higgs, etc.
- The LHC is a Top-factory → wealth of data to be explored.

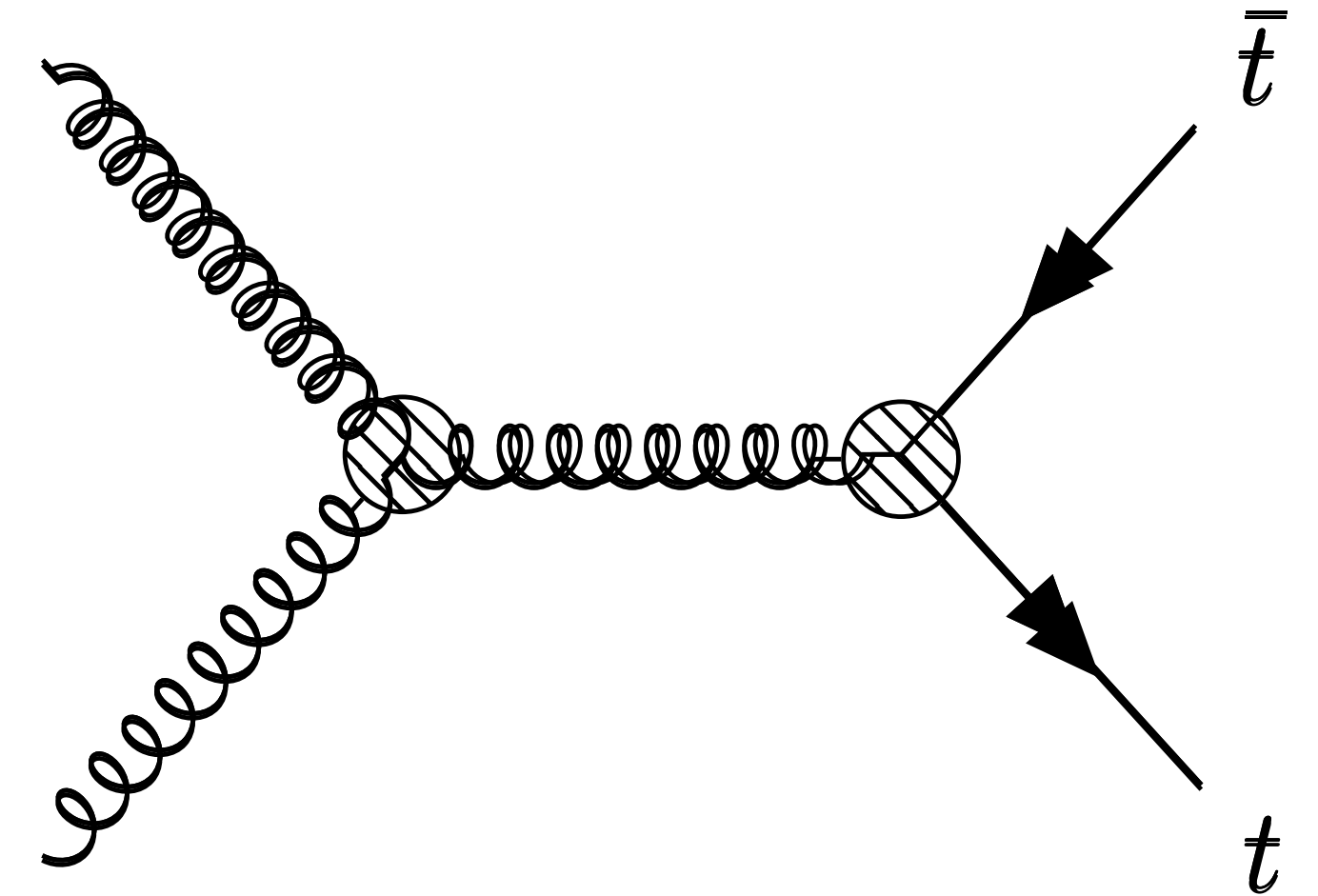
Why now?

- Existing top-fits:
 - TopFitter collaboration: 1512.03360, 1612.02294, 1901.03164
 - SMEFiT collaboration: 1901.05965
- What we can do:
 - Including ATLAS and CMS 13 TeV differential measurements
 - Studying the impact of NLO predictions
 - Detailed study of degeneracies
 - More honest treatment of theoretical uncertainties
 - Systematic study of the impact of uncertainties

The EFT

$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i^{d=6} + \cancel{\sum_j \frac{c_j}{\Lambda^4} \mathcal{O}_j^{d=8}} + \dots$$

- All relevant (including at least one Top) operators up to d=6
- Assumptions:
 - $U(2)_u \times U(2)_d \times U(2)_Q$ flavour symmetry
 - Light quarks considered massless (Yukawa=0)
 - Diagonal CKM
 - CP conserving
 - Warsaw-basis operators or linear combinations thereof
 - Total of 22 operators considered



The Operators

- Four-fermion vs. Two-fermion
- Gauge structure or Up vs. Down
- Chirality of the top
- ...

parameter	$t\bar{t}$	single t	tW	tZ	t decay	$t\bar{t}Z$	$t\bar{t}W$
$C_{Qq}^{1,8}$	Λ^{-2}	—	—	—	—	Λ^{-2}	Λ^{-2}
$C_{Qq}^{3,8}$	Λ^{-2}	$\Lambda^{-4} [\Lambda^{-2}]$	—	$\Lambda^{-4} [\Lambda^{-2}]$	$\Lambda^{-4} [\Lambda^{-2}]$	Λ^{-2}	Λ^{-2}
C_{tu}^8, C_{td}^8	Λ^{-2}	—	—	—	—	Λ^{-2}	—
$C_{Qq}^{1,1}$	$\Lambda^{-4} [\Lambda^{-2}]$	—	—	—	—	$\Lambda^{-4} [\Lambda^{-2}]$	$\Lambda^{-4} [\Lambda^{-2}]$
$C_{Qq}^{3,1}$	$\Lambda^{-4} [\Lambda^{-2}]$	Λ^{-2}	—	Λ^{-2}	Λ^{-2}	$\Lambda^{-4} [\Lambda^{-2}]$	$\Lambda^{-4} [\Lambda^{-2}]$
C_{tu}^1, C_{td}^1	$\Lambda^{-4} [\Lambda^{-2}]$	—	—	—	—	$\Lambda^{-4} [\Lambda^{-2}]$	—
C_{Qu}^8, C_{Qd}^8	Λ^{-2}	—	—	—	—	Λ^{-2}	—
C_{tq}^8	Λ^{-2}	—	—	—	—	Λ^{-2}	Λ^{-2}
C_{Qu}^1, C_{Qd}^1	$\Lambda^{-4} [\Lambda^{-2}]$	—	—	—	—	$\Lambda^{-4} [\Lambda^{-2}]$	—
C_{tq}^1	$\Lambda^{-4} [\Lambda^{-2}]$	—	—	—	—	$\Lambda^{-4} [\Lambda^{-2}]$	$\Lambda^{-4} [\Lambda^{-2}]$
$C_{\phi Q}^-$	—	—	—	Λ^{-2}	—	Λ^{-2}	—
$C_{\phi Q}^3$	—	Λ^{-2}	Λ^{-2}	Λ^{-2}	Λ^{-2}	Λ^{-2}	—
$C_{\phi t}$	—	—	—	Λ^{-2}	—	Λ^{-2}	—
$C_{\phi tb}$	—	Λ^{-4}	Λ^{-4}	Λ^{-4}	Λ^{-4}	—	—
C_{tZ}	—	—	—	Λ^{-2}	—	Λ^{-2}	—
C_{tW}	—	Λ^{-2}	Λ^{-2}	Λ^{-2}	Λ^{-2}	—	—
C_{bW}	—	Λ^{-4}	Λ^{-4}	Λ^{-4}	Λ^{-4}	—	—
C_{tG}	Λ^{-2}	$[\Lambda^{-2}]$	Λ^{-2}	—	$[\Lambda^{-2}]$	Λ^{-2}	Λ^{-2}

Event Kinematics

Or Rates vs. Tails

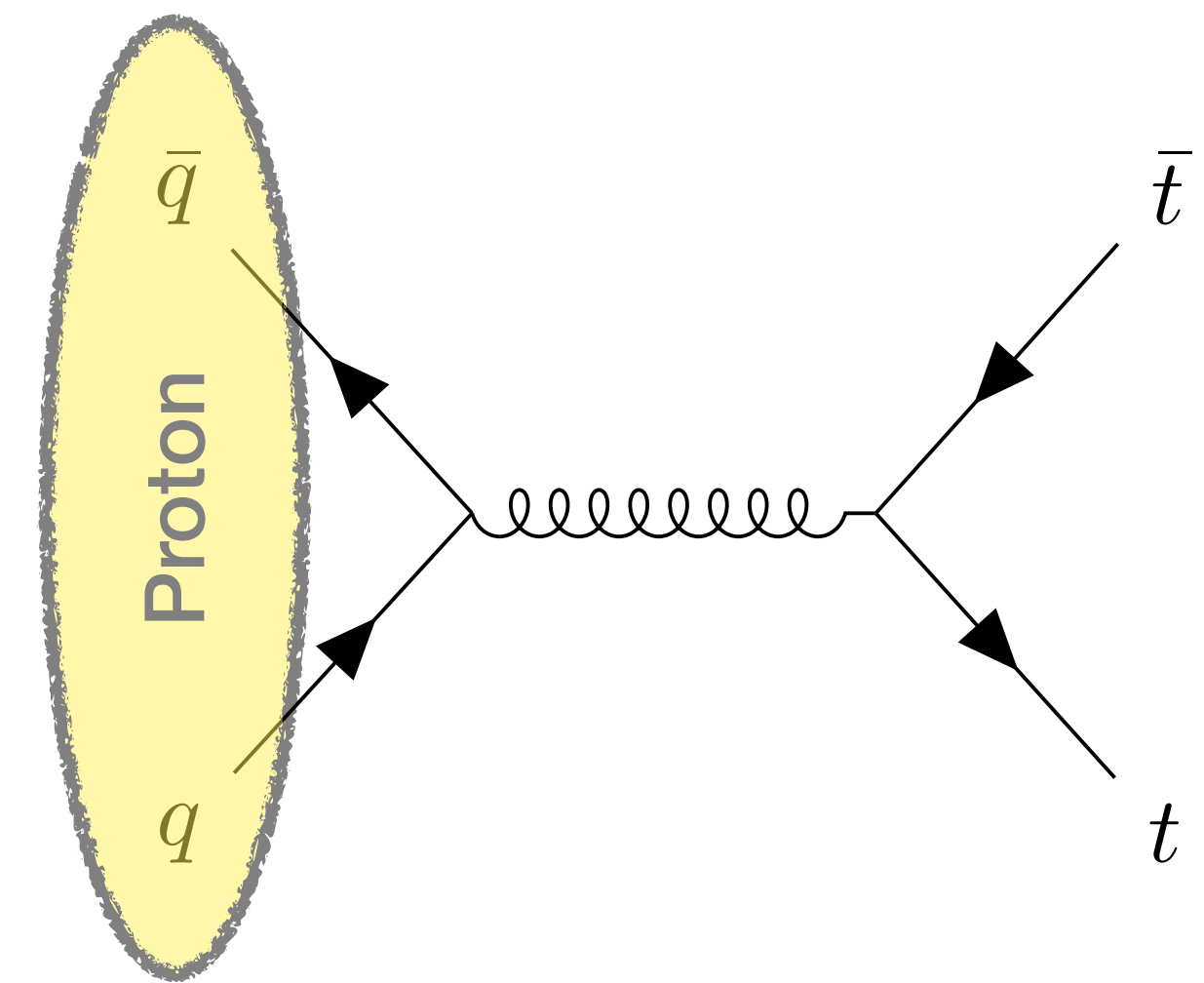
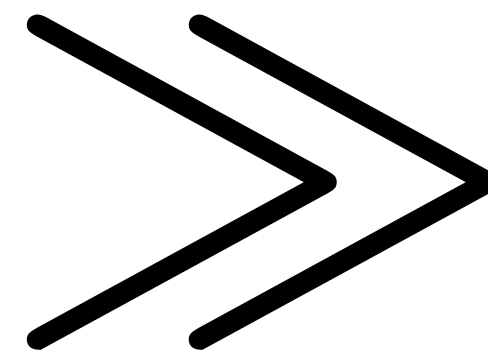
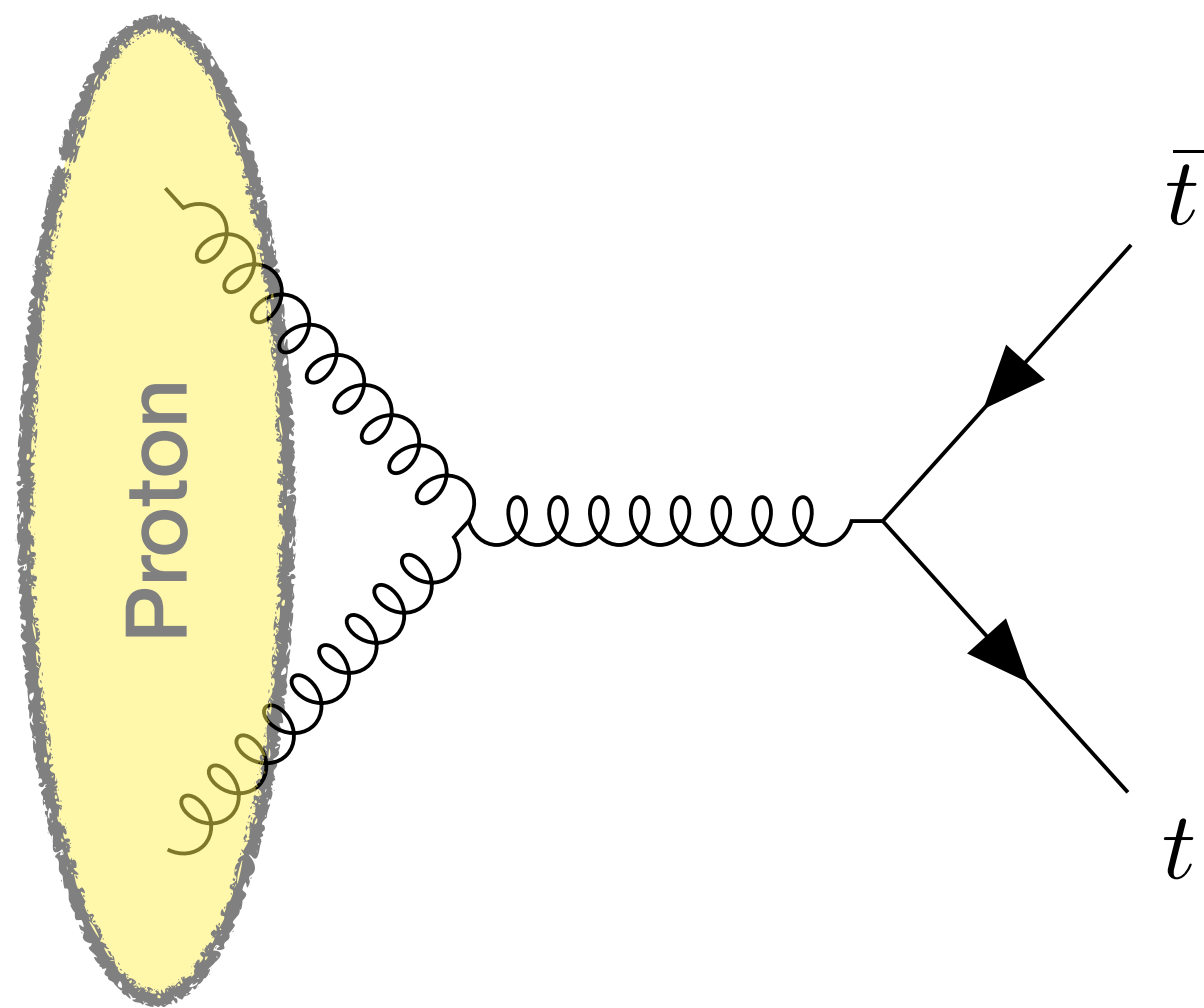
$$O_{tG} = C_{tG}(\bar{q}_3 \sigma^{\mu\nu} T^A u_3) \tilde{\phi} G_{\mu\nu}^A$$

$$O_{tu}^8 = C_{tu}^8(\bar{u}_3 \gamma^\mu T^A u_3)(\bar{u}_i \gamma_\mu T^A u_i)$$

$$\frac{1}{\sigma} \frac{d\sigma}{dm_{t\bar{t}}} \approx \frac{\sigma_{SM}(m_{t\bar{t}})}{\sigma_{SM}(2m_t)} \left(1 + \mathcal{O}(1) \frac{C_{tG}}{\Lambda^2} + \mathcal{O}(m_{t\bar{t}}^2) \frac{|C_{tG}|^2}{\Lambda^4} + \mathcal{O}(m_{t\bar{t}}^2) \frac{C_{tu}^8}{\Lambda^2} + \mathcal{O}(m_{t\bar{t}}^4) \frac{|C_{tu}^8|^2}{\Lambda^4} \right)$$

Very small

Very large
For high bins



Event Kinematics

Or Rates vs. Tails

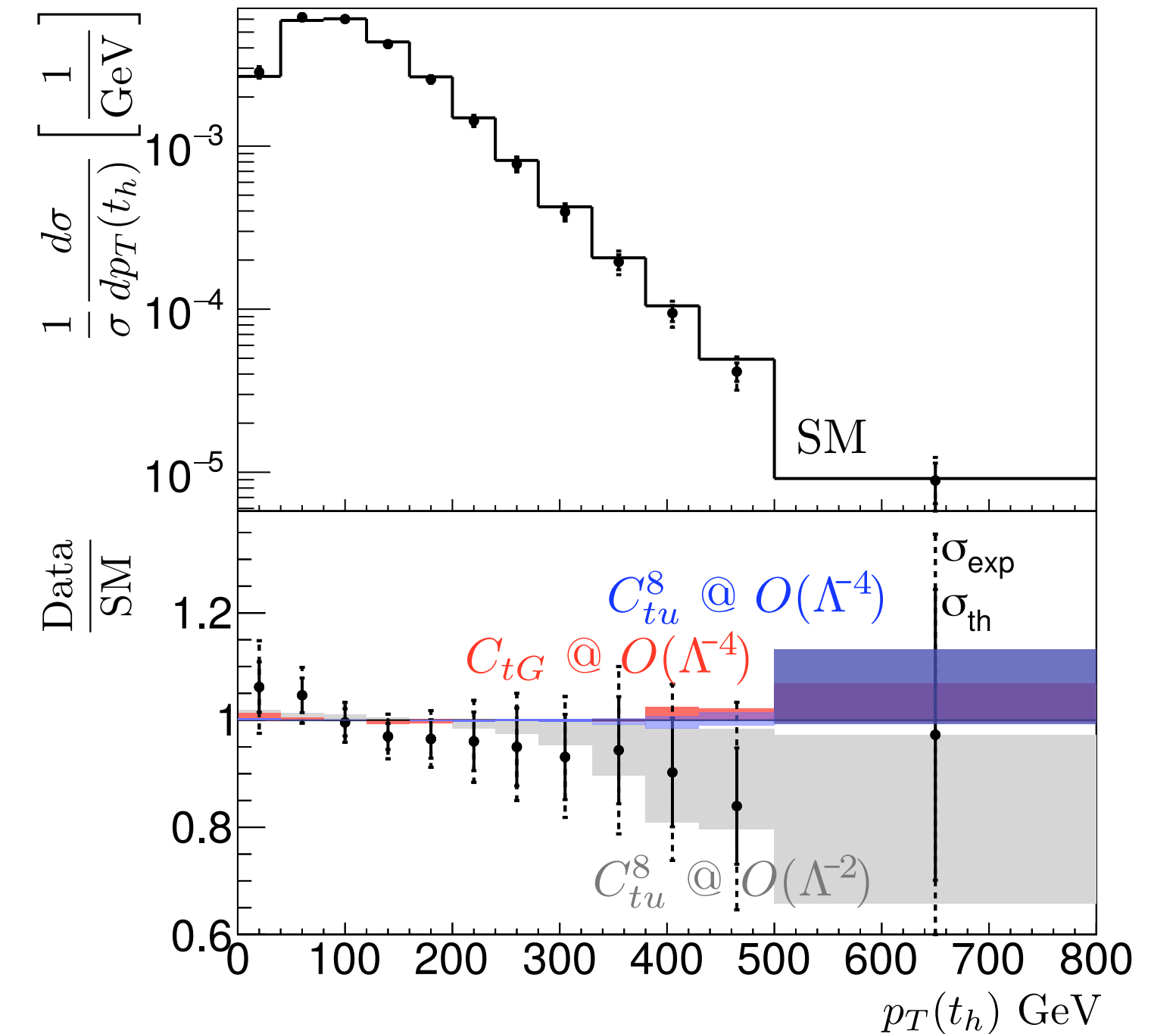
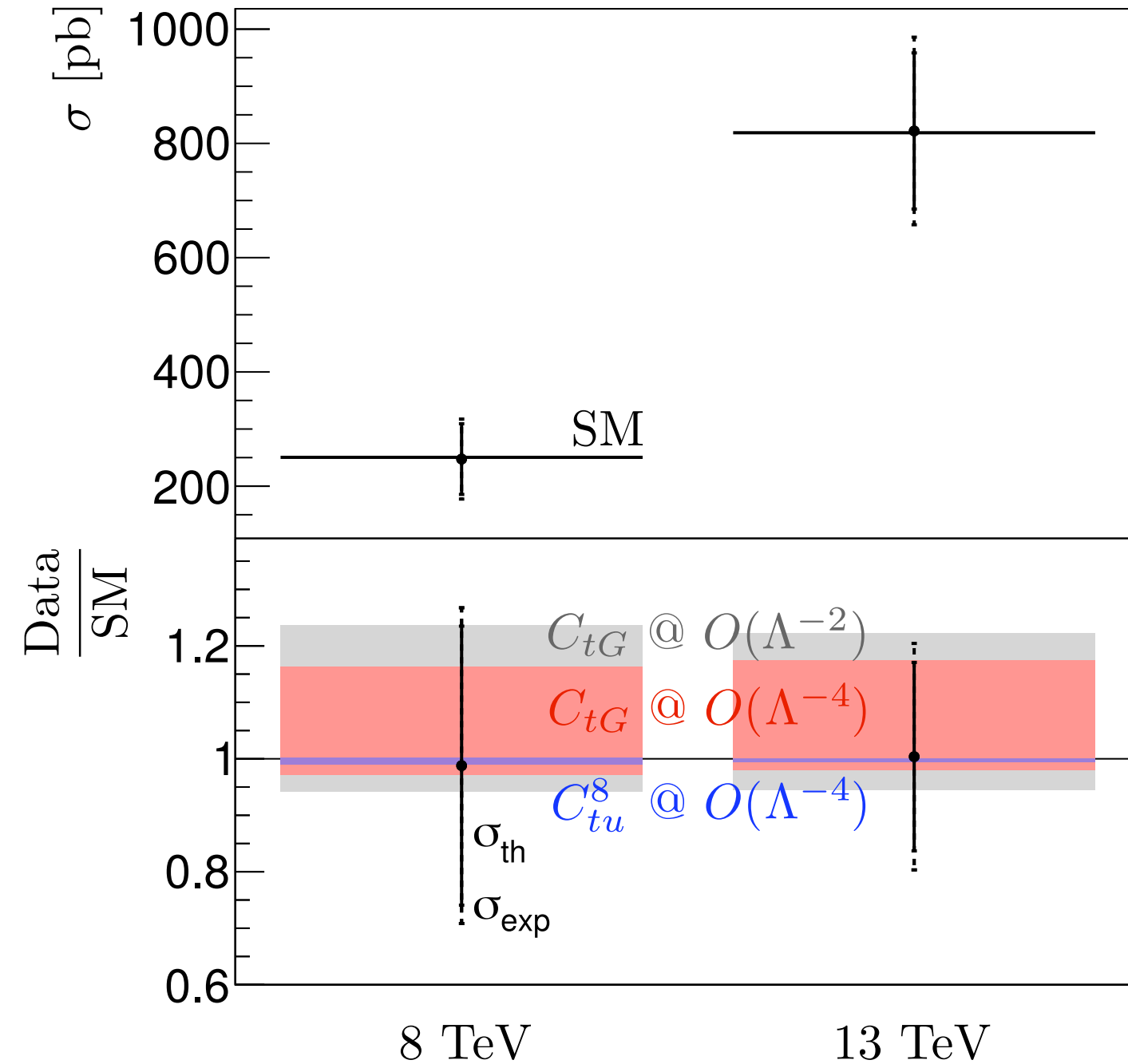
$$O_{tG} = C_{tG}(\bar{q}_3\sigma^{\mu\nu}T^A u_3)\tilde{\phi}G_{\mu\nu}^A$$

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Very small

Very large
For high bins



Up vs. Down

And Gauge Structure

$$O_{Qq}^{1,8} = C_{Qq}^{1,8} (\bar{q}_3 \gamma^\mu T^A q_3) (\bar{q}_i \gamma_\mu T^A q_i)$$

$$O_{Qq}^{3,8} = C_{Qq}^{3,8} (\bar{q}_3 \gamma^\mu T^A \tau^I q_3) (\bar{q}_i \gamma_\mu T^A \tau^I q_i)$$

$$r(x) = \frac{f_u(x) f_{\bar{u}}(s/xS)}{f_d(x) f_{\bar{d}}(s/xS)}$$

$$\sigma_{t\bar{t}} = \sigma_{t\bar{t}}^{SM} + \sigma_{t\bar{t}}^{NP} \left[3 C_{Qq}^{1,8} + 1 C_{Qq}^{3,8} \right]$$

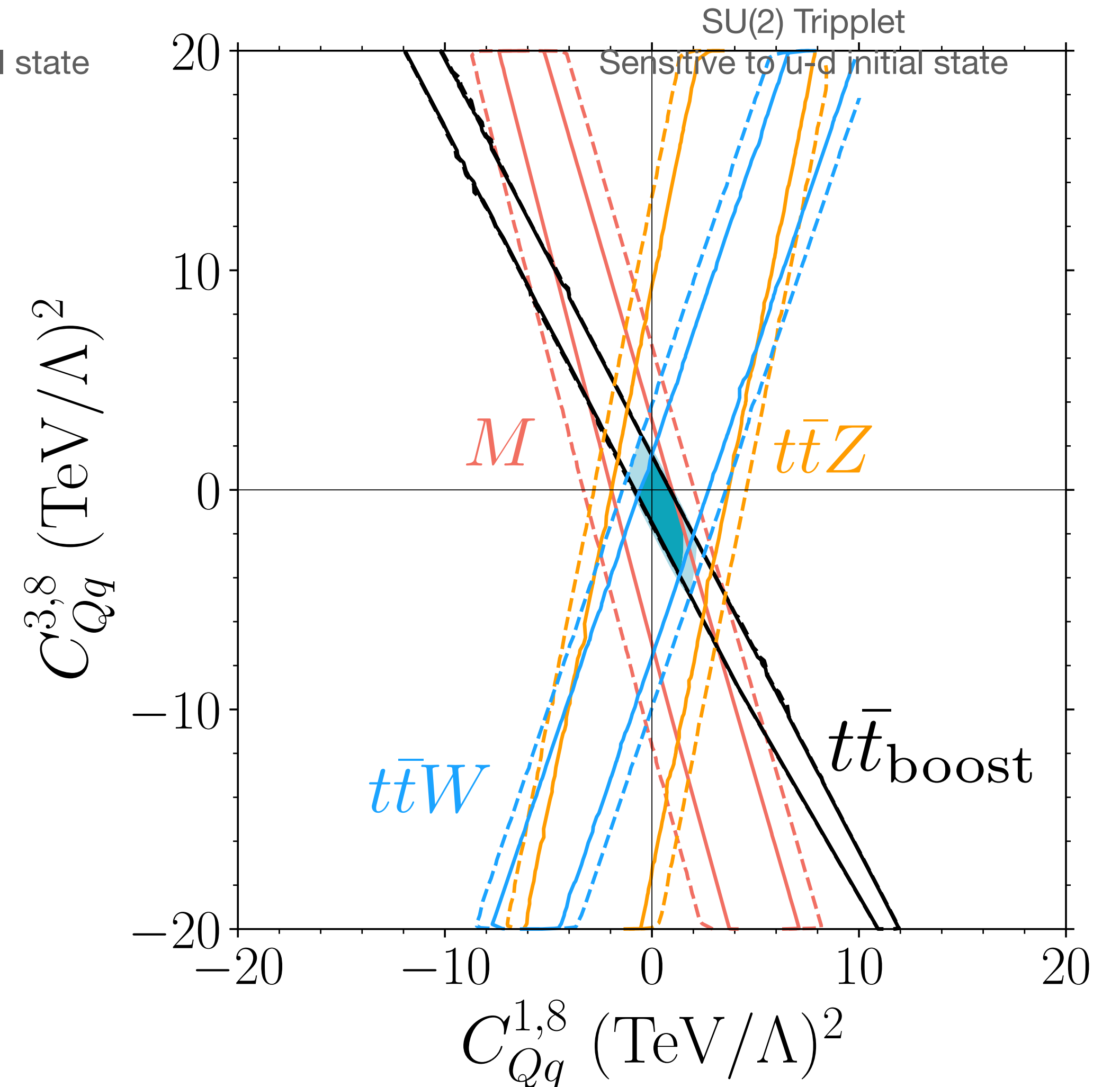
Probe different $r(x)$ in boosted regime

Include $t\bar{t}Z$ and $t\bar{t}W$ measurements

Valence quark maximum

SU(2) Singlet
Sensitive to u+d initial state

SU(2) Triplet
Sensitive to u-d initial state



Chirality

Or Left vs. Right

$$O_{tq}^8 = C_{tq}^8 (\bar{u}_3 \gamma^\mu T^A u_3) (\bar{q}_i \gamma_\mu T^A q_i)$$

Right handed Top
in final state

$$O_{Qq}^{1,8} = C_{Qq}^{1,8} (\bar{q}_3 \gamma^\mu T^A q_3) (\bar{q}_i \gamma_\mu T^A q_i)$$

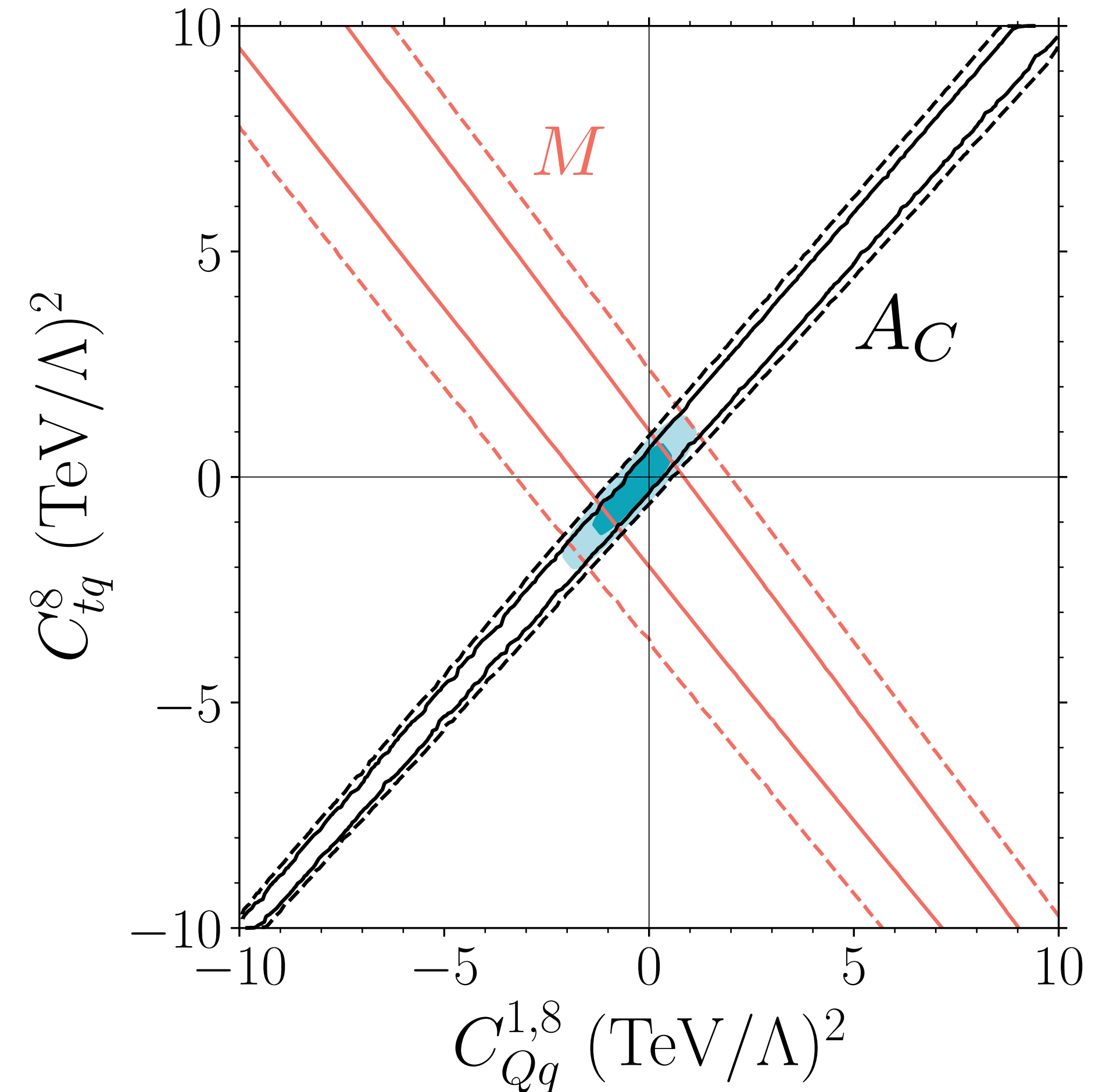
Left handed Top
in final state

$$\sigma_{t\bar{t}} = \sigma_{t\bar{t}}^{SM} + \sigma_{t\bar{t}}^{NP,int} \left(\frac{C_{tq}^8}{\Lambda^2} + \frac{C_{Qq}^{1,8}}{\Lambda^2} \right)$$

$$A_C = \frac{\sigma(\Delta|y| > 0) - \sigma(\Delta|y| < 0)}{\sigma(\Delta|y| > 0) + \sigma(\Delta|y| < 0)} \quad \text{with} \quad \Delta|y| = |y_t| - |y_{\bar{t}}|$$

$$A_C = \frac{\sigma_{t\bar{t}}^{SM,A} + \sigma_{t\bar{t}}^{NP,A} \left(\frac{C_{Qq}^{1,8}}{\Lambda^2} - \frac{C_{tq}^8}{\Lambda^2} \right)}{\sigma_{t\bar{t}}}$$

NLO Corrections can also help to resolve this degeneracy

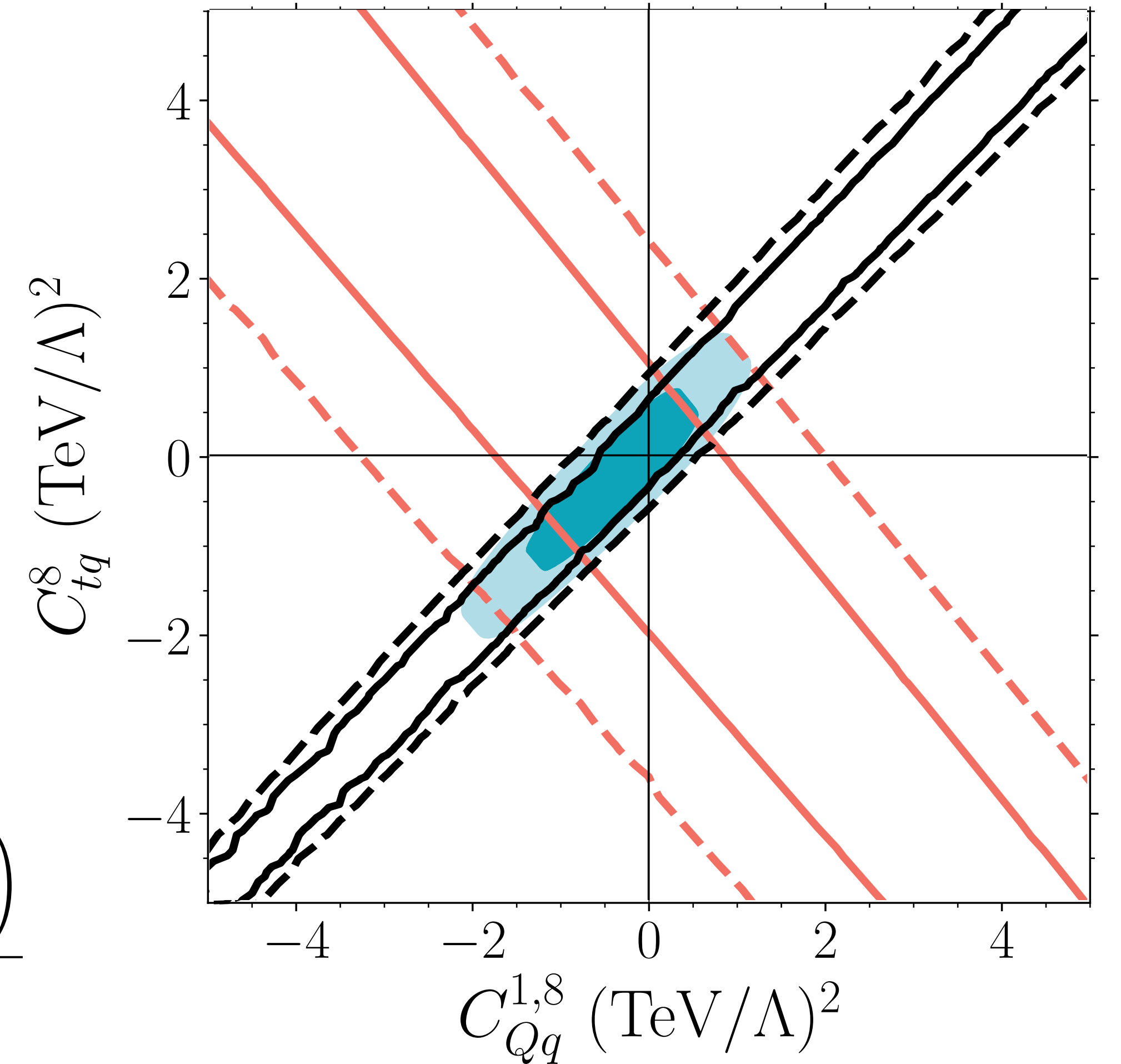


Quadratic Terms

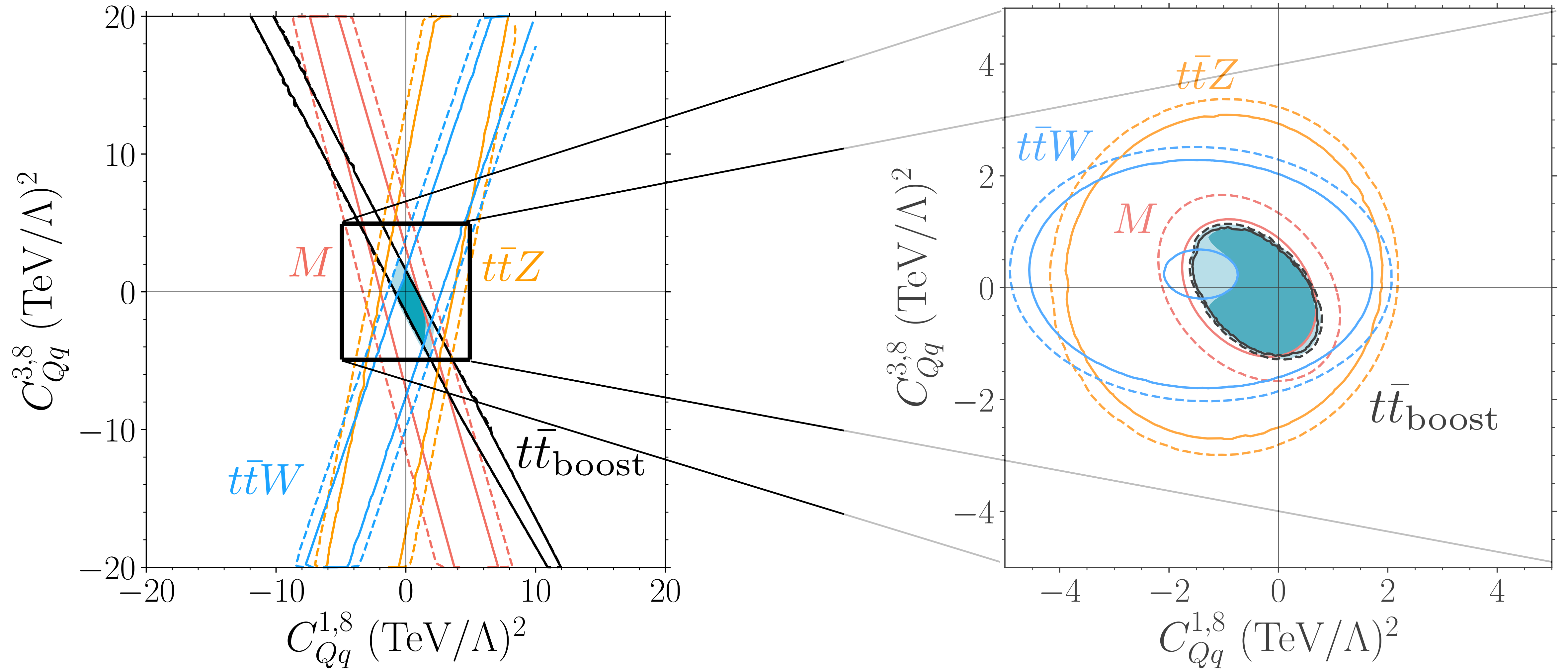
$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \dots$$

$$\begin{aligned} \sigma_{t\bar{t}} = & \sigma_{t\bar{t}}^{SM} + \sigma_{t\bar{t}}^{NP,int} \left(C_{tq}^8 + C_{Qq}^{1,8} \right) \\ & + \sigma_{t\bar{t}}^{NP,quad,1} \left(\left(C_{tq}^8 \right)^2 + \left(C_{Qq}^{1,8} \right)^2 \right) + \dots \end{aligned}$$

$$A_C = \frac{\sigma_{t\bar{t}}^{SM,A} + \sigma_{t\bar{t}}^{NP,A} \left(C_{Qq}^{1,8} - C_{tq}^8 \right) + \sigma_{t\bar{t}}^{NP^2,A} \left(\left(C_{Qq}^{1,8} \right)^2 - \left(C_{tq}^8 \right)^2 \right)}{\sigma_{t\bar{t}}}$$

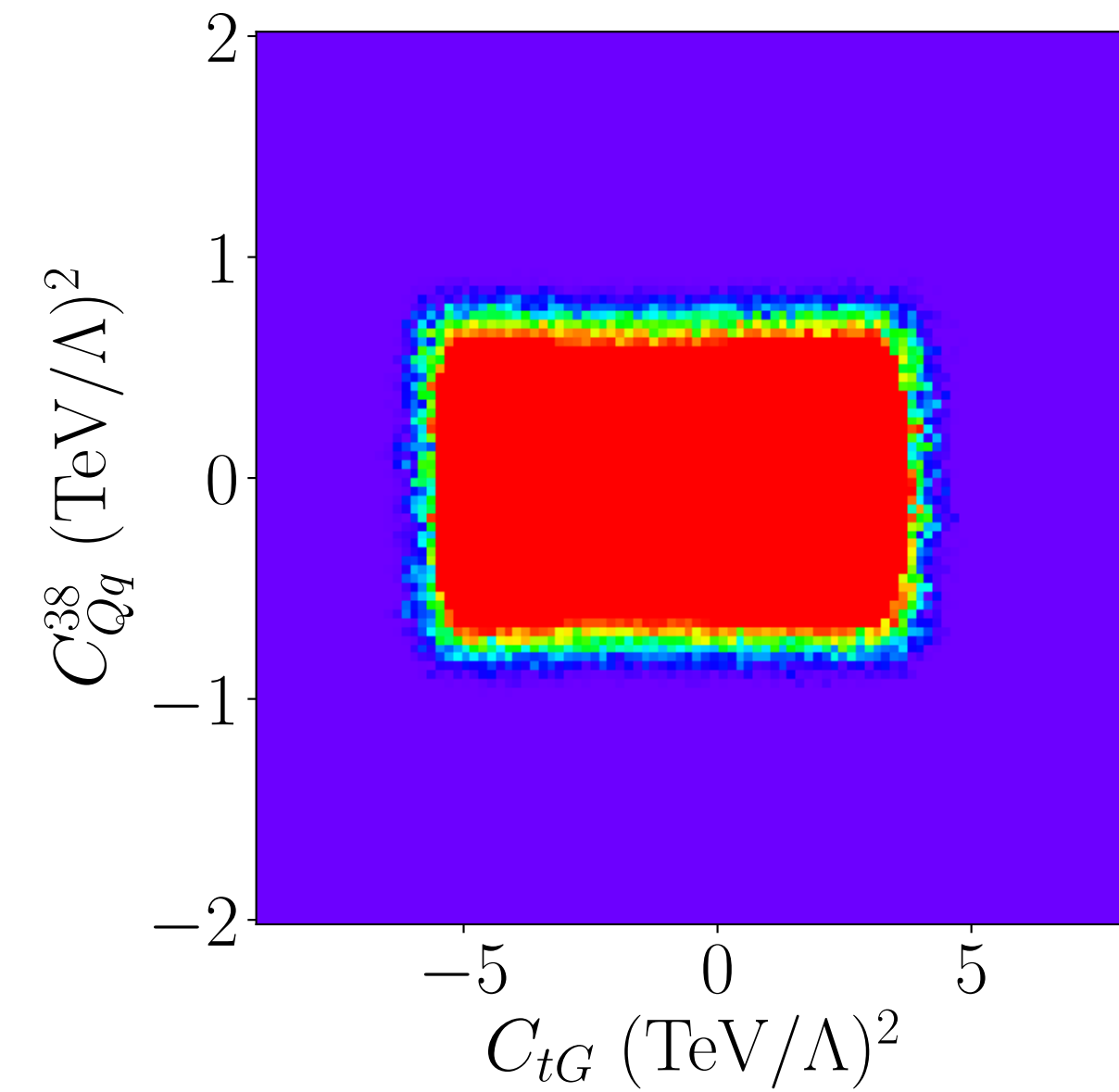


Quadratic Terms

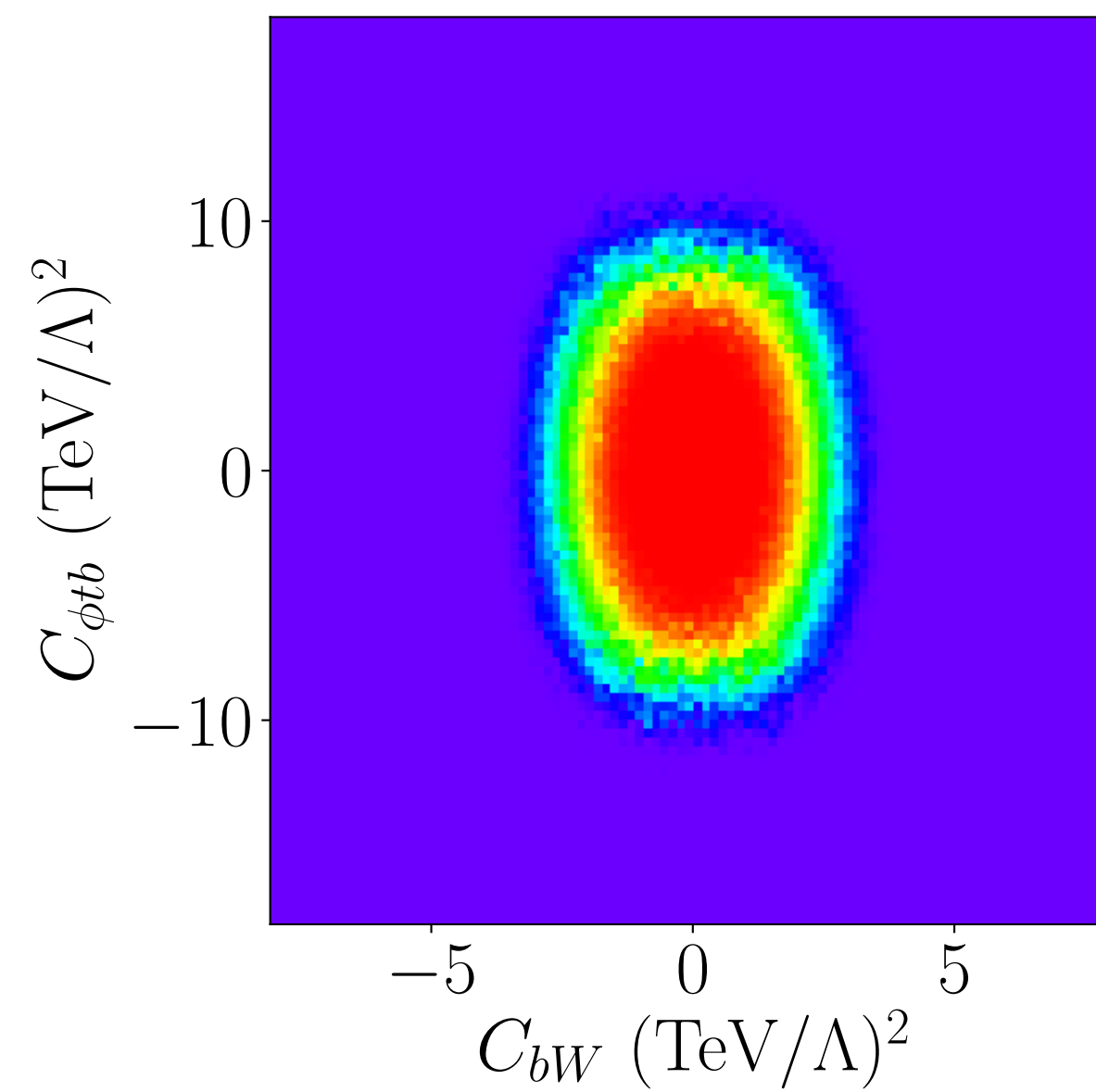


Single Top

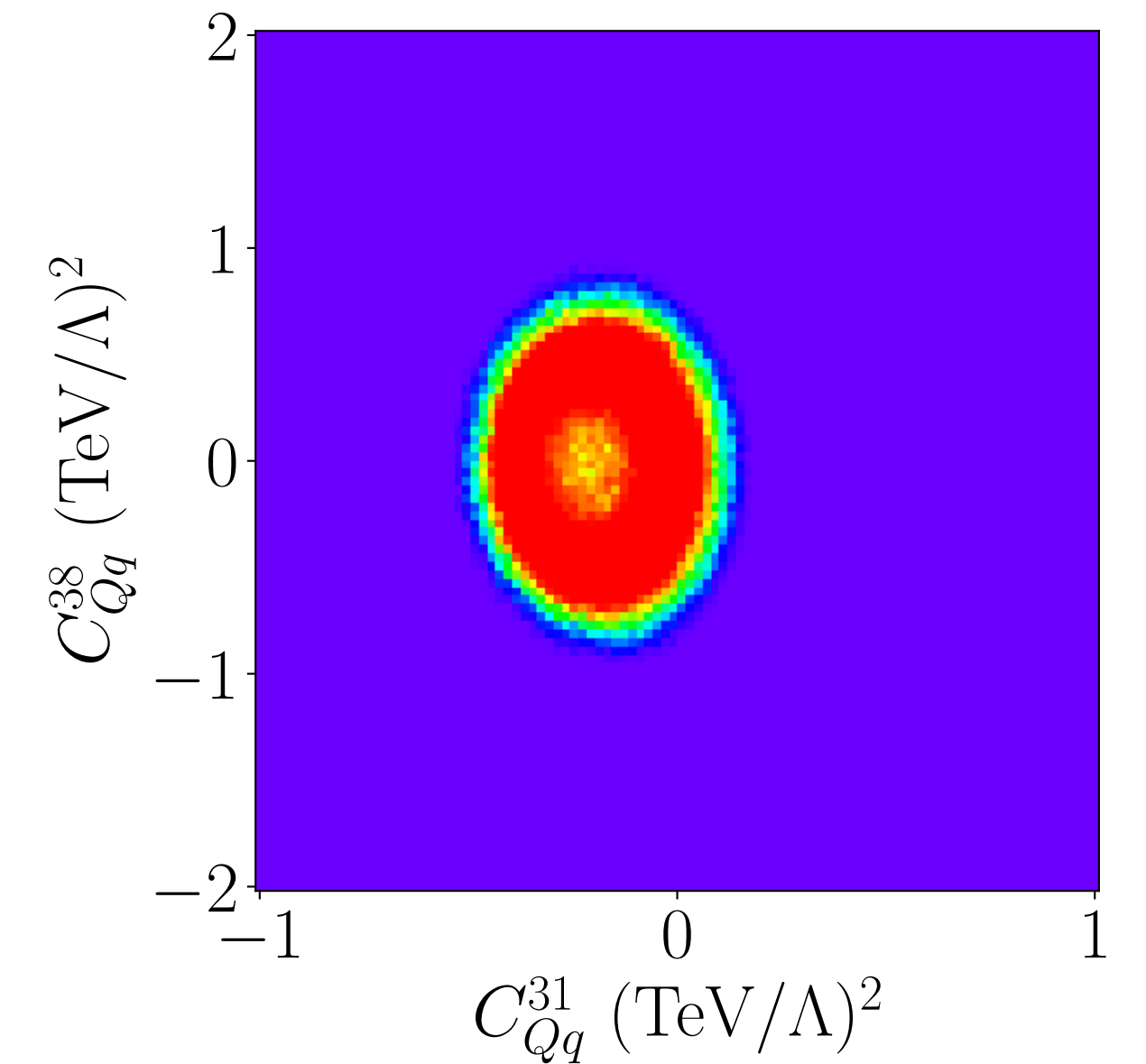
Some aspects



- Constrained by independent measurements



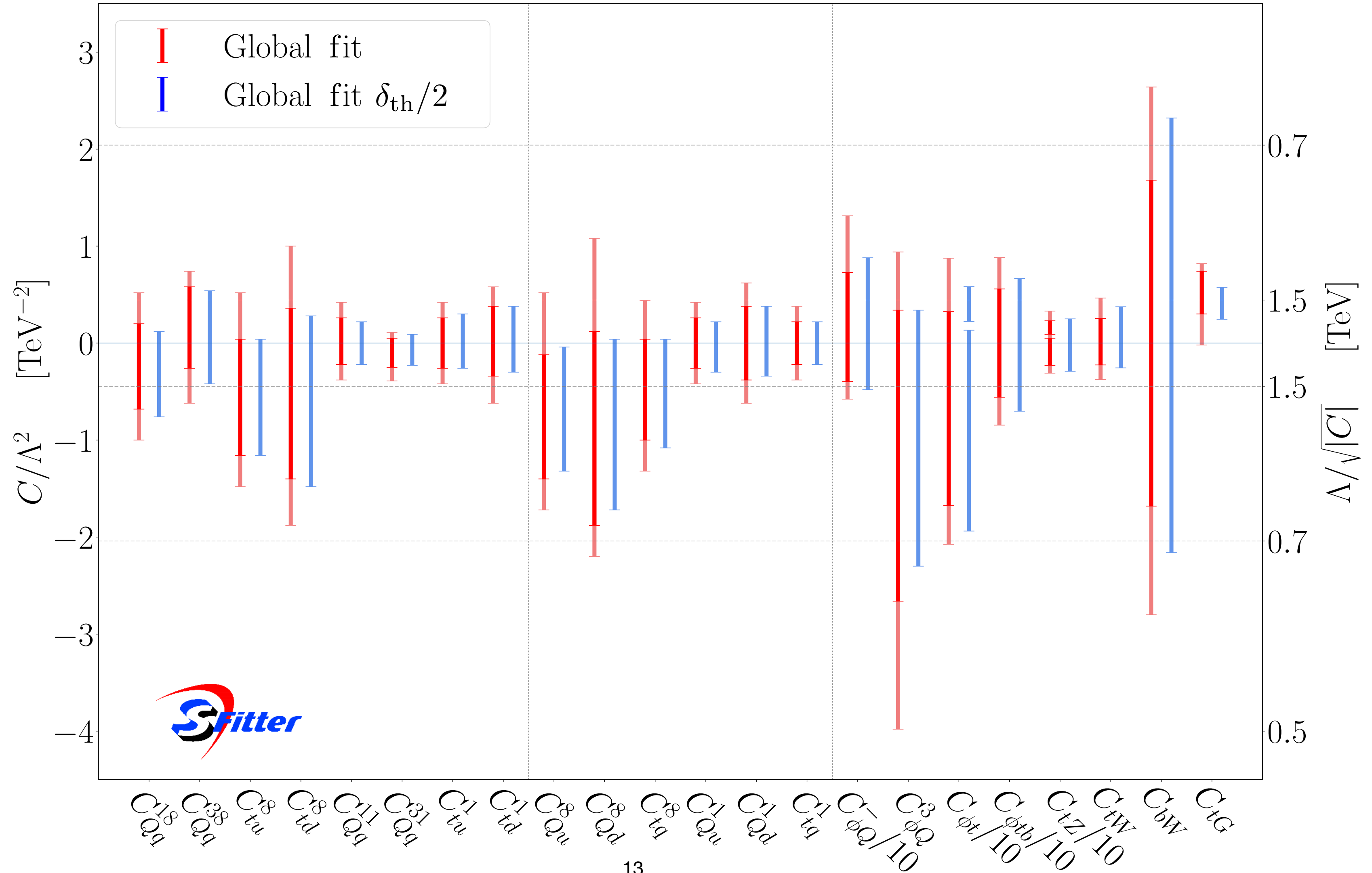
- Both contribute quadratically



- One linear and one quadratic contribution

Results

Run II, ATLAS+CMS, 68% and 95% C.L.



New Directions

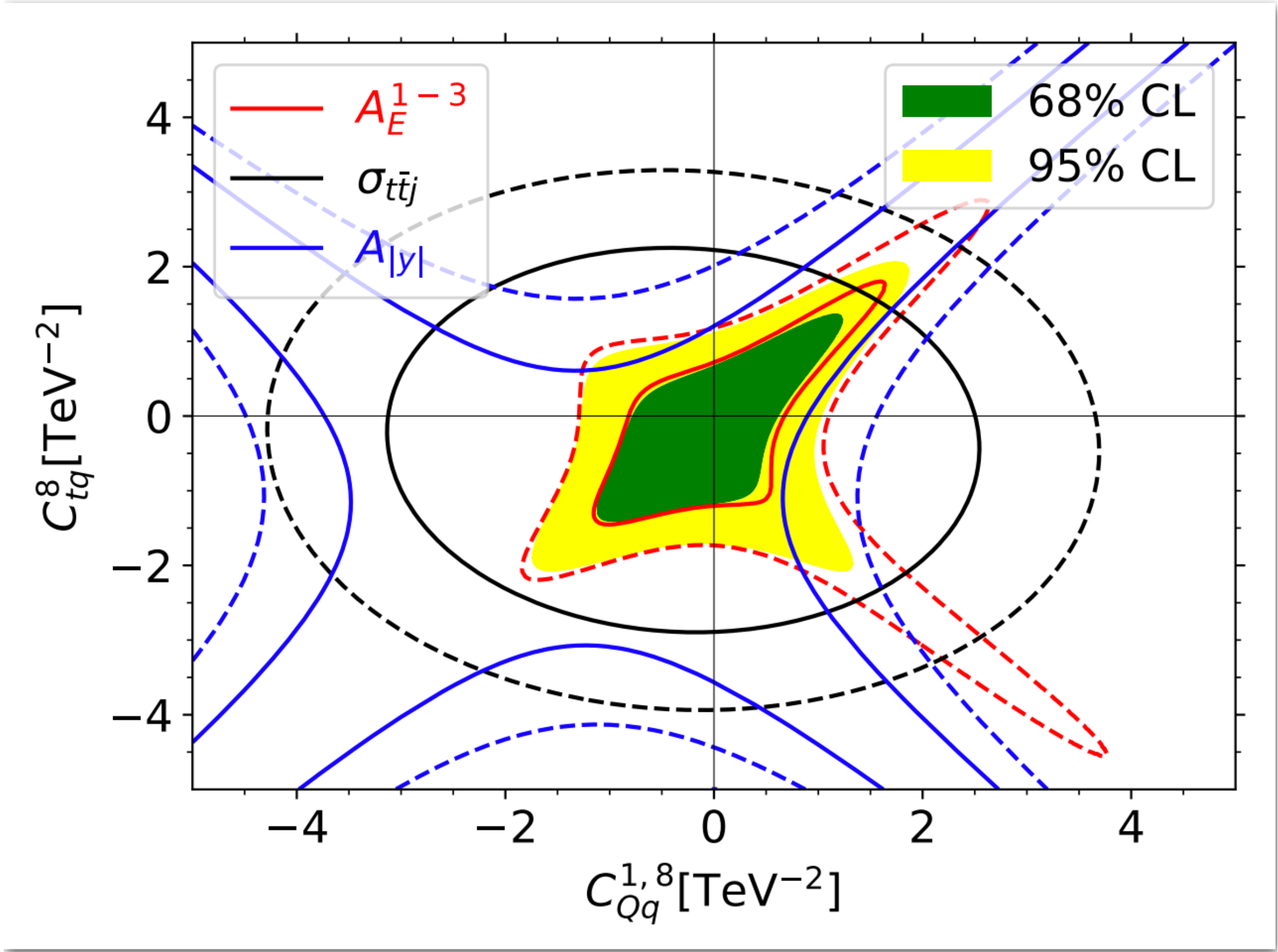
New Observables (e.g. Energy asymmetry)

$$A_E(\theta_j) = \frac{\sigma_{t\bar{t}j}(\theta_j, \Delta E > 0) - \sigma_{t\bar{t}j}(\theta_j, \Delta E < 0)}{\sigma_{t\bar{t}j}(\theta_j, \Delta E > 0) + \sigma_{t\bar{t}j}(\theta_j, \Delta E < 0)}$$

$$\Delta E = E_t - E_{\bar{t}}$$

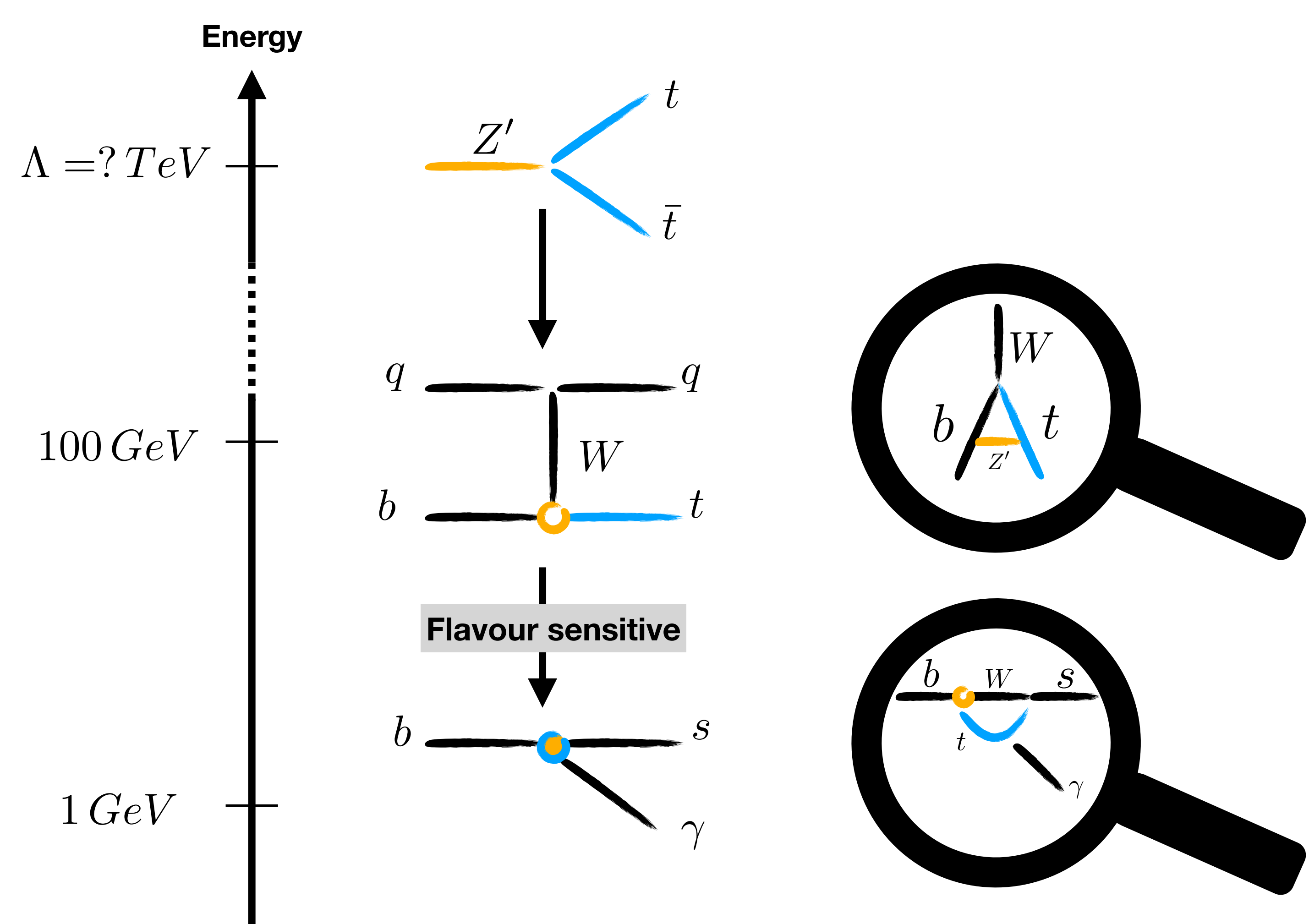
	$A_{ y }$	A_E
Order	NLO	LO
Process	tt	ttj

Resolve new directions in parameter space!



New Directions

Flavour (Work in Progress)

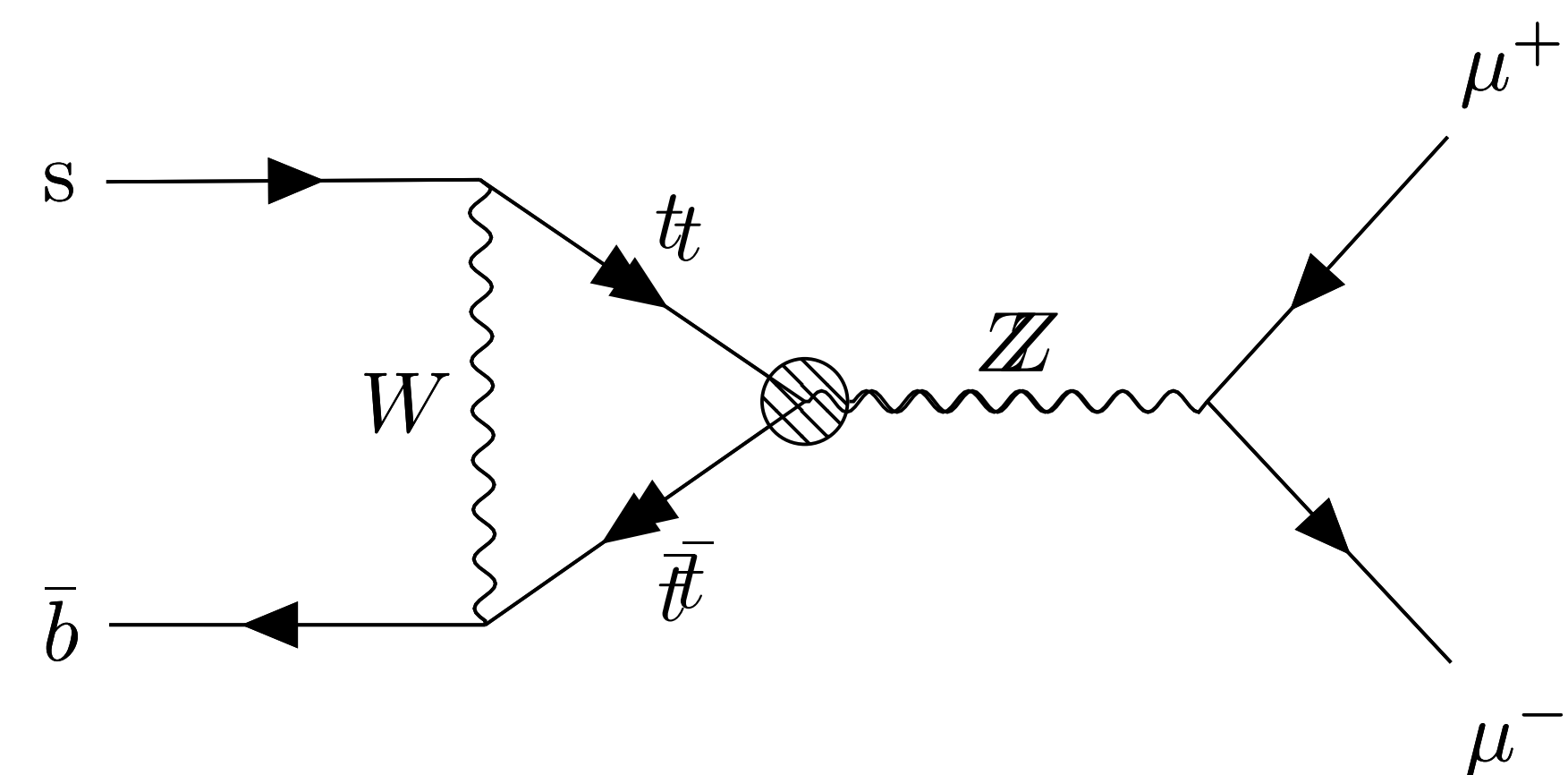
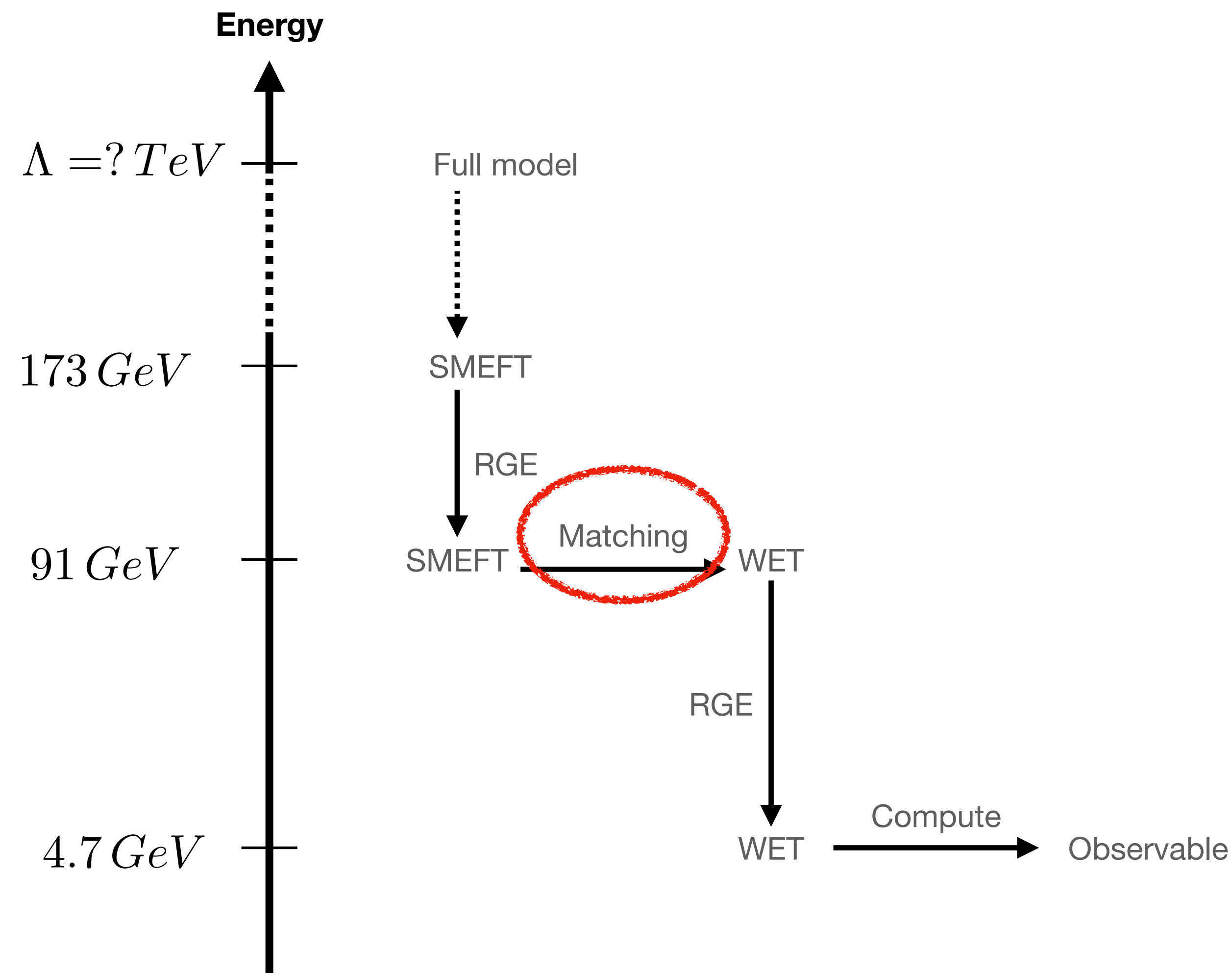


$$\dagger \mathcal{O}_{uW}^{ij} = \left(\bar{q}^i \sigma^{\mu\nu} u^j \right) \tau^I \tilde{\phi} W_{\mu\nu}^I$$

MFV	C_{uW}^{ij}
ii	0
33	$(a + y_t^2 b) y_t$
kj	0
$k3$	$c y_b^2 y_t V_{kb} V_{tb}^*$

New Directions

Flavour (Work in Progress)

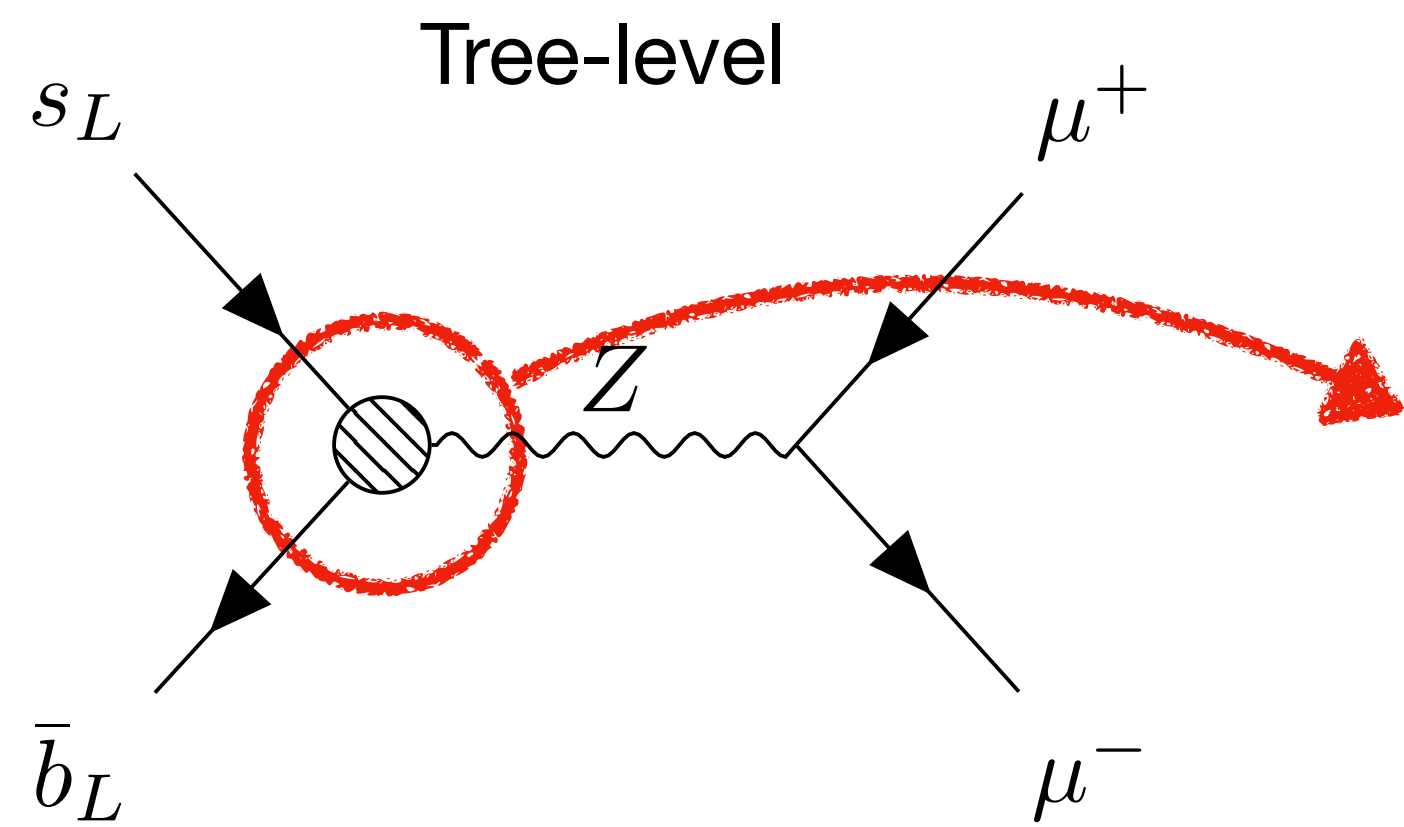
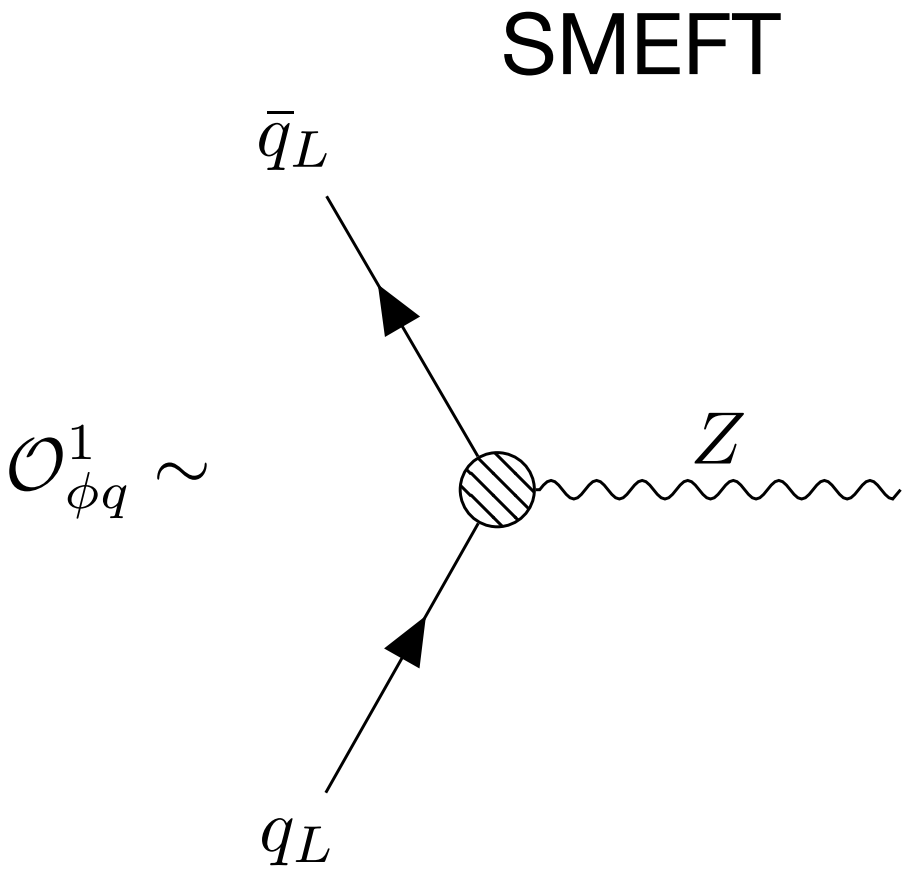
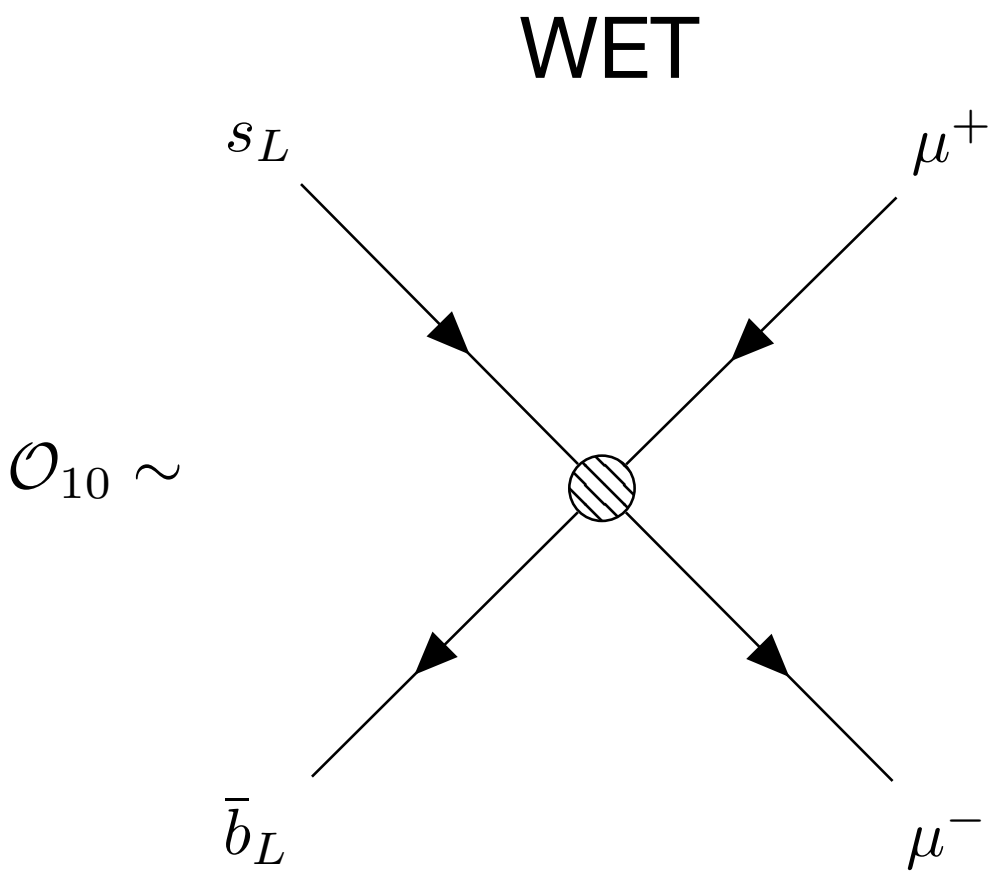


$$[\mathcal{O}_{\phi q}^1]_{ij} = \left(\phi^\dagger i \overleftrightarrow{D}_\mu \phi \right) (\bar{q}_i \gamma^\mu q_j)$$

MFV	$\mathcal{C}_{\phi q}^1$
ii	a
33	$(a + y_t^2 b) y_t$

New Directions

Flavour (Work in Progress)

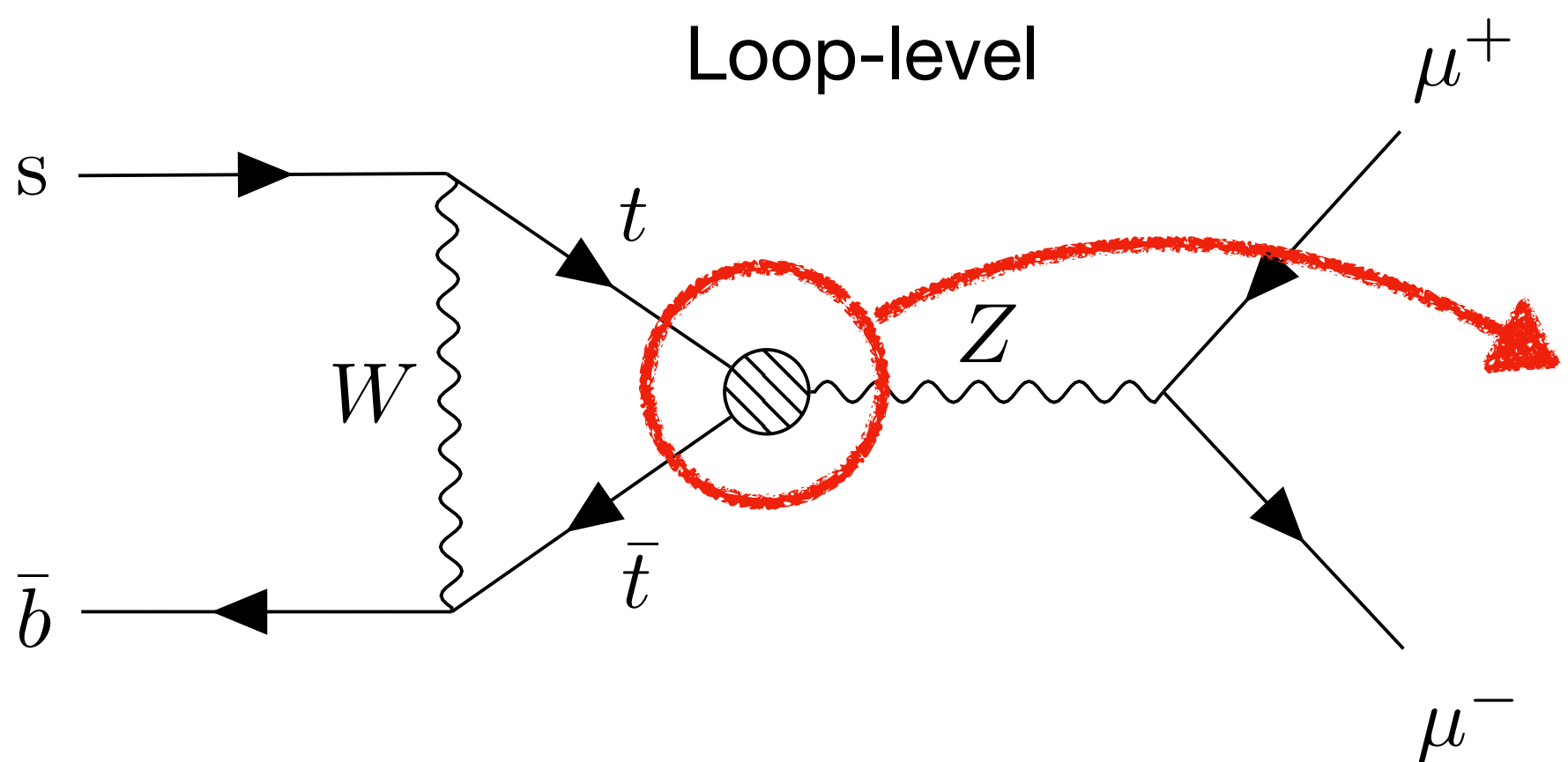


$$\sim [\mathcal{C}_{\phi q}^1]_{ij} V_{is}^* V_{jb} (\bar{s}_L \gamma_\mu b_L) Z^\mu$$

$\mathcal{C}_{10}$

b

$$\mathcal{C}_{\phi q}^1 \begin{cases} \text{Tree-level} \\ \text{Loop-level} \end{cases} \begin{matrix} b \\ a + y_t^2 b \end{matrix}$$

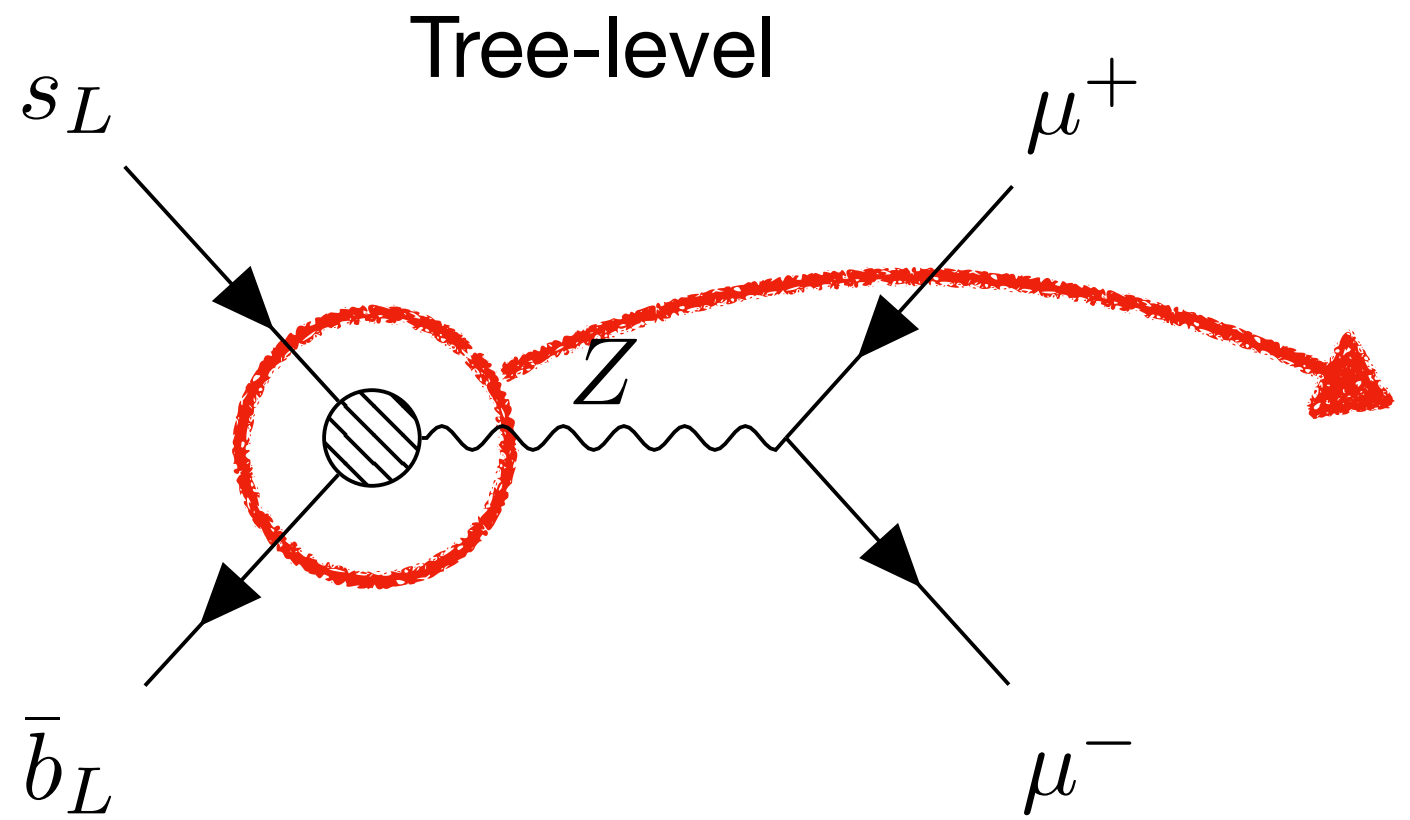
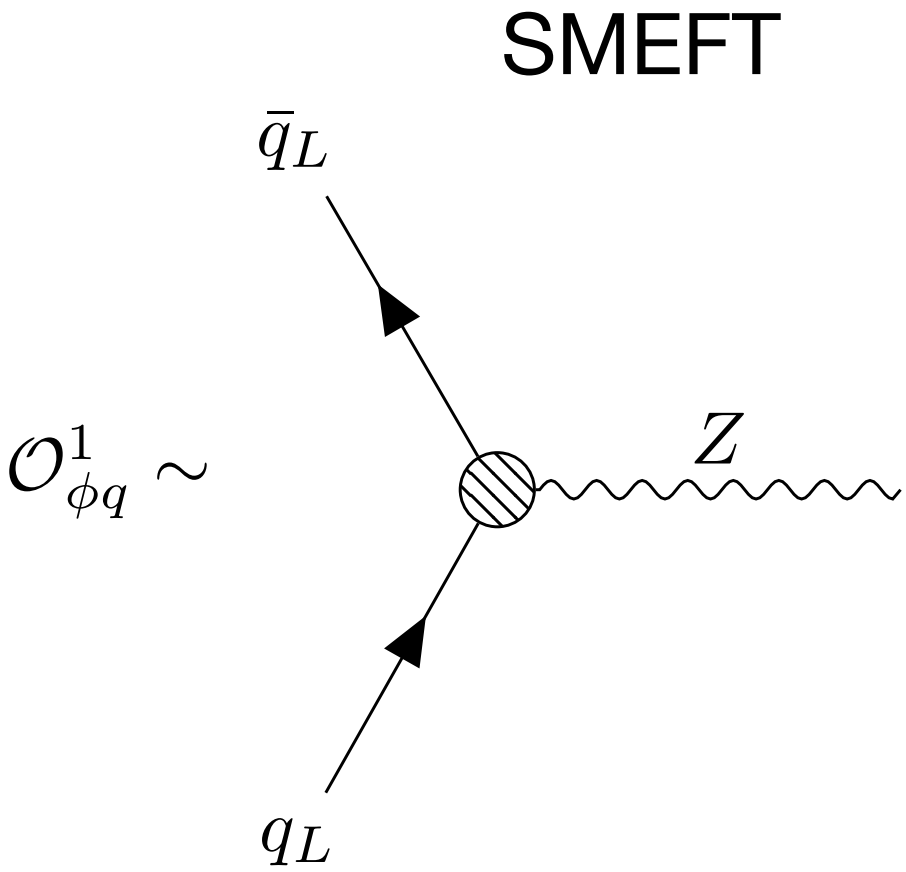
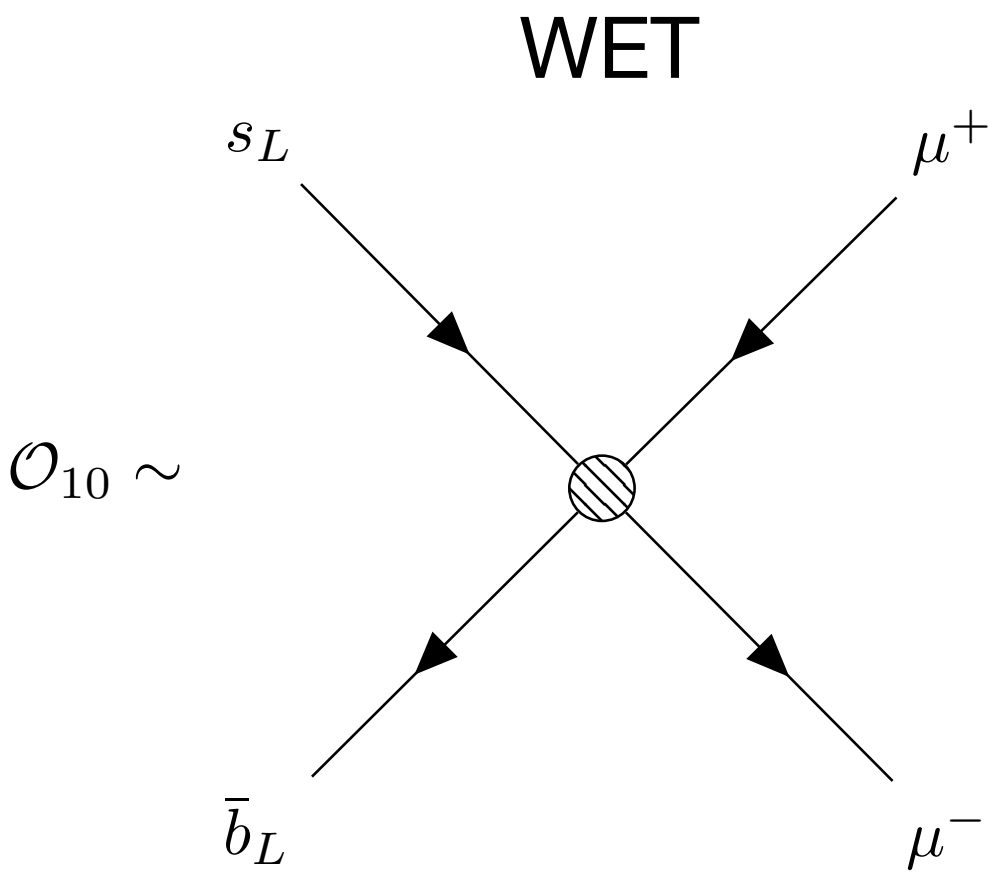


$$\sim [\mathcal{C}_{\phi q}^1]_{33} V_{ts}^* V_{tb} (\bar{t}_L \gamma_\mu t_L) Z^\mu$$

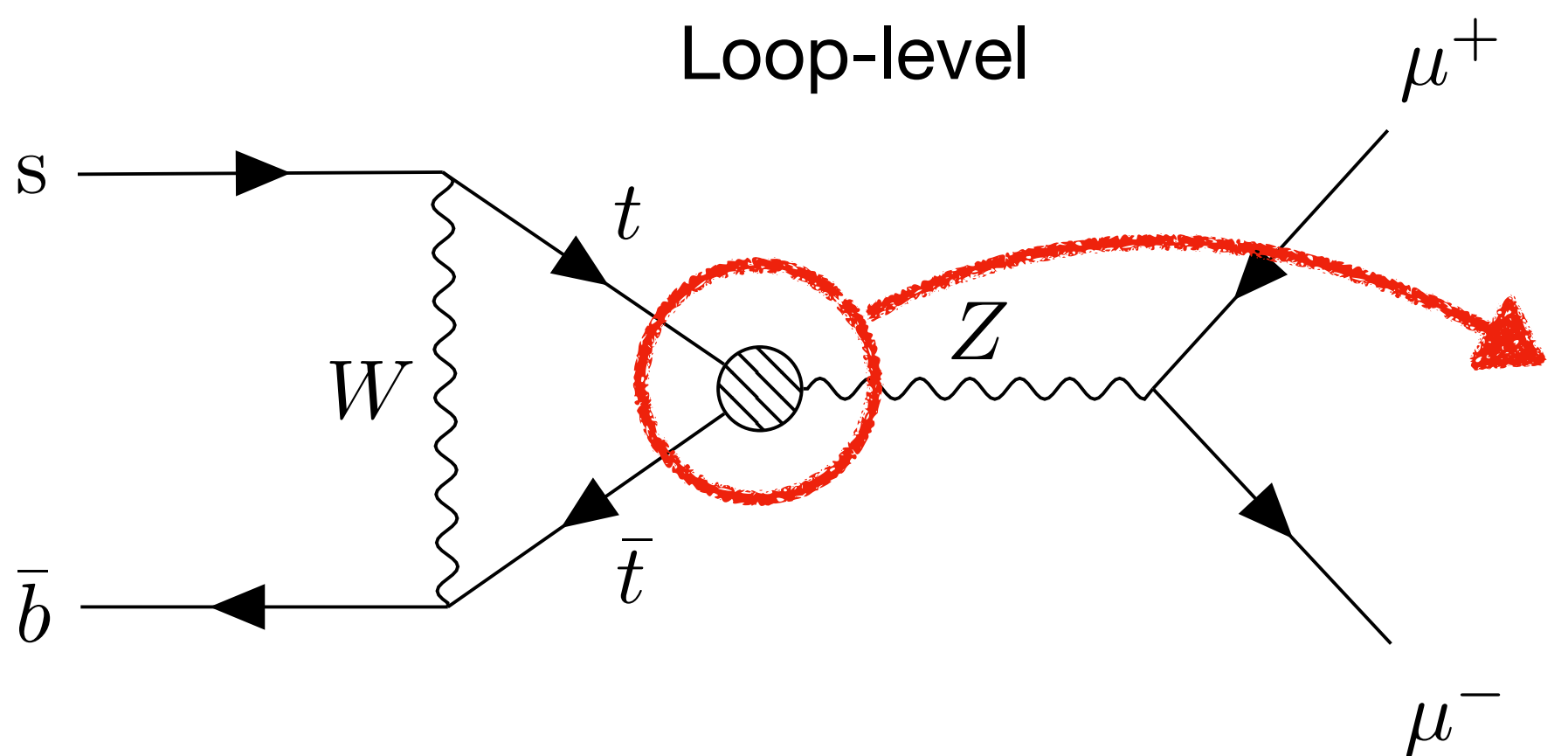
MFV	$\mathcal{C}_{\phi q}^1$
ii	a
33	$(a - y_t^2 b)y_t$

New Directions

Flavour (Work in Progress)



$$\sim [\mathcal{C}_{\phi q}^1]_{ij} V_{is}^* V_{jb} (\bar{s}_L \gamma_\mu b_L) Z^\mu$$



$$\sim [\mathcal{C}_{\phi q}^1]_{33} V_{ts}^* V_{tb} (\bar{t}_L \gamma_\mu t_L) Z^\mu$$

Matching & Running	$\mathcal{C}_{10}$
$a_{\phi q}^1$	0.1
$b_{\phi q}^1$	24.73

Conclusion

- Now is the time for Top-BSM
- Resolve degeneracies with different measurements
- Theorists can achieve a lot by lowering the uncertainties
- Still opportunities for new observables
- Interesting playground for connecting high- and low-energy physics

The Operators

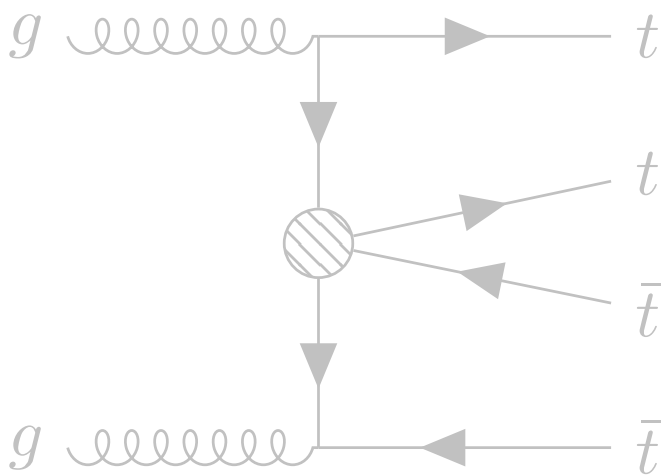
Examples
Processes

4 heavy fermions

$$C_{qu}^{8(3333)}(\bar{q}_3\gamma^\mu T^A q_3)(\bar{u}_3\gamma_\mu T^A u_3)$$

$t\bar{t}t\bar{t}$

$t\bar{t}b\bar{b}$



2 heavy - 2 light fermions

$$O_{qt8} = C_{qu}^{8(ii33)}(\bar{q}_i\gamma^\mu T^A q_i)(\bar{u}_3\gamma_\mu T^A u_3)$$

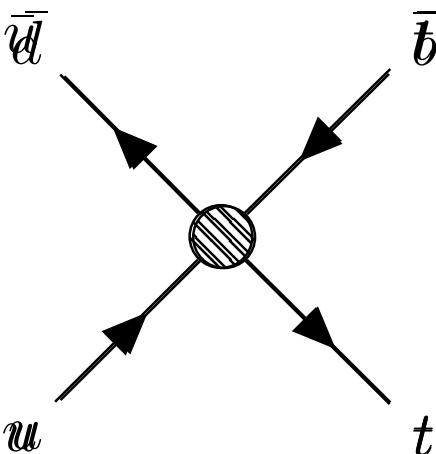
$$O_{qt1} = C_{qu}^{1(ii33)}(\bar{q}_i\gamma^\mu q_i)(\bar{u}_3\gamma_\mu u_3)$$

$t\bar{t}$ production

single t production

tZ

$t\bar{t}Z/W/H$



2 heavy fermions - boson

$$O_{tG} = \text{Re}\left\{C_{uG}^{(33)}\right\}(\bar{q}_3\sigma^{\mu\nu}T^A u_3)\tilde{\phi}G_{\mu\nu}^A$$

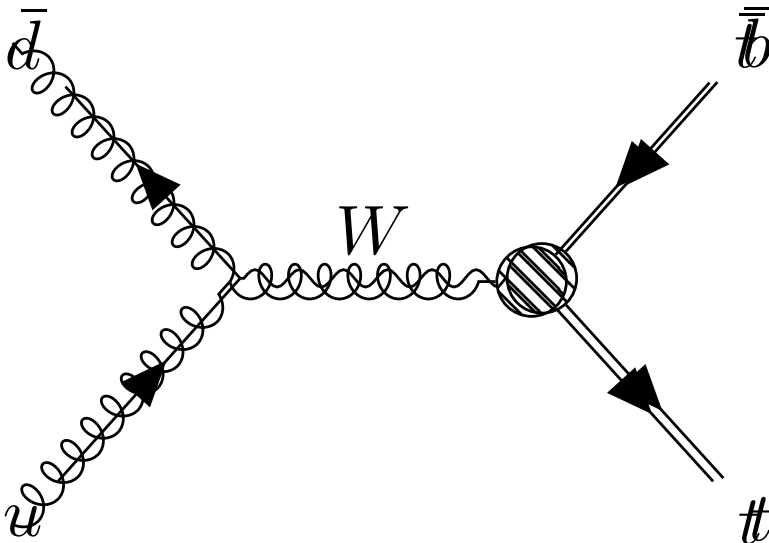
$$O_{tW} = \text{Re}\left\{C_{uW}^{(33)}\right\}(\bar{q}_3\sigma^{\mu\nu}\tau^I u_3)\tilde{\phi}W_{\mu\nu}^I$$

$t\bar{t}$ production

single t production

tZ/W

$t\bar{t}Z/W/H$



Event Kinematics

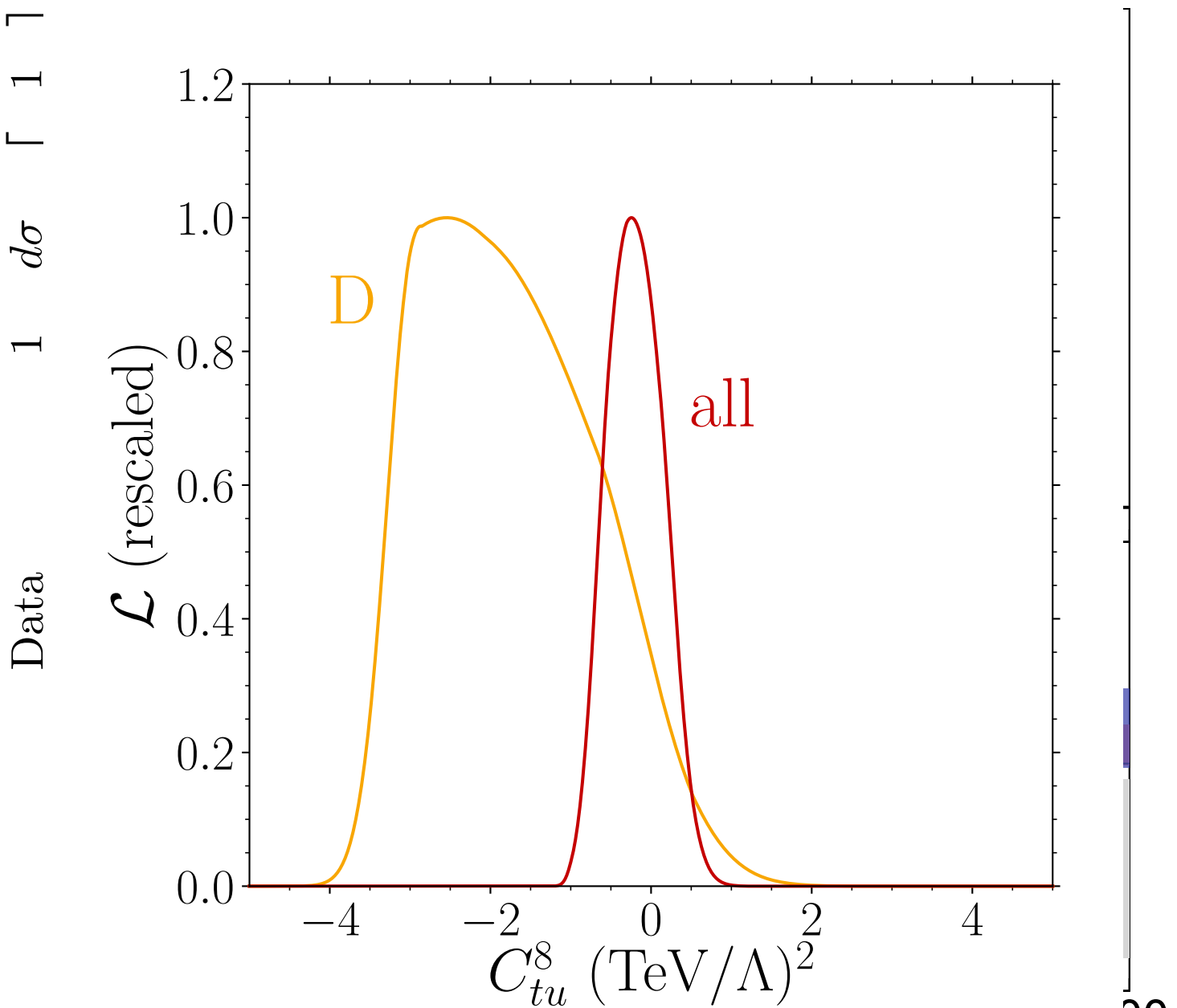
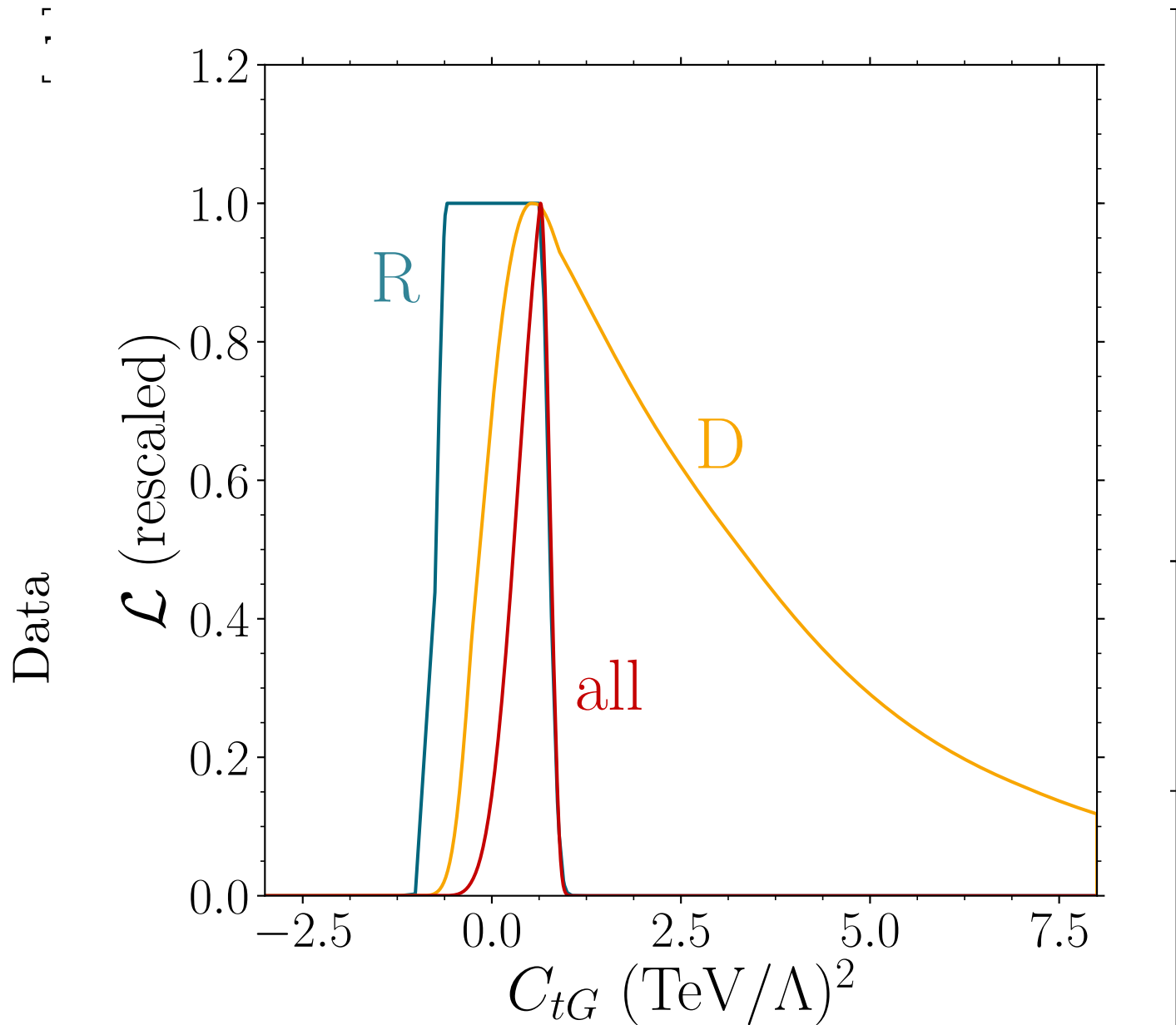
Or Rates vs. Tails

$$O_{tG} = C_{tG}(\bar{q}_3\sigma^{\mu\nu}T^A u_3)\tilde{\phi}G^A_{\mu\nu} \qquad O_{tu}^8 = C_{tu}^8(\bar{u}_3\gamma^\mu T^A u_3)(\bar{u}_i\gamma_\mu T^A u_i)$$

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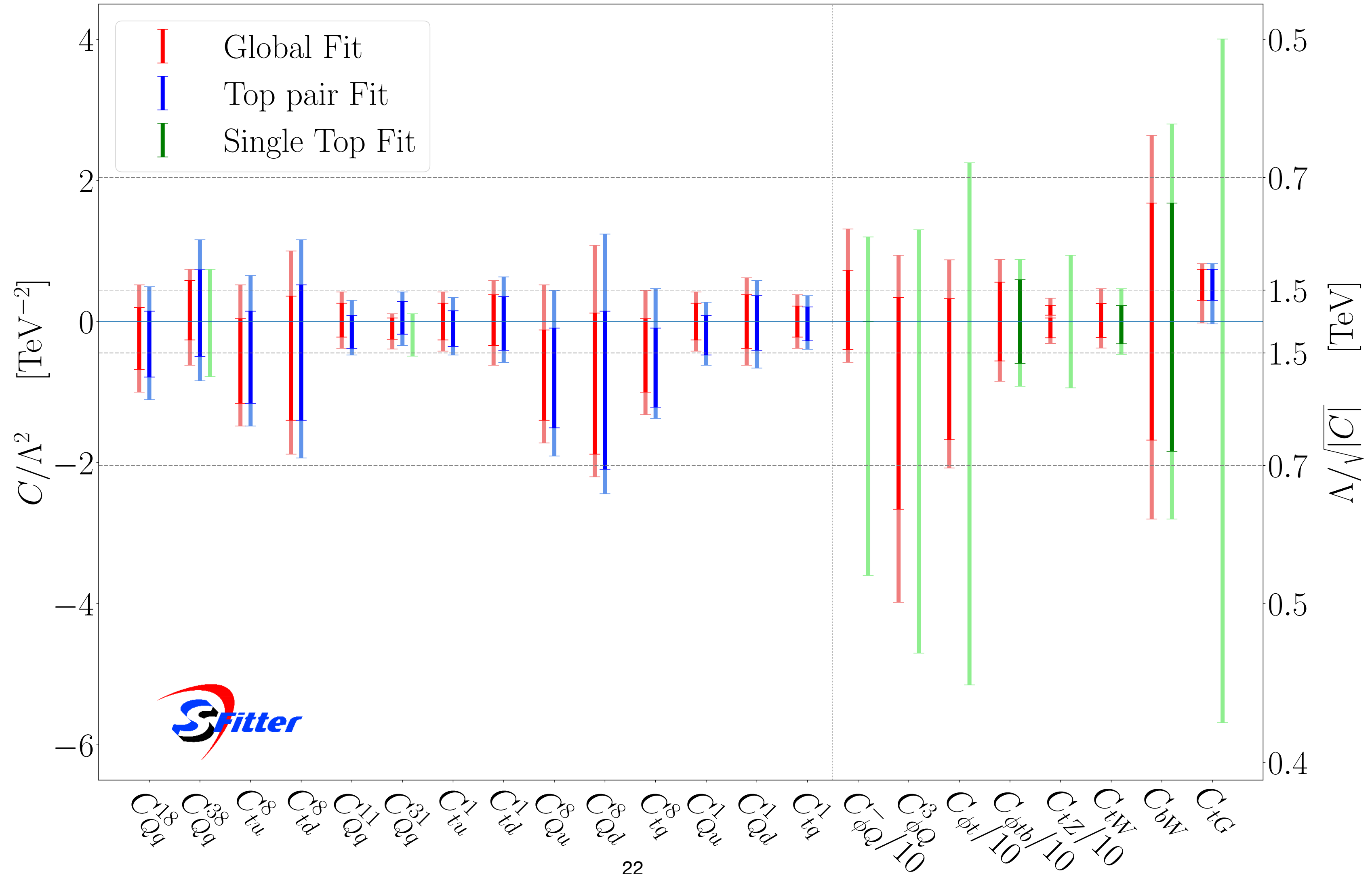
Very small

Very large
For high bins



Results

Run II, ATLAS+CMS, 68% and 95% C.L.



NLO Effects for chirality

